

Understanding Flavor at the LHC

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Thanks to my low-energy-physics collaborators:

Gudrun Hiller, YN

JHEP 0803 (2008) 046 [arXiv:0802.0916]

Gudrun Hiller, Yonit Hochberg, YN

JHEP 0903 (2009) 115 [arXiv:0812.0511]

Kfir Blum, Yuval Grossman, YN, Gilad Perez

Phys. Rev. Lett. 102 (2009) 211802 [arXiv:0903.2118]

Yuval Grossman, YN, Gilad Perez

arXiv:0904.0305

Oram Gedalia, Yuval Grossman, YN, Gilad Perez

arXiv:0906.1879

Helen Quinn + YN

‘The Mystery of the Missing Antimatter’ (PUP)

Thanks to my high- p_t -physics collaborators:

Yuval Grossman, YN, Jesse Thaler, Tomer Volansky, Jure Zupan
Phys. Rev. D76 (2007) 096006 [arXiv:0706.1845]

Jonathan Feng, Christopher Lester, YN, Yael Shadmi
Phys. Rev. D77 (2008) 076002 [arXiv:0712.0674]

Jonathan Feng, Sky French, Christopher Lester, YN, Yael Shadmi
arXiv:0906.4215

Feng, French, Galon, Lester, YN, Shadmi, Sanford, Yu
‘Identifying the SUSY theory of flavor at the LHC’

Why is flavor physics interesting?

- Flavor physics is sensitive to new physics at $\Lambda_{\text{NP}} \gg E_{\text{experiment}}$
FCNC suppressed within the SM by $\alpha_W^n, |V_{ij}|, m_f$
- The Standard Model flavor puzzle:
Why are the flavor parameters small and hierarchical?
(Why) are the neutrino flavor parameters different?
- The New Physics flavor puzzle:
If there is NP at the TeV scale, why are FCNC so small?
The solution \implies Clues for the subtle structure of the NP

Plan of Talk

1. Introduction
 - The LHC era
2. Open questions
 - The NP flavor puzzle
 - The SM flavor puzzle
3. What will we learn?
 - Flavor@ATLAS/CMS

The LHC Era

The LHC will explore the unknown

Energy $0.6 \rightarrow 4 \text{ TeV}$

Distance $10^{-19} \rightarrow 10^{-20} \text{ m}$

“Time” $10^{-11} \rightarrow 10^{-13} \text{ s}$

The LHC questions

- What is the mechanism of electroweak symmetry breaking?
- What separates the electroweak scale from the Planck scale?
- What are the dark matter particles?
- What happened at the electroweak phase transition (10^{-11} second after the big bang)?
- Was the baryon asymmetry generated by TeV scale physics?

Flavor at the LHC?

- The scale of flavor dynamics is unknown
- Very likely, it is well above the LHC direct reach

But...

- If new particles that couple to the SM fermions are discovered –
New flavor parameters can be measured
- Spectrum, flavor decomposition...
- New insights on flavor puzzles are likely

The NP Flavor Puzzle

What have we learned?

- The KM phase is different from zero (SM violates CP)
- The KM mechanism is the dominant source of the CP violation observed in meson decays
- Complete alternatives to the KM mechanism are excluded (Superweak, Approximate CP)
- No evidence for corrections to CKM
- NP contributions to the observed FCNC are at most comparable to the CKM contributions
- NP contributions are very small in $s \rightarrow d$, $c \rightarrow u$, $b \rightarrow d$, $b \rightarrow s$

The SM = Low energy effective theory

1. Gravity $\implies \Lambda_{\text{Planck}} \sim 10^{19} \text{ GeV}$
2. $m_\nu \neq 0 \implies \Lambda_{\text{Seesaw}} \leq 10^{15} \text{ GeV}$
3. m_H^2 -fine tuning; Dark matter $\implies \Lambda_{\text{NP}} \sim \text{TeV}$



- The SM = Low energy effective theory
- Must write non-renormalizable terms suppressed by $\Lambda_{\text{NP}}^{d-4}$
- $\mathcal{L}_{d=5} = \frac{y_{ij}^\nu}{\Lambda_{\text{seesaw}}} L_i L_j \phi \phi$
- $\mathcal{L}_{d=6}$ contains many flavor changing operators

New Physics

- The effects of new physics at a high energy scale Λ_{NP} can be presented as higher dimension operators

- For example, we expect the following dimension-six operators:

$$\frac{z_{sd}}{\Lambda_{\text{NP}}^2} (\overline{d_L} \gamma_\mu s_L)^2 + \frac{z_{cu}}{\Lambda_{\text{NP}}^2} (\overline{c_L} \gamma_\mu u_L)^2 + \frac{z_{bd}}{\Lambda_{\text{NP}}^2} (\overline{d_L} \gamma_\mu b_L)^2 + \frac{z_{bs}}{\Lambda_{\text{NP}}^2} (\overline{s_L} \gamma_\mu b_L)^2$$

- New contribution to neutral meson mixing, *e.g.*

$$\frac{\Delta m_B}{m_B} \sim \frac{f_B^2}{3} \times \frac{|z_{bd}|}{\Lambda_{\text{NP}}^2}$$

- Generic flavor structure $\equiv z_{ij} \sim 1$ or, perhaps, loop – factor

Some data

$\Delta m_K/m_K$	7.0×10^{-15}
$\Delta m_D/m_D$	8.7×10^{-15}
$\Delta m_B/m_B$	6.3×10^{-14}
$\Delta m_{B_s}/m_{B_s}$	2.1×10^{-12}
ϵ_K	2.3×10^{-3}
A_Γ/y_{CP}	≤ 0.2
$S_{\psi K_S}$	0.67 ± 0.02
$S_{\psi\phi}$	≤ 1

High Scale?

- For $z_{ij} \sim 1$ (and $\mathcal{I}m(z_{ij}) \sim 1$), $\Lambda_{\text{NP}} \gtrsim \frac{10^{-4}}{\sqrt{\Delta m/m}} \text{ TeV}$

Mixing	$\Lambda_{\text{NP}}^{\text{CPC}} \gtrsim$	$\Lambda_{\text{NP}}^{\text{CPV}} \gtrsim$
$K - \bar{K}$	1000 TeV	20000 TeV
$D - \bar{D}$	1000 TeV	3000 TeV
$B - \bar{B}$	400 TeV	800 TeV
$B_s - \bar{B}_s$	70 TeV	70 TeV

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Did we misinterpret the Higgs fine tuning problem?

Did we misinterpret the dark matter puzzle?

Small (hierachical?) flavor parameters?

- For $\Lambda_{\text{NP}} \sim 1 \text{ TeV}$, $z_{ij} \lesssim 10^8 (\Delta m_{ij}/m)$

Mixing	$ z_{ij} \lesssim$	$\text{Im}(z_{ij}) \lesssim$
$K - \bar{K}$	8×10^{-7}	6×10^{-9}
$D - \bar{D}$	5×10^{-7}	1×10^{-7}
$B - \bar{B}$	5×10^{-6}	1×10^{-6}
$B_s - \bar{B}_s$	2×10^{-4}	2×10^{-4}

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The flavor structure of NP@TeV must be highly non-generic

How? Why? = The NP flavor puzzle

The SM Flavor Puzzle

What have we learned?

- $\Delta m_{21}^2 = (7.9 \pm 0.3) \times 10^{-5} \text{ eV}^2$, $|\Delta m_{32}^2| = (2.6 \pm 0.2) \times 10^{-3} \text{ eV}^2$
- $\sin^2 \theta_{12} = 0.31 \pm 0.02$, $\sin^2 \theta_{23} = 0.47 \pm 0.07$, $\sin^2 \theta_{13} = 0_{-0.0}^{+0.08}$
- Neutrino-flavor is different

Smallness and Hierarchy

$$\begin{aligned} Y_t &\sim 1, & Y_c &\sim 10^{-2}, & Y_u &\sim 10^{-5} \\ Y_b &\sim 10^{-2}, & Y_s &\sim 10^{-3}, & Y_d &\sim 10^{-4} \\ Y_\tau &\sim 10^{-2}, & Y_\mu &\sim 10^{-3}, & Y_e &\sim 10^{-6} \\ |V_{us}| &\sim 0.2, & |V_{cb}| &\sim 0.04, & |V_{ub}| &\sim 0.004, & \delta_{\text{KM}} &\sim 1 \end{aligned}$$

- For comparison: $g_s \sim 1$, $g \sim 0.6$, $g' \sim 0.3$, $\lambda \sim 1$
- The SM flavor parameters have structure:
smallness and hierarchy
- Why? = The SM flavor puzzle

The Froggatt-Nielsen (FN) mechanism

- Approximate “horizontal” symmetry (e.g. $U(1)_H$)
- Small breaking parameter $\epsilon = \langle S_{-1} \rangle / \Lambda \ll 1$
- $\mathbf{10}(2, 1, 0)$, $\bar{\mathbf{5}}(0, 0, 0)$



$$Y_t : Y_c : Y_u \sim 1 : \epsilon^2 : \epsilon^4$$

$$Y_b : Y_s : Y_d \sim 1 : \epsilon : \epsilon^2$$

$$Y_\tau : Y_\mu : Y_e \sim 1 : \epsilon : \epsilon^2$$

$$|V_{us}| \sim |V_{cb}| \sim \epsilon, \quad |V_{ub}| \sim \epsilon^2, \quad \delta_{\text{KM}} \sim 1$$

+

$$m_3 : m_2 : m_1 \sim 1 : 1 : 1$$

$$|U_{e2}| \sim 1, \quad |U_{\mu 3}| \sim 1, \quad |U_{e3}| \sim 1$$

Testing FN with Neutrinos

- The data:
 - $\Delta m_{21}^2 = (7.9 \pm 0.3) \times 10^{-5} \text{ eV}^2$, $|\Delta m_{32}^2| = (2.6 \pm 0.2) \times 10^{-3} \text{ eV}^2$
 - $\sin^2 \theta_{12} = 0.31 \pm 0.02$, $\sin^2 \theta_{23} = 0.47 \pm 0.07$, $\sin^2 \theta_{13} = 0_{-0.0}^{+0.08}$
- The tests:
 - $s_{23} \sim 1$, $m_2/m_3 \sim \epsilon^x$?
Inconsistent with FN
 - $s_{23} \sim 1$, $s_{12} \sim 1$, $s_{13} \sim \epsilon^x$?
Inconsistent with FN
 - $\sin^2 2\theta_{23} = 1 - \epsilon^x$?
Inconsistent with FN

Neutrino Mass Anarchy

- Facts:

- $\sin \theta_{23} \sim 0.70 > \text{any } |V_{ij}|$
- $\sin \theta_{12} \sim 0.56 > \text{any } |V_{ij}|$
- $m_2/m_3 \gtrsim 1/6 > \text{any } m_i/m_j$ for charged fermions
- $\sin \theta_{13} \sim 0.1$ is still possible

- Possible interpretation:

- Neutrino parameters are all of $O(1)$ (no structure):
Neutrino mass anarchy
- Consistent with FN
- Close to GUT+FN predictions:

$$s_{23} \sim \frac{m_s/m_b}{|V_{cb}|} \sim 1; \quad s_{12} \sim \frac{m_d/m_s}{|V_{us}|} \sim 0.2; \quad s_{13} \sim \frac{m_d/m_b}{|V_{ub}|} \sim 0.5$$

Structure is in the eye of the beholder

$$|U|_{3\sigma} = \begin{pmatrix} 0.79 - 0.86 & 0.50 - 0.61 & 0.0 - 0.2 \\ 0.25 - 0.53 & 0.47 - 0.73 & 0.56 - 0.79 \\ 0.21 - 0.51 & 0.42 - 0.69 & 0.61 - 0.83 \end{pmatrix}$$

- Tribimaximal-ists:

$$|U|_{\text{TBM}} = \begin{pmatrix} \sqrt{2/3} & \sqrt{1/3} & 0 \\ \sqrt{1/6} & \sqrt{1/3} & \sqrt{1/2} \\ \sqrt{1/6} & \sqrt{1/3} & \sqrt{1/2} \end{pmatrix}$$

- Anarch-ists:

$$|U|_{\text{anarchy}} = \begin{pmatrix} \mathcal{O}(0.6) & \mathcal{O}(0.6) & \mathcal{O}(0.6) \\ \mathcal{O}(0.6) & \mathcal{O}(0.6) & \mathcal{O}(0.6) \\ \mathcal{O}(0.6) & \mathcal{O}(0.6) & \mathcal{O}(0.6) \end{pmatrix}$$

What will we learn?

Flavor Physics at the LHC era

- If ATLAS/CMS observe no NP...
- and flavor factories observe no NP...

Flavor Physics at the LHC era

- If ATLAS/CMS observe no NP...
- but flavor factories observe NP...
- We may have misinterpreted the fine-tuning problem
- We may have misinterpreted the dark matter puzzle
- Flavor will provide the only clue for an accessible scale of NP

Flavor Physics at the LHC era

ATLAS/CMS will, hopefully, observe NP at $\Lambda_{\text{NP}} \lesssim TeV$;

In combination with flavor factories, we may...

- Understand how the NP flavor puzzle is (not) solved
 \implies Probe NP at $\Lambda_{\text{NP}} \gg TeV$
- Get hints about the solution to the SM flavor puzzle

Gauge+Gravity Mediation

- Example: High (but not too high) scale gauge mediation
 - Gravity mediation sub-dominant but non-negligible
 - $r = \frac{\text{gravity-med}}{\text{gauge-med}} \sim \left(\frac{m_M}{m_{\text{Pl}}}\right)^2 \left(\frac{4\pi}{\alpha_3(m_M)}\right)^2 \frac{3}{8n_M}$
 - $\widetilde{M}_{\tilde{Q}_L}^2(m_M) = \tilde{m}_{\tilde{Q}_L}^2 (\mathbf{1} + r X_{\tilde{Q}_L})$
 - Degeneracy depends on r

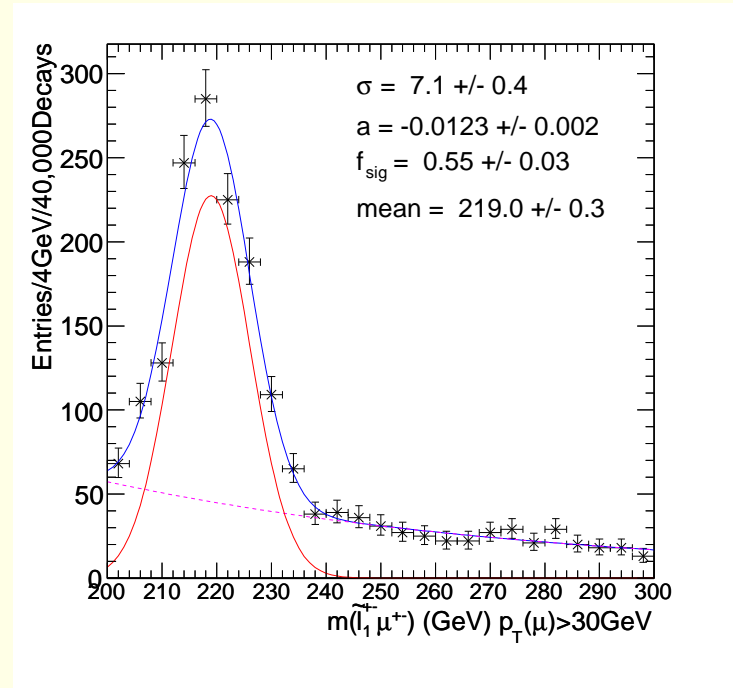
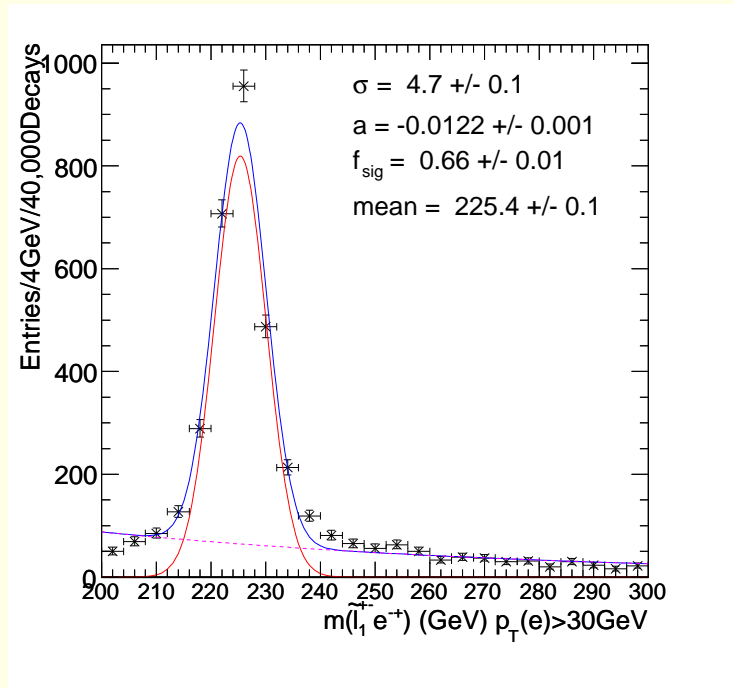
Assume: The flavor structure of X determined by FN:

- $X_{\tilde{Q}_L} \sim \begin{pmatrix} 1 & V_{us} & V_{ub} \\ \cdot & 1 & V_{cb} \\ \cdot & \cdot & 1 \end{pmatrix}; \quad X_{\tilde{D}_R} \sim \begin{pmatrix} 1 & \frac{m_d/m_s}{V_{us}} & \frac{m_d/m_b}{V_{ub}} \\ \cdot & 1 & \frac{m_s/m_b}{V_{cb}} \\ \cdot & \cdot & 1 \end{pmatrix}$

- Mixing depends only on X which is related to the SM flavor

Measuring small mass splitting + mixing

$$\chi_1^0 \rightarrow \tilde{\ell}_1^\pm e^\mp \text{ or } \tilde{\ell}_2^\pm \mu^\mp; \quad \tilde{\ell}_2^\pm \rightarrow \tilde{\ell}_1^\mp (\ell^\pm \ell^\pm)_{\text{soft}} \text{ or } \tilde{\ell}_2^\pm \rightarrow \tilde{\ell}_1^\pm (\ell^\pm \ell^\mp)_{\text{soft}}$$



$$m(\tilde{\ell}_1^\pm e^\mp) = m_{\chi_1^0}$$

$$m(\tilde{\ell}_1^\pm \mu^\pm) = m_{\chi_1^0} - E_{\text{shift}}$$

$$\Delta m = \frac{2m_{\chi_1^0} m_{\tilde{\ell}}}{m_{\chi_1^0}^2 + m_{\tilde{\ell}}^2} E_{\text{shift}}$$

Solving the NP Flavor Puzzle

If ATLAS/CMS observe squarks and sleptons...

- Determine the sfermion mass scale (\tilde{m})
- Determine the sfermion mass splitting ($m_{\tilde{f}_j} - m_{\tilde{f}_i}$)
- Determine the sfermion flavor decomposition (K_{ij})



Learn how the SUSY flavor suppression is obtained

What will we learn?

Physics at $\Lambda_{\text{NP}} \gg \Lambda_{\text{LHC}}$

If ATLAS/CMS determine sfermion mass splittings...

- Find the ratio between gravity- and gauge-mediated contributions (r)
- Determine the messenger scale of gauge mediation (m_M)
- Find the hierarchy between the GMSB and see-saw scales



Probe physics at $m_M \sim 10^{15} \text{ GeV}$

Solving the SM Flavor Puzzle?

If ATLAS/CMS determine sfermion flavor decomposition...

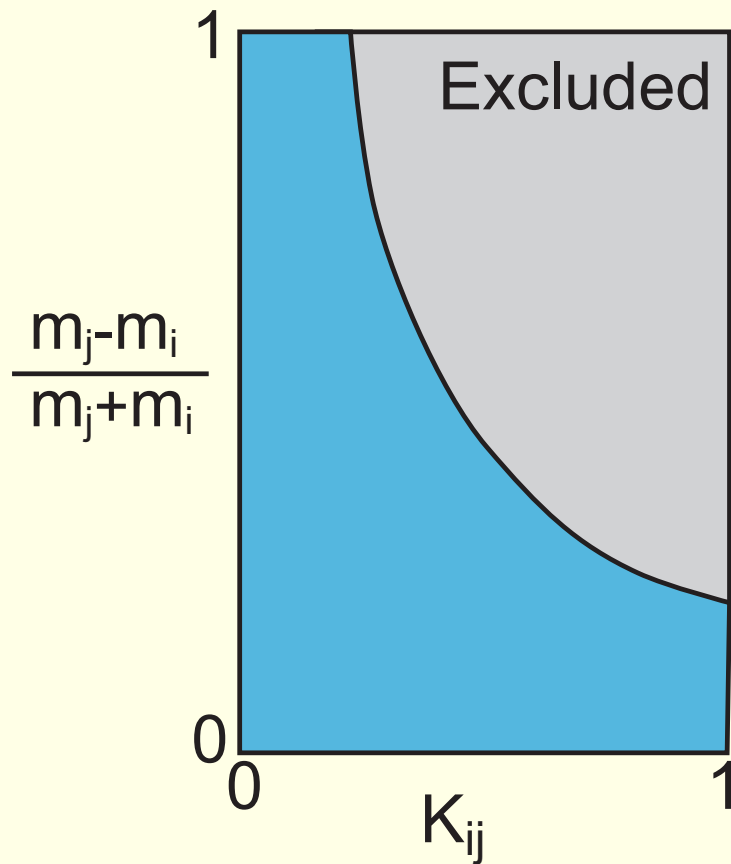
- Determine X of $\tilde{M}^2 = \tilde{m}^2(\mathbf{1} + rX)$
- Does X have the FN-predicted structure?



Test theories that explain the SM flavor structure

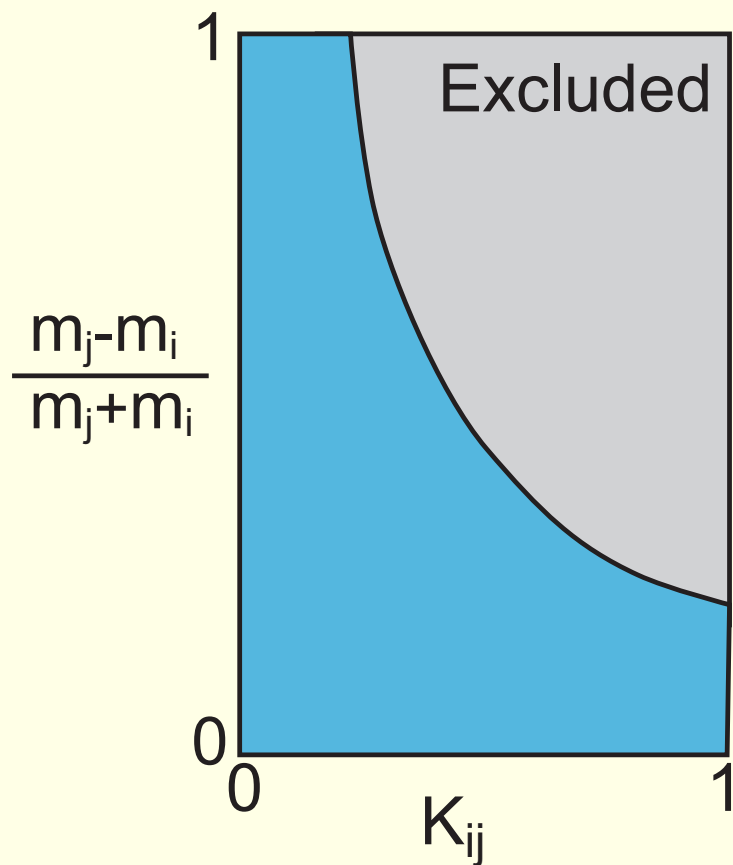
What will we learn?

The NP flavor plane

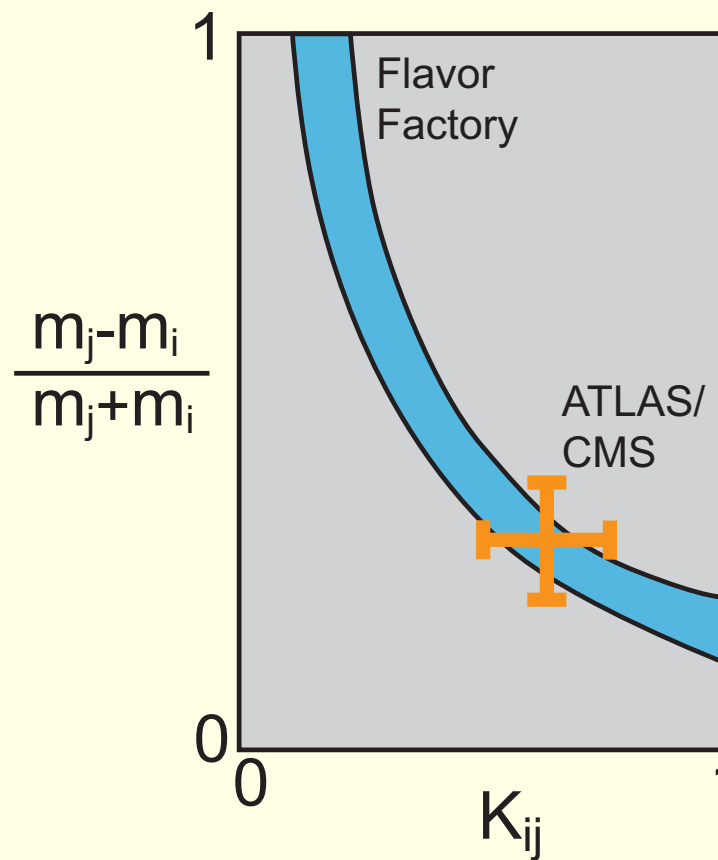


Flavor Factories

The NP flavor plane



Flavor Factories



FF+ATLAS/CMS

What will we learn?

- ATLAS/CMS and flavor factories give complementary information
- In the absence of NP at ATLAS/CMS – flavor factories will be crucial to find Λ_{NP}
- With NP at ATLAS/CMS –
The NP flavor puzzle is likely to be understood
 \implies A probe of physics at $\Lambda_{\text{NP}} \gg \Lambda_{\text{LHC}}$
- With supersymmetry –
The SM flavor puzzle may be solved

Backup Transparencies

A brief history of FV

- $\Gamma(K \rightarrow \mu\mu) \ll \Gamma(K \rightarrow \mu\nu) \implies \text{Charm}$ [GIM, 1970]
- $\Delta m_K \implies m_c \sim 1.5 \text{ GeV}$ [Gaillard-Lee, 1974]
- $\varepsilon_K \neq 0 \implies \text{Third generation}$ [KM, 1973]
- $\Delta m_B \implies m_t \gg m_W$ [Various, 1986]

Why is CPV interesting?

- Within the SM, a single CP violating parameter η :
In addition, QCD = CP invariant (θ_{QCD} irrelevant)
Strong predictive power (correlations + zeros)
Excellent tests of the flavor sector
- η cannot explain the baryon asymmetry – a puzzle:
There must exist new sources of CPV
Electroweak baryogenesis? (Testable at the LHC)
Leptogenesis? (Window to Λ_{seesaw})

A brief history of CPV

- 1964 – 2000

- $|\varepsilon| = (2.284 \pm 0.014) \times 10^{-3}$; $\mathcal{R}e(\varepsilon'/\varepsilon) = (1.67 \pm 0.26) \times 10^{-3}$

A brief history of CPV

- 1964 – 2000

- $|\varepsilon| = (2.284 \pm 0.014) \times 10^{-3}$; $\mathcal{R}e(\varepsilon'/\varepsilon) = (1.67 \pm 0.26) \times 10^{-3}$

- 2000 – 2009

- $S_{\psi K_S} = +0.67 \pm 0.02$

- $S_{\eta' K_S} = +0.61 \pm 0.07$, $S_{\pi^0 K_S} = +0.57 \pm 0.17$,
 $S_{\rho^0 K_S} = +0.63 \pm 0.17$, $S_{f_0 K_S} = +0.62 \pm 0.11$

- $S_{K^+ K^- K_S} = -0.74 \pm 0.11$

- $S_{\pi^+ \pi^-} = -0.61 \pm 0.08$, $C_{\pi^+ \pi^-} = -0.38 \pm 0.06$

- $S_{\psi \pi^0} = -0.93 \pm 0.15$, $S_{D^+ D^-} = -0.89 \pm 0.26$

- $\mathcal{A}_{K^\mp \rho^0} = +0.37 \pm 0.11$, $\mathcal{A}_{\eta K^\mp} = -0.27 \pm 0.09$, $\mathcal{A}_{f_2 K^\mp} = -0.68 \pm 0.20$

- $\mathcal{A}_{K^\mp \pi^\pm} = -0.098 \pm 0.012$, $\mathcal{A}_{\eta K^{*0}} = +0.19 \pm 0.05$

- ...

Flavor@GeV \implies NP@TeV

A recent example [Blum et al, 0903.2118, PRL in press]

- $\frac{\Delta m_K}{m_K} = (7.01 \pm 0.01) \times 10^{-15}; \quad \epsilon_K = (2.23 \pm 0.01) \times 10^{-3}$
- $\frac{\Delta m_D}{m_D} = (8.6 \pm 2.1) \times 10^{-15}; \quad A_\Gamma = (1.2 \pm 2.5) \times 10^{-3}$
- Consider $\frac{1}{\text{TeV}^2} [\overline{Q_{Li}}(X_Q)_{ij}\gamma_\mu Q_{Lj}]^2$
- Take $Y_d = \lambda_d, \quad Y_u = V^\dagger \lambda_u, \quad X_Q = V_d^\dagger \text{diag}(\lambda_1, \lambda_2)V_d$
- $K + D \implies$ Degeneracy: $\lambda_2 - \lambda_1 \leq 0.004 - 0.0005$
 - Supersymmetry: $\frac{m_{\tilde{Q}_2} - m_{\tilde{Q}_1}}{m_{\tilde{Q}_2} + m_{\tilde{Q}_1}} \leq 0.27 - 0.034$
 - RS-I: $\sqrt{\frac{\text{TeV}}{m_{\text{KK}}}} f_{Q_2} \lesssim 0.06 - 0.02.$

A beautiful relation

- Assume no direct CP violation

- A surprising relation: $y \tan \phi = x(1 - |q/p|)$

Grossman et al., arXiv:0904.0305



$$K \quad \arg(\epsilon) = \arctan(-x/y)$$

$$B_s \quad A_{\text{SL}}^s = -2|y/x|S_{\psi\phi}/(1 - S_{\psi\phi}^2)^{1/2}$$

$$D \quad (1 - |q/p|)/\tan \phi = y/x$$

Flavor Violation (FV)

- $\mathcal{L}_{\text{kinetic+gauge}} + \mathcal{L}_{\text{Higgs}}$ has a large global symmetry:
 $G_{\text{global}} = [U(3)]^5$
- $\mathcal{L}_{\text{Yukawa}} = \overline{Q}_{Li} Y_{ij}^u \tilde{\phi} U_{Rj} + \overline{Q}_{Li} Y_{ij}^d \phi D_{Rj} + \overline{L}_{Li} Y_{ij}^e \phi E_{Rj}$
 breaks $G_{\text{global}} \rightarrow U(1)_B \times U(1)_e \times U(1)_\mu \times U(1)_\tau$
- Flavor physics:
 interactions that break the $[SU(3)]^5$ symmetry



- $Q_L \rightarrow V_Q Q_L, \quad U_R \rightarrow V_U U_R, \quad D_R \rightarrow V_D D_R$
 = Change of interaction basis
- Can be used to reduce the number of parameters in Y^u, Y^d

Kobayashi and Maskawa

The number of real and imaginary quark flavor parameters:

- With two generations:

$$2 \times (4_R + 4_I) - 3 \times (1_R + 3_I) + 1_I = 5_R + 0_I$$

- With three generations:

$$2 \times (9_R + 9_I) - 3 \times (3_R + 6_I) + 1_I = 9_R + 1_I$$

- The two generation SM is CP conserving

The three generation SM is CP violating

CP violation = a single imaginary parameter in the CKM matrix:

- V unitary with 3 real (λ, A, ρ) and 1 imaginary (η) parameters:

$$V \simeq \begin{pmatrix} 1 & \lambda & A\lambda^3(\rho + i\eta) \\ -\lambda & 1 & A\lambda^2 \\ A\lambda^3(1 - \rho + i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

What have we learned?

$S_{\psi K_S}$

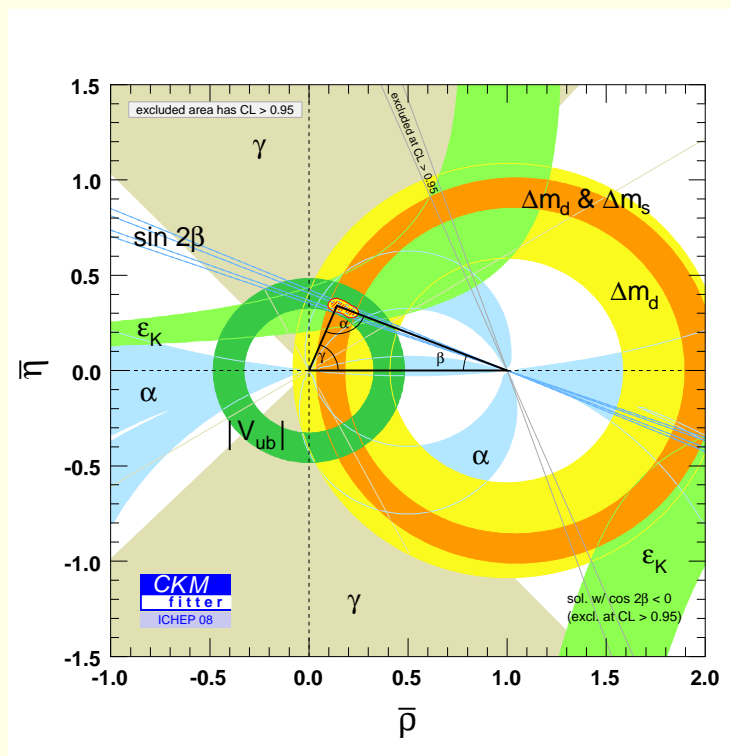
- Babar/Belle: $A_{\psi K_S}(t) = \frac{\frac{d\Gamma}{dt}[\overline{B^0_{\text{phys}}}(t) \rightarrow \psi K_S] - \frac{d\Gamma}{dt}[B^0_{\text{phys}}(t) \rightarrow \psi K_S]}{\frac{d\Gamma}{dt}[\overline{B^0_{\text{phys}}}(t) \rightarrow \psi K_S] + \frac{d\Gamma}{dt}[B^0_{\text{phys}}(t) \rightarrow \psi K_S]}$
- Theory: $A_{\psi K_S}(t)$ dominated by interference between $A(B^0 \rightarrow \psi K_S)$ and $A(B^0 \rightarrow \overline{B^0} \rightarrow \psi K_S)$
- $\implies A_{\psi K_S}(t) = S_{\psi K_S} \sin(\Delta m_B t)$
 $\implies S_{\psi K_S} = \frac{1}{|A(B^0 \rightarrow \overline{B^0})|} \mathcal{I}m \left[\frac{A(B^0 \rightarrow \overline{B^0}) A(\overline{B^0} \rightarrow \psi K_S)}{A(B^0 \rightarrow \psi K_S)} \right]$
- SM: $S_{\psi K_S} = \mathcal{I}m \left[\frac{V_{tb}^* V_{td}}{V_{tb} V_{td}^*} \frac{V_{cb} V_{cd}^*}{V_{cb}^* V_{cd}} \right] = \frac{2\eta(1-\rho)}{\eta^2 + (1-\rho)^2}$
- The approximations involved are better than one percent!
- Experiments: $S_{\psi K_S} = 0.671 \pm 0.024$

Testing CKM – Take I

- Assume: CKM matrix is the only source of FV and CPV
- λ known from $K \rightarrow \pi l \nu$
 A known from $b \rightarrow c l \nu$
- Many observables are $f(\rho, \eta)$:
 - $b \rightarrow u l \nu \implies \propto |V_{ub}/V_{cb}|^2 \propto \rho^2 + \eta^2$
 - $\Delta m_{B_d}/\Delta m_{B_s} \implies \propto |V_{td}/V_{ts}|^2 \propto (1 - \rho)^2 + \eta^2$
 - $S_{\psi K_S} \implies \frac{2\eta(1-\rho)}{(1-\rho)^2 + \eta^2}$
 - $S_{\rho\rho}(\alpha)$
 - $\mathcal{A}_{DK}(\gamma)$
 - ϵ_K

What have we learned?

The flavor-factories plot



CKMFitter

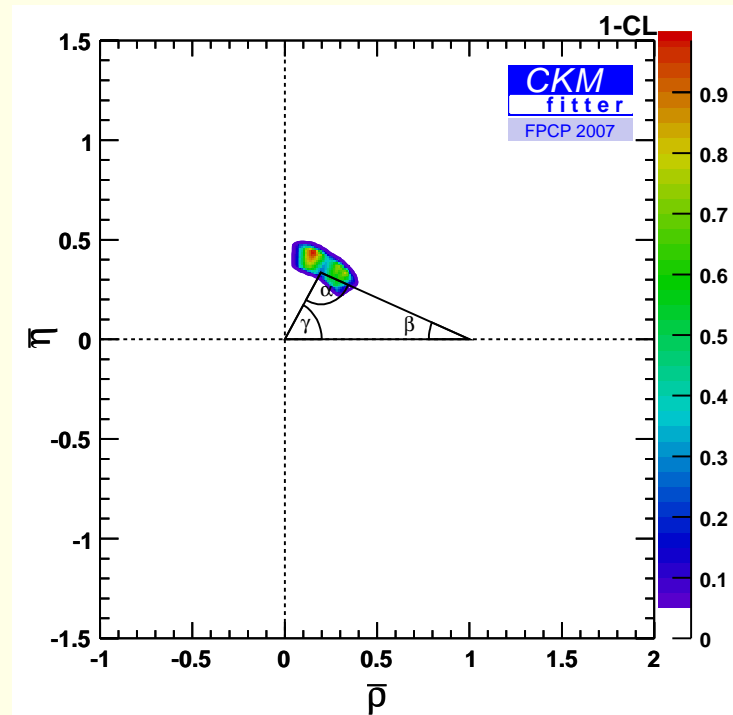
Very likely, the CKM mechanism dominates FV and CPV

Testing CKM - take II

- Assume: New Physics in leading tree decays - negligible
- Allow arbitrary new physics in loop processes
- Use only tree decays and $B^0 - \bar{B}^0$ mixing
- Define $h_d e^{2i\sigma_d} = \frac{A^{\text{NP}}(B^0 \rightarrow \bar{B})}{A^{\text{SM}}(B^0 \rightarrow \bar{B})}$
- Use $|V_{ub}/V_{cb}|$, \mathcal{A}_{DK} , $S_{\psi K}$, $S_{\rho\rho}$, Δm_{B_d} , $\mathcal{A}_{\text{SL}}^d$
- Fit to $\boxed{\eta}$, ρ , $\boxed{h_d}$, σ_d
- Find whether $\eta = 0$ is allowed
If not \implies The KM mechanism is at work
- Find whether $h_d \gg 1$ is allowed
If not \implies The KM mechanism is dominant

What have we learned?

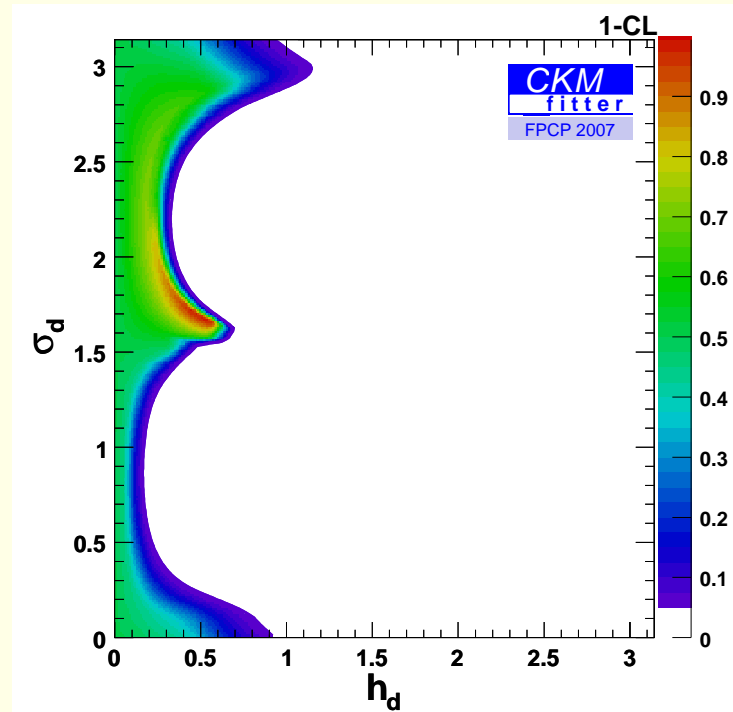
$\eta \neq 0$?



- The KM mechanism is at work

What have we learned?

$$\underline{h_d \ll 1?}$$



- The KM mechanism dominates CP violation
- The CKM mechanism is a major player in flavor violation

Minimal flavor violation (MFV)

- MFV = the only source of FV are the SM Yukawa matrices
- MFV \implies NP@TeV scale is consistent with FCNC constraints
- Most likely, an approximation
- Predictions:
 - Spectrum: often MFV implies degeneracies
 - Mixing: the third generation is approximately decoupled
- Example: Gauge mediated supersymmetry breaking
 - Squark spectrum: $2 + 1$
 - Squark decays: $\tilde{q}_{1,2} \rightarrow q_{1,2}, \quad \tilde{q}_3 \rightarrow q_3$
- In principle, testable in ATLAS/CMS

The FN mechanism: Predictions (quarks)

- In the quark sector: 8 FN charges, 9 observables
- One prediction that is independent of charge assignments:

$$|V_{ub}| \sim |V_{us}V_{cb}|$$

Experimentally correct to within a factor of 2

- In addition, six inequalities:

$$|V_{us}| \gtrsim \frac{m_d}{m_s}, \frac{m_u}{m_c}; \quad |V_{ub}| \gtrsim \frac{m_d}{m_b}, \frac{m_u}{m_t}; \quad |V_{cb}| \gtrsim \frac{m_s}{m_b}, \frac{m_c}{m_t}$$

Experimentally fulfilled

- When ordering the quarks by mass:

$$V_{CKM} \sim \mathbf{1} \text{ (diagonal terms not suppressed parameterically)}$$

Experimentally fulfilled

The FN mechanism: Predictions (leptons)

- In the lepton sector: 5 FN charges, 9 observables
- Four predictions that are independent of charge assignments:

$$m_{\nu_i}/m_{\nu_j} \sim |U_{ij}|^2$$

$$|U_{e3}| \sim |U_{e2}U_{\mu 3}|$$

- In addition, three inequalities:

$$|U_{e2}| \gtrsim \frac{m_e}{m_\mu}; \quad |U_{e3}| \gtrsim \frac{m_e}{m_\tau}; \quad |U_{\mu 3}| \gtrsim \frac{m_\mu}{m_\tau}$$

- When ordering the leptons by mass:

$$U \sim \mathbf{1}$$