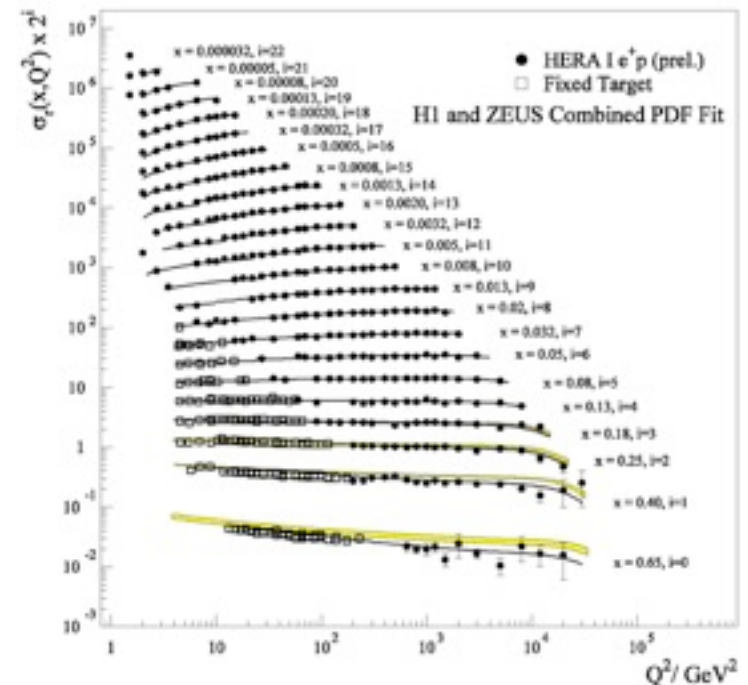
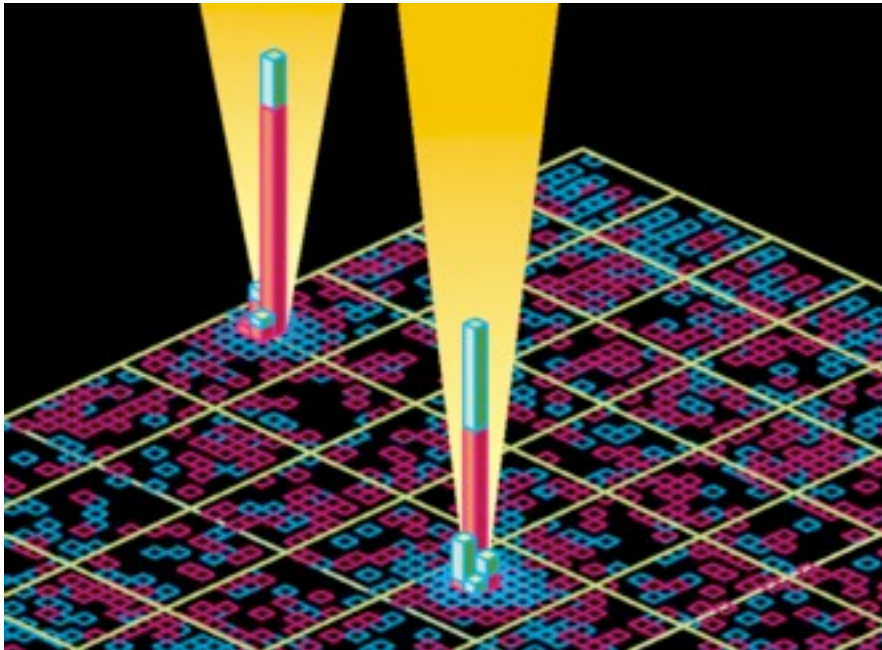


The Symmetries of QCD

R. Sekhar Chivukula
Michigan State University

Mar. 11, 2009

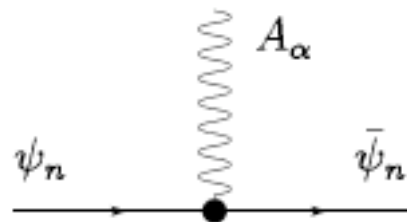
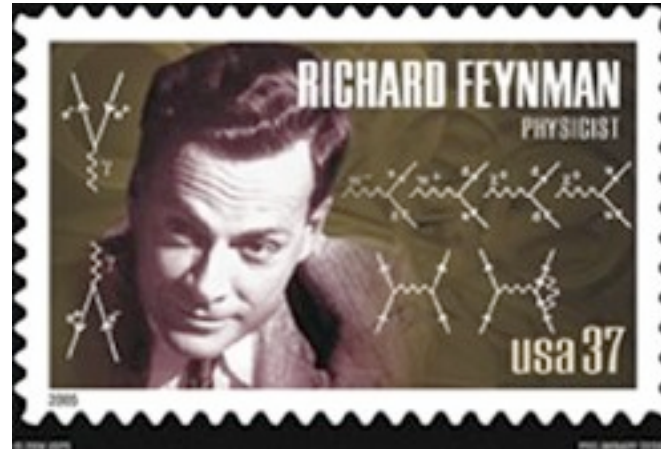
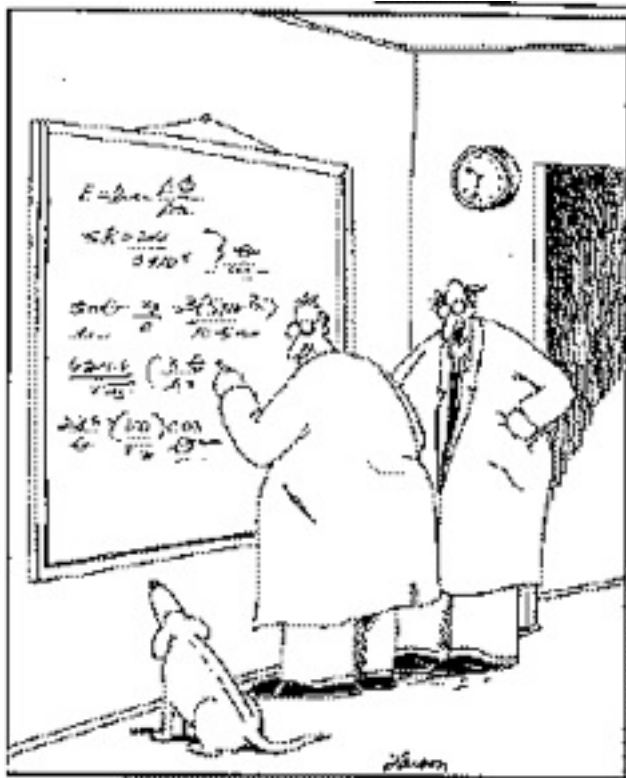
Fermilab Colloquium



The Symmetries of QCD

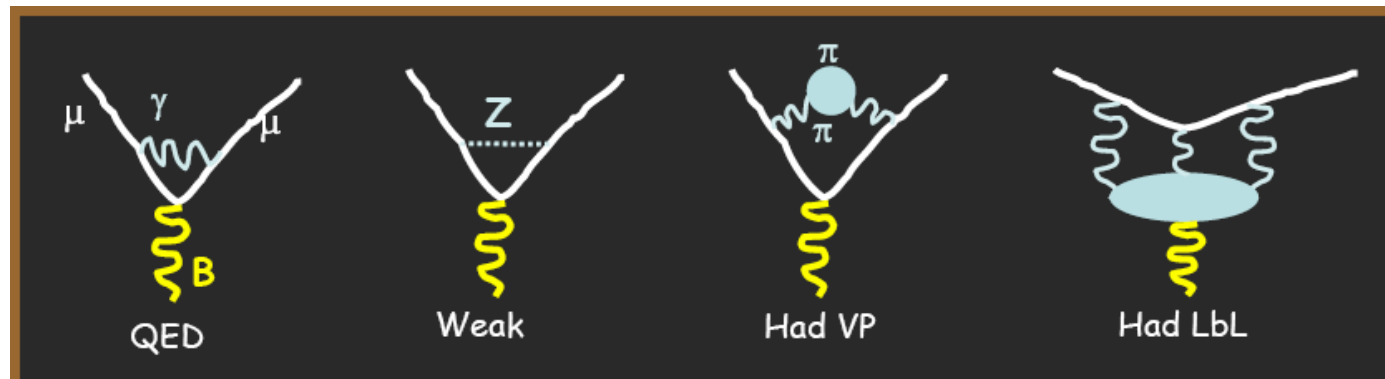
- What is QFT? Gauge Theory? QCD?
- Classical and Quantum Symmetries
- No Quantum Scale Symmetry!
- Chiral Symmetries are Broken
- Symmetry Summary
- Applications and Connections:
 - Morals for the LHC Era

Quantum Field Theory: QED



$$g = 2 + \frac{\alpha}{\pi} + \dots$$

“They act so cute when they try to understand Quantum Field Theory”



QED: Gauge Theory

Gauge “Symmetry”*

$$\psi(x) \rightarrow \exp(i e \theta(x)) \psi(x)$$

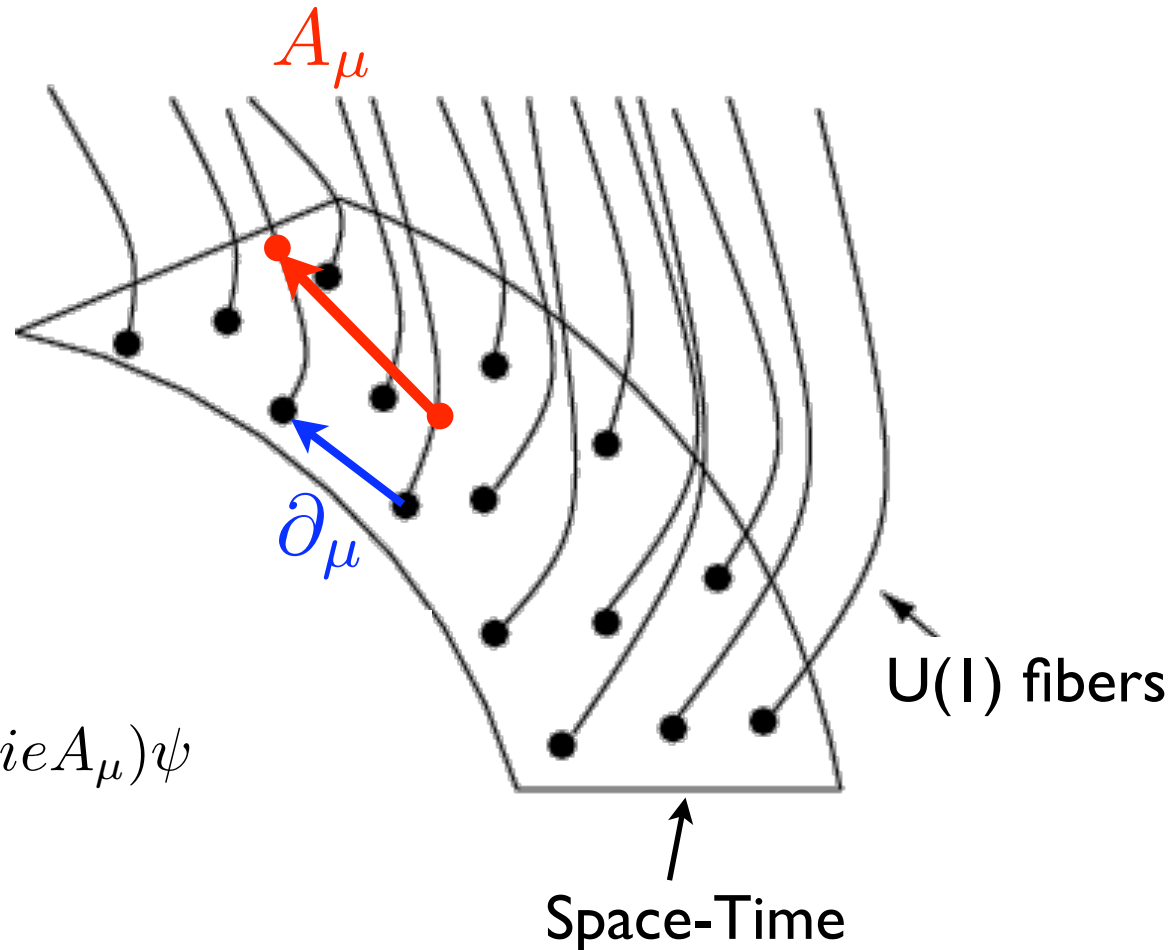
$$A_\mu(x) \rightarrow A_\mu(x) + \partial_\mu \theta(x)$$

$$\vec{E}^2 - \vec{B}^2$$

$$\vec{J} \cdot \vec{A}$$

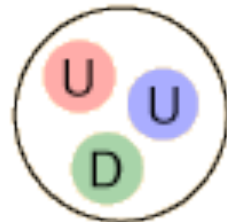
$$L_{QED} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + i \bar{\psi} \gamma^\mu (\partial_\mu - i e A_\mu) \psi$$

Weak Coupling

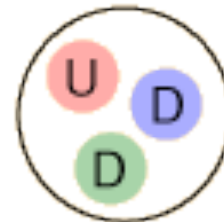


* Not a symmetry in the usual sense: rather it is a redundancy in the description of the theory!

Quantum Chromodynamics: Hadrons Contain Quarks



Proton
938 MeV



Neutron
940 MeV



π^+
140 MeV



ρ^+
770 MeV

Light quarks have “masses” of 300 MeV?

Why is the pion so light?

QCD Binds Quarks into Hadrons



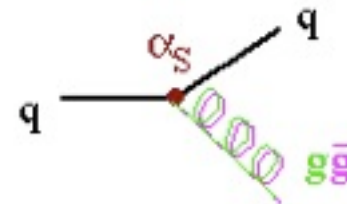
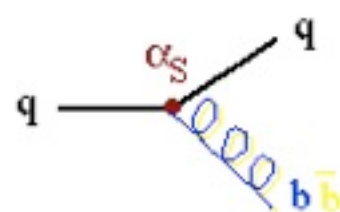
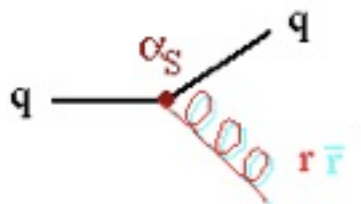
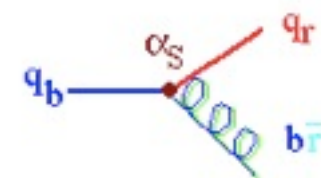
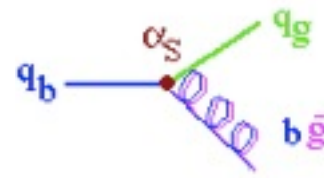
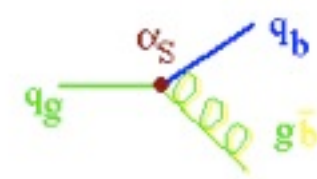
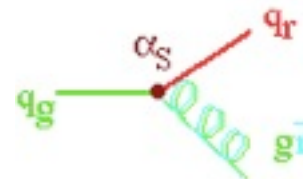
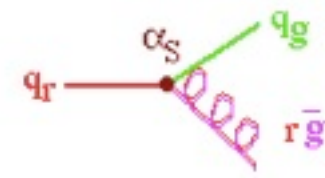
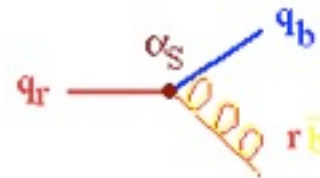
Quarks carry a color



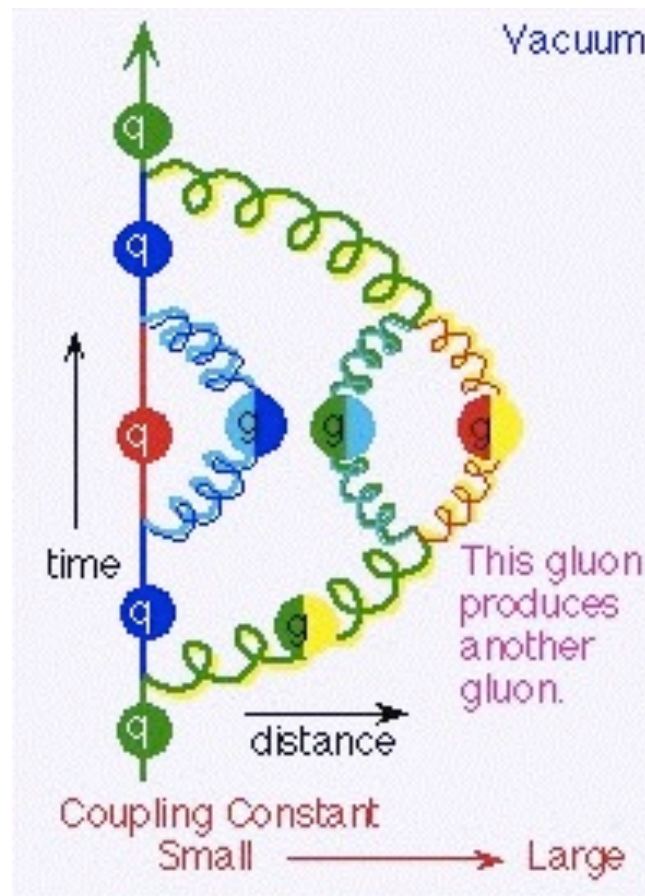
Anti-quarks carry an anti-color



Gluons carry a color and an anti-color

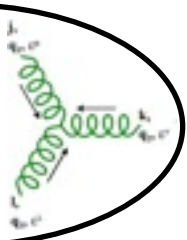


Gluons Couple to Gluons



QCD Lagrangian

$$\vec{E}_a^2 - \vec{B}_a^2$$

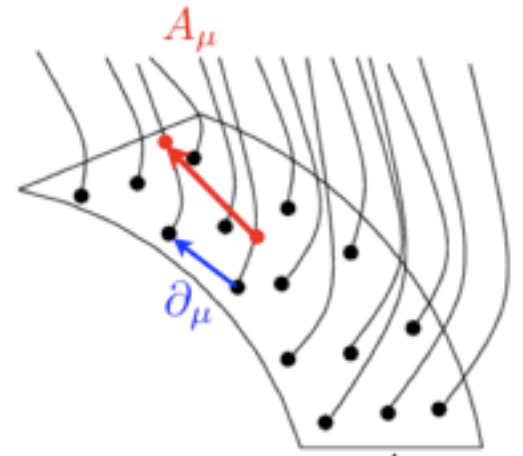


$$J_\mu^a \cdot A^{a\mu}$$

$$L_{\text{QCD}} = -\frac{1}{4} F_{\mu\nu}^{(a)} F^{(a)\mu\nu} + i \sum_q \bar{\psi}_q^i \gamma^\mu (D_\mu)_{ij} \psi_q^j - \sum_q m_q \bar{\psi}_q^i \psi_{qi}$$

$$F_{\mu\nu}^{(a)} = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - g_s f_{abc} A_\mu^b A_\nu^c$$

$$(D_\mu)_{ij} = \delta_{ij} \partial_\mu + ig_s \sum_a \frac{\lambda_{i,j}^a}{2} A_\mu^a$$



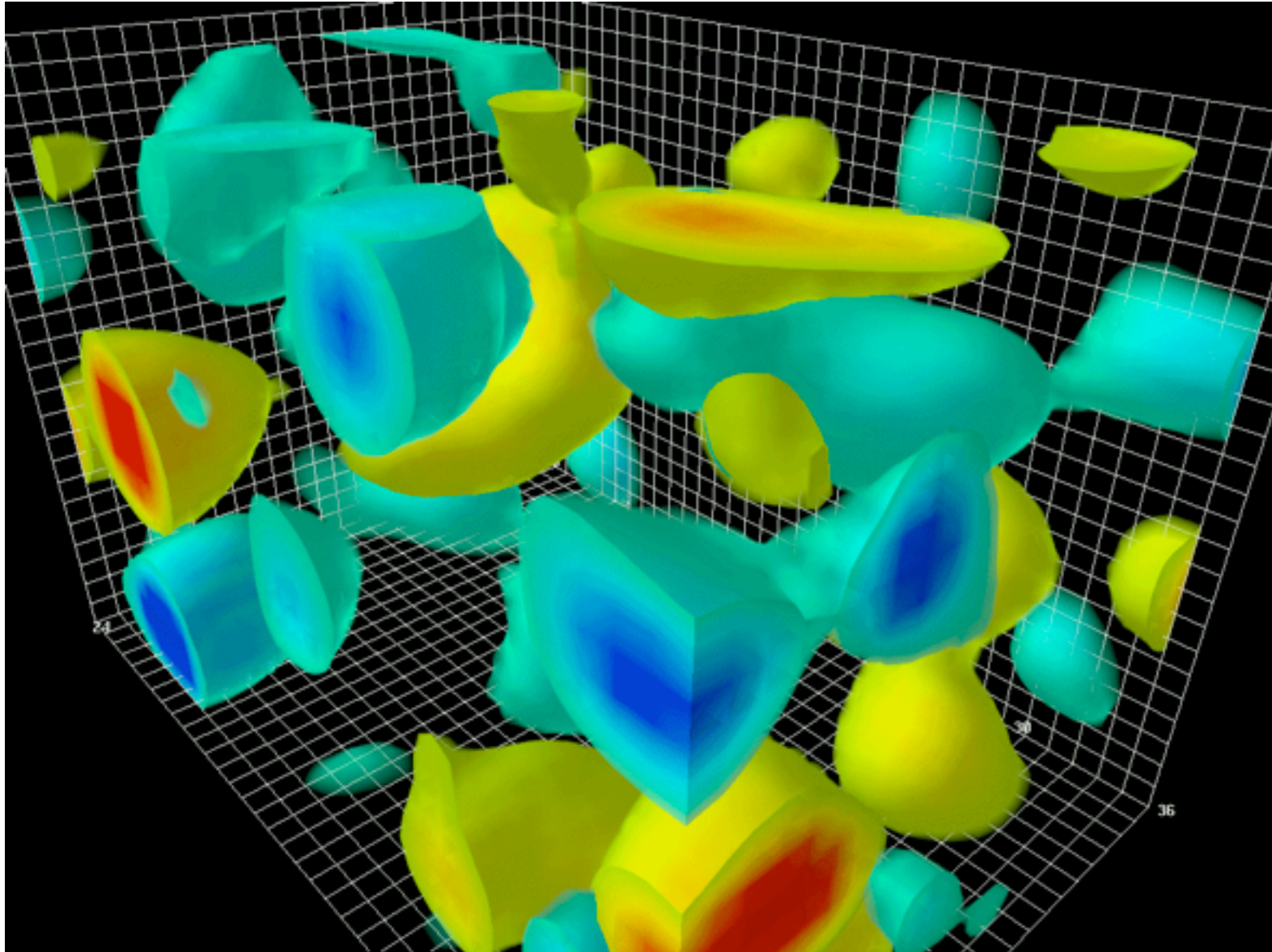
SU(3) gauge “symmetry”:

$$\psi_q(x) \rightarrow U(x)\psi_q(x)$$

$$A^\mu(x) \rightarrow U(x)A^\mu(x)U^\dagger(x) - \frac{1}{g_s}U(x)\partial^\mu U^\dagger(x)$$

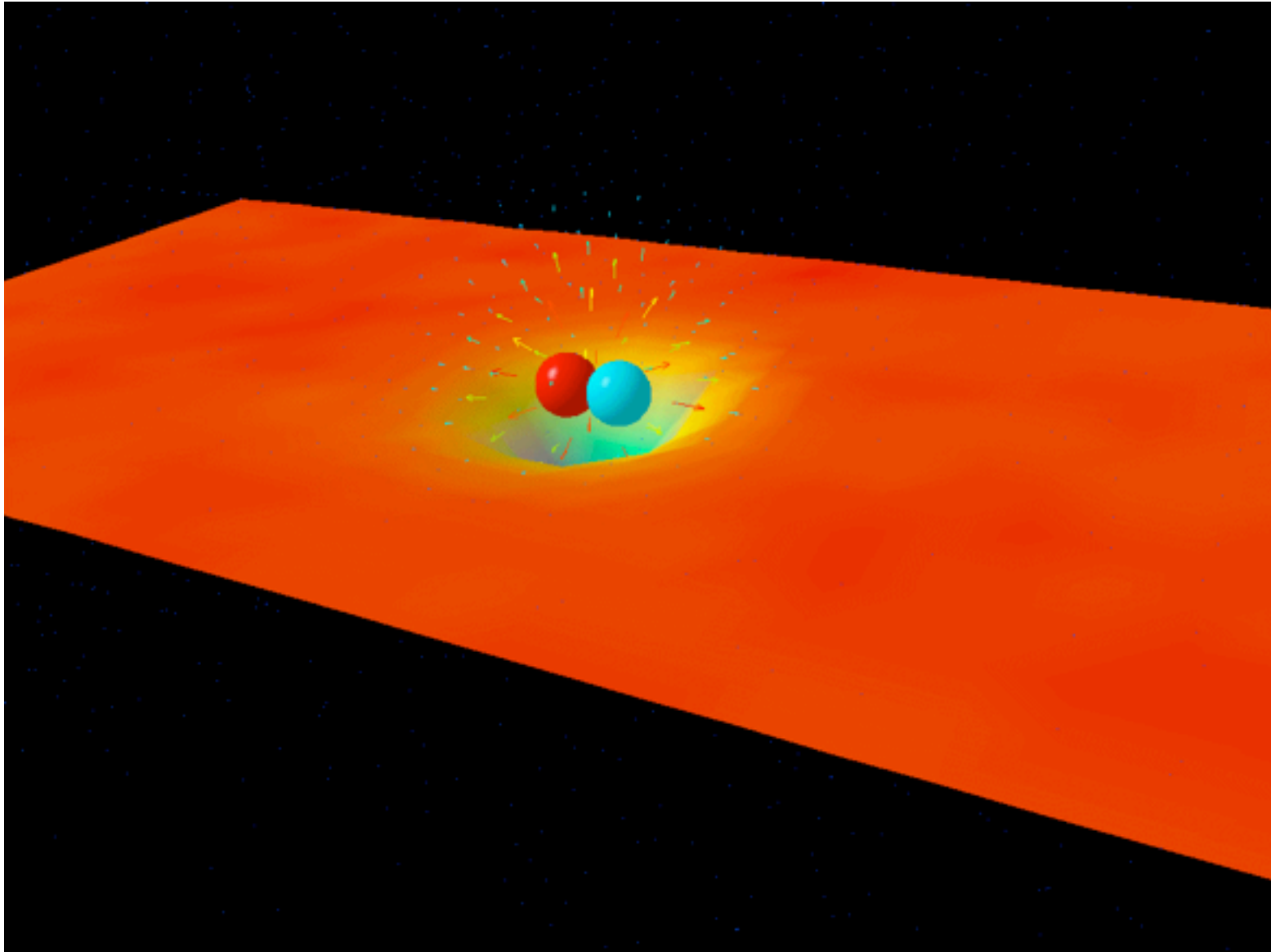
Parameters: Quark masses and QCD coupling

Lattice Simulation



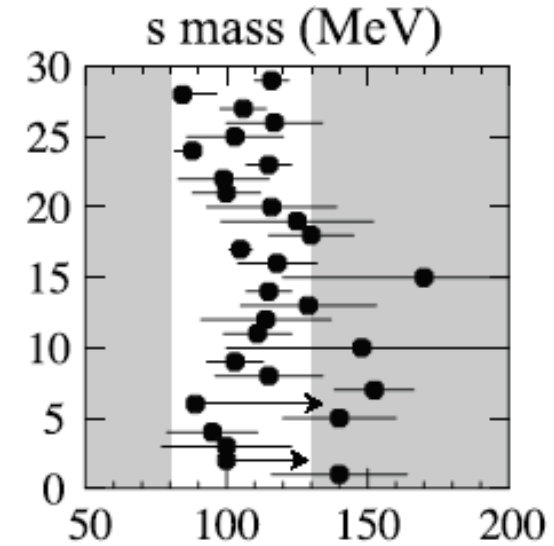
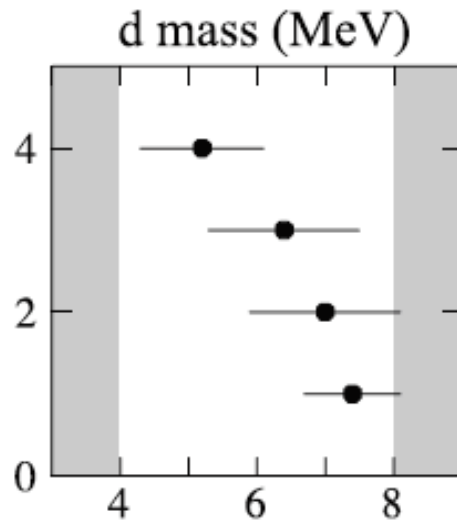
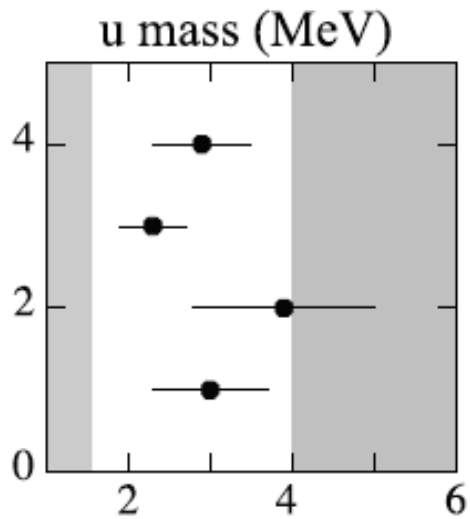
Derek Leinweber: www.physics.adelaide.edu.au/theory/staff/leinweber/

Quarks are Confined!



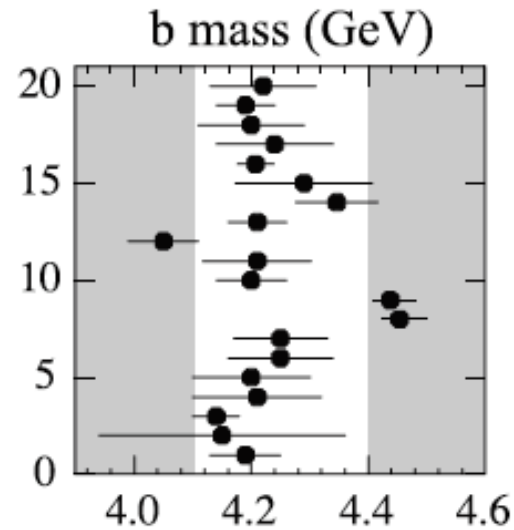
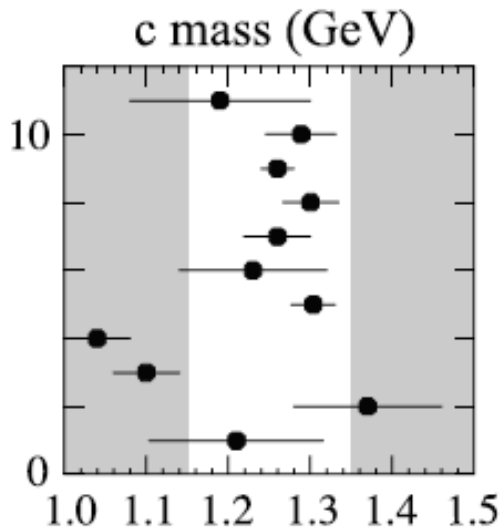
Derek Leinweber: www.physics.adelaide.edu.au/theory/staff/leinweber/

Lagrangian Quark Masses



Light?

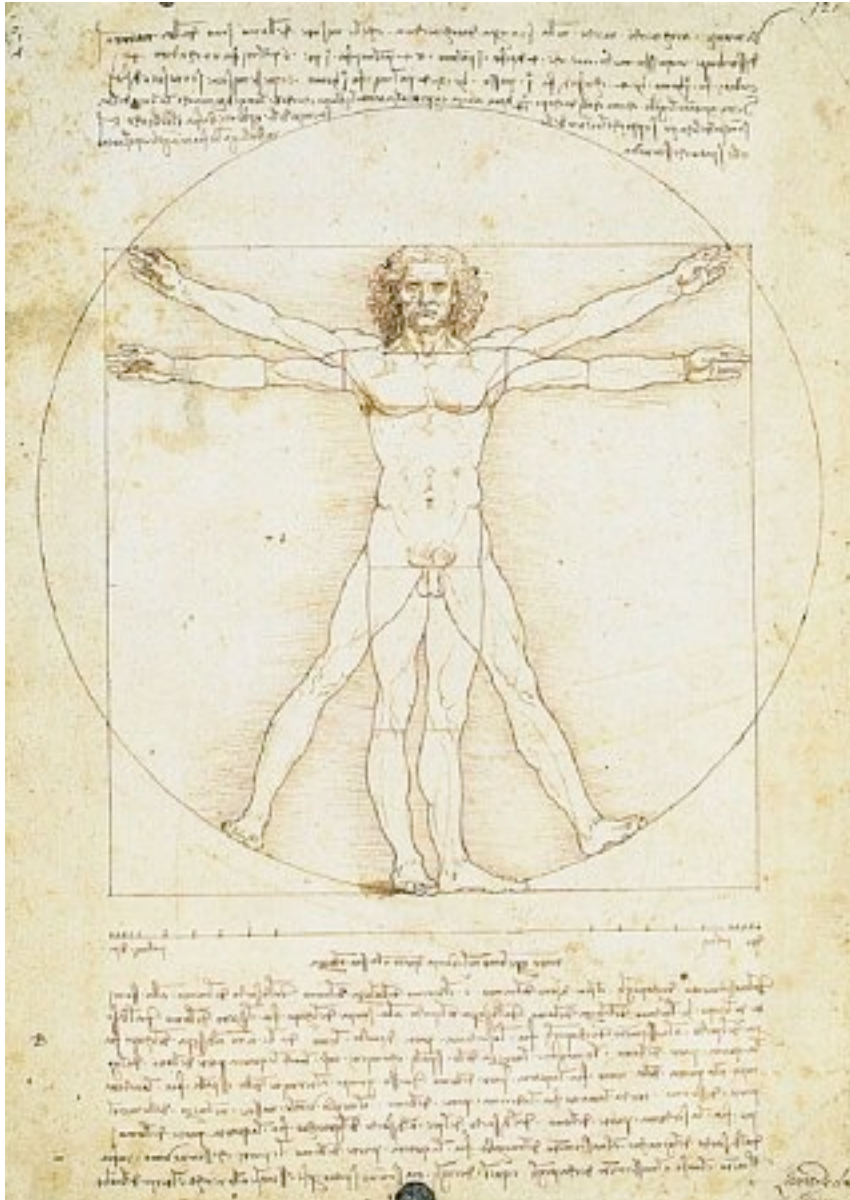
Heavy!



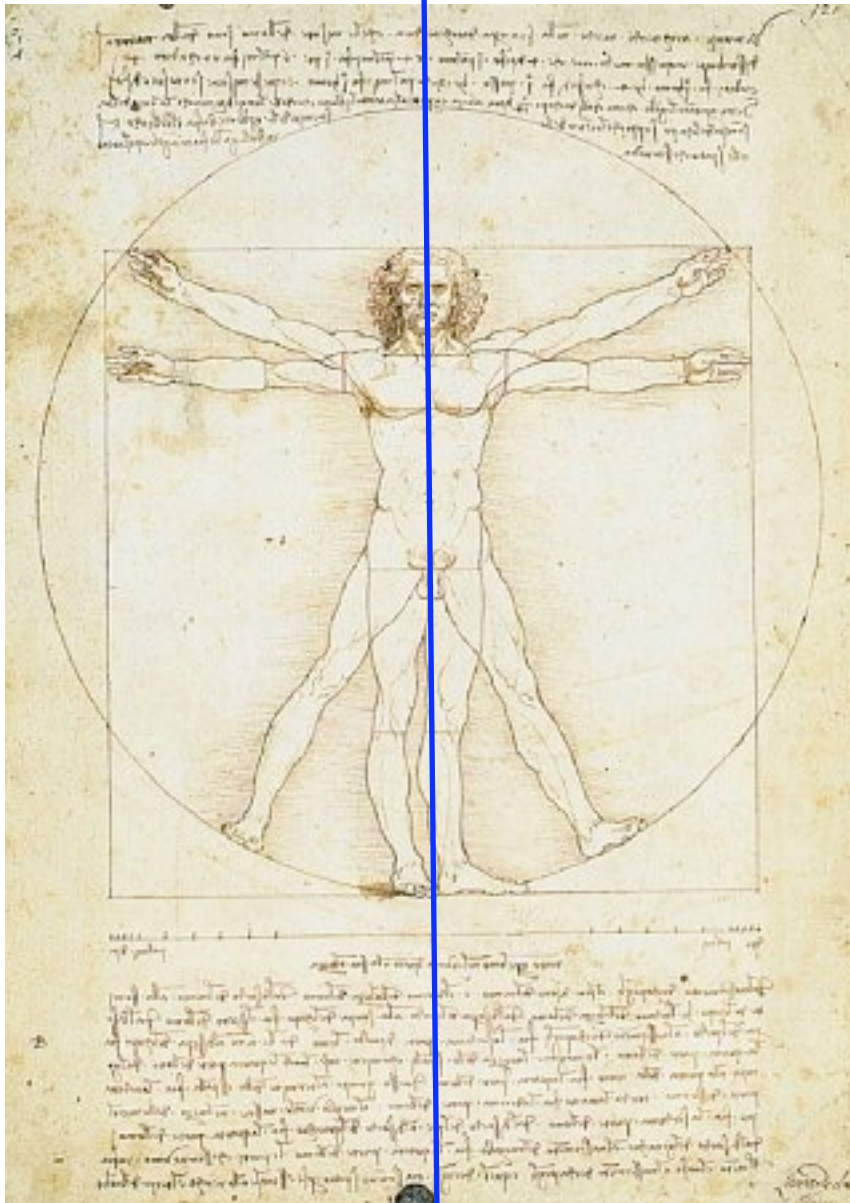
Explaining these masses is the “flavor problem”

RPP2004

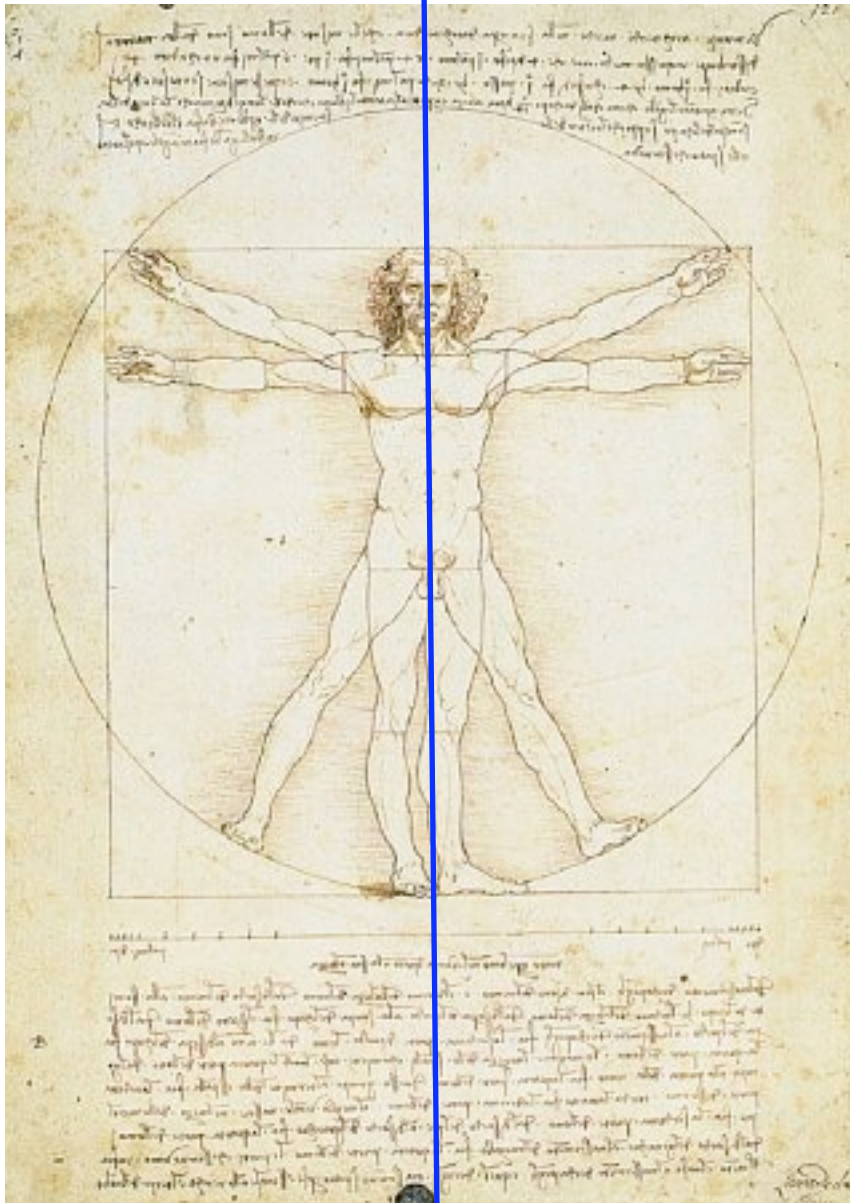
Symmetries



Symmetries



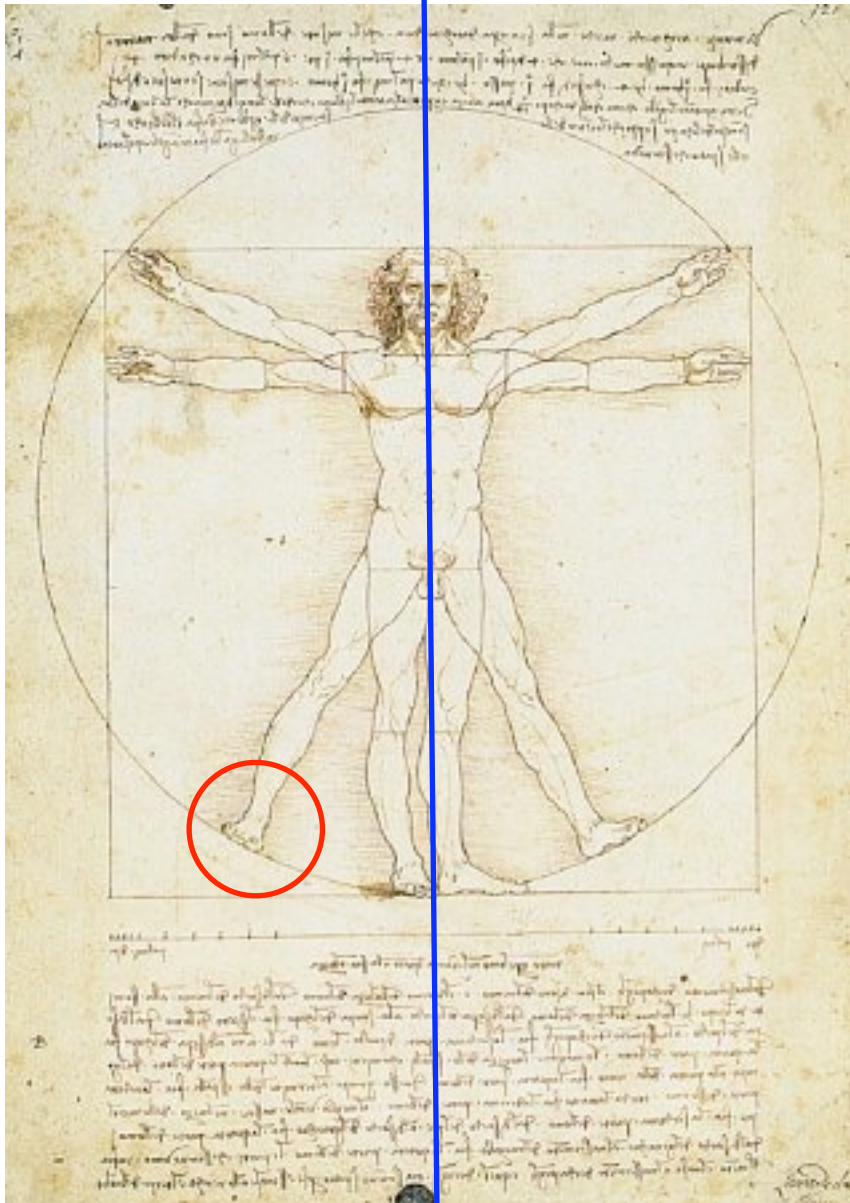
Symmetries



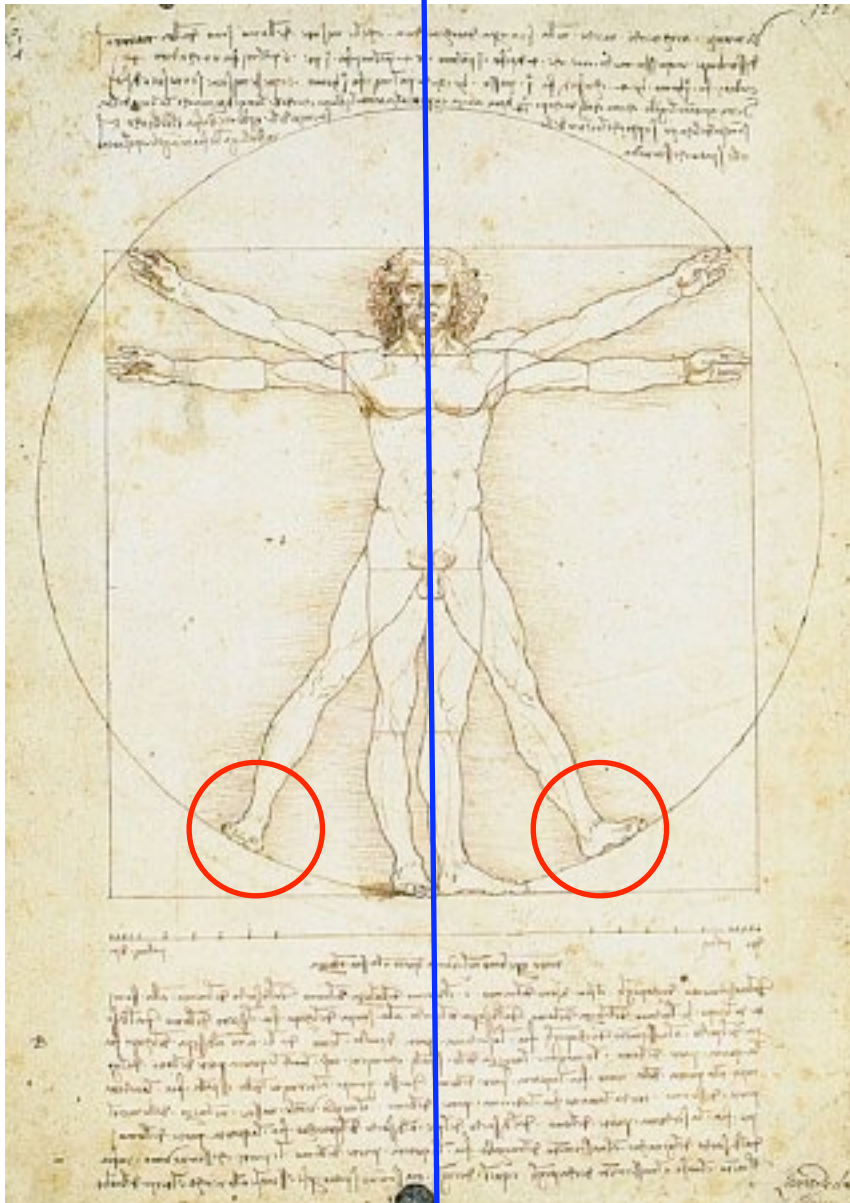
Bilateral Symmetry

Symmetries

Bilateral Symmetry

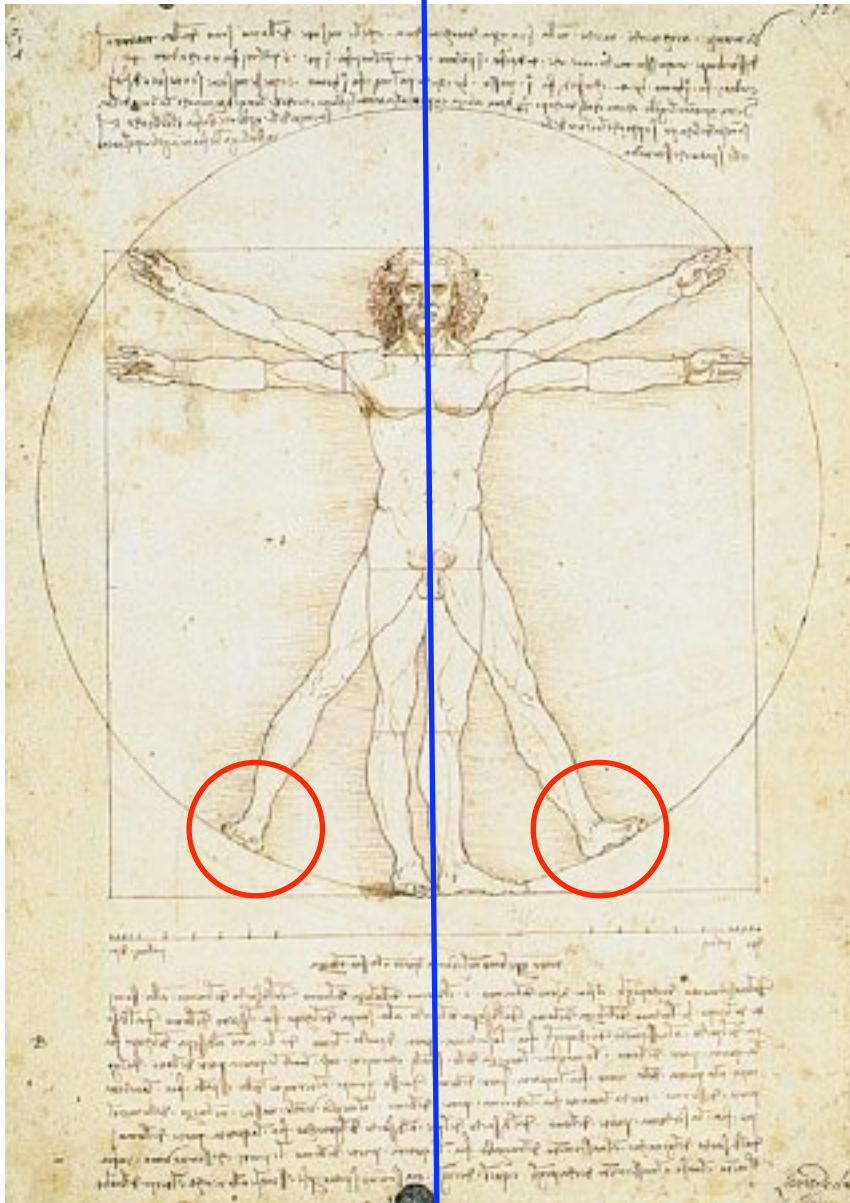


Symmetries



Bilateral Symmetry

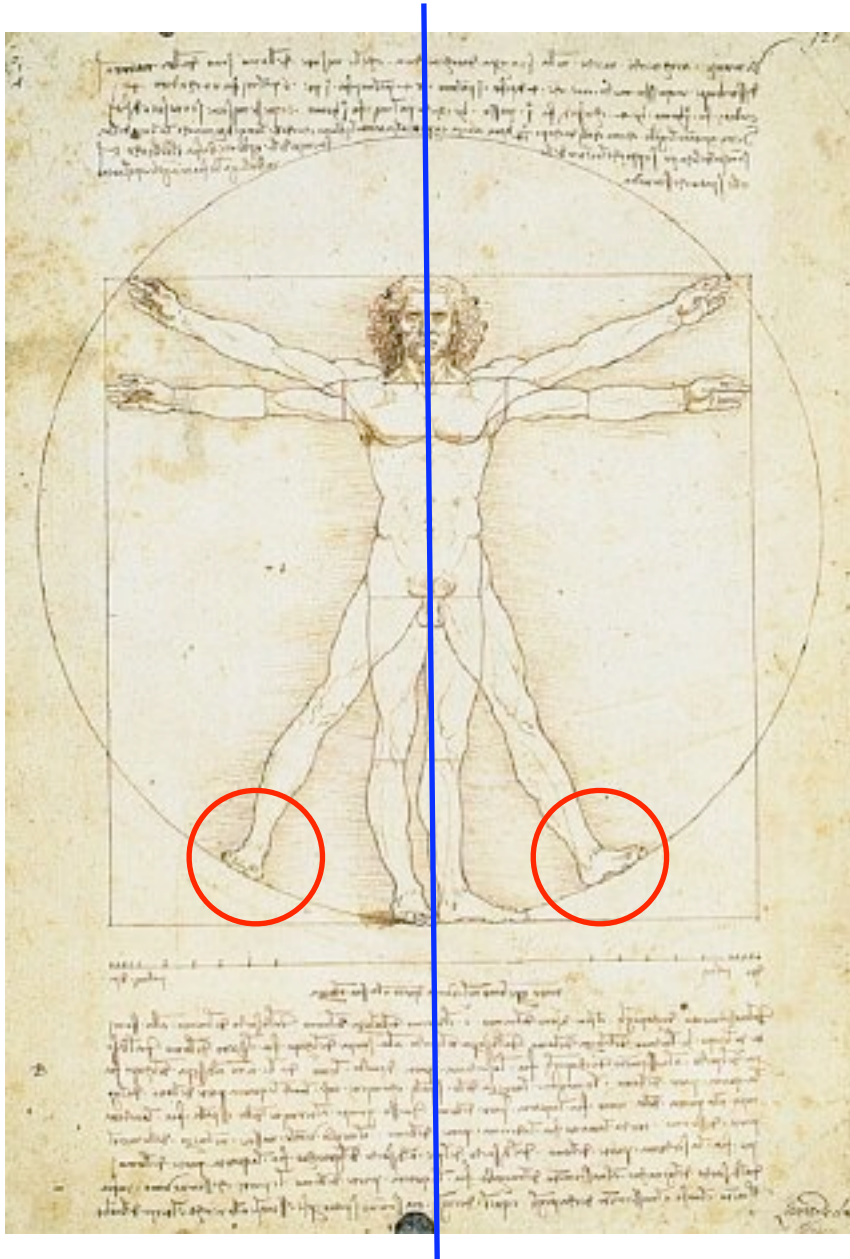
Symmetries



Bilateral Symmetry

Approximate
Symmetry

Symmetries

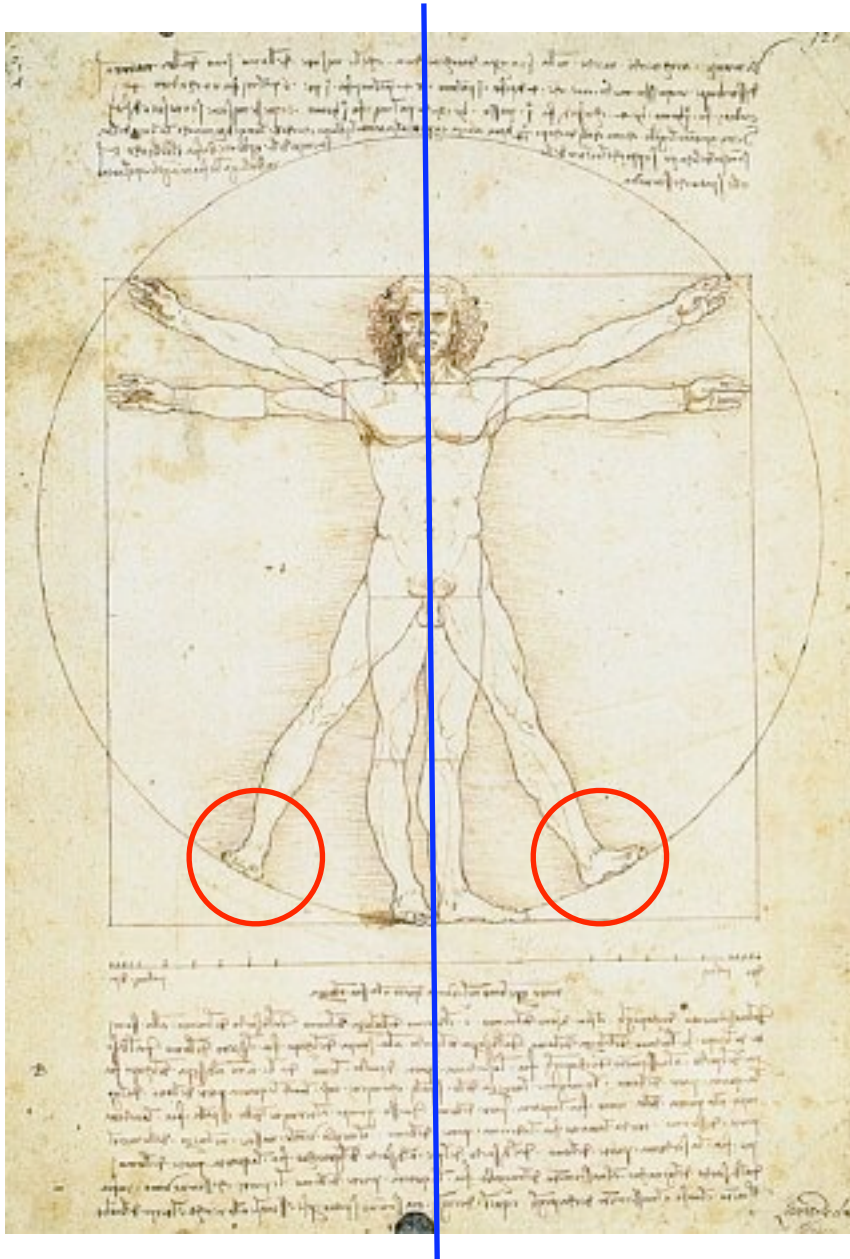


Bilateral Symmetry

Approximate
Symmetry

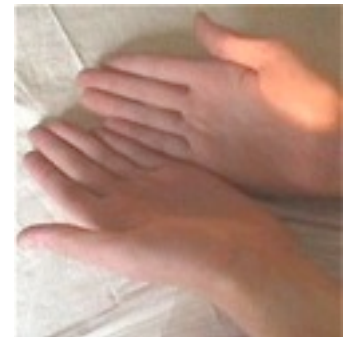
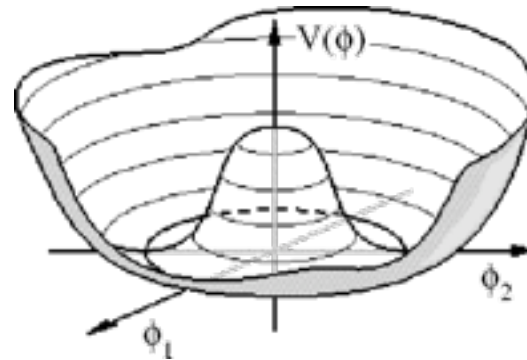


Symmetries

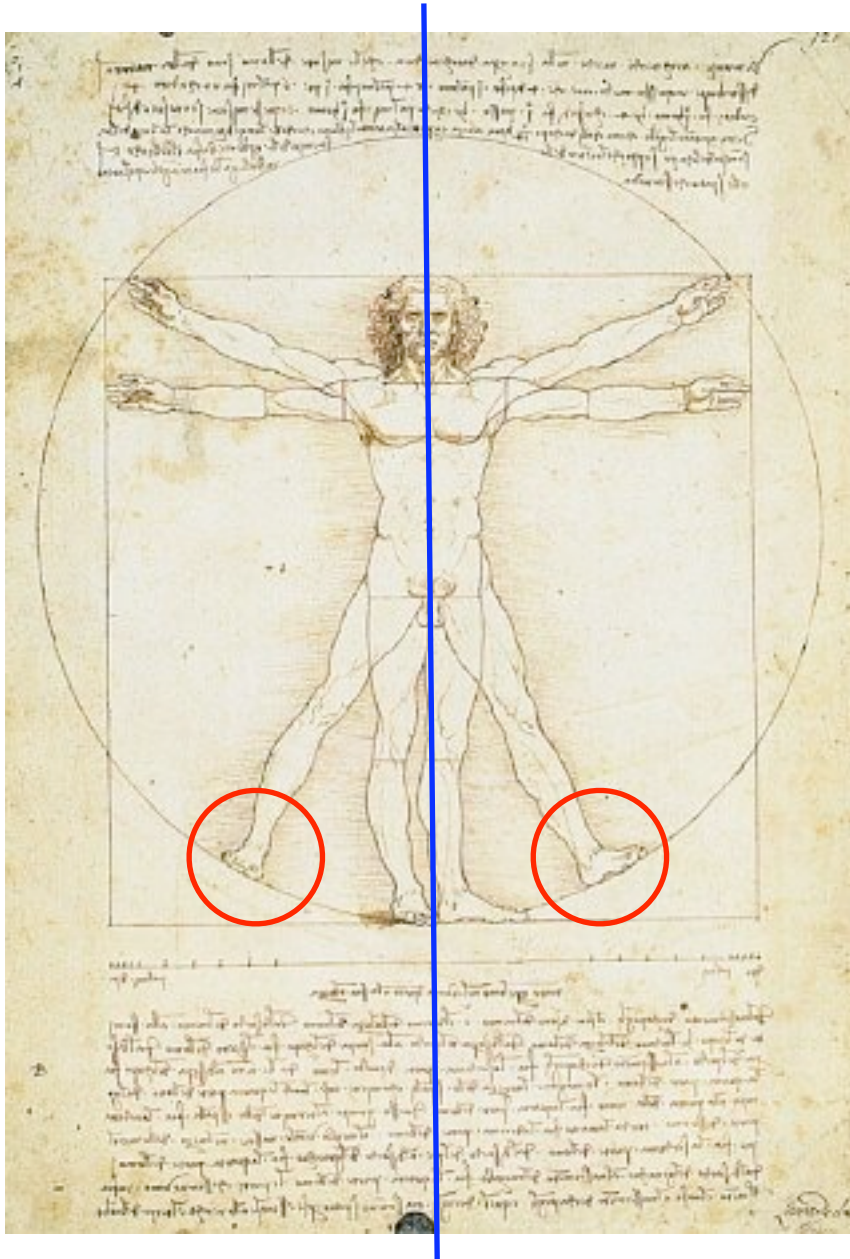


Bilateral Symmetry

Approximate
Symmetry

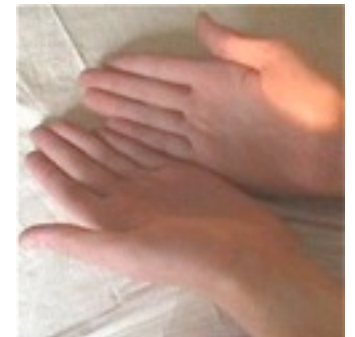
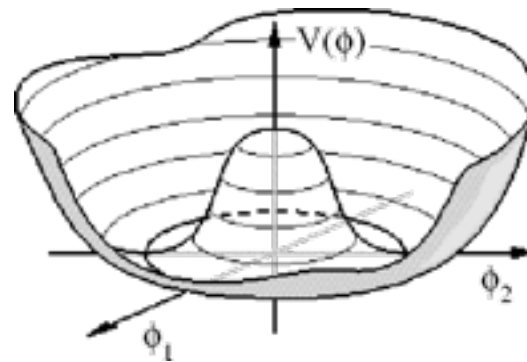


Symmetries



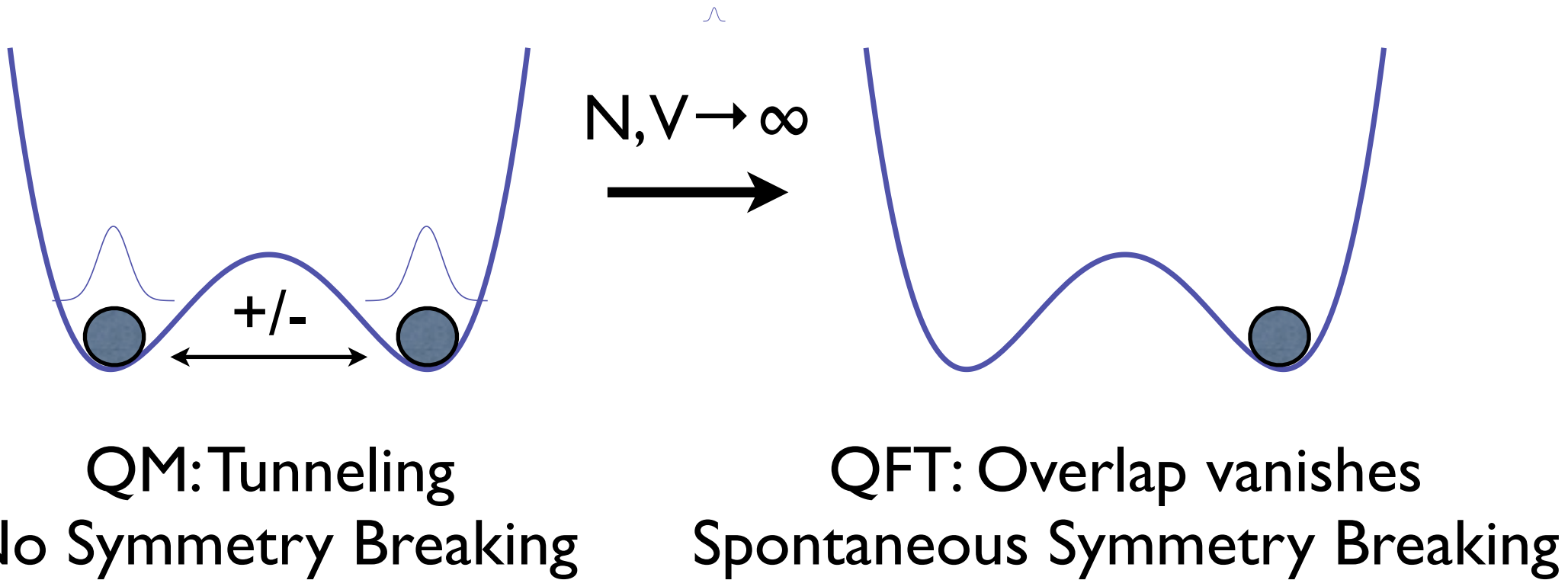
Bilateral Symmetry

Approximate
Symmetry

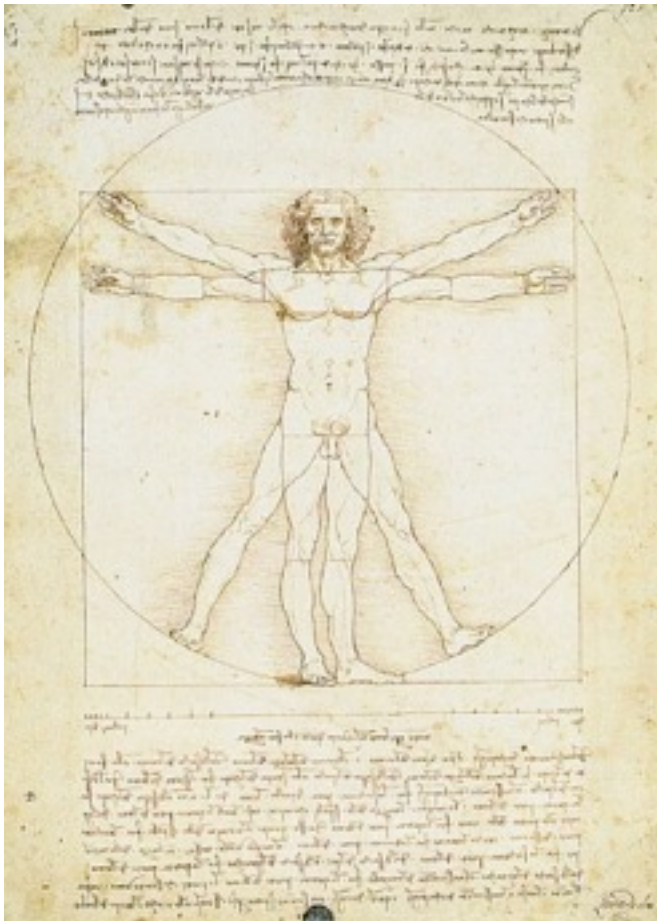


Spontaneously
Broken Symmetry -
By dynamics!

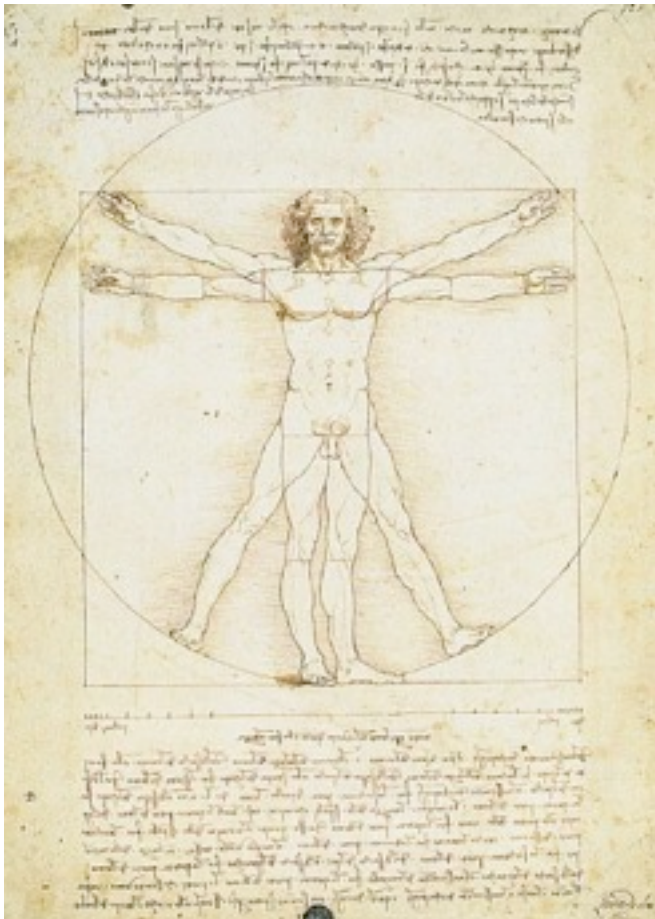
QFT differs from QM!



Anomalous Symmetry

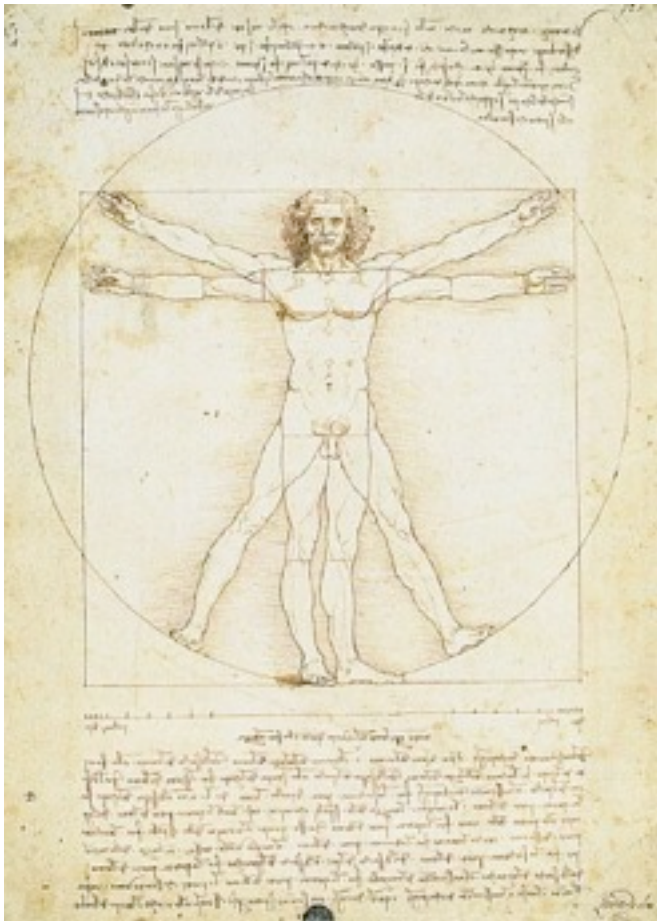


Anomalous Symmetry



Sometimes,
quantum
dynamics
eliminates
the symmetries!

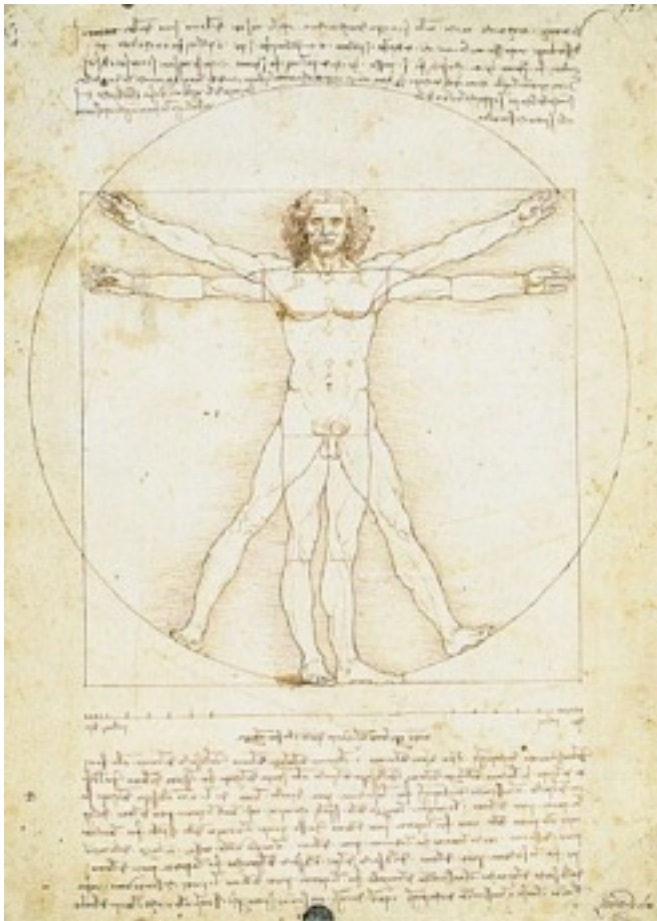
Anomalous Symmetry



Sometimes,
quantum
dynamics
eliminates
the symmetries!



Anomalous Symmetry

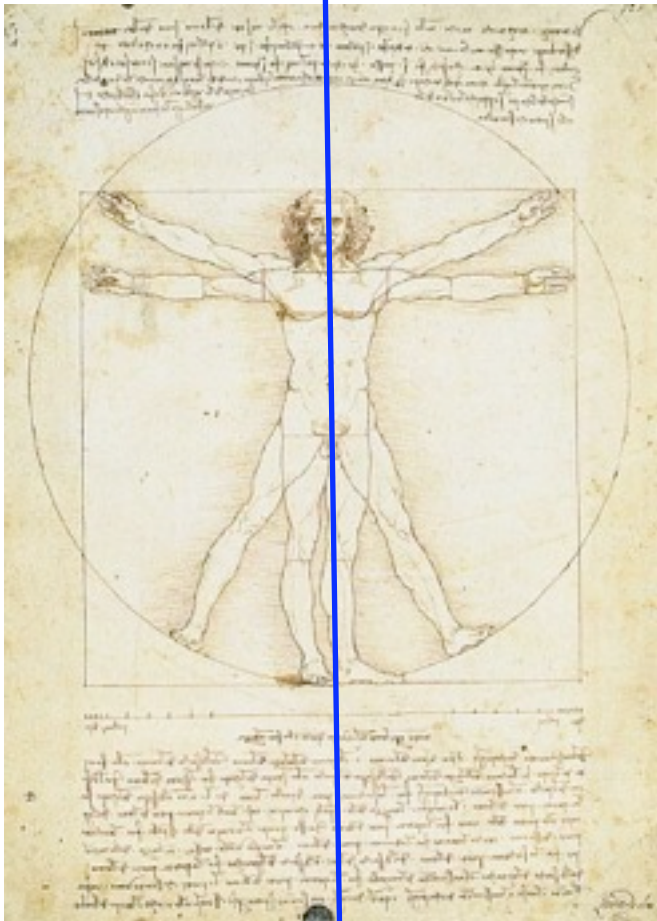


Sometimes,
quantum
dynamics
eliminates
the symmetries!



Resistance is futile!

Anomalous Symmetry

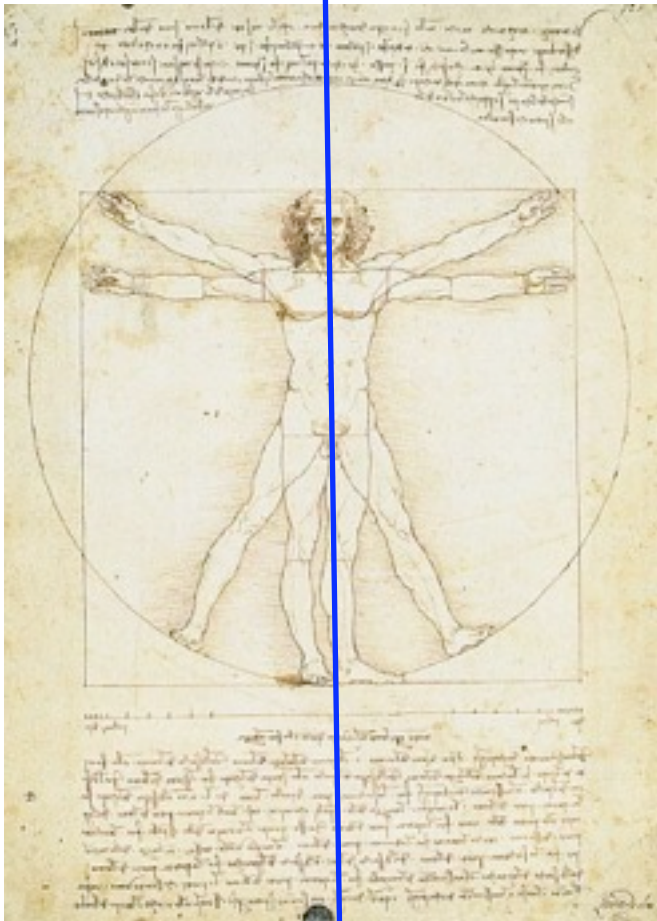


Sometimes,
quantum
dynamics
eliminates
the symmetries!

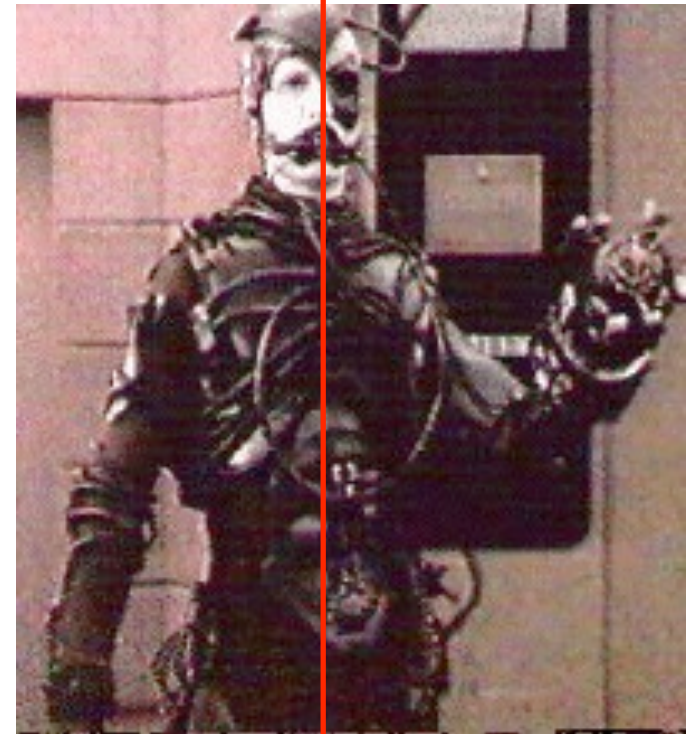


Resistance is futile!

Anomalous Symmetry



Sometimes,
quantum
dynamics
eliminates
the symmetries!



Resistance is futile!

Space-time Symmetries of Classical Chromodynamics (“~~Q~~CD”)

- Poincare Invariance
- C, P, and T (as written...)
- (Approximate) Scale Invariance*

A scale transformation:



$$x^\mu \rightarrow \lambda \cdot x^\mu$$

$$\psi_q(x) \rightarrow \lambda^{3/2} \psi_q(\lambda x) \quad A_\mu^a(x) \rightarrow \lambda A_\mu^a(\lambda x)$$

NB: Broken by all mass terms...

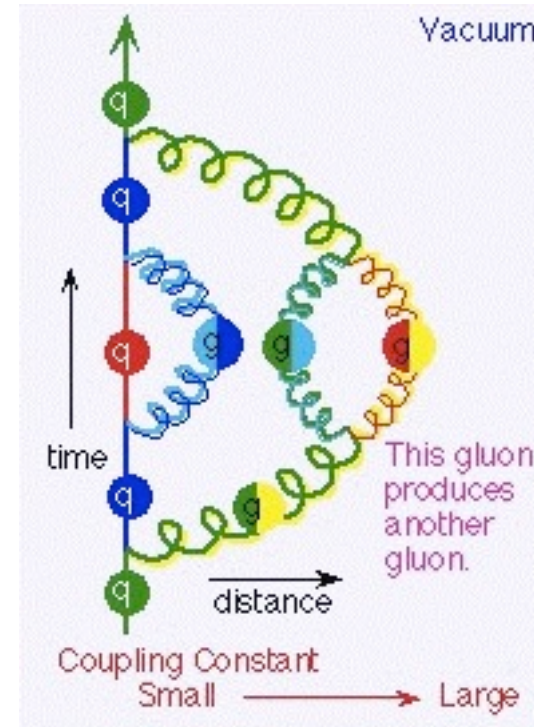
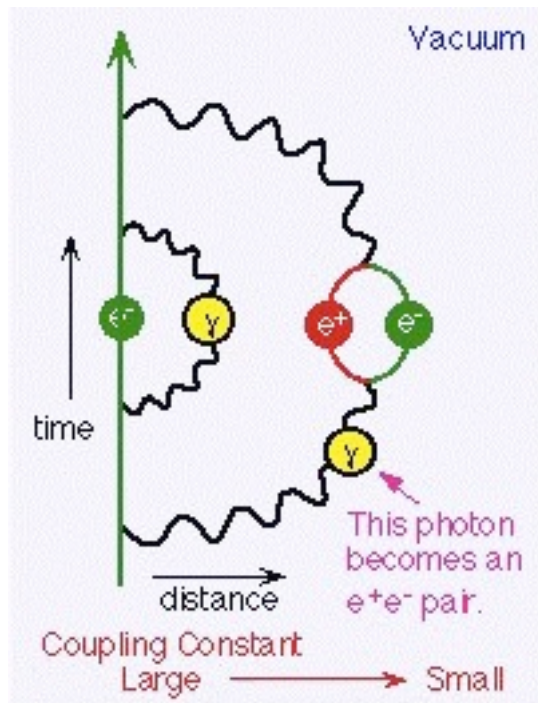
*(& Proper conformal trans.)

NB: Local scale invariance the reason for the name “gauge transformation” (Weyl)

No Quantum Scale Invariance!

- Quantum vacuum is a polarizable medium

QED →



← QCD

Fermions and Gluons make opposite contributions!

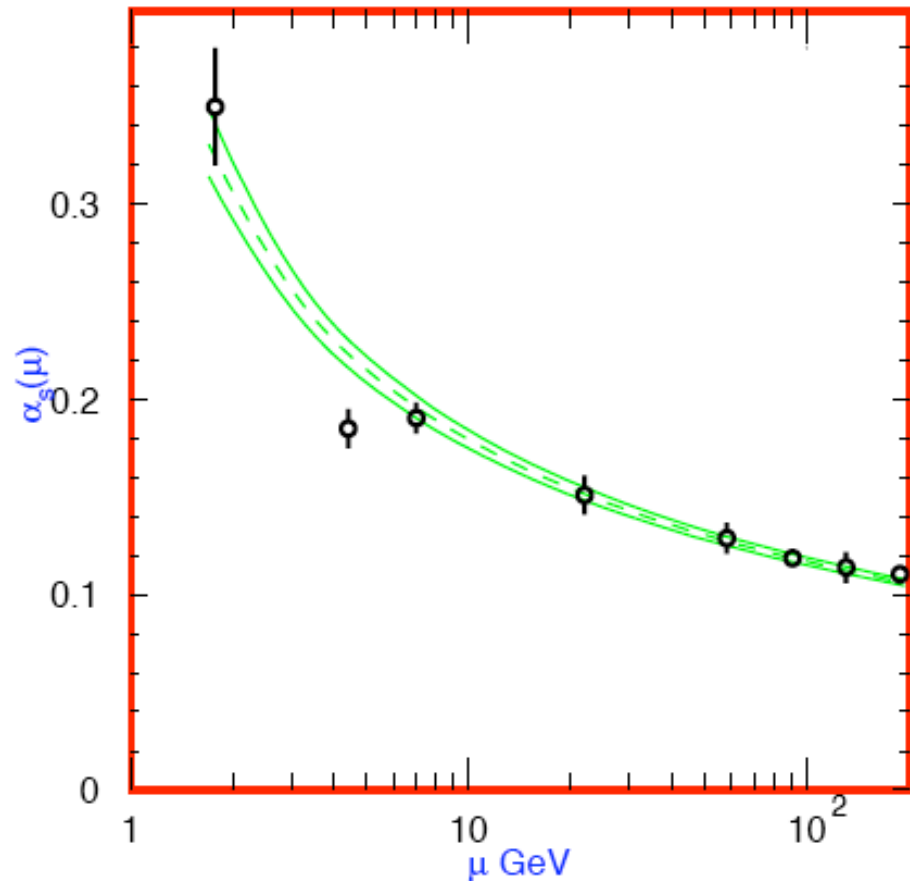
QCD Beta Function

$$\mu \frac{\partial \alpha_s}{\partial \mu} = 2\beta(\alpha_s) = -\frac{\beta_0}{2\pi} \alpha_s^2 - \frac{\beta_1}{4\pi^2} \alpha_s^3$$

$$\beta_0 = 11 - \frac{2}{3}n_f,$$

$$\beta_1 = 51 - \frac{19}{3}n_f,$$

Asymptotic Freedom
vs.
Infrared Slavery



2004 Nobel Prize:
Gross, Politzer, Wilczek

RPP2004

Dimensional Transmutation

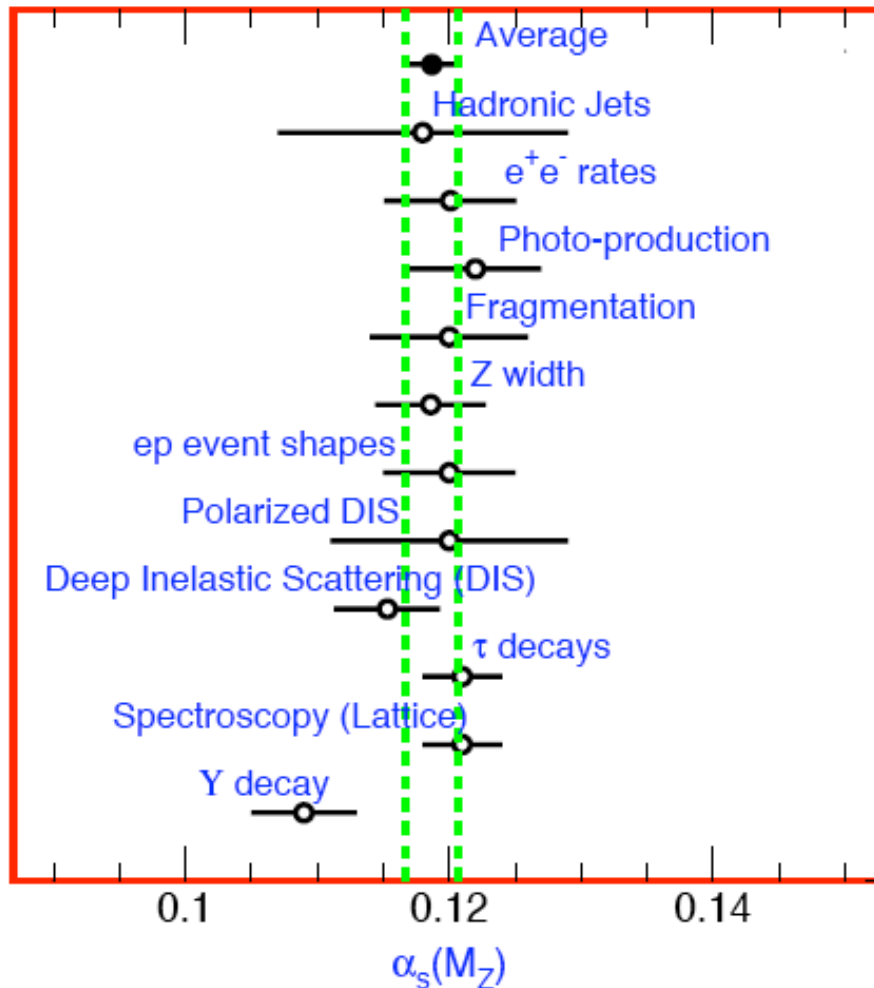
- Value of strong-coupling fixed at one scale, predicted at others.
- Alternatively, given by energy scale: Λ_{QCD}

$$\alpha_s(\mu) = \frac{4\pi}{\beta_0 \ln(\mu^2/\Lambda^2)} \left[1 - \frac{2\beta_1}{\beta_0^2} \frac{\ln[\ln(\mu^2/\Lambda^2)]}{\ln(\mu^2/\Lambda^2)} + \frac{4\beta_1^2}{\beta_0^4 \ln^2(\mu^2/\Lambda^2)} \right. \\ \left. \times \left(\left(\ln[\ln(\mu^2/\Lambda^2)] - \frac{1}{2} \right)^2 + \frac{\beta_2\beta_0}{8\beta_1^2} - \frac{5}{4} \right) \right].$$

RPP2004

Strong-interaction scale set by Λ_{QCD}

Strong Coupling



$$\Lambda_{\text{QCD}}^{(5)} = 217^{+25}_{-23} \text{ MeV}$$

$$\Lambda_{\text{QCD}}^{(3)} \simeq 350 \text{ MeV}$$

Λ_{QCD} defines “light”
and “heavy”

Why are nuclear binding energies so much smaller?

Scale Current

- Any continuous transformation defines an associated current
- A continuous symmetry has a conserved current

$$\partial_\mu s^\mu = T_\lambda^\lambda = -\frac{\beta}{2g} F_{\mu\nu}^{(a)} F^{(a)\mu\nu} + \sum_q m_q \bar{\psi}_q \psi_q + \dots$$



Quantum Effect!

Scale Symmetry is Anomalous!

Global Quark Flavor Symmetries of “CD” (I)

- Baryon Number: $\psi_q \rightarrow e^{i\alpha} \psi_q$
- (Approximate*) Gell-Mann $SU(3)_v$:

$$\begin{pmatrix} u \\ d \\ s \end{pmatrix} \rightarrow U \begin{pmatrix} u \\ d \\ s \end{pmatrix}$$

Together: $U(3)_v$

Realized “As Is” in Quantum Theory!

*Requires $m_u \approx m_d \approx m_s$

Global Quark Flavor Symmetries of “CD” (2)

- (Approximate*) Chiral $SU(3)_L \times SU(3)_R$:

$$\bar{\psi}_q i\not{D}\psi_q \equiv \bar{\psi}_q^L i\not{D}\psi_q^L + \bar{\psi}_q^R i\not{D}\psi_q^R$$

$$\begin{pmatrix} u^{L,R} \\ d^{L,R} \\ s^{L,R} \end{pmatrix} \rightarrow U_{L,R} \begin{pmatrix} u^{L,R} \\ d^{L,R} \\ s^{L,R} \end{pmatrix}$$

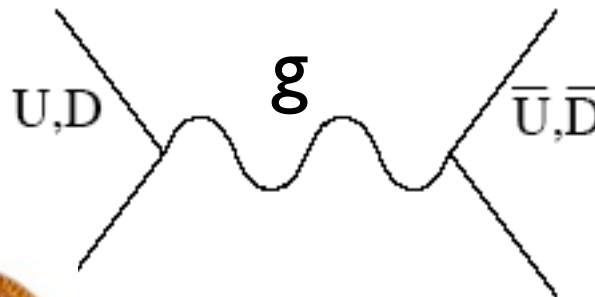
*For “light” quarks: u, d, and s(?)

$$m_q \bar{\psi}_q \psi_q \equiv m_q \bar{\psi}_q^L \psi_q^R + m_q \bar{\psi}_q^R \psi_q^L$$

Quantum Theory different...

Chiral Symmetries are Spontaneously Broken

- Chiral symmetries are good quantum symmetries, but ...
- ... strong low-energy QCD dynamics rearranges the vacuum


$$\rightarrow \langle \bar{U}_L U_R \rangle = \langle \bar{D}_L D_R \rangle \neq 0.$$



2008 Nobel Prize: Nambu

Goldstone Bosons

- $SU(3)_L \times SU(3)_R$ breaks to Gell-Mann $SU(3)_V$
- Implies existence of eight (massless) Goldstone Bosons

$$j_{A\mu}^i = -f_\pi \partial_\mu \pi^i + \dots$$

$$\partial^\mu j_{A\mu}^i = -f_\pi \partial^\mu \partial_\mu \pi^i + \dots = 0 \Leftrightarrow m_\pi^2 = 0$$

Pseudo Goldstone Bosons

- Relevant for light quarks only: $m_q \leq \Lambda_{QCD}$
- u, d, and s are not massless,
- Therefore, pions are light*:

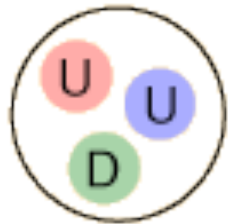
$$m_\pi^2 \propto (m_u + m_d)\Lambda_{QCD}$$

$$m_K^2 \propto (m_s + m_{u,d})\Lambda_{QCD}$$

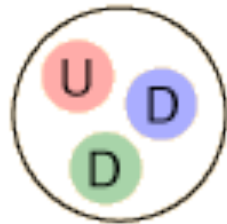
$$m_\eta^2 \propto \frac{1}{3}(m_u + m_d + 4m_s)\Lambda_{QCD}$$

*NB: QCD Lagrangian Quark masses appear here!

Constituent Quarks



Proton
938 MeV



Neutron
940 MeV



π^+
140 MeV



ρ^+
770 MeV

~300 MeV quark-model mass is of a “dressed” or “constituent” quark

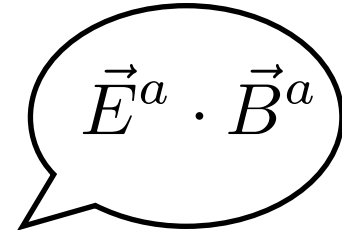
~99% of proton and neutron “mass” due to QCD field energy!



What about $U(1)_A$?

- $U(1)_A$: $\psi_q^L \rightarrow e^{i\theta} \psi_q^L$ & $\psi_q^R \rightarrow e^{-i\theta} \psi_q^R$
- Broken by chiral condensate
- Where is the ninth Goldstone Boson?
- Another Quantum “Anomaly”:

$$\partial_\mu j_A^\mu \propto \frac{g^2}{16\pi^2} \varepsilon^{\mu\nu\alpha\beta} F_{\mu\nu}^{(a)} F_{\alpha\beta}^{(a)}$$



Adler, Bell, Jackiw

't Hooft

Strong-CP Problem?

- But the same term could be in the QCD Lagrangian...
- ... and violates CP!
- Yields a neutron electric dipole moment
- Motivates the Axion...

't Hooft

Peccei and Quinn

QCD is the culmination almost a century of effort in QFT...

- Heisenberg, Pauli, Dirac, Jordan, Oppenheimer,...
- Schwinger, Feynman, Tomonaga, Bethe,...
- Nambu, Goldstone, ...
- Yang & Mills
- 't Hooft, Veltman, Lee, Ward,...
- Gross, Politzer, Wilczek

Shelter Island, 1947

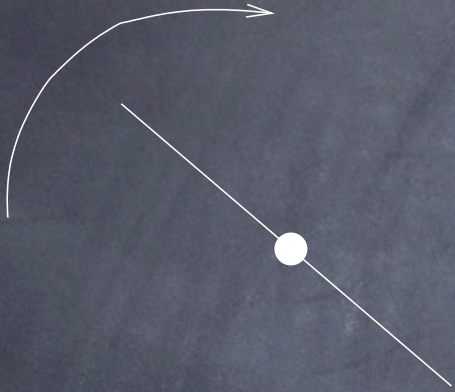


Symmetries in QCD

Symmetry	Anomaly?	Fate
Color (gauged)	no	unbroken (confined)
<u>Scale Invariance</u>	yes	not a symmetry
$SU(3)_V \times U(1)_B$	no	unbroken
<u>$SU(3)_A$</u>	no	spontaneously broken
$U(1)_A$	yes	not a symmetry
CP	no?!	not a symmetry!?

New Theoretical Developments...

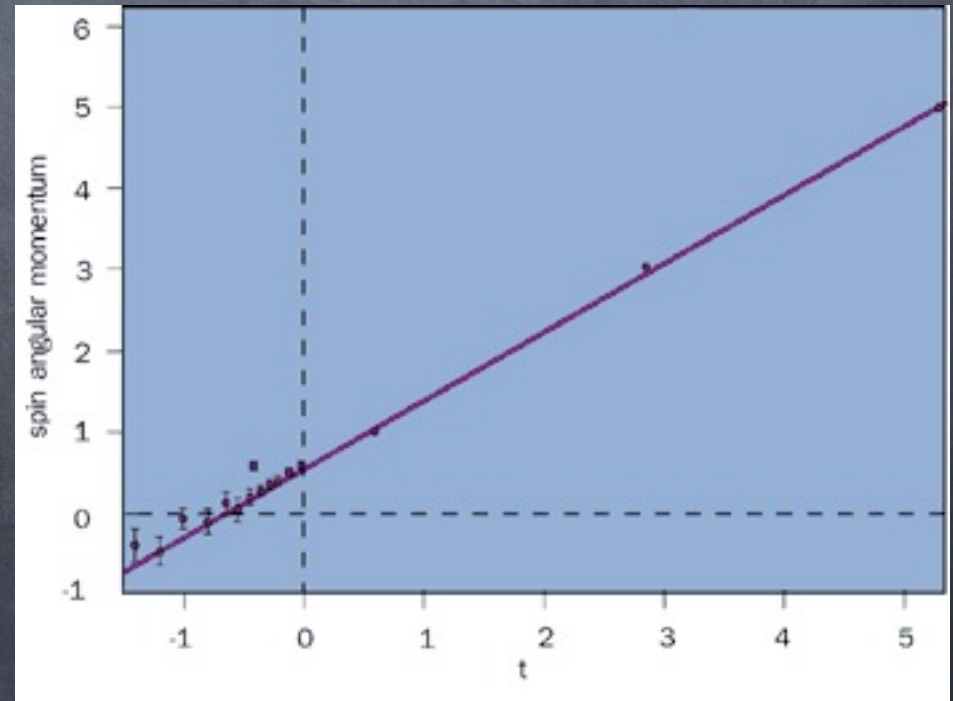
QCD and String Theory: Relativistic Classical String



$$E \propto \frac{J^2}{2I} \quad I \propto ML^2$$

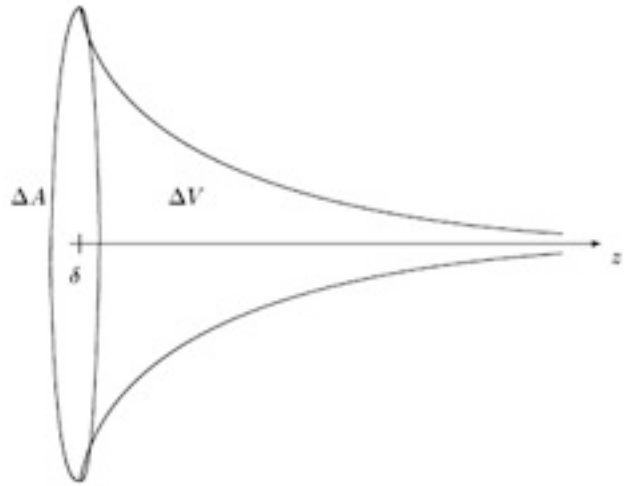
$$M, L \propto E \quad \Rightarrow \quad J \propto E^2$$

Regge Trajectories



AdS/CFT Duality

Conjecture: Equivalence of 5D theory in AdS and 4D CFT



$$ds^2 = \left(\frac{R}{z}\right)^2 [\eta_{\mu\nu} dx^\mu dx^\nu - dz^2]$$

$$R < z < R'$$

UV \longrightarrow IR

NB: Rescaling Invariance!

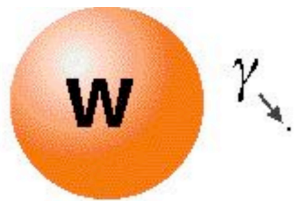
Strong evidence for N=4 SUSY YM string theory on AdS

Strongly-coupled CFT \Leftrightarrow Weakly-coupled 5D Theory!

Is there an 5D/QCD duality?

Morals for the LHC Era

What is the origin of the Higgs Mechanism?

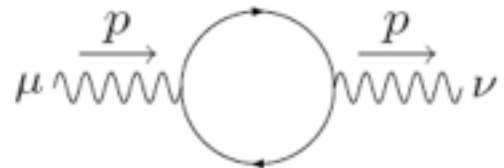


Why do some particles have large masses while the photon and neutrinos have puny masses?



The Higgs Mechanism

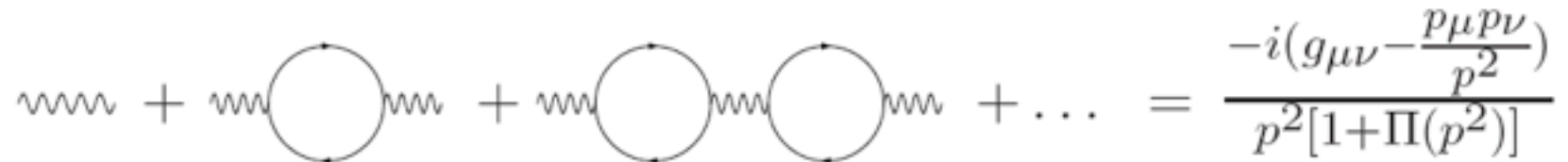
The polarization tensor $\Pi_{\mu\nu}(p)$ is defined as:



$$i \Pi_{\mu\nu}(p) \equiv i(p_\mu p_\nu - p^2 g_{\mu\nu}) \Pi(p^2)$$

where the form of $\Pi_{\mu\nu}(p)$ is governed by gauge invariance, i.e. it satisfies $p^\mu \Pi_{\mu\nu}(p) = p^\nu \Pi_{\mu\nu}(p) = 0$.

The renormalized propagator is the sum of a geometric series



$$\text{wavy line} + \text{wavy line with loop} + \text{wavy line with two loops} + \dots = \frac{-i(g_{\mu\nu} - \frac{p_\mu p_\nu}{p^2})}{p^2 [1 + \Pi(p^2)]}$$

The pole at $p^2 = 0$ is shifted to a non-zero value if:

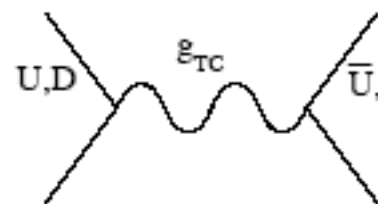
$$\Pi(p^2) \underset{p^2 \rightarrow 0}{\simeq} \frac{-g^2 v^2}{p^2}$$


“Eaten” Goldstone Boson

Then $p^2 [1 + \Pi(p^2)] = p^2 - g^2 v^2$, yielding a gauge boson mass of gv .

Technicolor

- Use scaled-up QCD to break electroweak symmetry

$$SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$$


$$\rightarrow \langle \bar{U}_L U_R \rangle = \langle \bar{D}_L D_R \rangle \neq 0$$

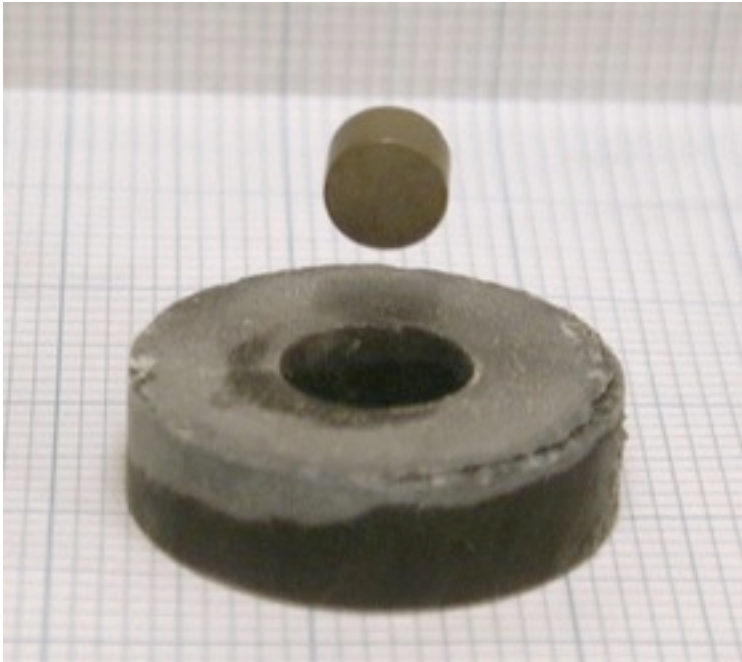
No hierarchy
problem!

$$\pi^\pm, \pi^0 \rightarrow W_L^\pm, Z_L$$

$$M_W = \frac{g F_{TC}}{2} \rightarrow F_{TC} \approx 250 \text{ GeV}$$

S?

(“Low-Energy” Analog)



$$\langle \phi^{--} \rangle \neq 0$$

“Abelian Higgs Model”



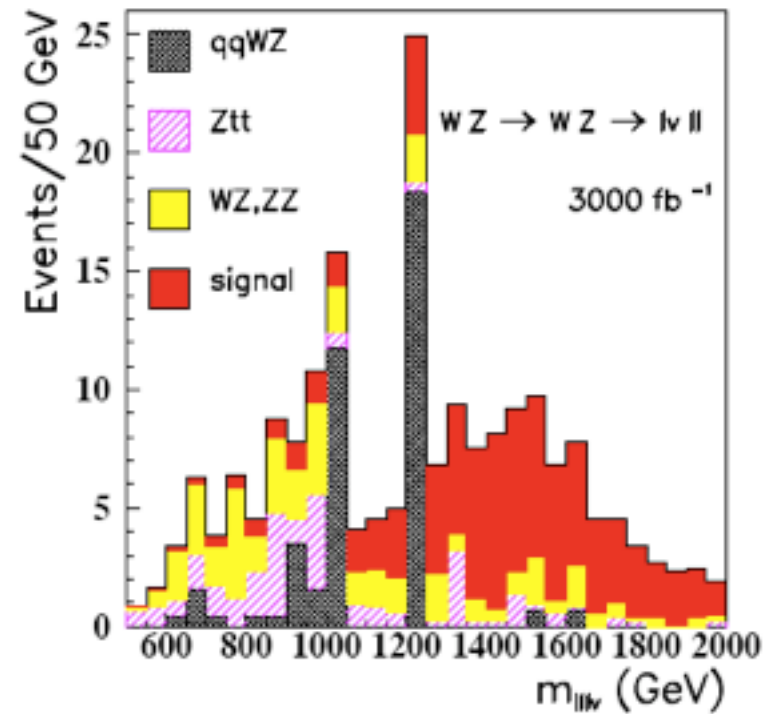
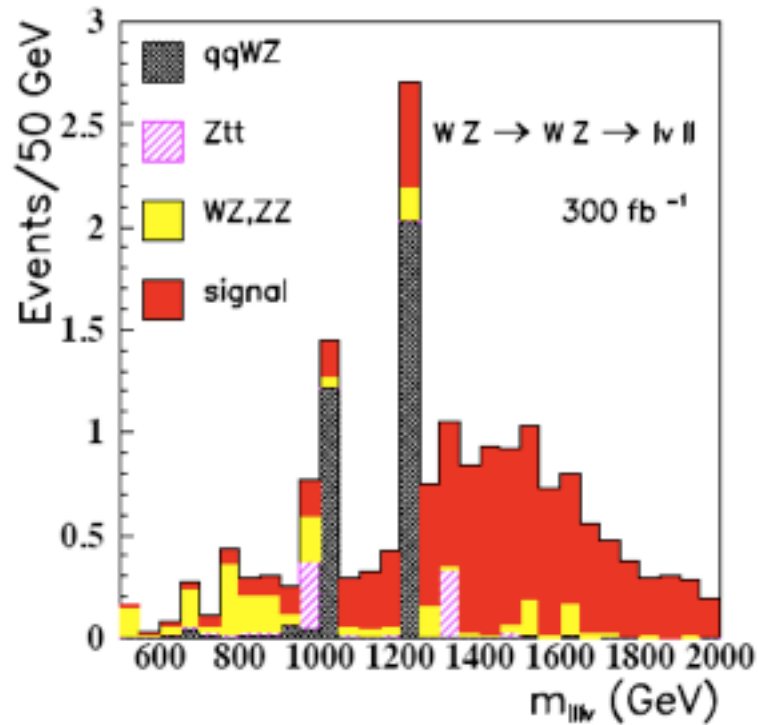
B

C

S

Weinberg: “Superconductivity for Particular Theorists”

WZ Scattering at SLHC



$$p_T(l_1) > 150 \text{ GeV}, \quad p_T(l_2) > 100 \text{ GeV}, \quad p_T(l_3) > 50 \text{ GeV}$$

$$|m(l_1 l_2) - m_Z| < 10 \text{ GeV}$$

$$E_T^{\text{miss}} > 75 \text{ GeV}$$

+ forward jets

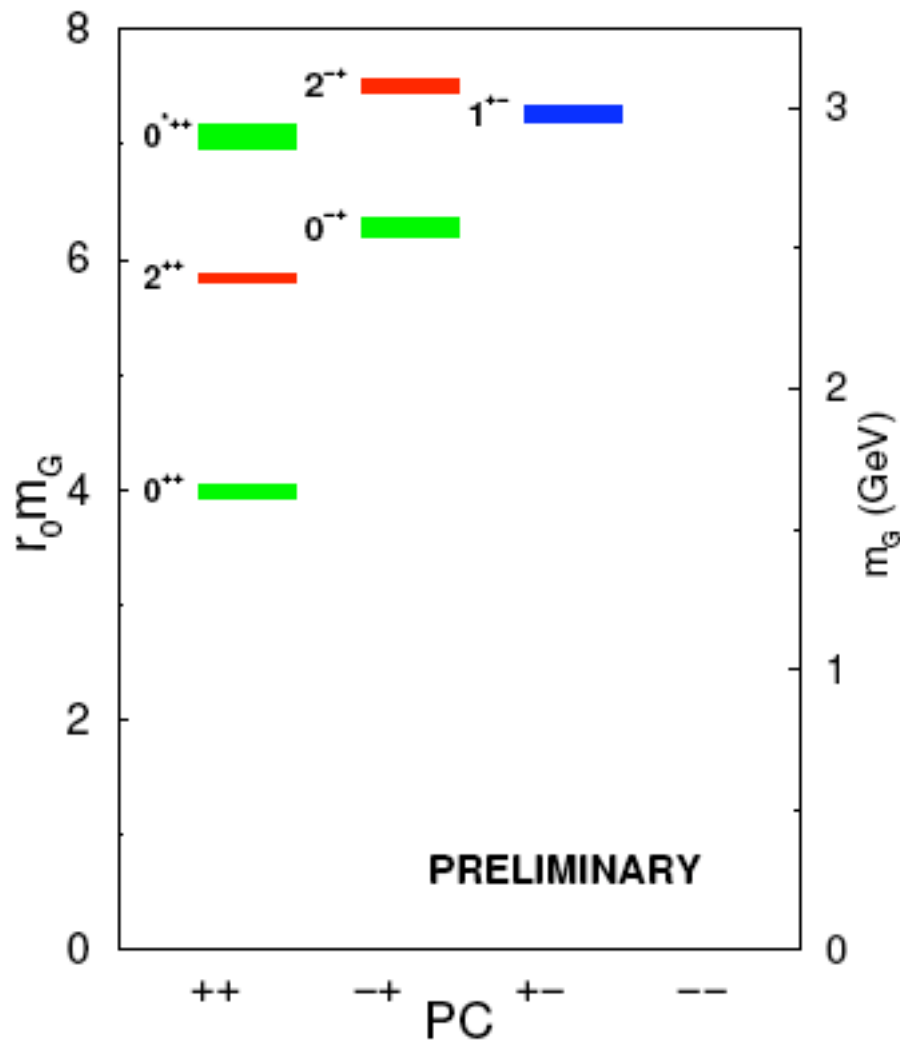
QCD

- Strong QCD dynamics breaks quark chiral symmetries ...
- ... yielding $\sim 99\%$ of the mass of ordinary matter ...
- and provides an example for models of dynamical electroweak symmetry breaking.
- QCD is the very model of a modern quantum field theory.



Extra Slides

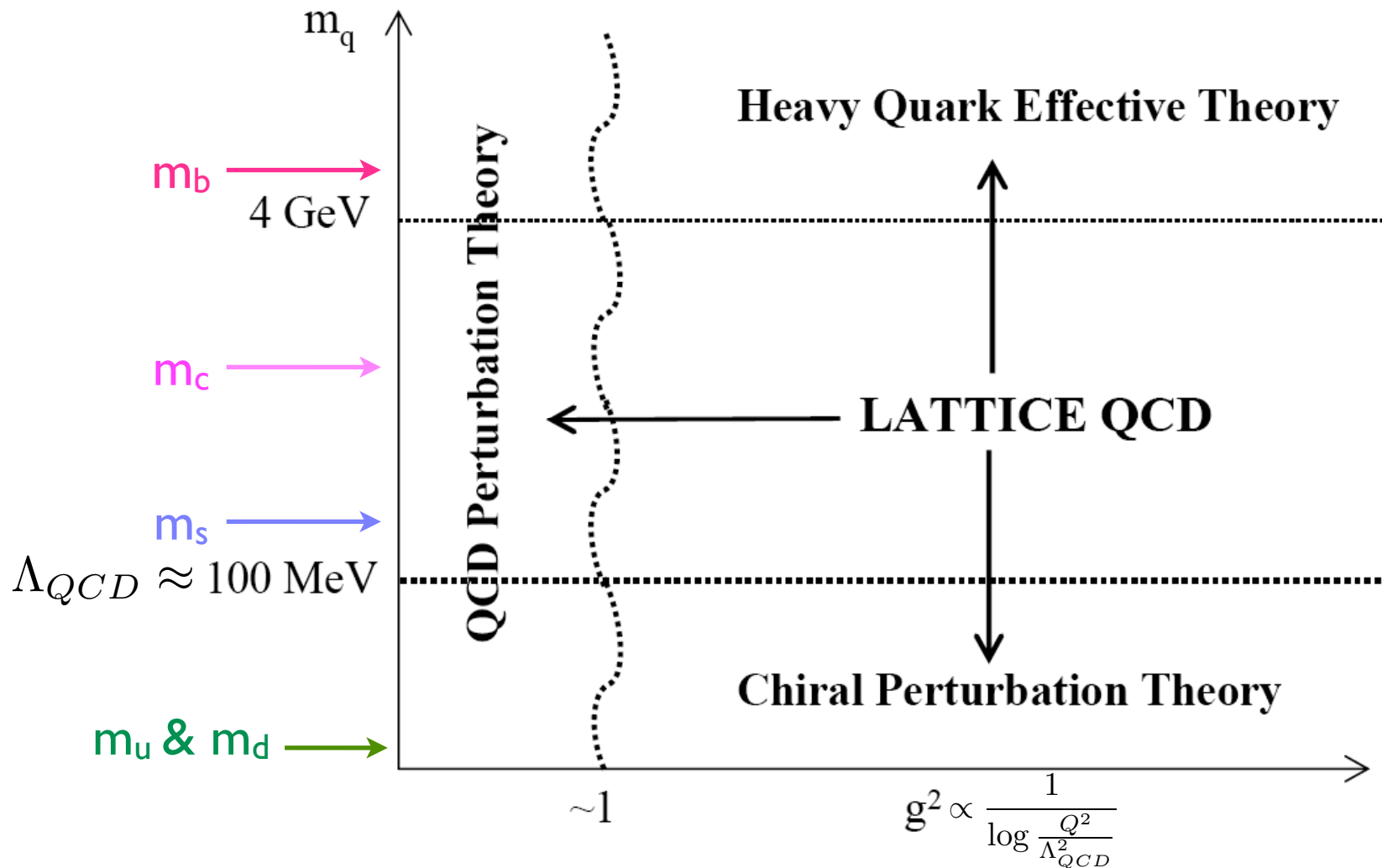
Mass-Gap in Pure Yang-Mills Theory



Glueball
Spectrum

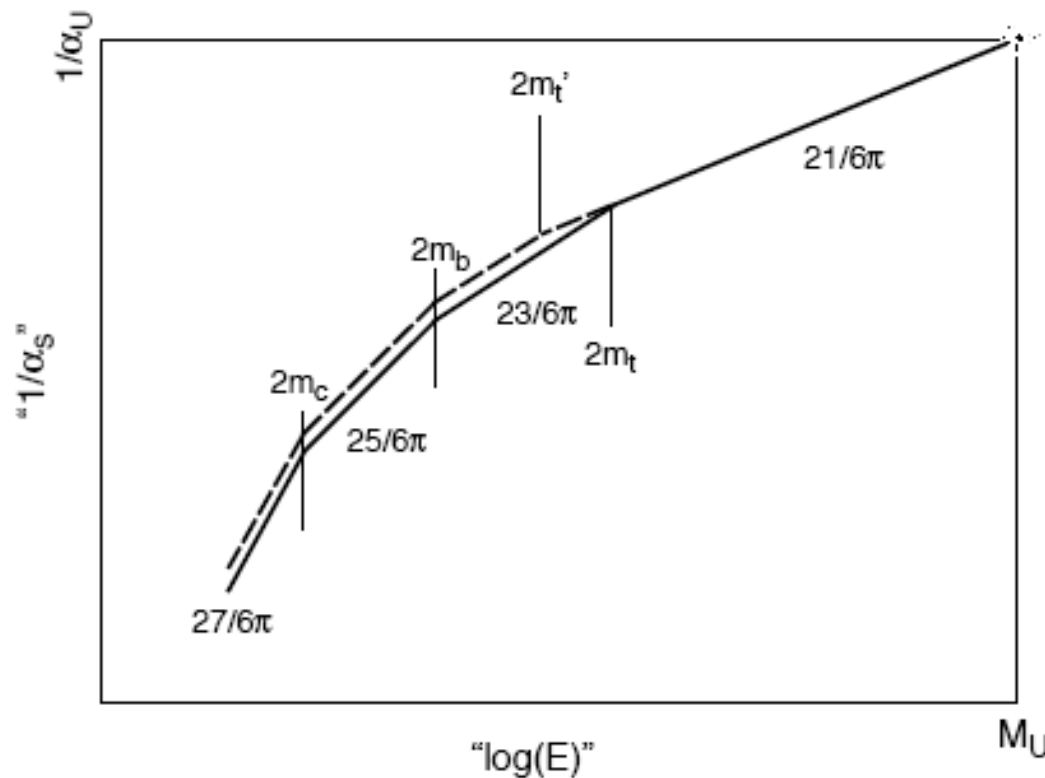
Morningstar and Peardon,
nucl-th/0309068

QCD: Modern Perspective



Rajan Gupta, LANL

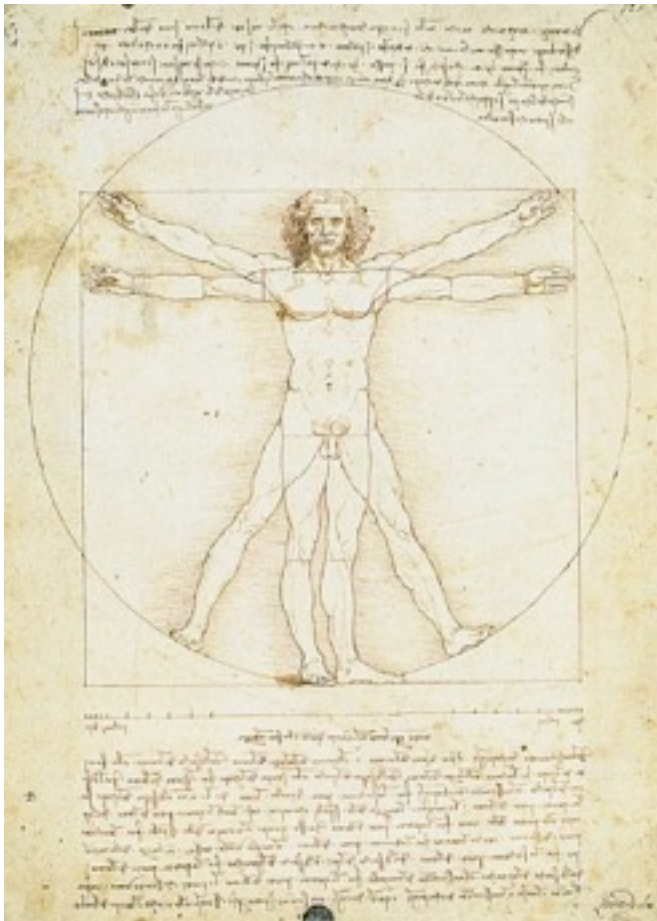
Top-Quark Matters



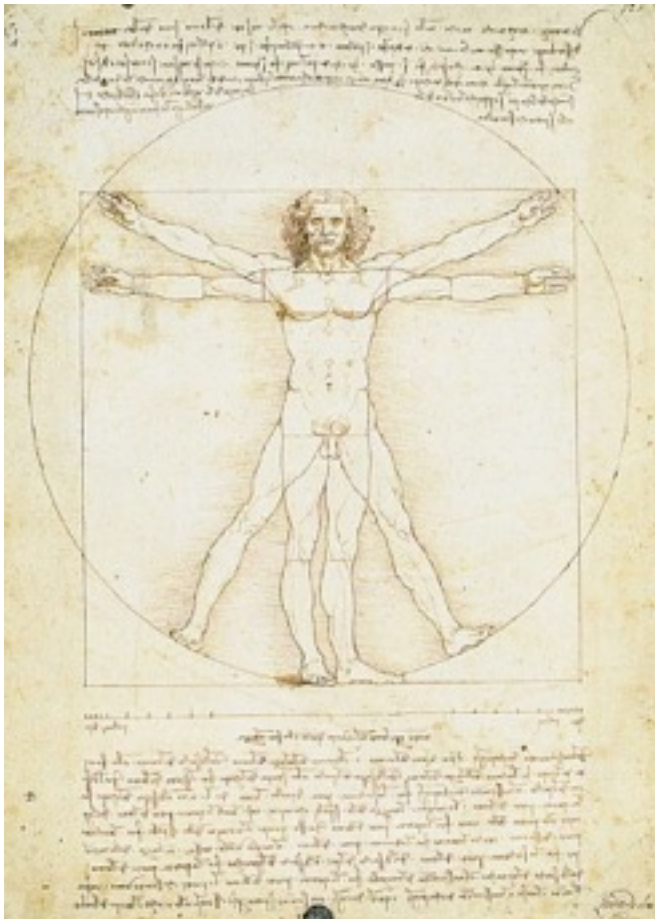
$$\frac{M_{\text{proton}}}{1 \text{ GeV}} \propto \left(\frac{m_t}{1 \text{ GeV}} \right)^{2/27}$$

Quigg, hep-ph/9707508

Anomalous Symmetry

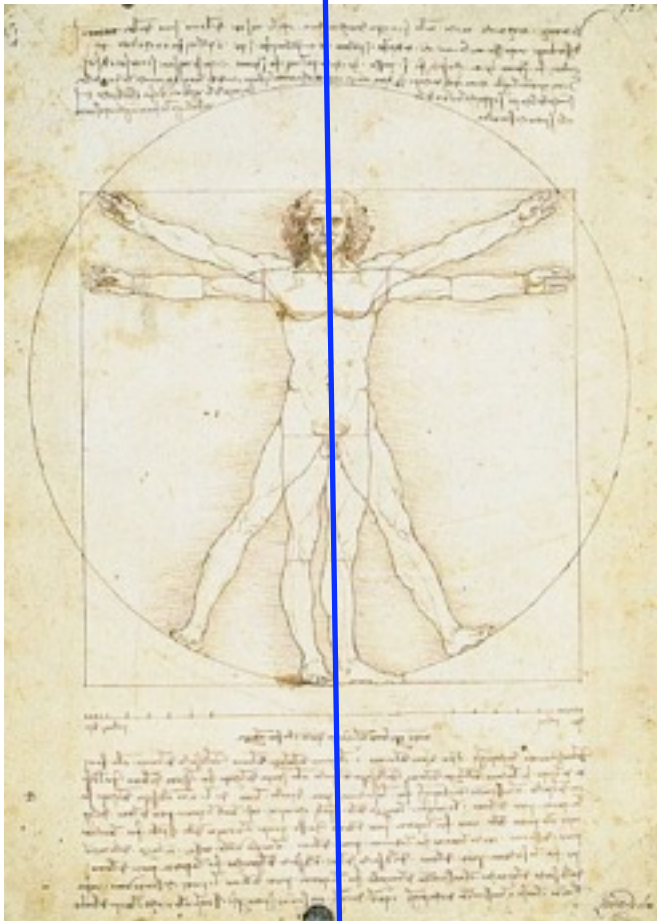


Anomalous Symmetry



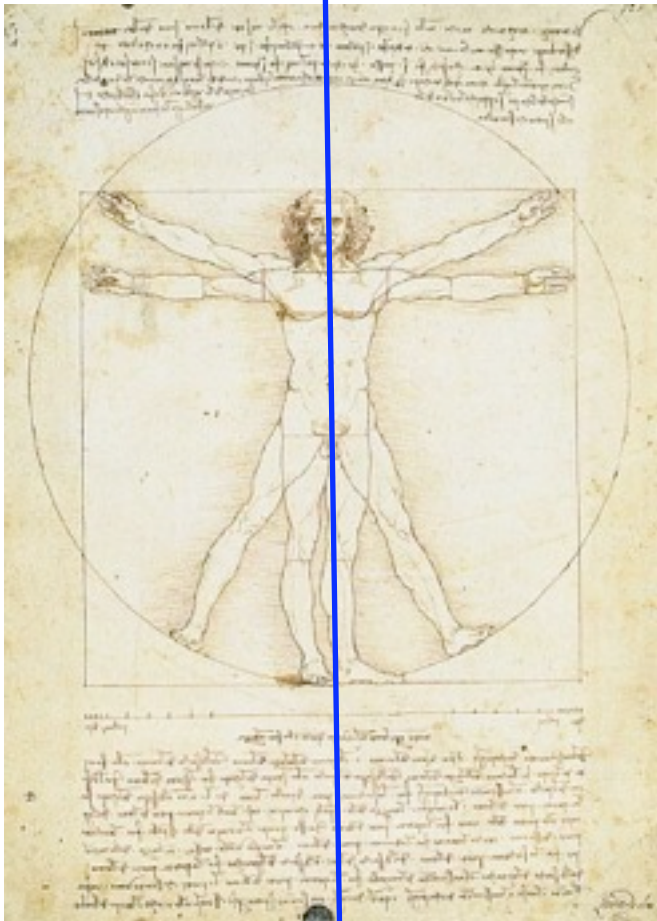
Sometimes,
quantum
dynamics
eliminates
the symmetries!

Anomalous Symmetry

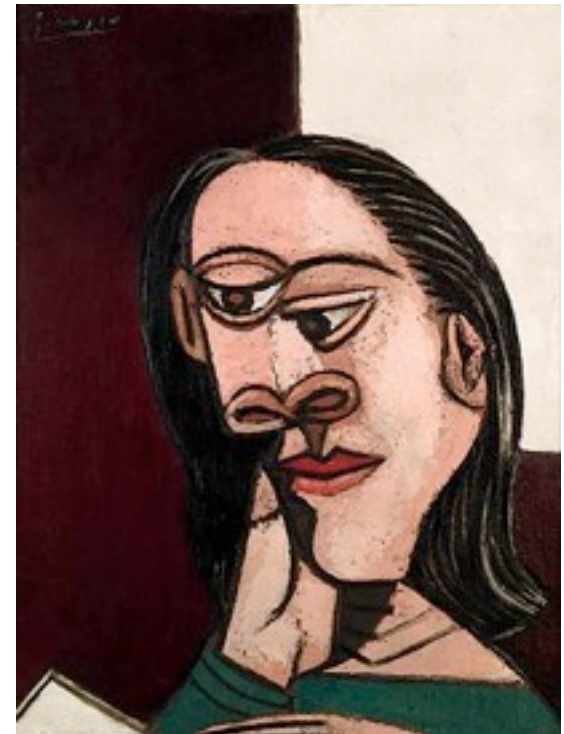


Sometimes,
quantum
dynamics
eliminates
the symmetries!

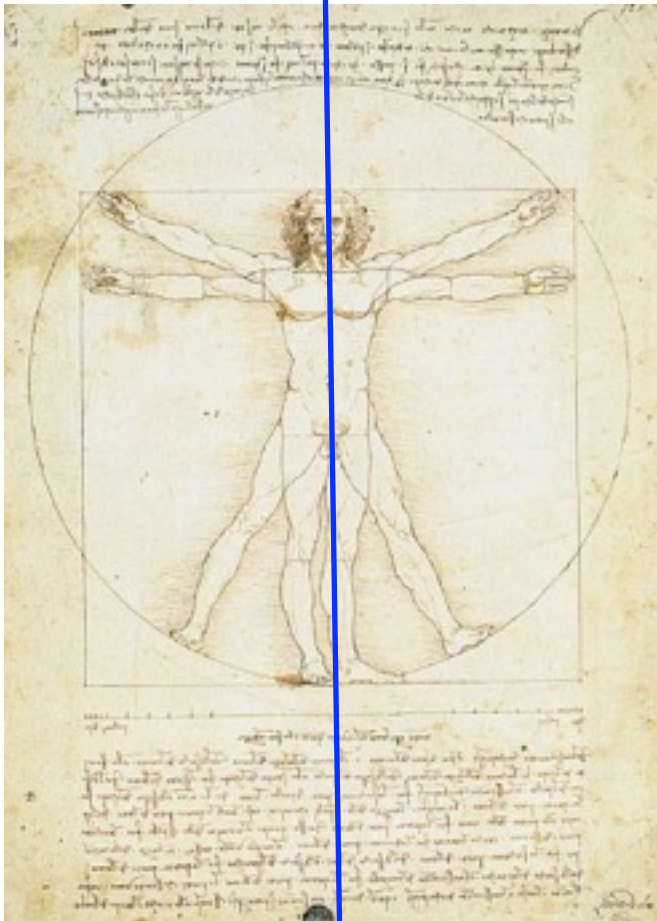
Anomalous Symmetry



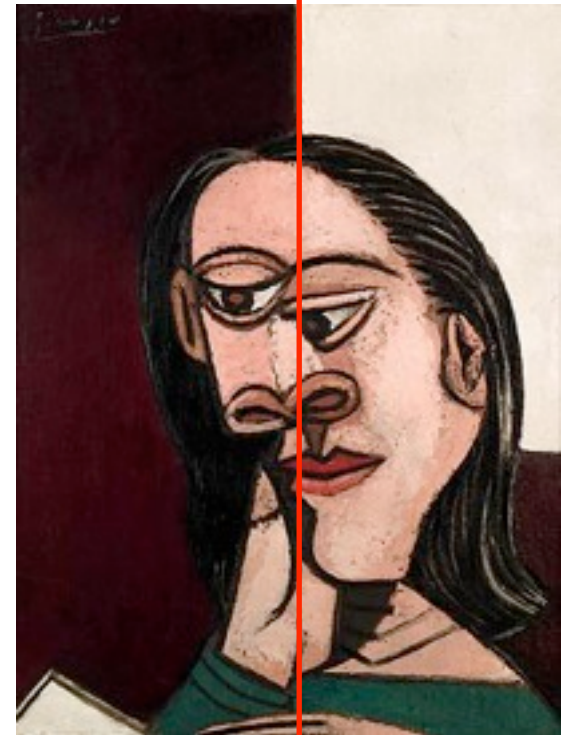
Sometimes,
quantum
dynamics
eliminates
the symmetries!



Anomalous Symmetry



Sometimes,
quantum
dynamics
eliminates
the symmetries!



QCD is Still Important

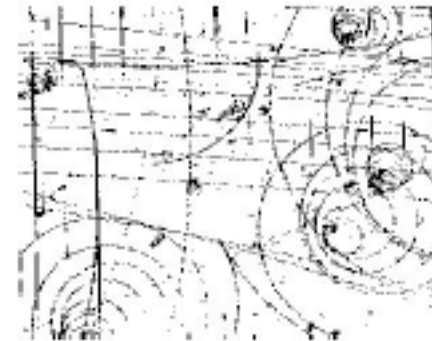
Clay Mathematics Institute

Dedicated to increasing and disseminating mathematical knowledge

Yang–Mills and Mass Gap

The laws of quantum physics stand to the world of elementary particles in the way that Newton's laws of classical mechanics stand to the macroscopic world. Almost half a century ago, Yang and Mills introduced a remarkable new framework to describe elementary particles using structures that also occur in geometry. Quantum Yang–Mills theory is now the foundation of most of elementary particle theory, and its predictions have been tested at many experimental laboratories, but its mathematical foundation is still unclear. The successful use of Yang–Mills theory to describe the strong interactions of elementary particles depends on a subtle quantum mechanical property called the "mass gap:" the quantum particles have positive masses, even though the classical waves travel at the speed of light. This property has been discovered by physicists from experiment and confirmed by computer simulations, but it still has not been understood from a theoretical point of view. Progress in establishing the existence of the Yang–Mills theory and a mass gap and will require the introduction of fundamental new ideas both in physics and in mathematics.

[The Millennium Problems](#)
[Official Problem Description](#)
[Lecture by Lorenzo Sadun at University of Texas \(video\)](#)



Especially at Tevatron/LHC...

Everything but
this, and possibly even
this, is QCD!

