CP-Violation or Nuclear Excitation?

The crucial role of neutrino-nucleus interaction modelling in neutrino oscillation measurements

Stephen Dolan

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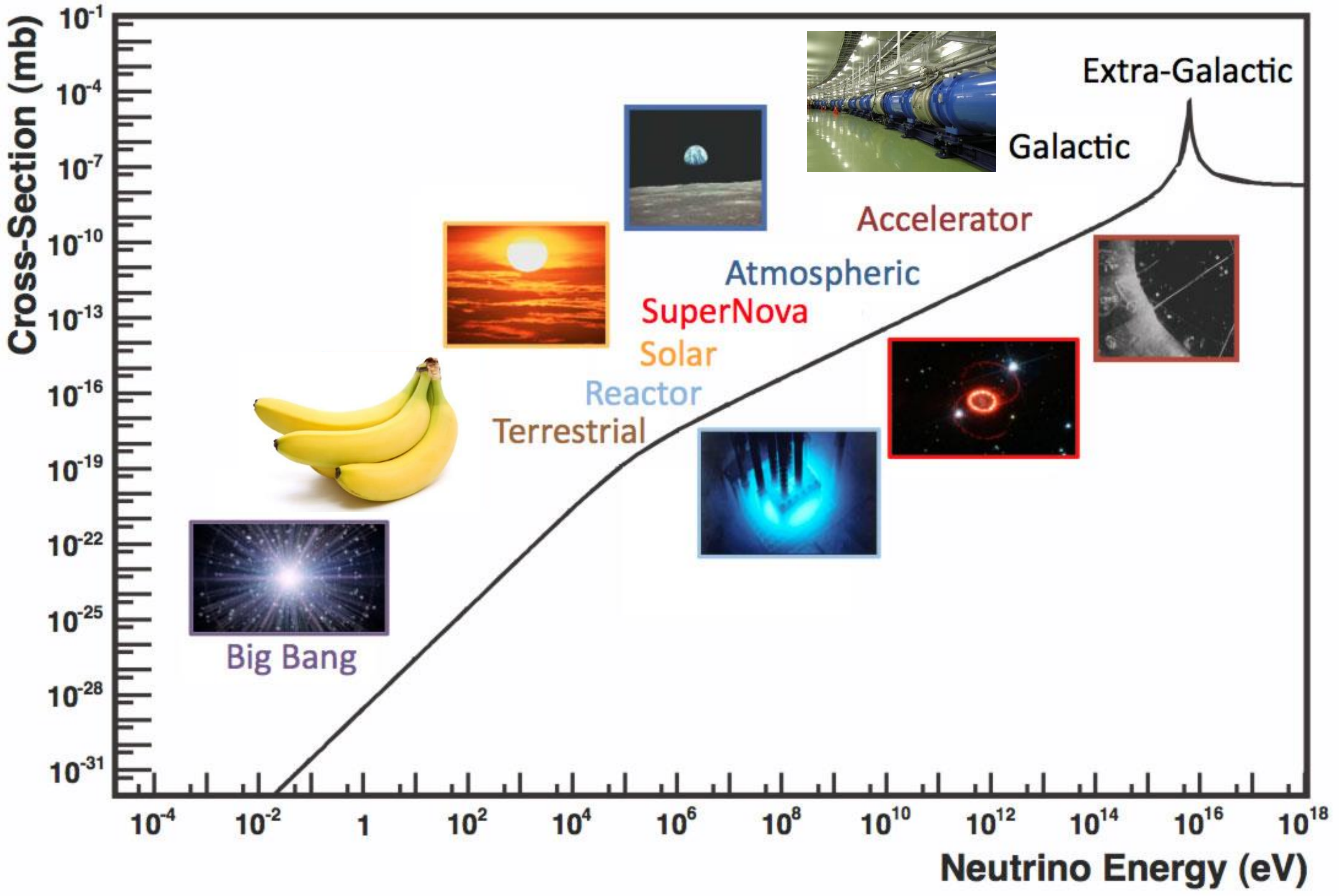
Overview

- **Neutrino Oscillations**
- Accelerator-Based Experiments
- $\boldsymbol{\nu}$ Interactions for $\boldsymbol{\nu}$ Oscillations
- Reconstructing Neutrino Energy
- The Path to Precision Measurements

Neutrino Sources



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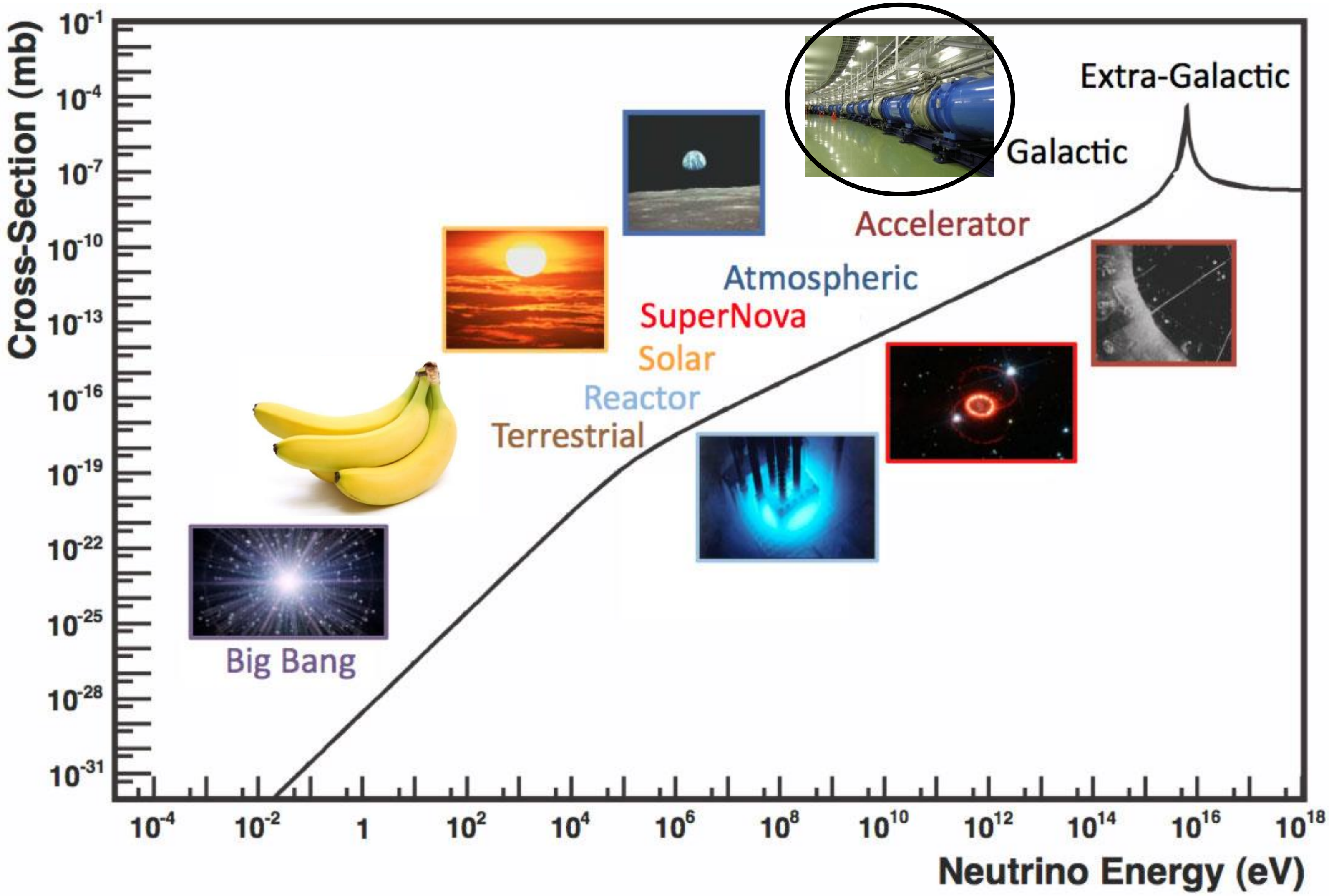
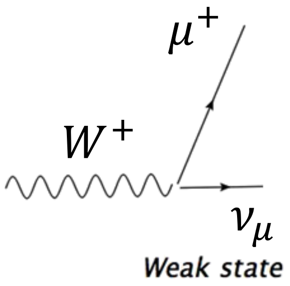


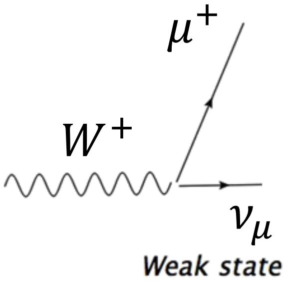
Neutrino Oscillations

- Neutrinos are **produced** in particular weak eigenstates $(\nu_e, \nu_\mu, \nu_\tau)$



Neutrino Oscillations

- Neutrinos are **produced** in particular weak eigenstates $(\nu_e, \nu_\mu, \nu_\tau)$
- These are linear combinations of mass eigenstates (ν_1, ν_2, ν_3) related by a unitary matrix, U_{PMNS}



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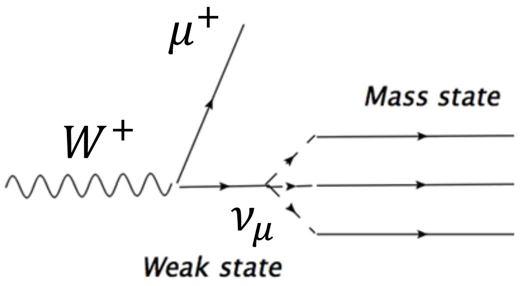

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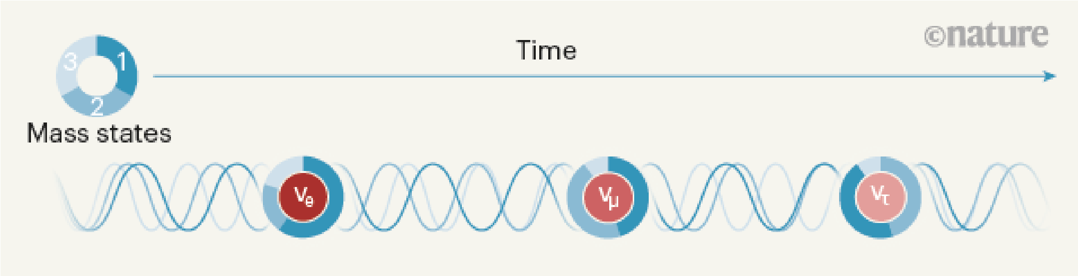

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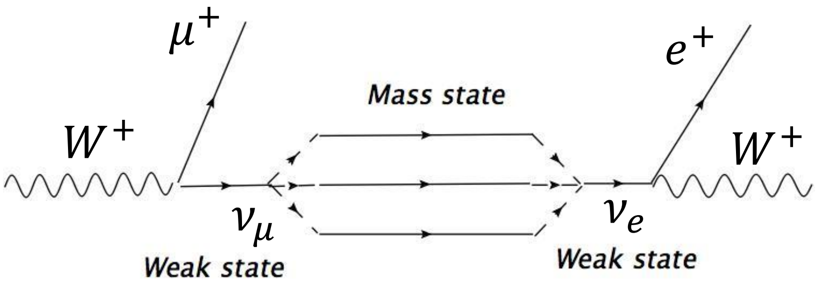

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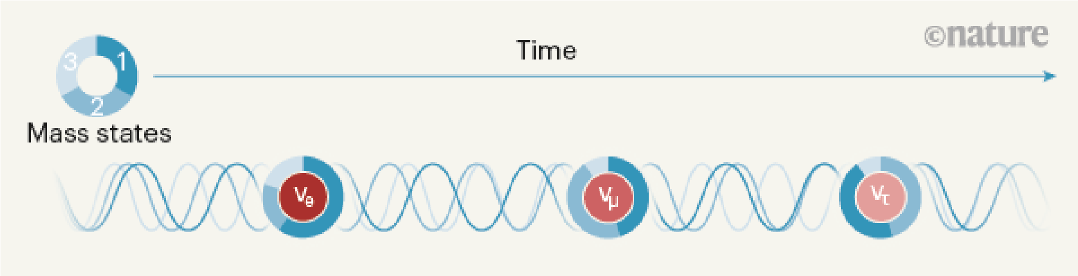

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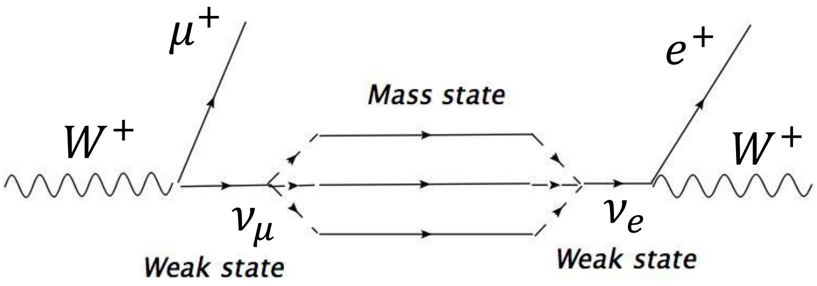

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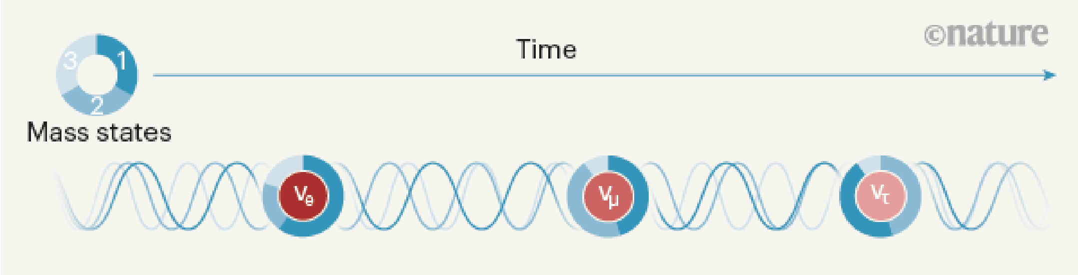
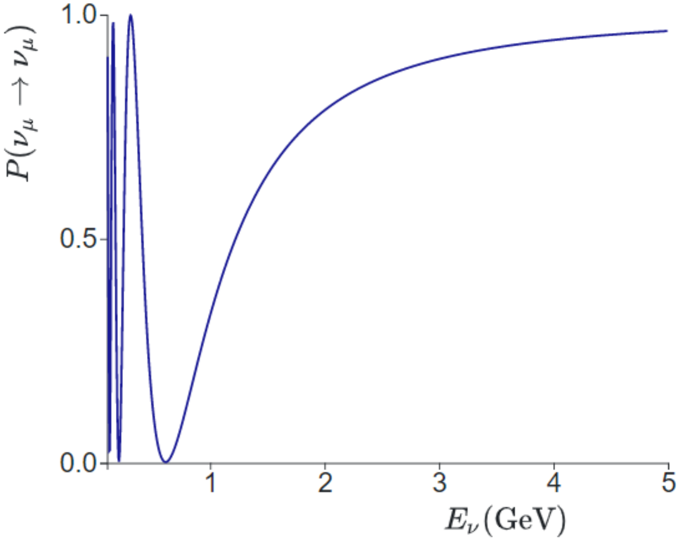
➡ The probability of finding a neutrino as a particular flavour "*oscillates*" as its mass states evolve



Neutrino Oscillations

- The oscillation probability depends on:
  - The neutrino energy



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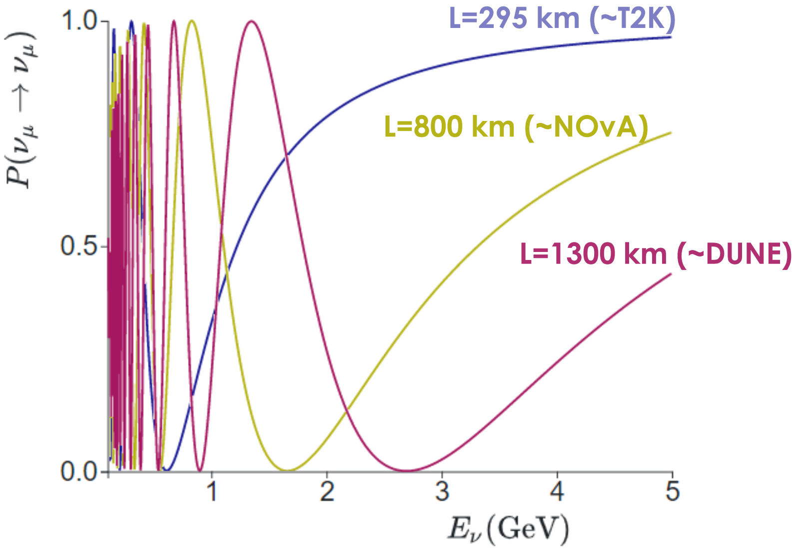

Neutrino Oscillations

- The oscillation probability depends on:
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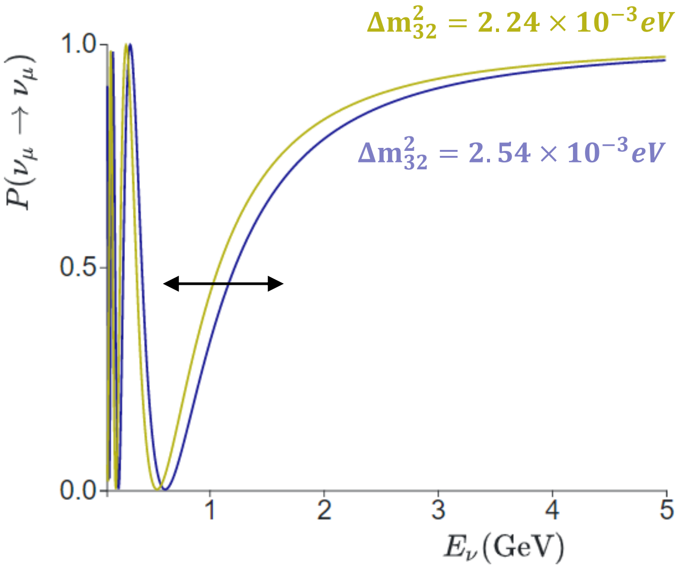
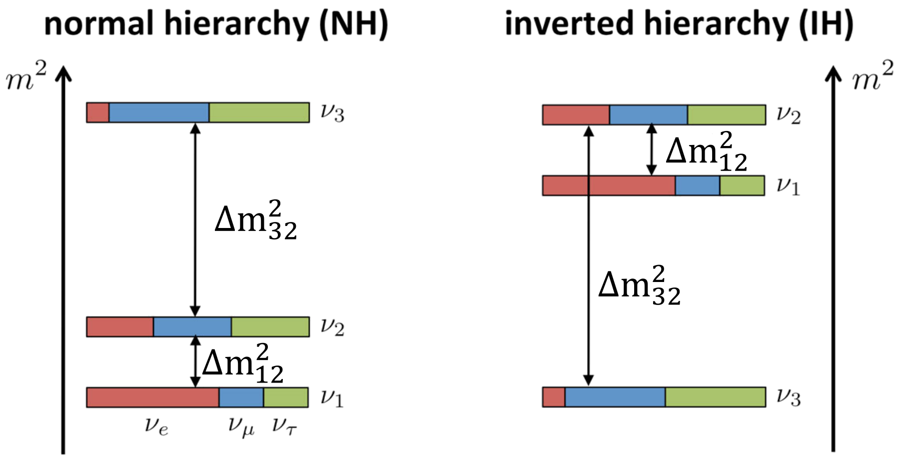
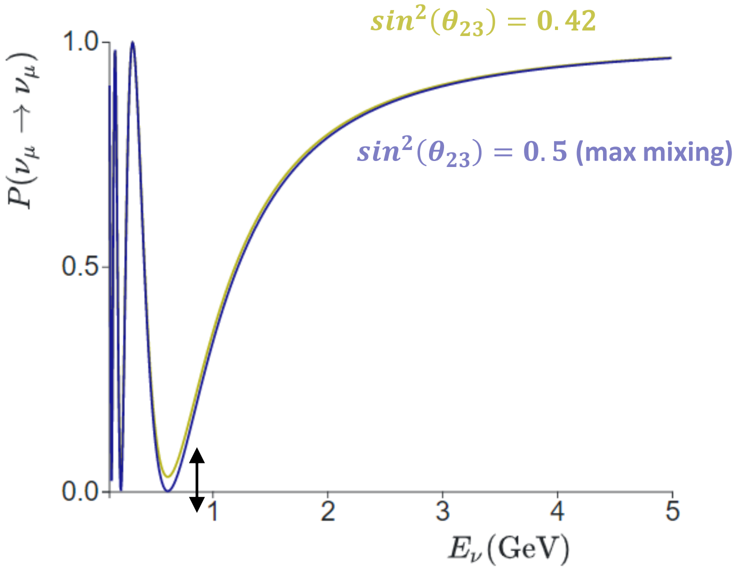




- Neutrino oscillations in a vacuum are sensitive only to the square of the mass splittings.
- "Matter effects" can give us the sign, but this is a challenging measurement.
- We don't yet know the right "hierarchy"

Neutrino Oscillations

- The oscillation probability depends on:
  - The neutrino energy
  - The travelled distance (“baseline”)
  - The difference in masses of ν_1, ν_2, ν_3
  - The PMNS mixing parameters



$$U = \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{+i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}$$

$s_{ij} = \sin \theta_{ij}$

$c_{ij} = \cos \theta_{ij}$

- Three mixing angles: $\theta_{12}, \theta_{13}, \theta_{23}$

Predominantly from KamLAND reactor neutrino experiment

From reactor experiments (e.g. Daya Bay) and from measuring $P(\nu_\mu \to \nu_e)$

Measuring $P(\nu_\mu \to \nu_\mu)$

Neutrino Oscillations

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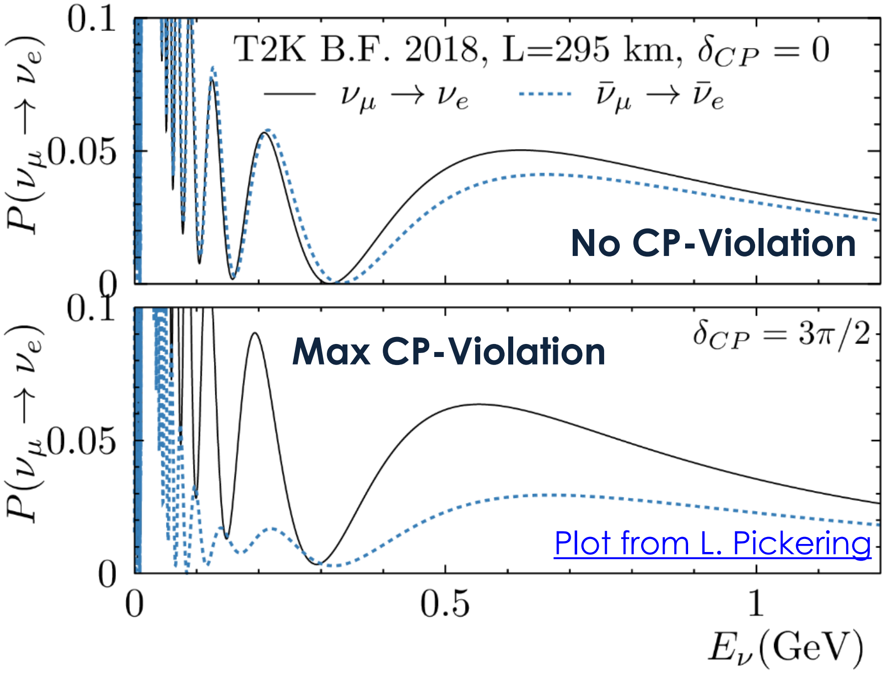

No CP-Violation

Max CP-Violation

Plot from L. Pickering

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$s_{ij} = \sin \theta_{ij}$
$c_{ij} = \cos \theta_{ij}$

- Three mixing angles: $\theta_{12}, \theta_{13}, \theta_{23}$
- One CP-Violating phase: δ_{CP}

Predominantly from KamLAND reactor neutrino experiment

From reactor experiments (e.g. Daya Bay) and from measuring $P(\nu_\mu \to \nu_e)$

Measuring $P(\nu_\mu \to \nu_\mu)$

Required to have a difference between neutrino and anti-neutrino vacuum oscillations

The Story So Far (2016)

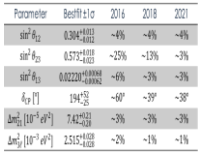
www.nu-fit.org

Parameter	Bestfit ±1σ	Precision
$\sin^2\theta_{12}$	$0.307^{+0.013}_{-0.012}$	$\sim 4\%$
$\sin^2\theta_{23}$	$0.574^{+0.026}_{-0.144}$	$\sim 25\%$
$\sin^2\theta_{13}$	$0.02217^{+0.0013}_{-0.0010}$	$\sim 6\%$
δ_{CP} [°]	272^{+61}_{-64}	$\sim 63°$
Δm^2_{21} $[10^{-5}\ eV^2]$	$7.49^{+0.19}_{-0.17}$	$\sim 3\%$
$\Delta m^2_{3\ell}$ $[10^{-3}\ eV^2]$	$2.484^{+0.045}_{-0.048}$	$\sim 2\%$

The Story So Far (2018)

Parameter	Bestfit ±1σ	2016	2018
$\sin^2\theta_{12}$	$0.307^{+0.013}_{-0.012}$	$\sim 4\%$	$\sim 4\%$
$\sin^2\theta_{23}$	$0.538^{+0.033}_{-0.069}$	$\sim 25\%$	$\sim 13\%$
$\sin^2\theta_{13}$	$0.02206^{+0.00075}_{-0.00075}$	$\sim 6\%$	$\sim 3\%$
δ_{CP} [°]	234^{+43}_{-31}	$\sim 63°$	$\sim 39°$
$\Delta m^2_{21}\ [10^{-5}\ eV^2]$	$7.40^{+0.21}_{-0.20}$	$\sim 3\%$	$\sim 3\%$
$\Delta m^2_{3\ell}\ [10^{-3}\ eV^2]$	$2.494^{+0.033}_{-0.031}$	$\sim 2\%$	$\sim 1\%$

The Story So Far (2021)

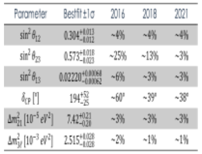


Parameter	Bestfit ±1σ	2016	2018	2021
$\sin^2\theta_{12}$	$0.304^{+0.013}_{-0.012}$	$\sim$4%	$\sim$4%	$\sim$4%
$\sin^2\theta_{23}$	$0.573^{+0.018}_{-0.023}$	$\sim$25%	$\sim$13%	$\sim$3%
$\sin^2\theta_{13}$	$0.02220^{+0.00068}_{-0.00062}$	$\sim$6%	$\sim$3%	$\sim$3%
δ_{CP} [°]	194^{+52}_{-25}	$\sim$60°	$\sim$39°	$\sim$38°
Δm^2_{21} $[10^{-5}\ eV^2]$	$7.42^{+0.21}_{-0.20}$	$\sim$3%	$\sim$3%	$\sim$3%
$\Delta m^2_{3\ell}$ $[10^{-3}\ eV^2]$	$2.515^{+0.028}_{-0.028}$	$\sim$2%	$\sim$1%	$\sim$1%

Precision neutrino-oscillation physics!

Nature **580**, 339-344

The Story So Far (2021)



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- Maximal θ_{23} mixing? (flavour symmetries?)

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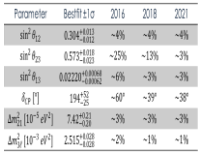

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- A new source of CP-violation? (implications for cosmology and leptogensis)

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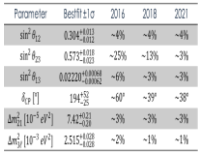

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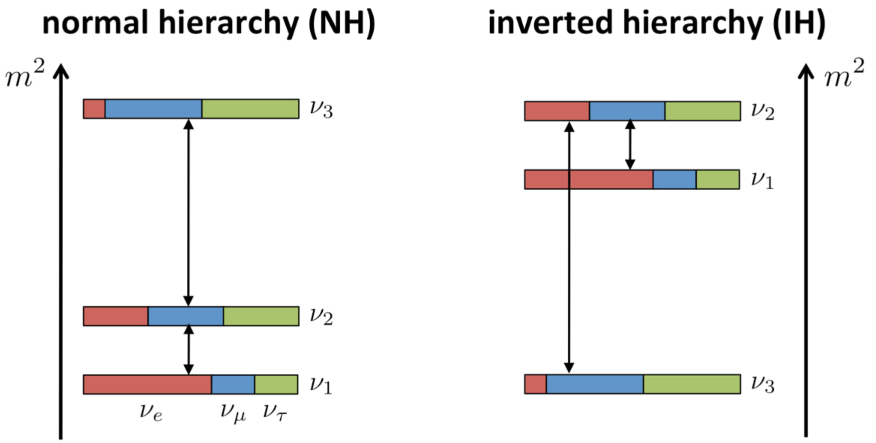

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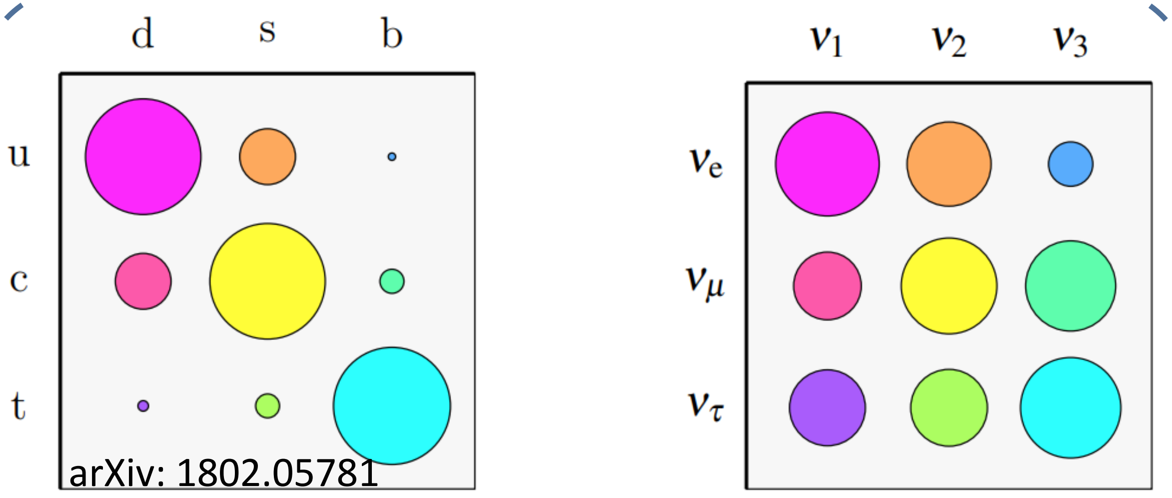
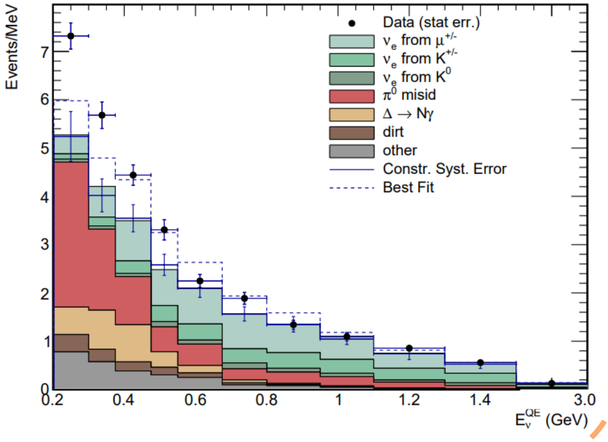


(a) Quark Mixing Elements.

(b) Lepton Mixing Elements.



Precision neutrino-oscillation physics!

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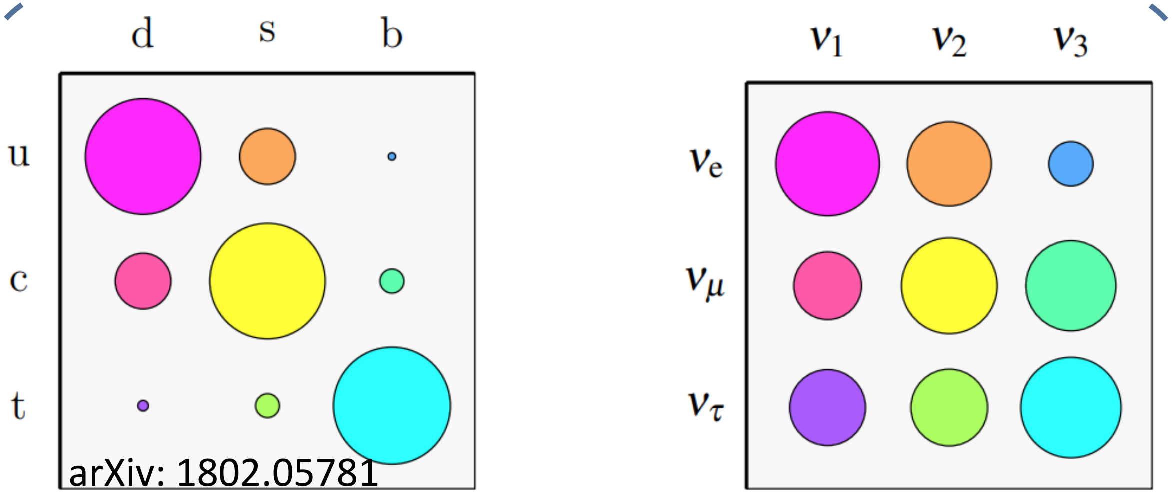
- A new source of CP-violation?
- What's the neutrino mass ordering?
- Why are CKM and PMNS mixing so different?
- Are there only three flavours? (are there sterile neutrinos?)

The Story So Far (2021)



arXiv: 1802.05781

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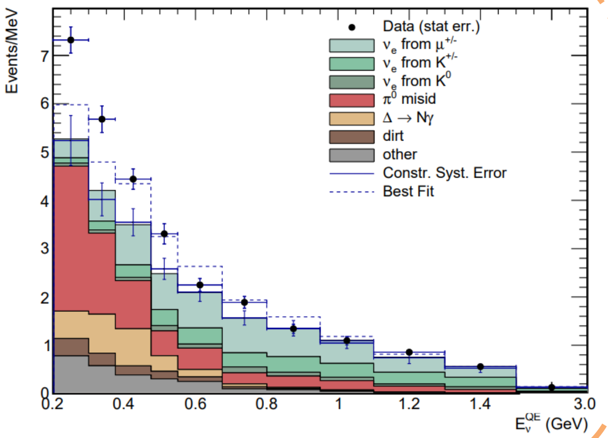

Facilities for exploring physics beyond the standard model + PMNS

- The next generation of experiments will offer unprecedented precision (10-50 times more statistics for the long-baseline program)
- Opportunities to see new physics feeding down to create deviations from PMNS behaviour (e.g. "NSIs")
- A complementary approach to pushing back the frontiers of particle physics

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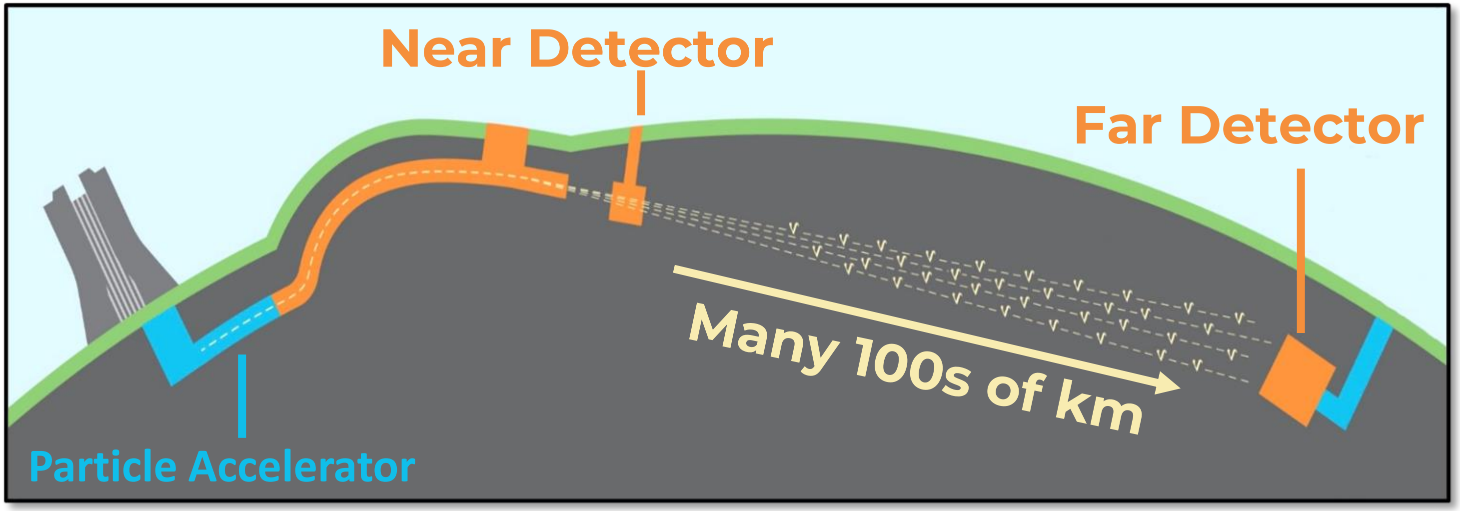

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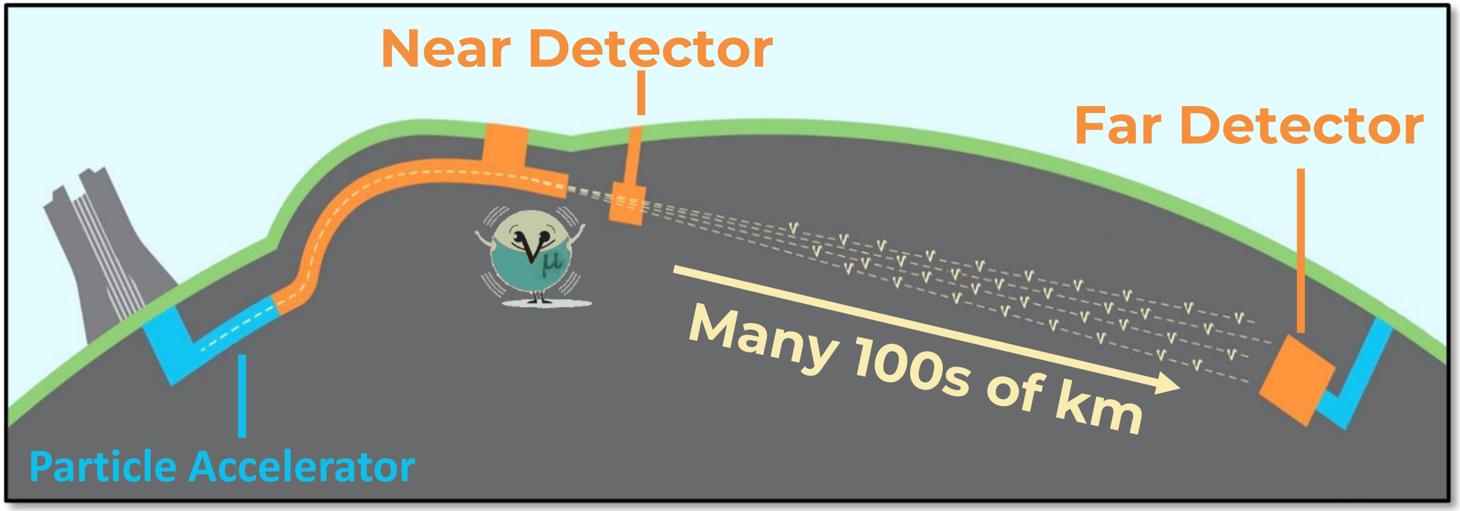

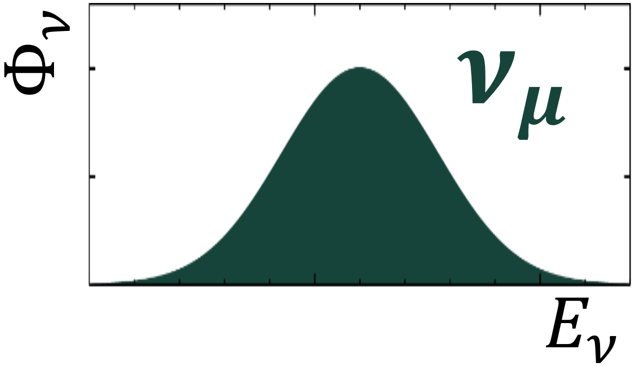

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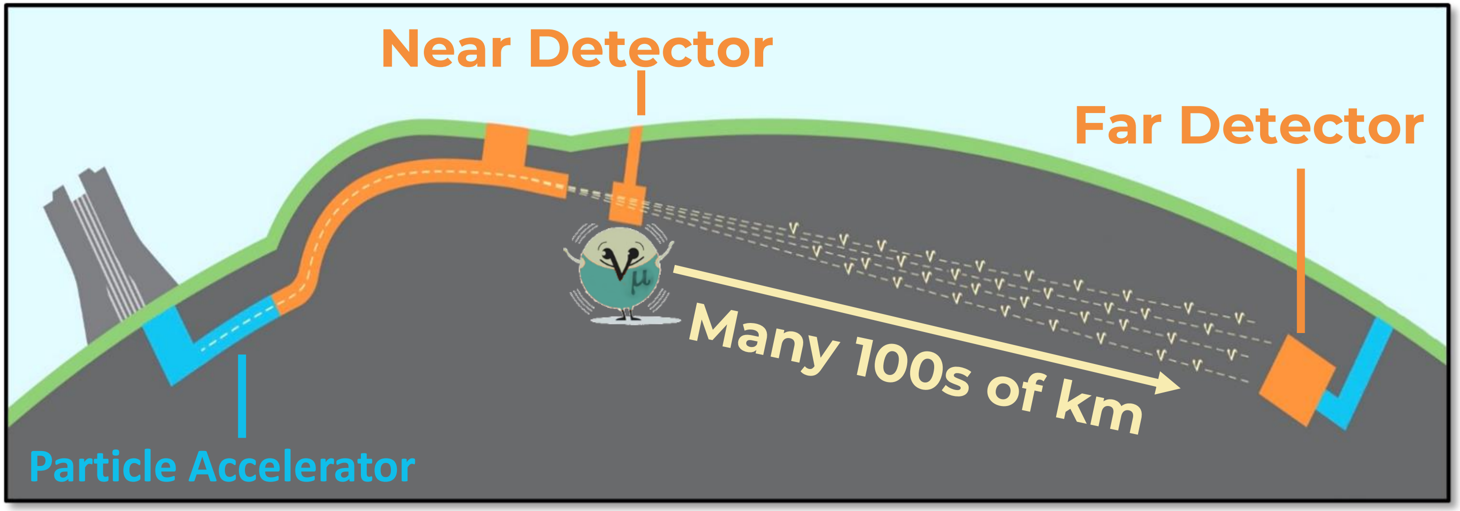
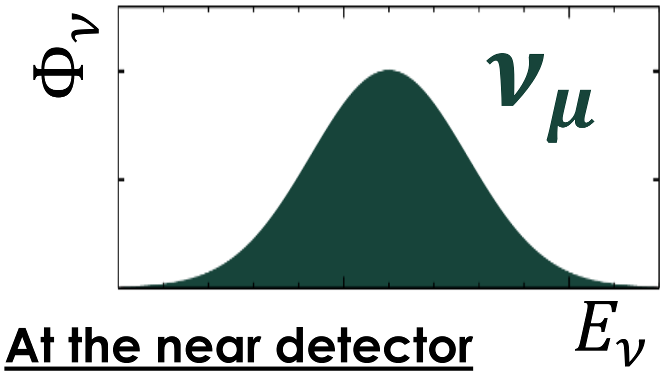




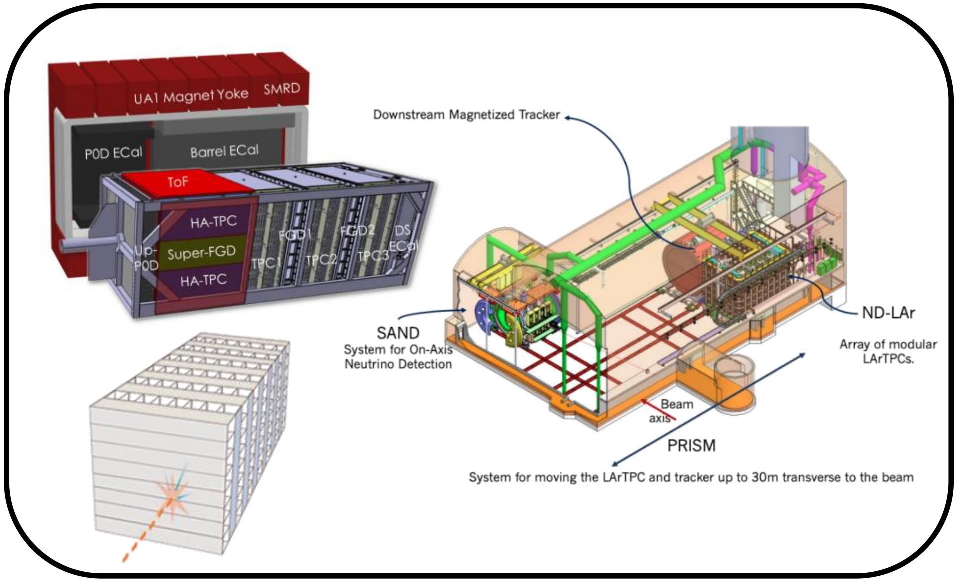
At the near detector

$$N_\mu(E_\nu) = \sigma(E_\nu)\Phi_\nu(E_\nu)\epsilon(E_\nu)$$

Interaction cross section

Neutrino flux

Detector effects



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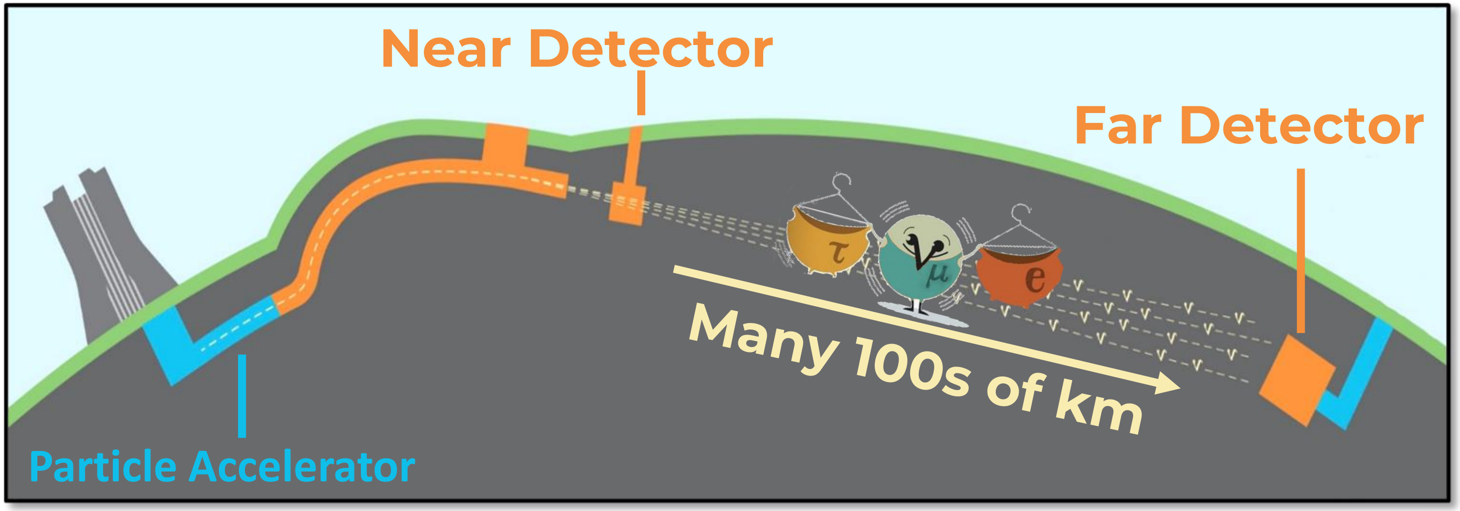

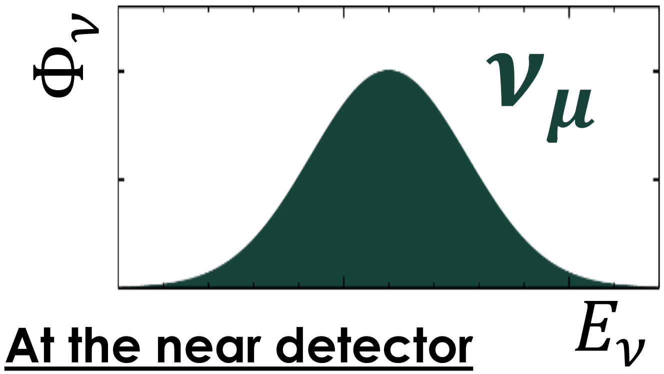

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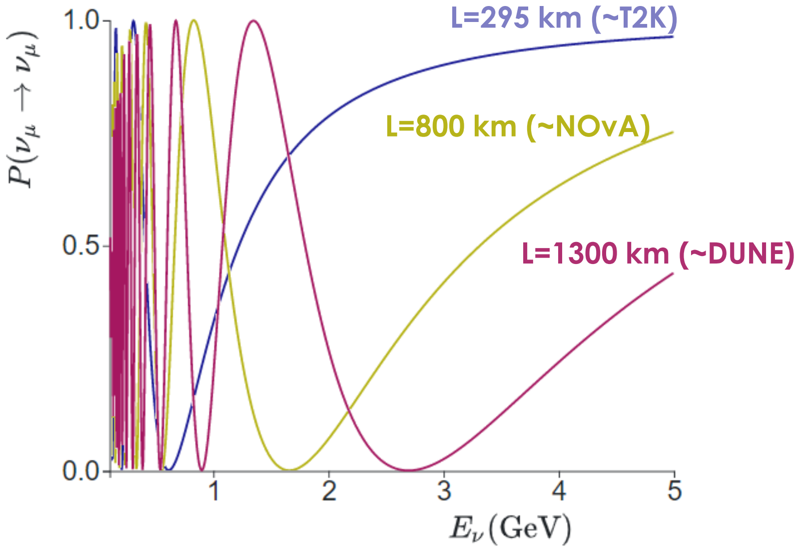

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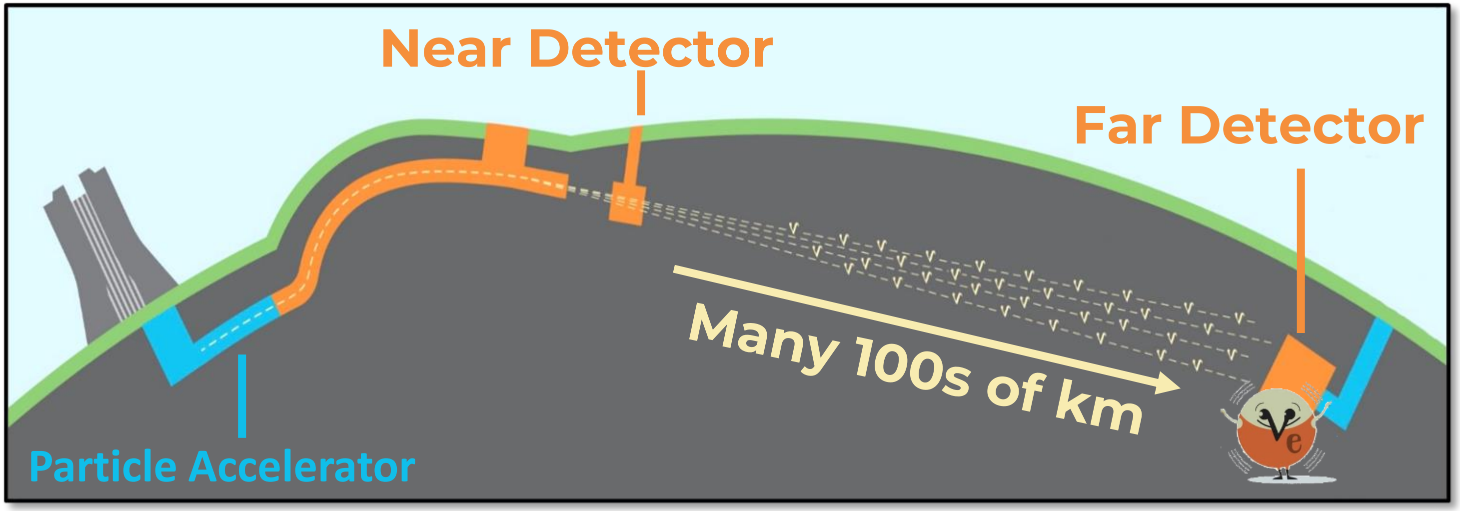

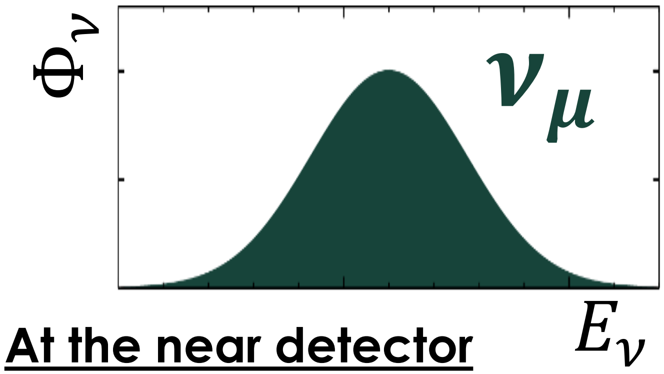

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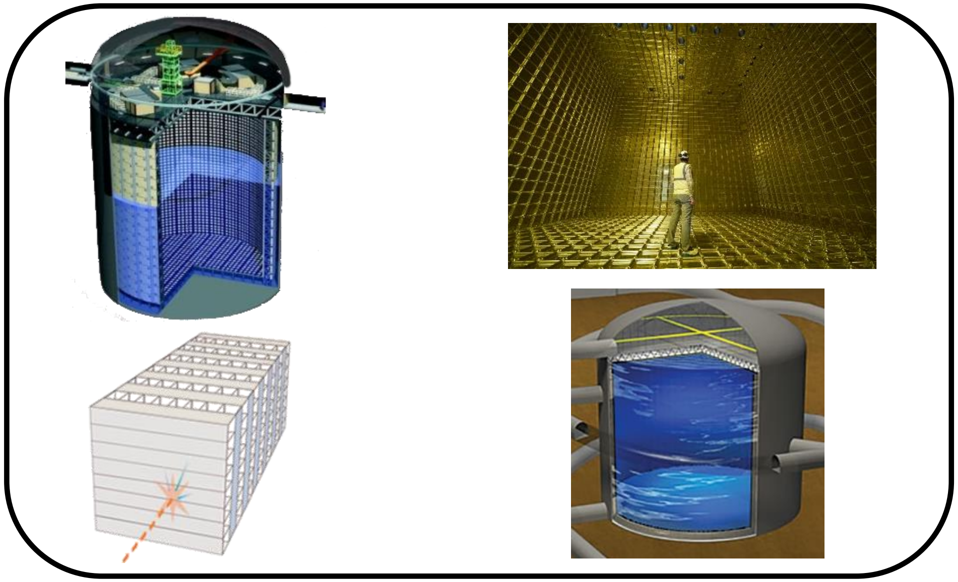
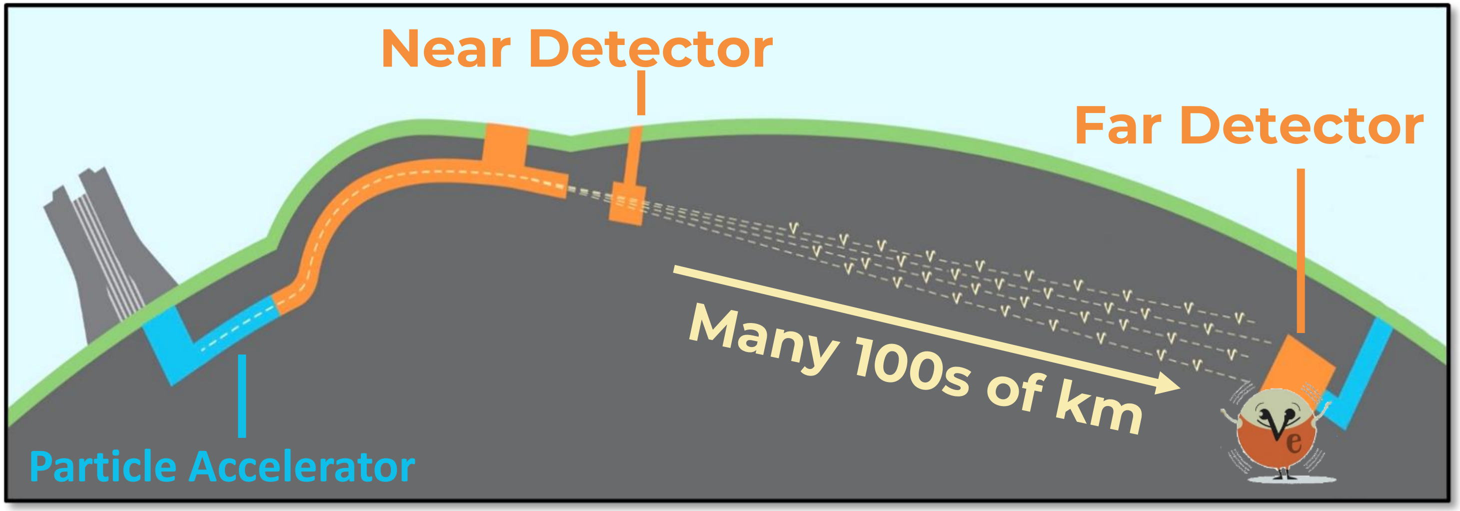
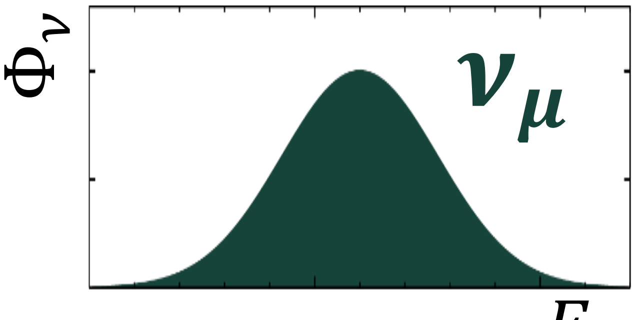
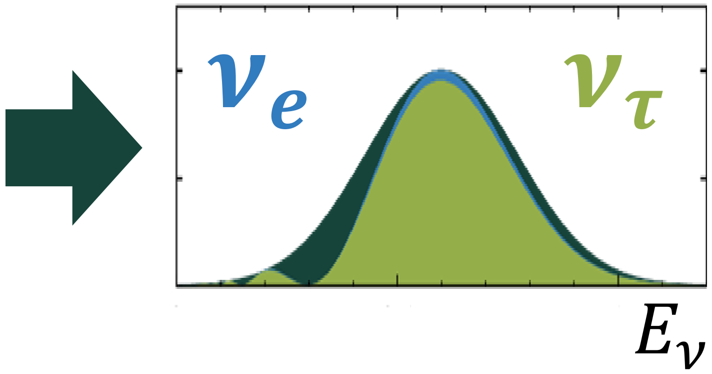


Accelerator-Based Experiments







At the near detector

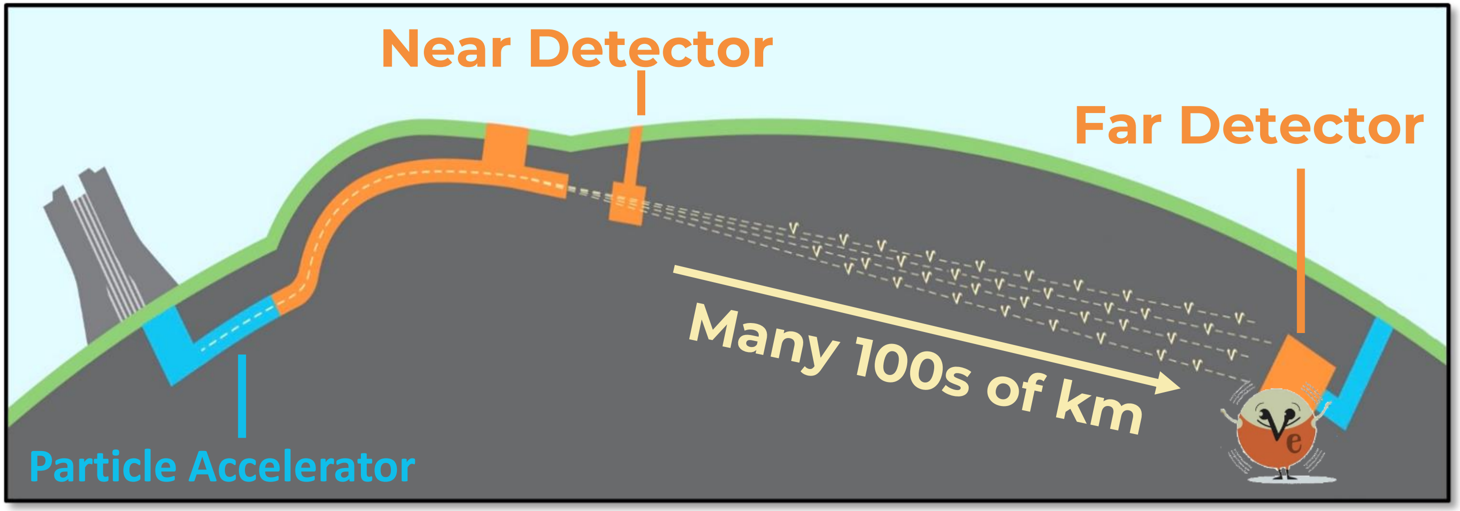
$$N_\mu(E_\nu) = \sigma(E_\nu)\Phi_\nu(E_\nu)\epsilon(E_\nu)$$

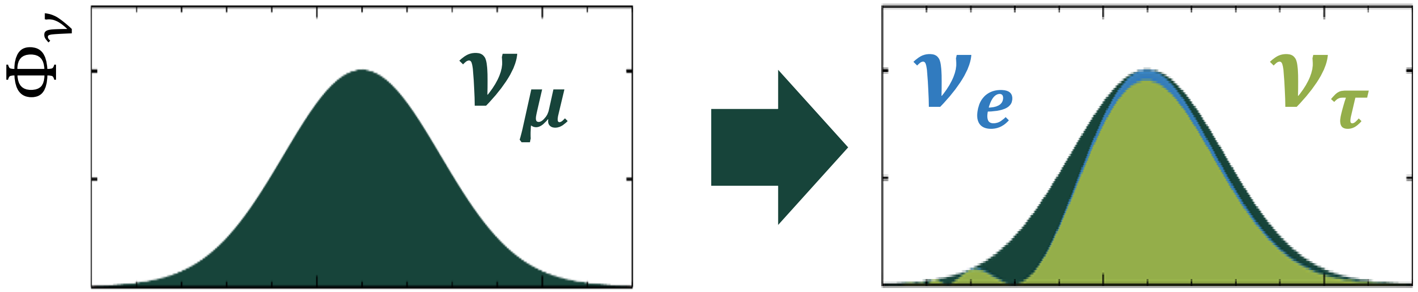
Interaction cross section

Neutrino flux

Detector effects

Accelerator-Based Experiments





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Oscillation probability

Accelerator-Based Experiments

Current long-baseline experiments

	T2K	NOvA
Baseline	295 km	800 km
N_μ^{rec} (ν-mode)	318	211
N_μ^{rec} ($\bar{\nu}$-mode)	137	105
N_e^{rec} ($\nu-$mode)	94	82
N_e^{rec} ($\bar{\nu}$-mode)	16	33

Reconstructed events in samples at the experiment's far detectors

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Current systematic uncertainties

Phys. Rev. D **98**, 032012

Source (NOvA)	$N(\nu_e)$
$\sigma_{\nu N}$ and FSI	7.7%
Total Syst.	**9.2%**

NEUTRINO 2022
XXX International Conference on Neutrino Physics and Astrophysics

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Accelerator-Based Experiments

Future long-baseline experiments

Coming 2027-2032

	HYPER K *arXiv:1805.04163*	DUNE *arXiv:2002.03005*
Baseline	295 km	1300 km
N_μ^{rec} (ν-mode)	~10000	~7000
N_μ^{rec} ($\bar{\nu}$-mode)	~14000	~3500
N_e^{rec} ($\nu-$mode)	~2000	~1500
N_e^{rec} ($\bar{\nu}$-mode)	~2000	~500

Approximate late-stage projections for reconstructed events in samples at the experiment's far detectors

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2002.03005

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Approximate late-stage projections for reconstructed events in samples at the experiment's far detectors

Crucial to reduce uncertainties related to neutrino interaction cross sections

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Overview

- Neutrino Oscillations
- Accelerator-Based Experiments
- **ν Interactions for ν Oscillations**
- Reconstructing Neutrino Energy
- The Path to Precision Measurements

Where are we so far?

- Current neutrino oscillation experiments are mostly **statistics limited**
- Systematic uncertainties related to neutrino-nucleus interactions are often dominant and **are unacceptably large for the next generation of experiments**
- Key questions:
  1. *Why is modelling neutrino interactions so difficult?*

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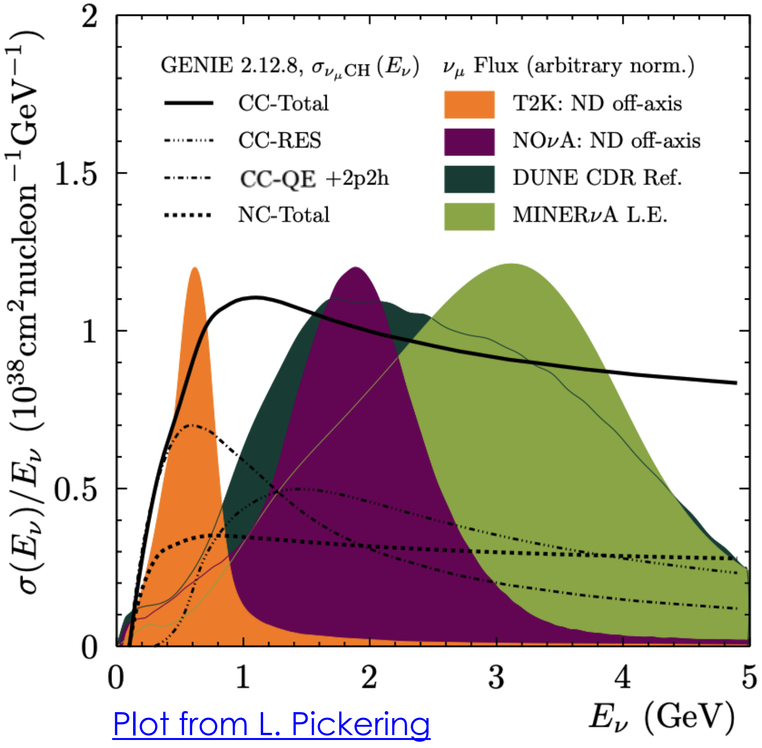
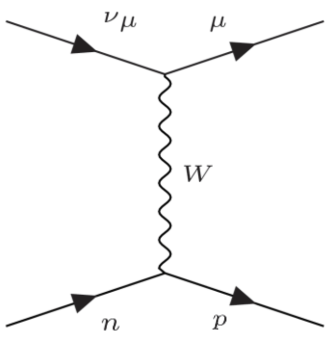
Neutrino-nucleus interactions



Plot from L. Pickering

CC-QE

(Charged-Current Quasi-Elastic)



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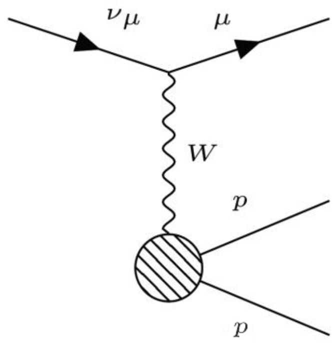
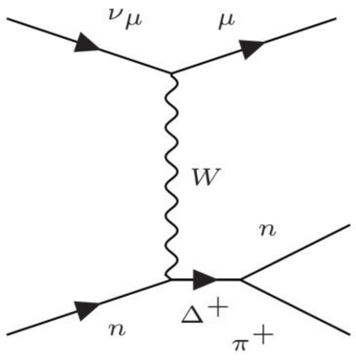


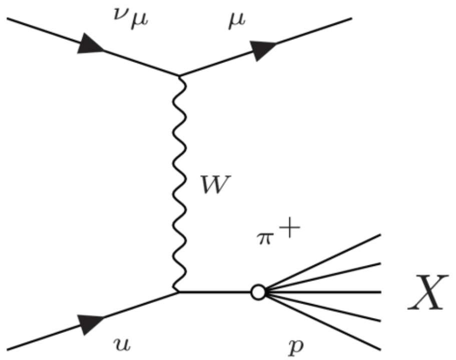
CC-SPP

(Single Pion Production)



CC-DIS

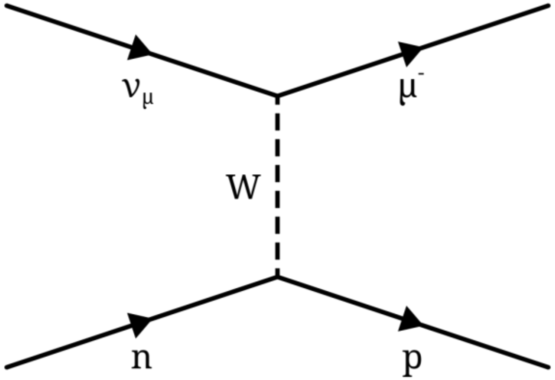
(Deep Inelastic Scattering)



Neutrino-nucleon scattering

- Even the most simple “CCQE” interaction is hard to describe as the target in an extended object

$$M \sim \frac{g_W^2}{8}\frac{1}{M_W^2}\left[\bar{u}_\mu \gamma_\mu (1-\gamma_5) u_\nu\right]\left[\bar{u}_p \underset{\overset{\uparrow}{???}}{(\ldots)} u_n\right]$$



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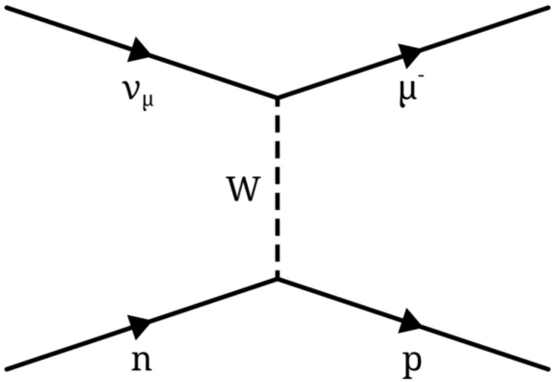

$$M \sim \frac{g_W^2}{8}\frac{1}{M_W^2}[\bar{u}_\mu \gamma_\mu (1-\gamma_5) u_\nu][\bar{u}_p(\dots)u_n]$$

↑

$$J_H^\beta = \bar{u}_p \left[f_{1V}\gamma^\beta + i\frac{\xi f_{2V}}{2M}\sigma^{\beta\delta}q_\delta + \frac{f_{3V}}{M}q^\beta + f_A\gamma^\beta\gamma_5 + \frac{f_p}{M}q^\beta\gamma_5 + \frac{f_{3A}}{M}\left(P_p^\beta + P_n^\beta\right)\gamma_5 \right] u_n$$

$$M = \left(M_p + M_n\right)/2 \qquad q = p_\nu - p_\mu = P_p - P_n \qquad \xi = \mu_p - \mu_n \qquad \sigma^{\mu\nu} = \frac{i}{2}\left[\gamma^\mu, \gamma^\nu\right]$$

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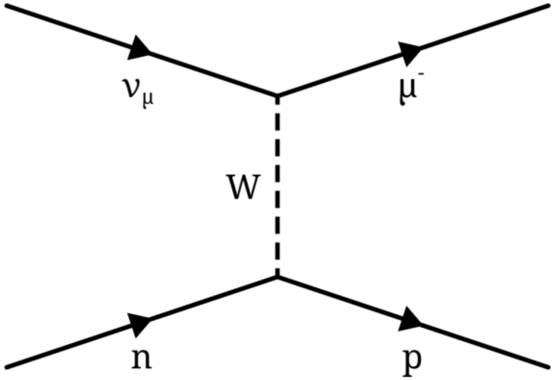

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$$\uparrow$$

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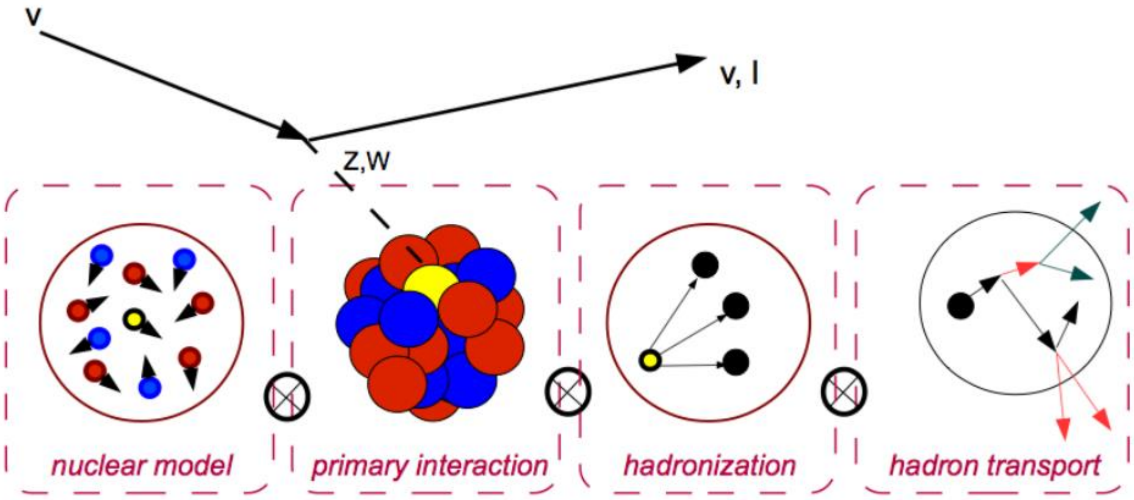
- The f factors are the "form factors" (read "fudge factors")
- Many of these can be extracted from electron scattering experiments
- f_A is the axial form factor, here we don't have much data to help us!
  - Usually we take a dipole form but **recent lattice QCD** calculations suggest this might not be a good idea Ann. Rev. Nucl. Part Vol. 72:205-232

Neutrino-nucleus scattering



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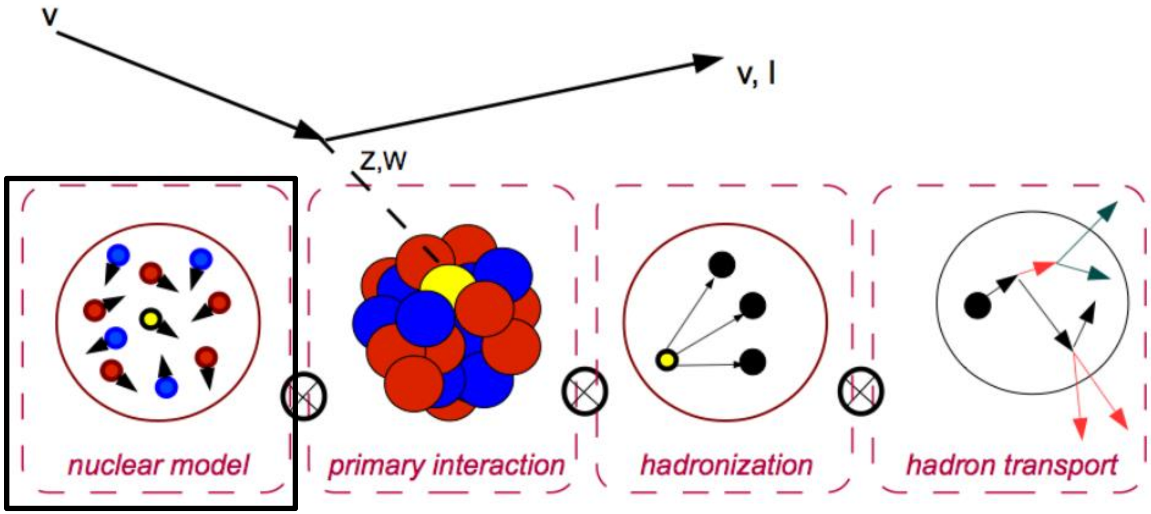

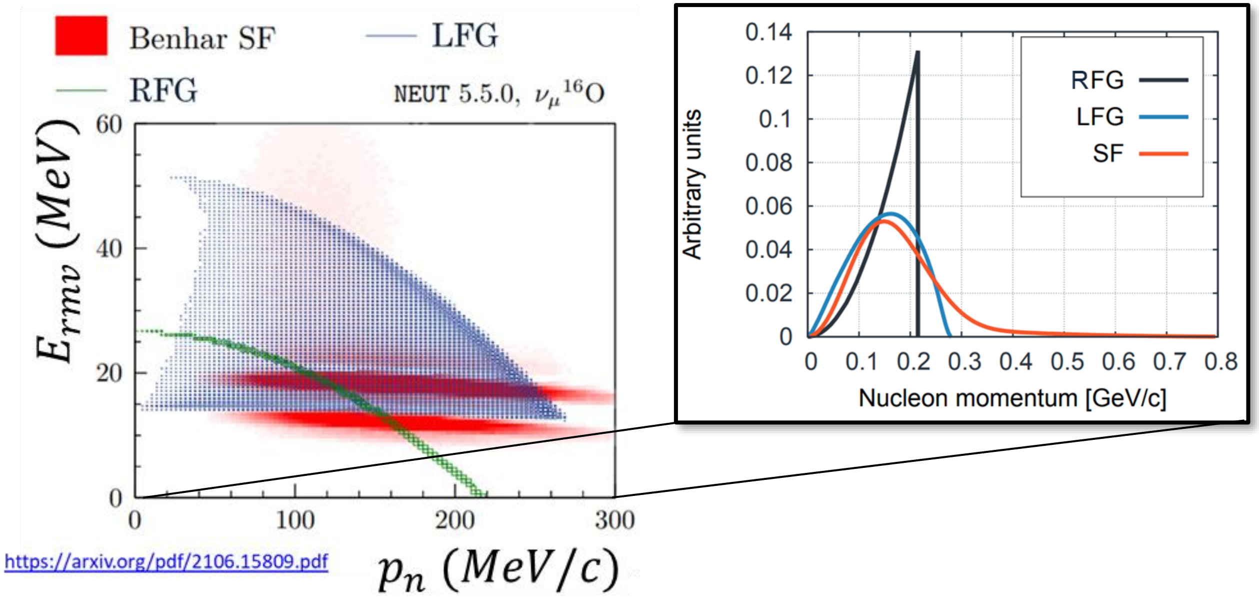

https://arxiv.org/pdf/2106.15809.pdf

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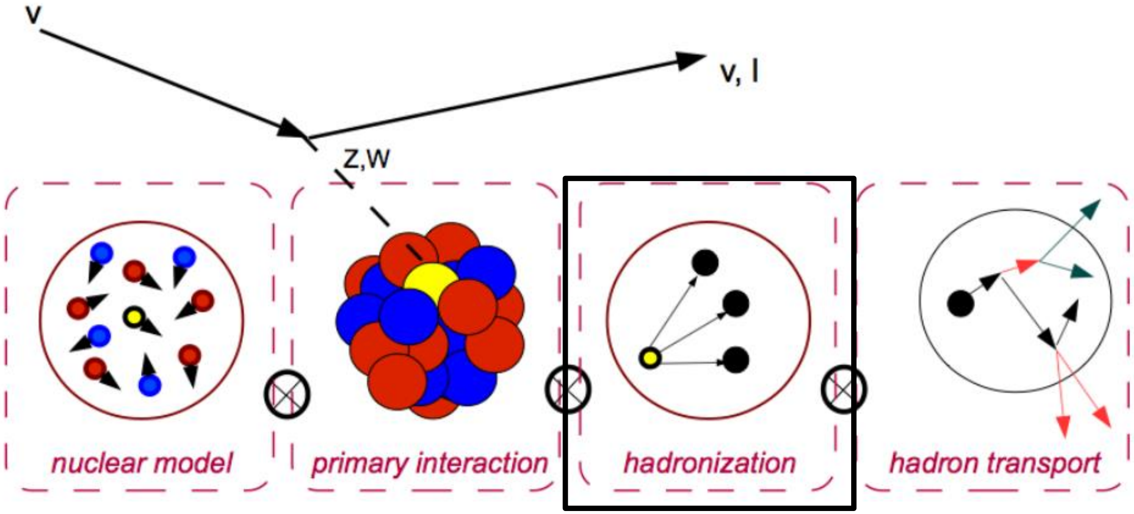
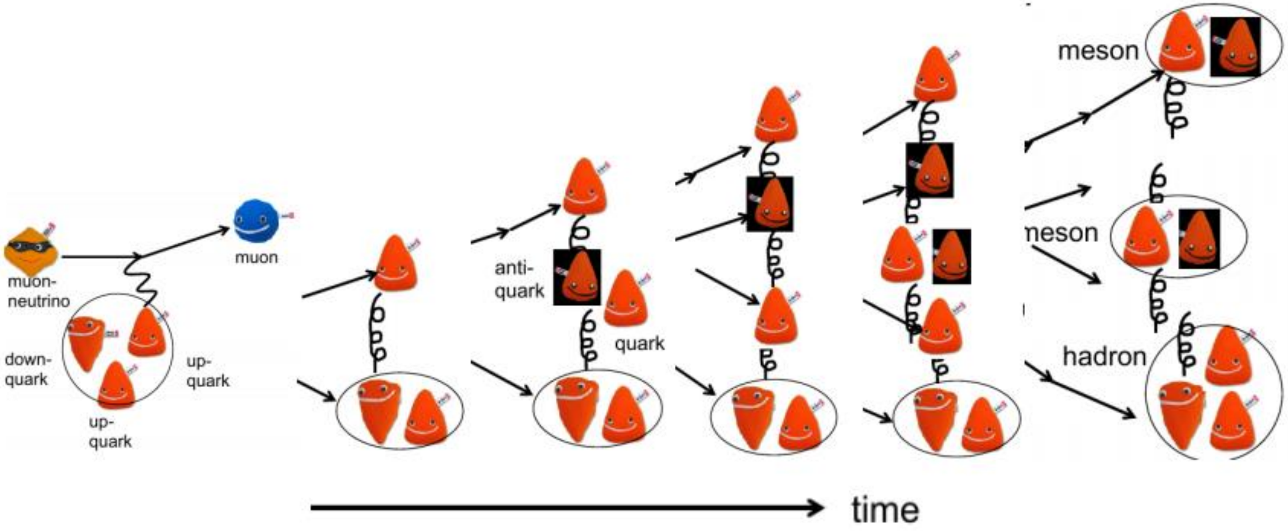
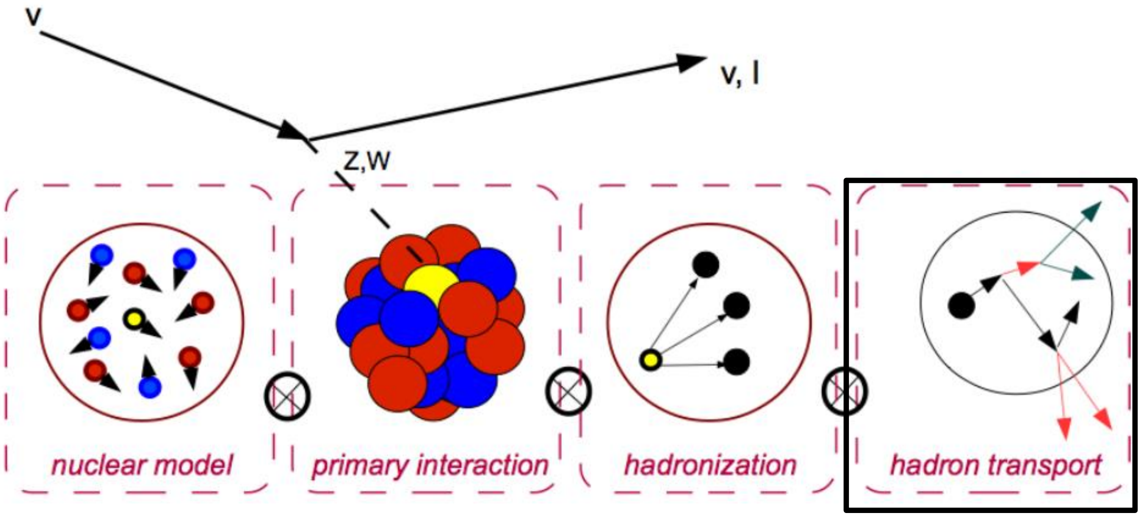




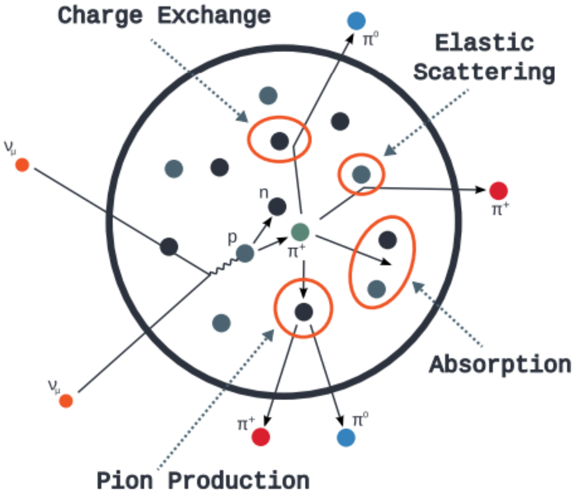
T. Katori

Neutrino-nucleus scattering



- Hadrons re-interact inside the nuclear medium: **Final State Interactions**

- Impractical to solve exactly, forced to use approximate methods



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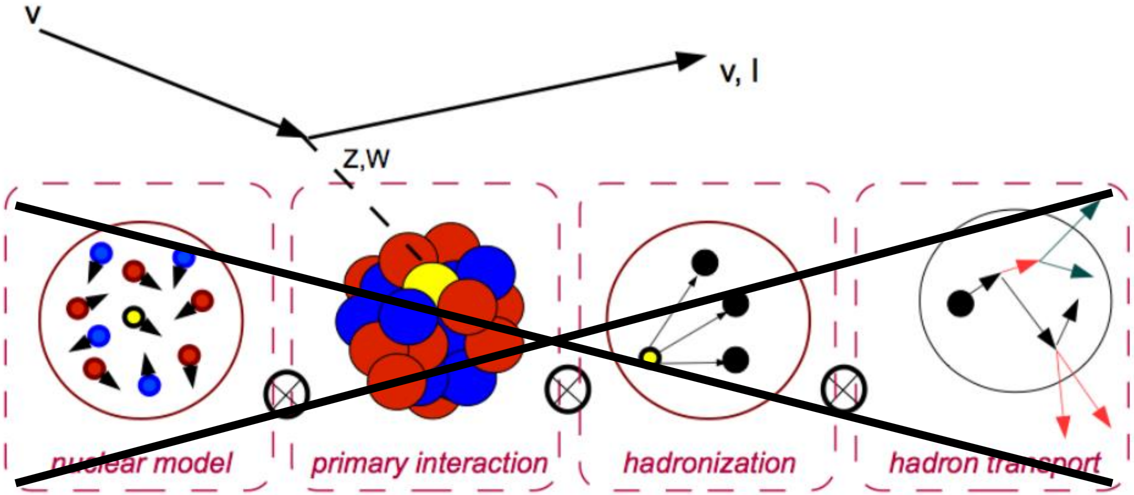
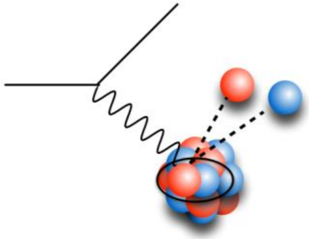
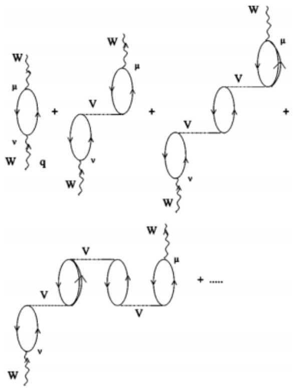


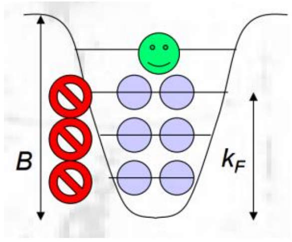
Multi-nucleon Interactions



Long range nuclear correlations



Pauli blocking



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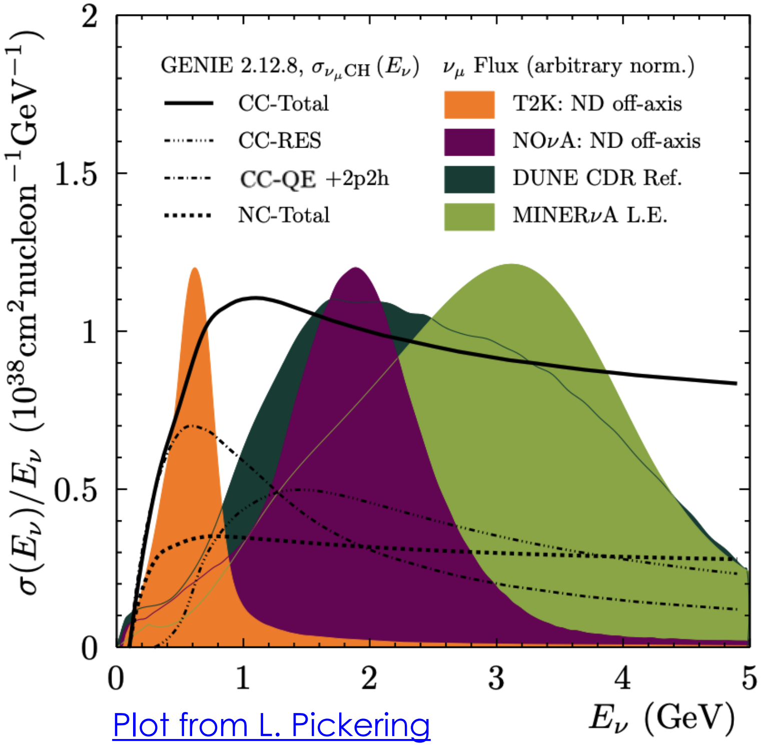
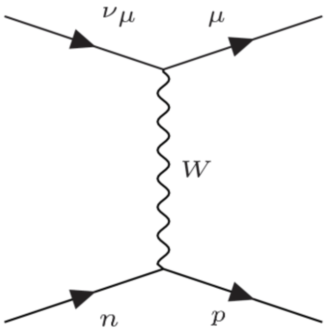
Neutrino-nucleus cross sections



Plot from L. Pickering

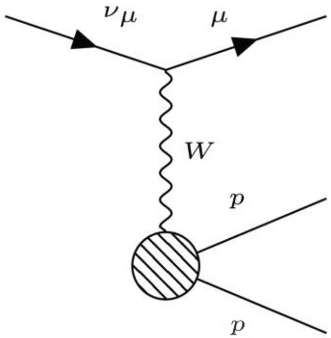
CC-QE

(Charged-Current Quasi-Elastic)



CC-2p2h

(Two-Particle-Two-Hole)



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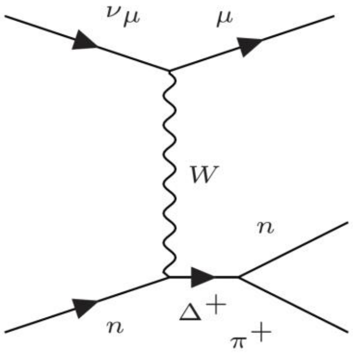

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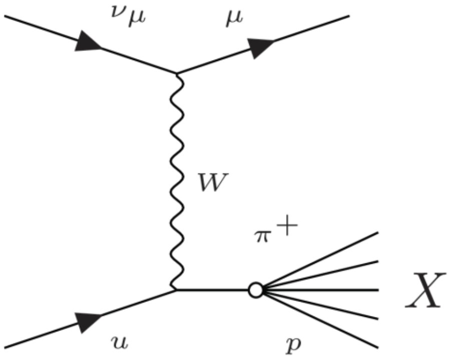

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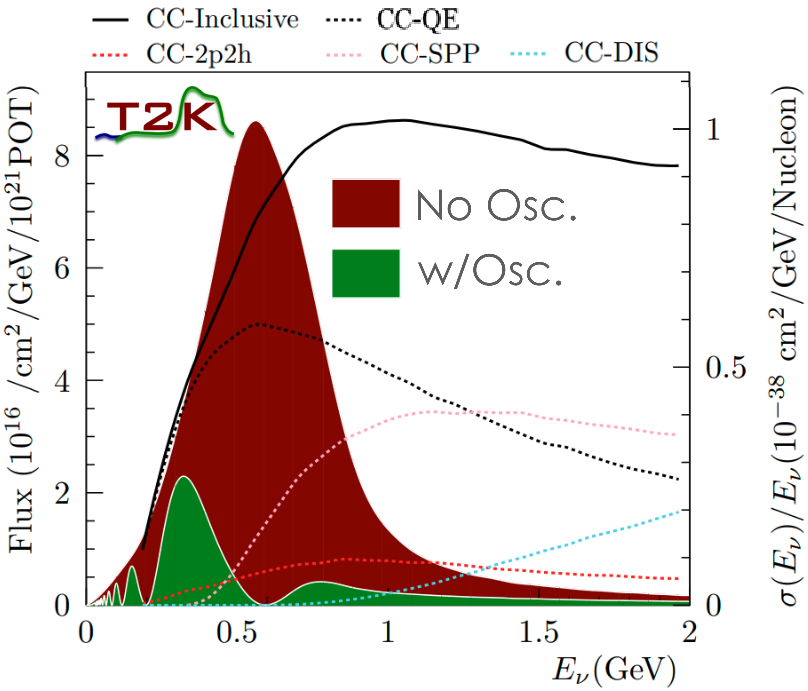

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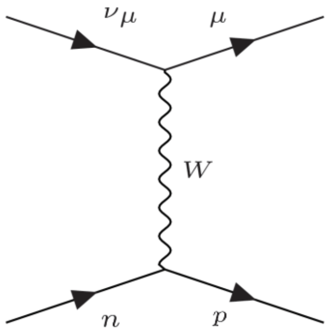

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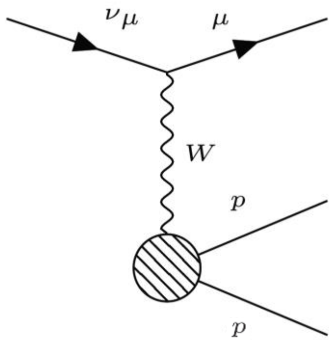

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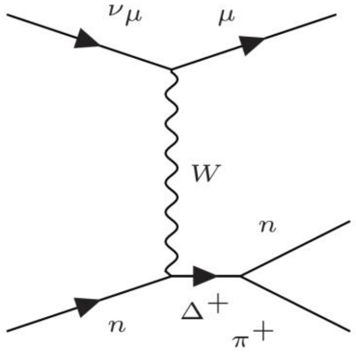

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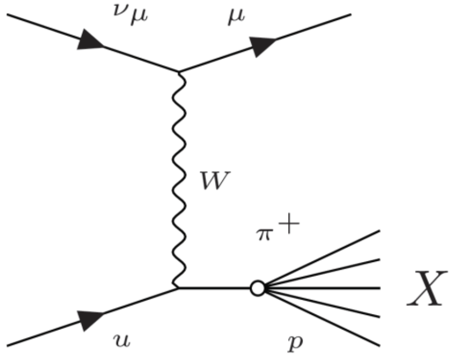
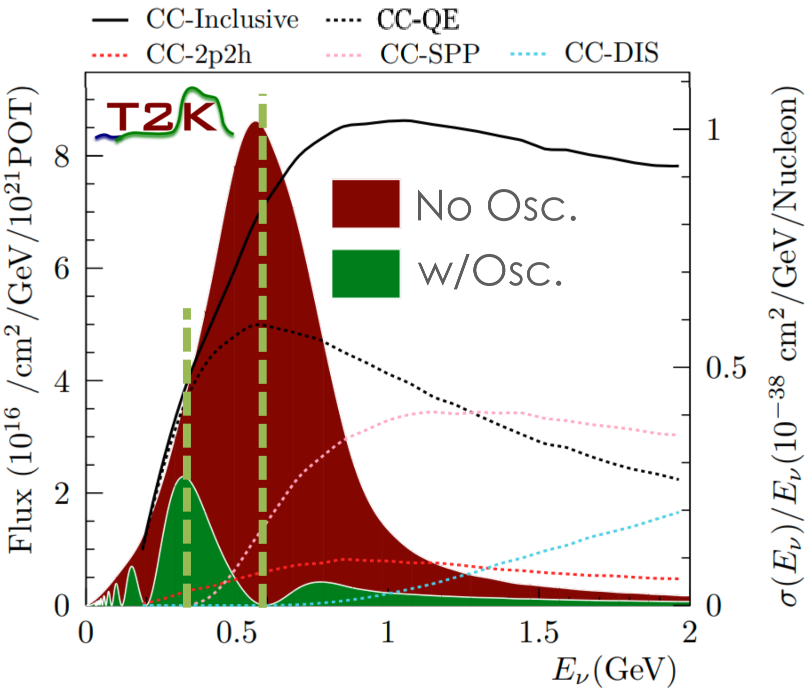


Event rates to oscillation parameters



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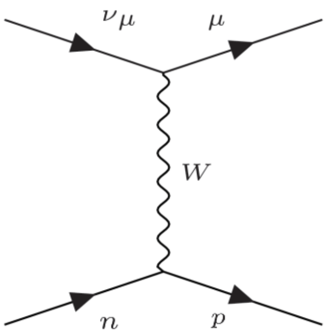

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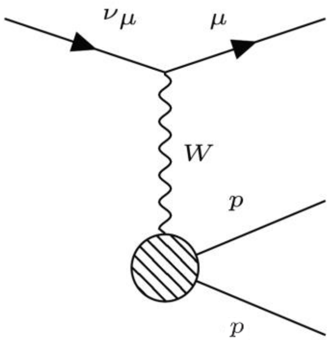

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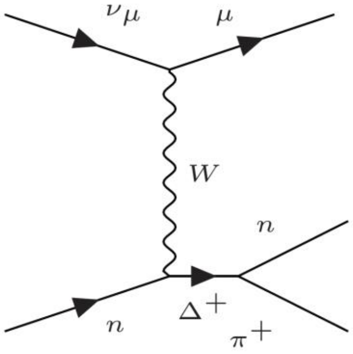

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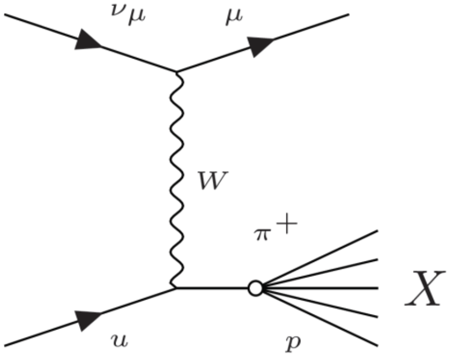

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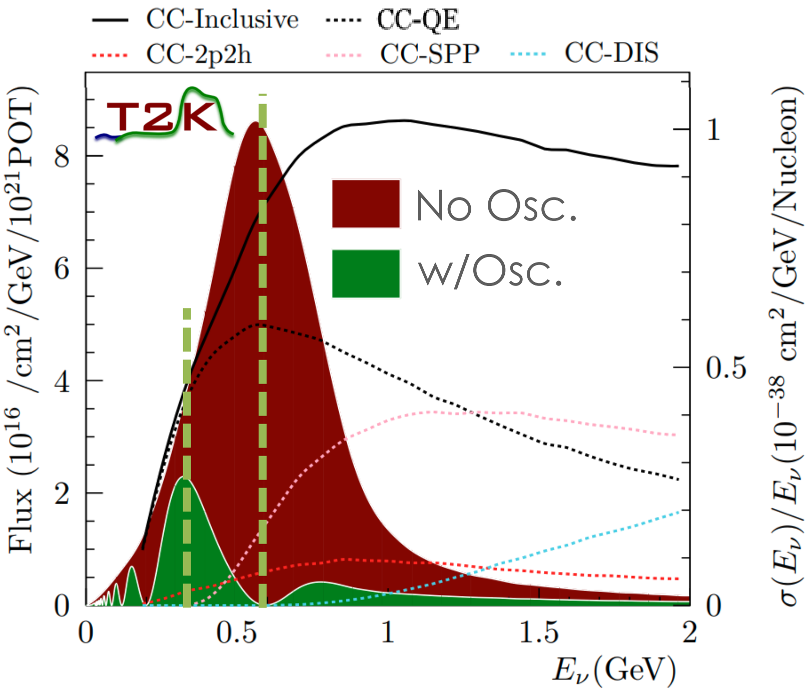

- Near / far ratios don't fully cancel systematics:
  - Dramatic change in E_ν distribution due to oscillations
  - ν_μ at ND vs ν_e at FD (for appearance)
  - Different ND/FD design, acceptance

At the near detector

$$N_\mu(E_\nu) = \sigma(E_\nu)\Phi_\nu(E_\nu)\epsilon(E_\nu)$$

Interaction cross section

Neutrino flux

Detector effects

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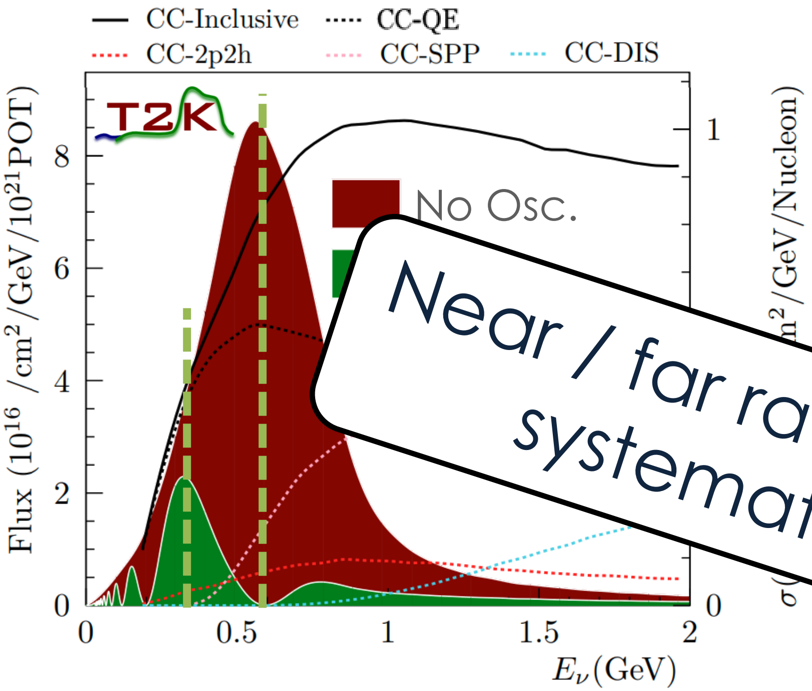
Oscillation probability

Event rates to oscillation parameters



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  - ν_μ at ND vs ν_e at FD (for
    arance)
  - design,

Near / far ratios don’t fully cancel systematic uncertainties!

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Interaction cross section

Neutrino flux

Detector effects

At the far detector

$$N_\mu(E_\nu) = P(\nu_\mu \to \nu_\mu)\sigma(E_\nu)\Phi_\nu(E_\nu)\epsilon(E_\nu)$$

$$N_e(E_\nu) = P(\nu_\mu \to \nu_e)\sigma(E_\nu)\Phi_\nu(E_\nu)\epsilon(E_\nu)$$

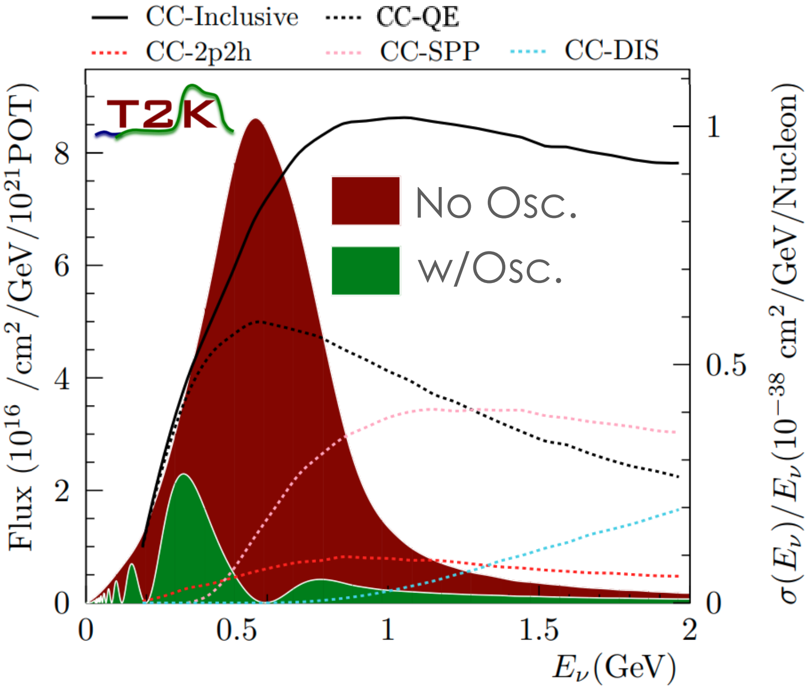
Oscillation probability

Where are we so far?

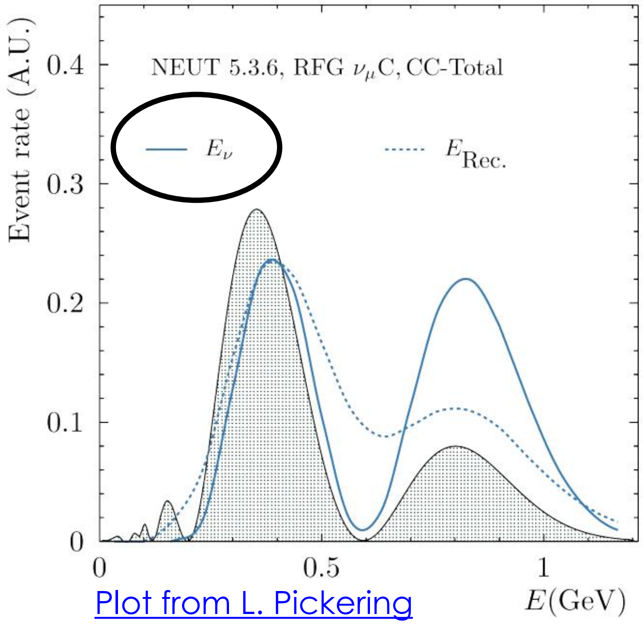
- Key questions:

  1. *Why is modelling neutrino interactions so difficult?*
  2. *Why does the near detector not allow a better cancellation of uncertainties?*
  3. ***What exactly do we need to understand in order to reduce uncertainties on oscillation measurements?***

Event rates to oscillation parameters



What we would like to measure



Plot from L. Pickering

$$N_\ell(E_\nu) = P(\nu_\mu \to \nu_\ell)(E_\nu)\ \sigma(E_\nu)\ \Phi_\nu(E_\nu)\ \epsilon(E_\nu)$$

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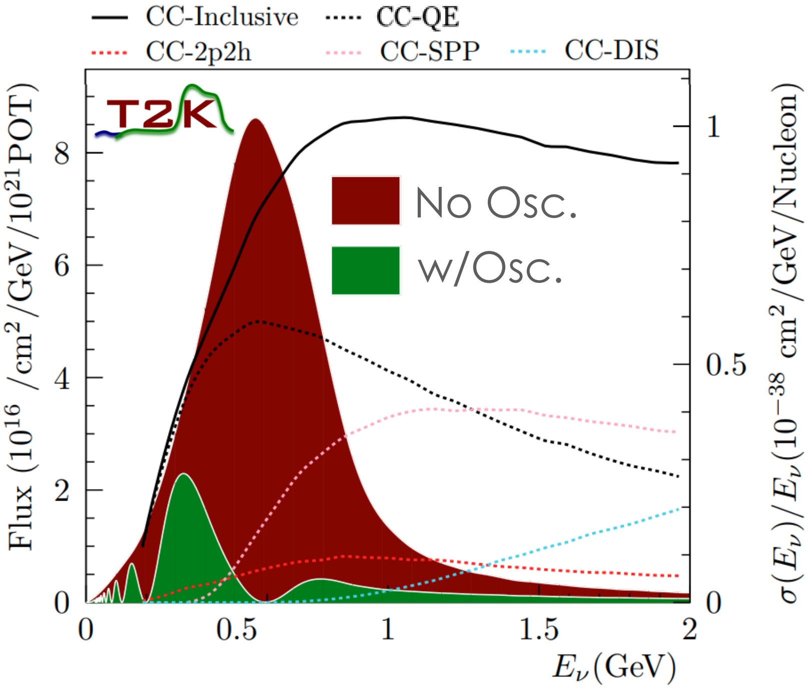
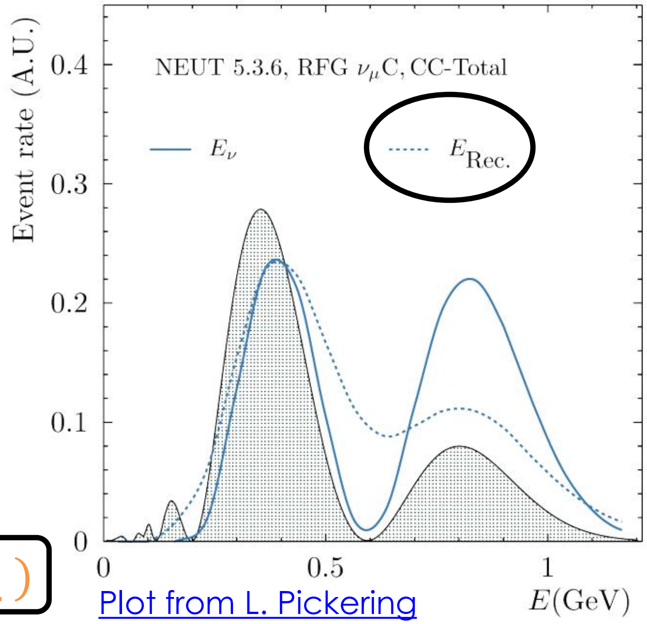


What we can actually measure



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$$N_\ell(E_{Rec.}) = P(\nu_\mu \to \nu_\ell)(E_\nu)\ \sigma(E_\nu)\ \Phi_\nu(E_\nu)\ \epsilon(E_\nu)\ S(E_\nu, E_{Rec.})$$

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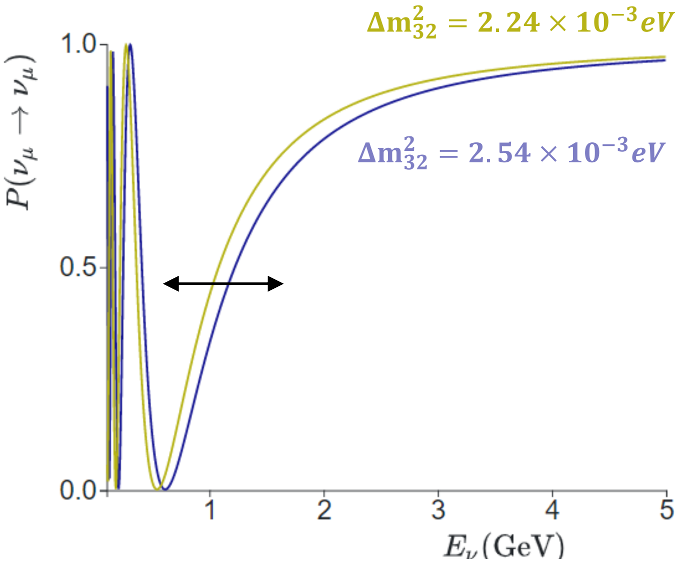
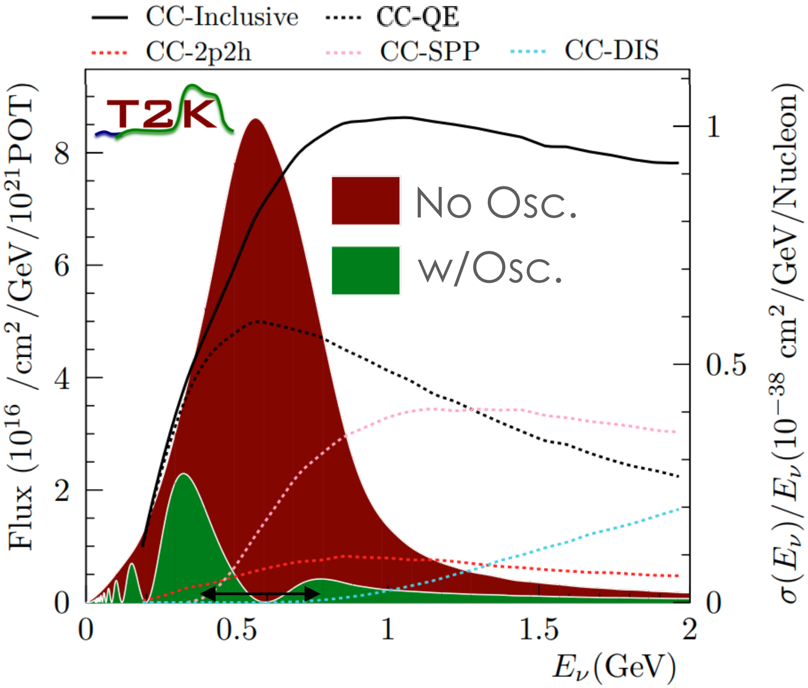
- For a precision probe of oscillation parameters, reconstructing the shape of the oscillated spectrum is crucial





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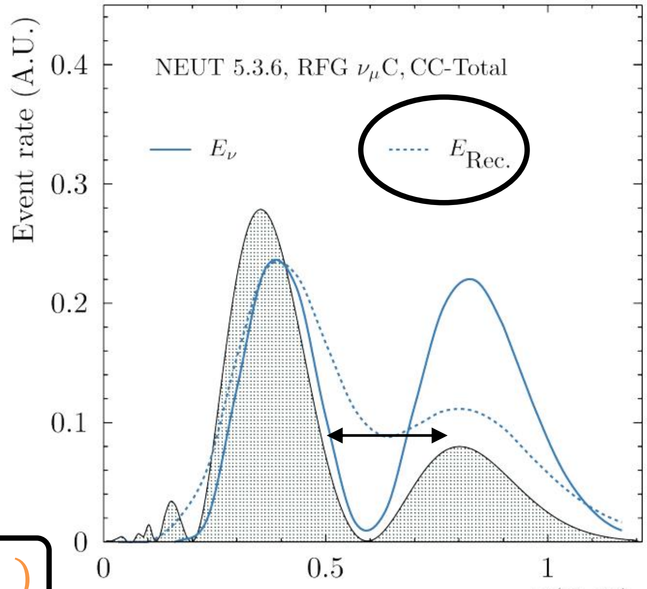

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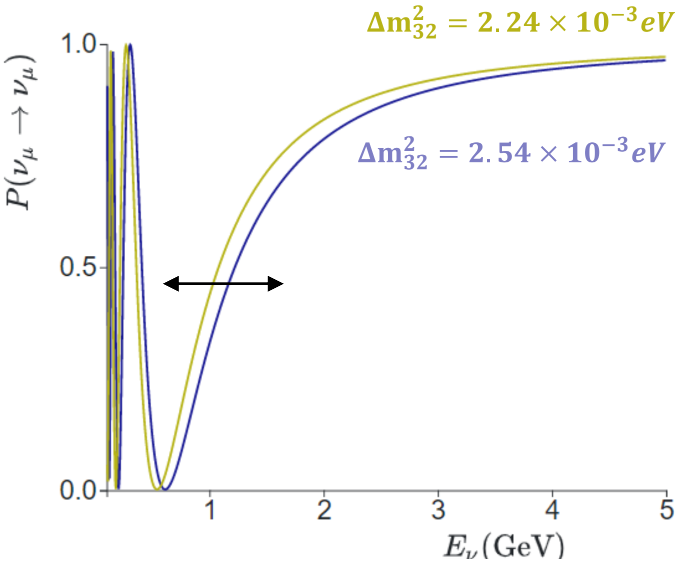

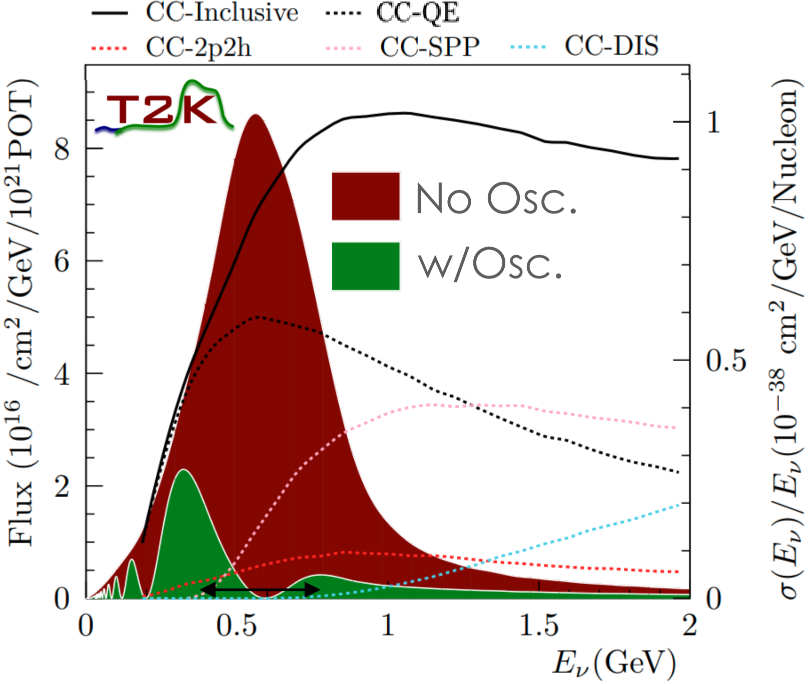

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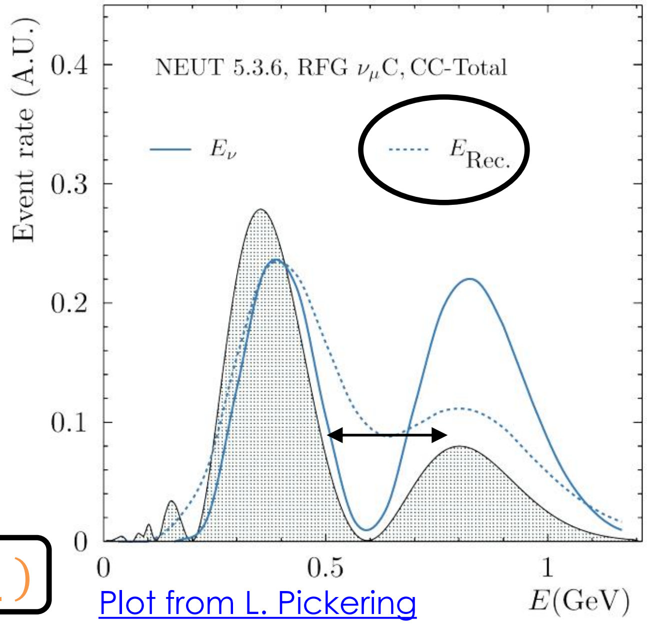

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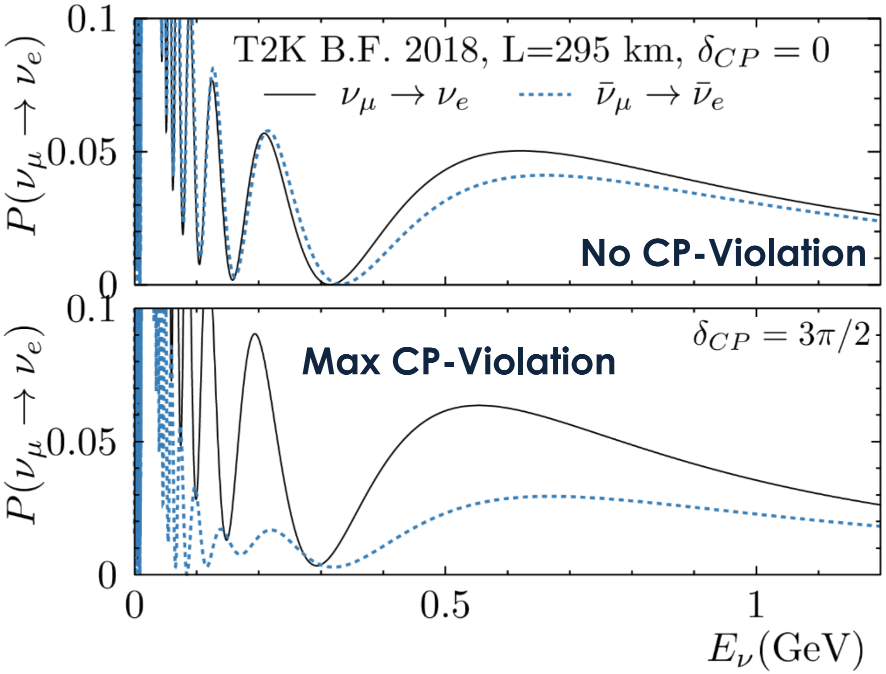

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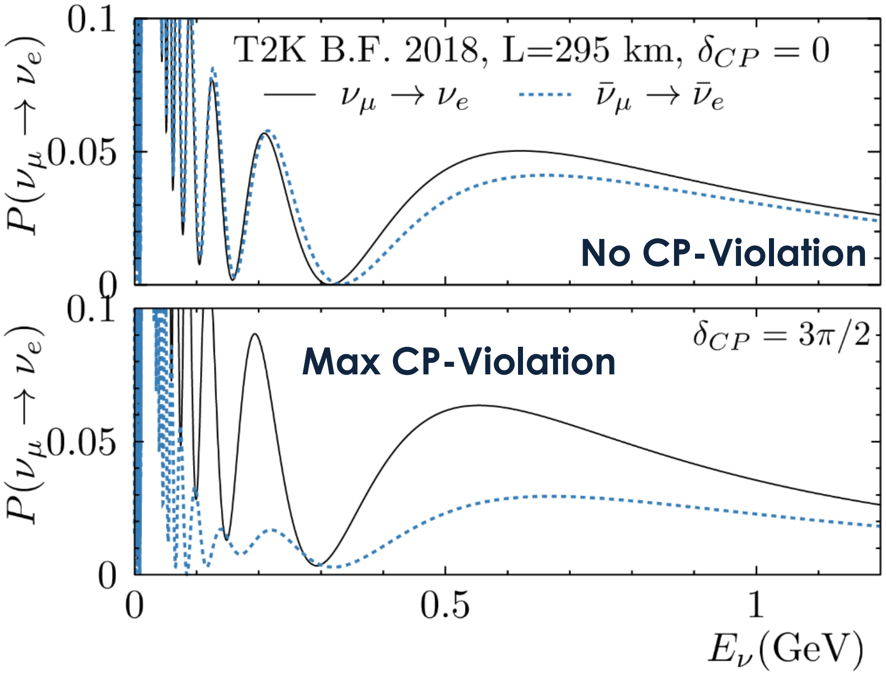

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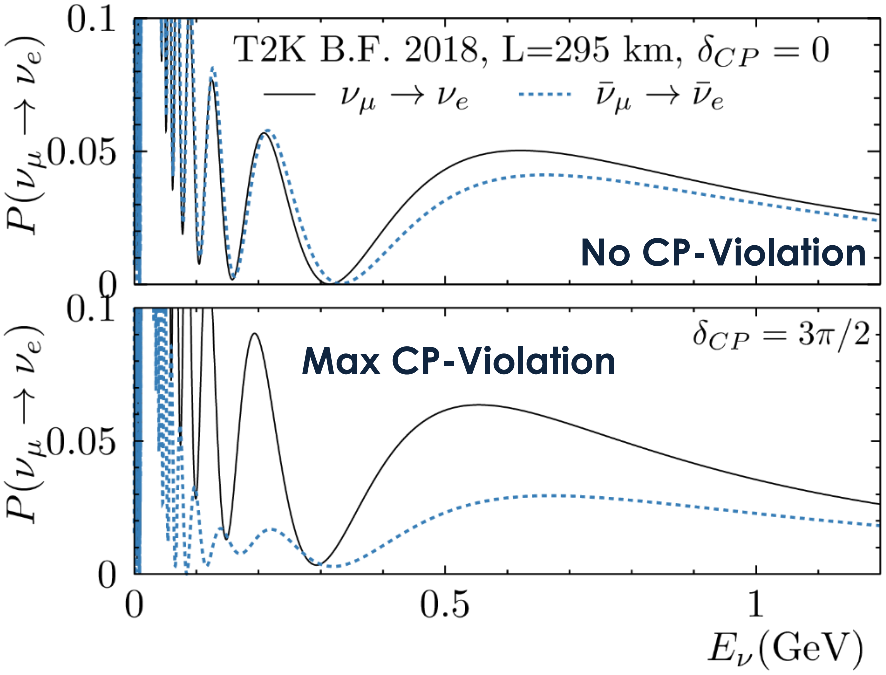
- But we mainly measure muon neutrino interactions at the near detector
- **A good modelling of ν_e/ν_μ cross section ratio is essential**



Three things we need to model

(a non exhaustive list)

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    - *So we know how to extrapolate from our near to far detectors*

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Overview

- Neutrino Oscillations
- Accelerator-Based Experiments
- $\boldsymbol{\nu}$ Interactions for $\boldsymbol{\nu}$ Oscillations
- **Reconstructing Neutrino Energy**
- The Path to Precision Measurements

Reconstructing E_ν

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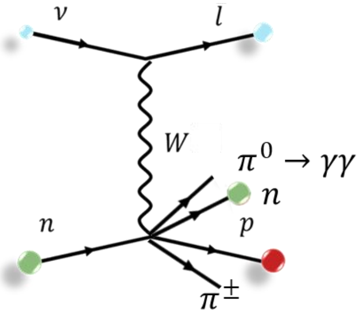
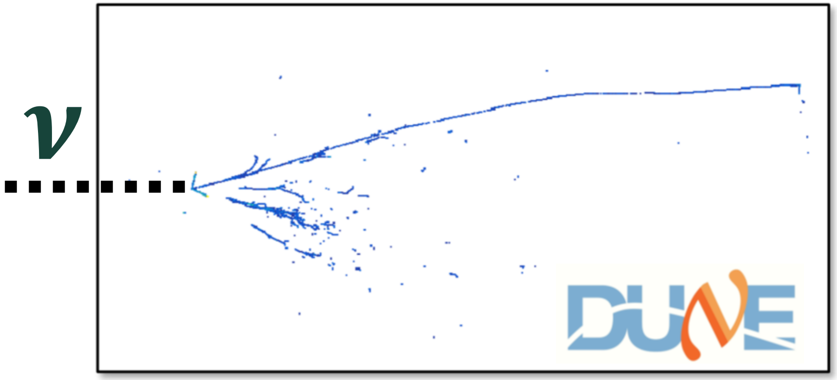
“Calorimetric method”

$$E_\nu = E_\ell + E_{had,vis}$$

- Add the lepton energy to the sum of all visible hadronic energy
- But not all hadrons deposit all their energy inside the detector





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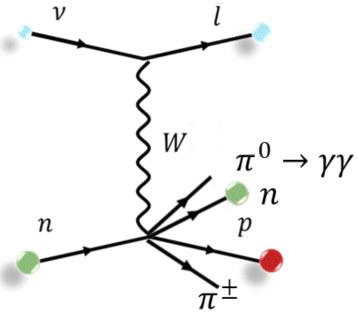
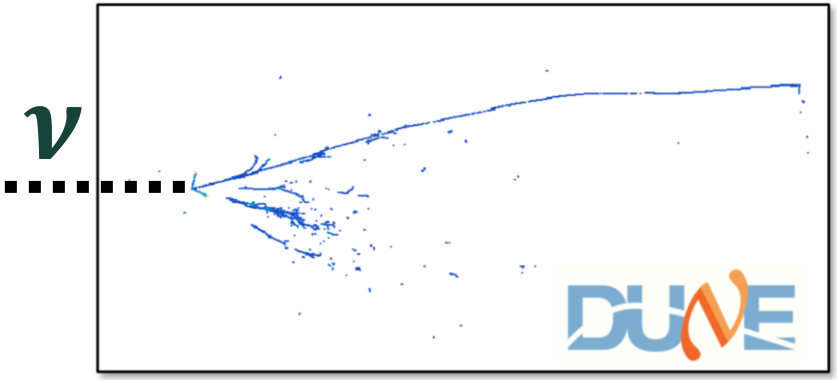
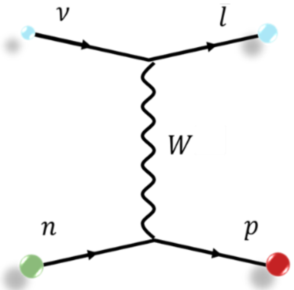


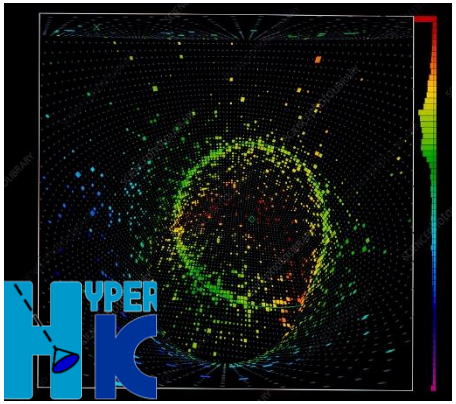


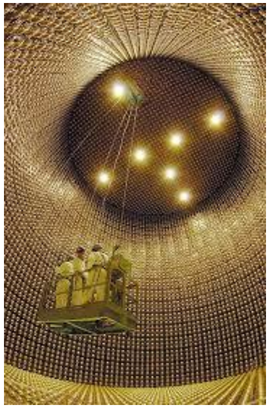
"Kinematic method"

$$E_\nu = \frac{m_p^2 - (m_n - E_B)^2 - m_\ell^2 + 2E_\ell(m_n - E_B)}{2(m_n - E_B - E_\ell + p_\ell \cos\theta_\ell)}$$

- Uses only the outgoing lepton kinematics
- Assume elastic scatter off a stationary nucleon

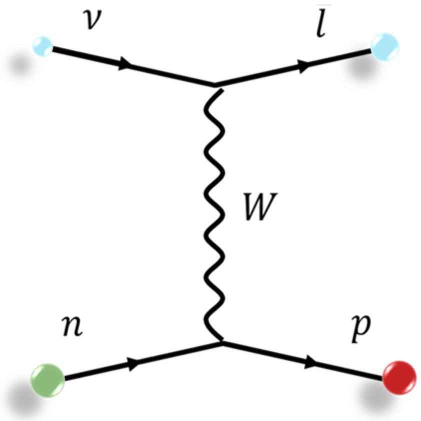


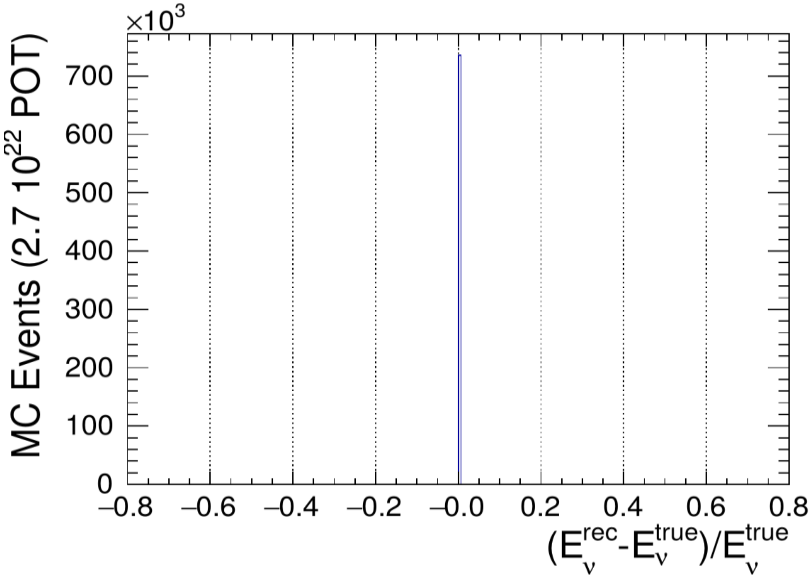




Nuclear effects and E_ν (T2K/HK)

CCQE (1p1h)





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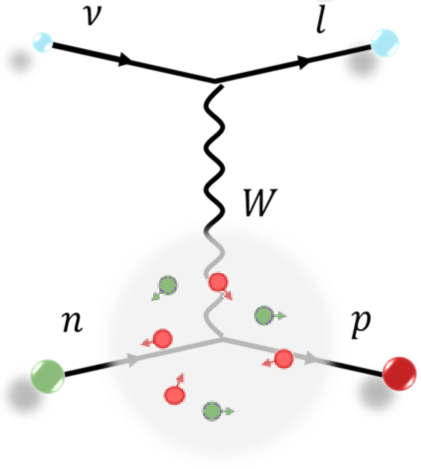
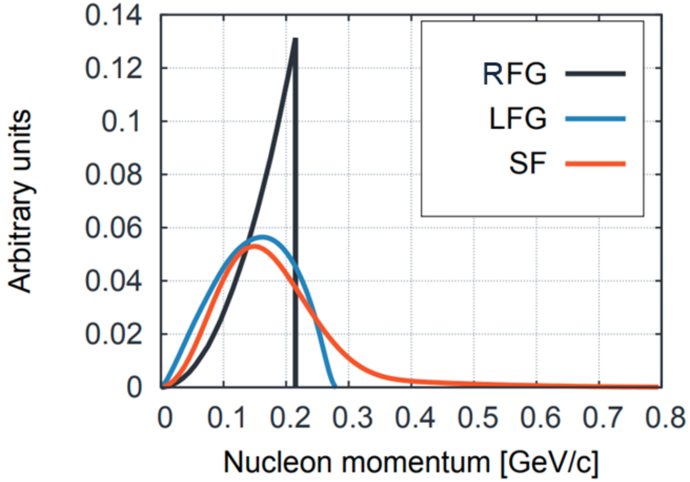
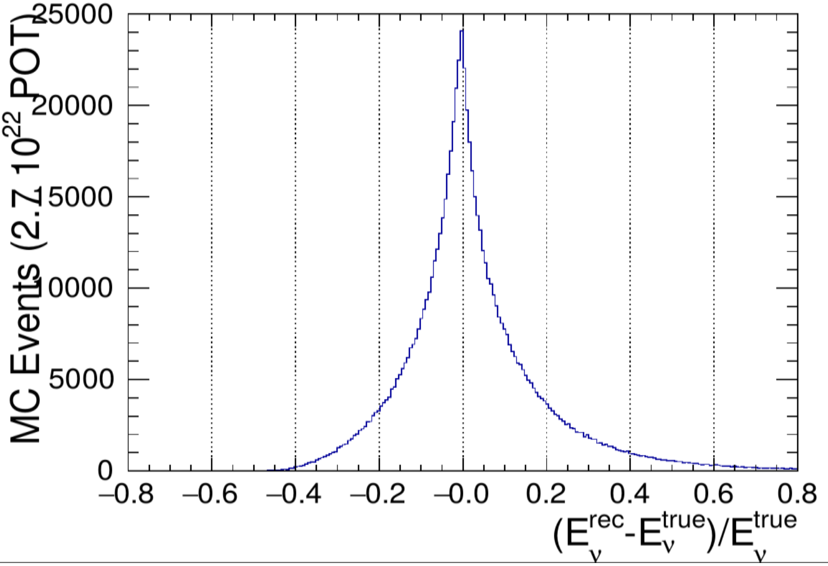
Proxy for E_ν from lepton kinematics is exact only for **CCQE elastic scattering** off a **stationary nucleon**

Nuclear effects and E_ν (T2K/HK)

CCQE (1p1h)







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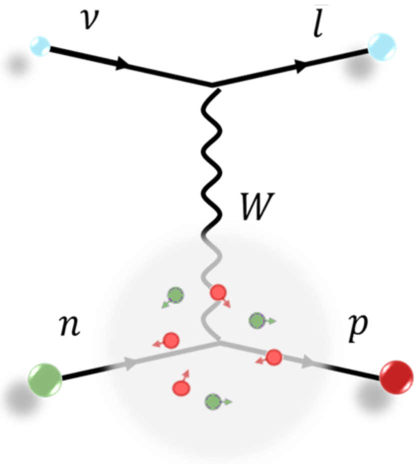
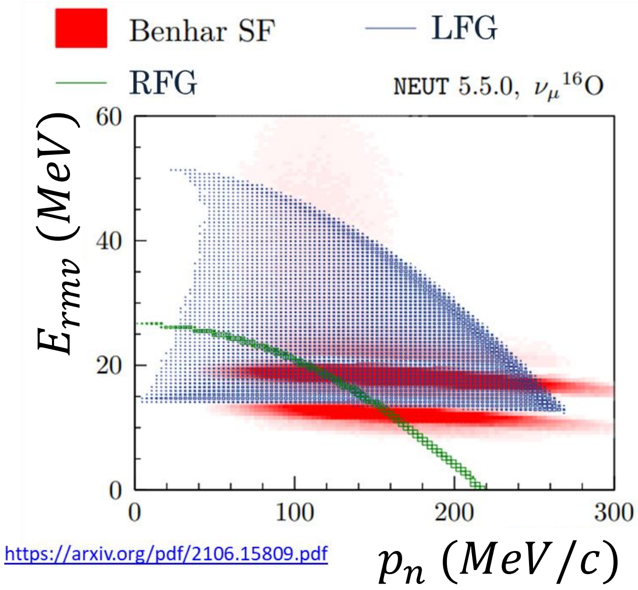
The motion of the nucleons inside the nucleus (*Fermi motion*) causes a **smearing** on E_ν

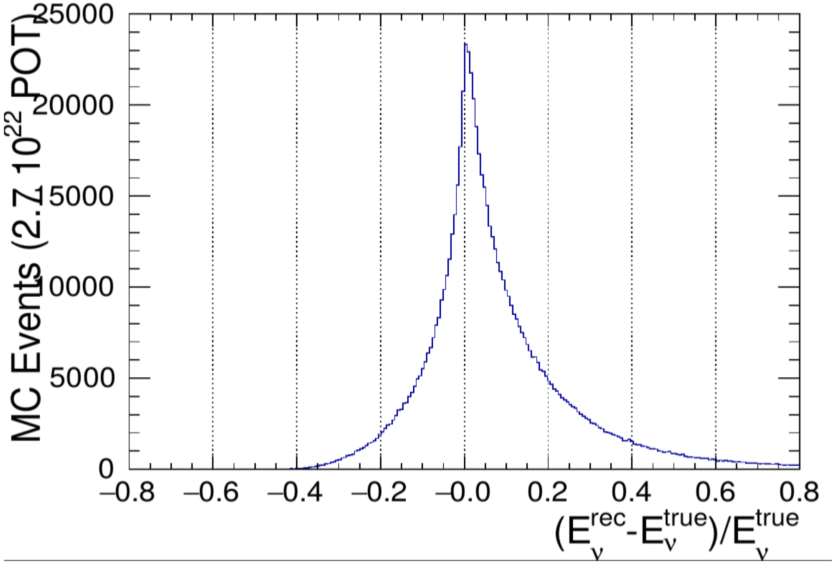
Nuclear effects and E_ν (T2K/HK)

CCQE (1p1h)





https://arxiv.org/pdf/2106.15809.pdf



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The energy loss in the nucleus (to extract the struck nucleon from its shell) introduces a **bias**

Nuclear effects and E_ν (T2K/HK)

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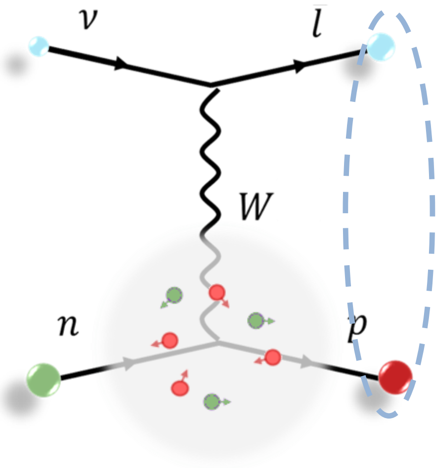
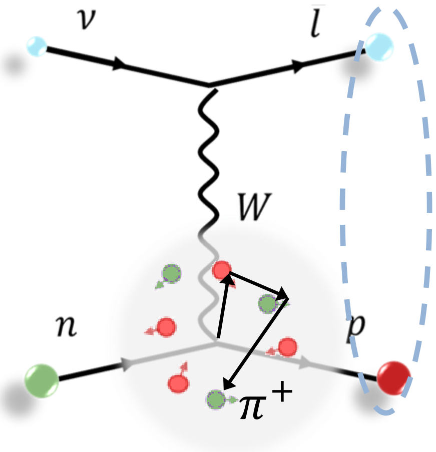
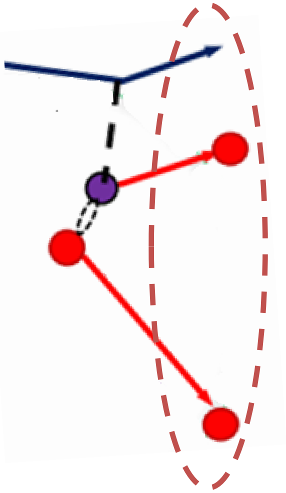


CCRES

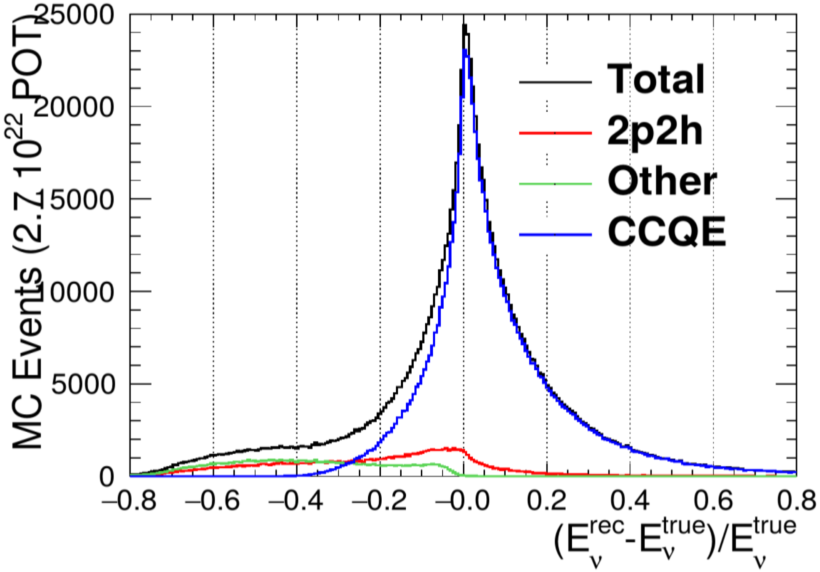


Final state interactions (FSI) can cause different interaction modes to have the same final state

2p2h



Interactions off a bound state of two nucleons can result in **2p2h** final states



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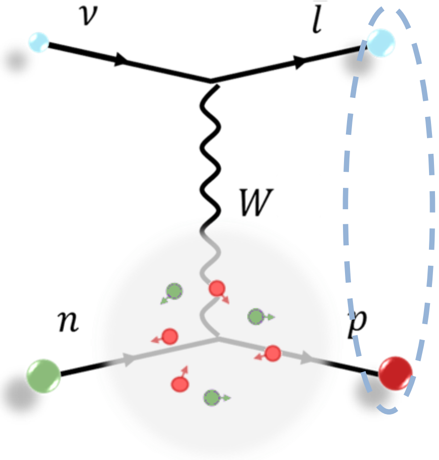
The motion of the nucleons inside the nucleus (*Fermi motion*) causes a **smearing** on E_ν

The energy loss in the nucleus (to extract the struck nucleon from its shell) introduces a **bias**

Not a good proxy for non-CCQE events: 2p2h and CC1π with pion abs. FSI

Nuclear effects and E_ν (T2K/HK)

CCQE (1p1h)



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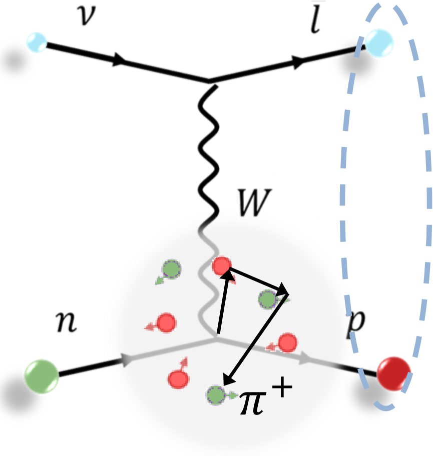

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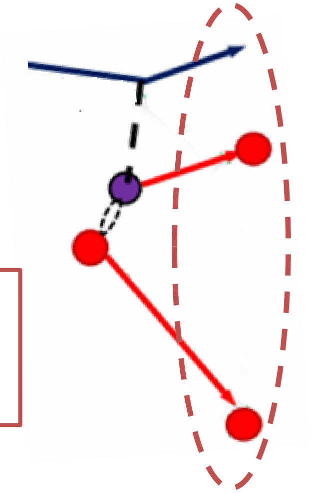

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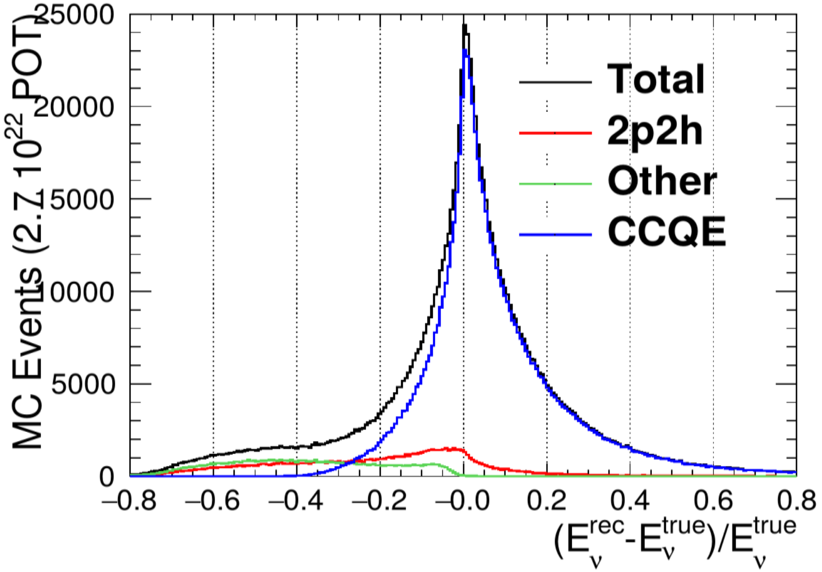

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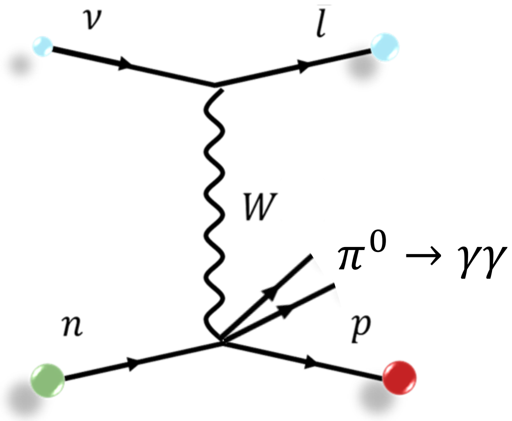
First-order effects

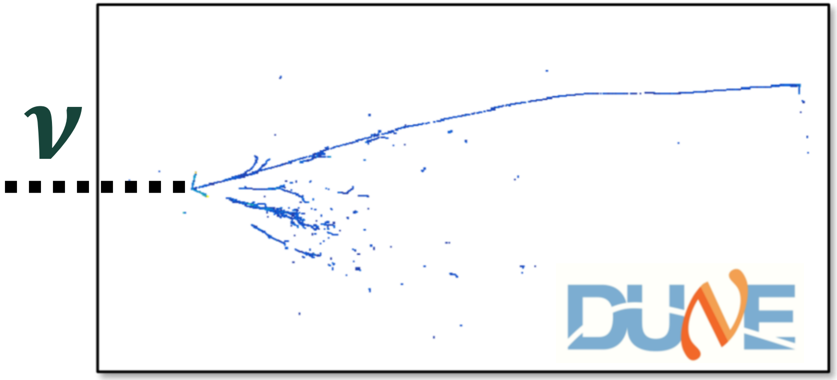
Fermi motion causes a **smearing** on E_ν^{QE}

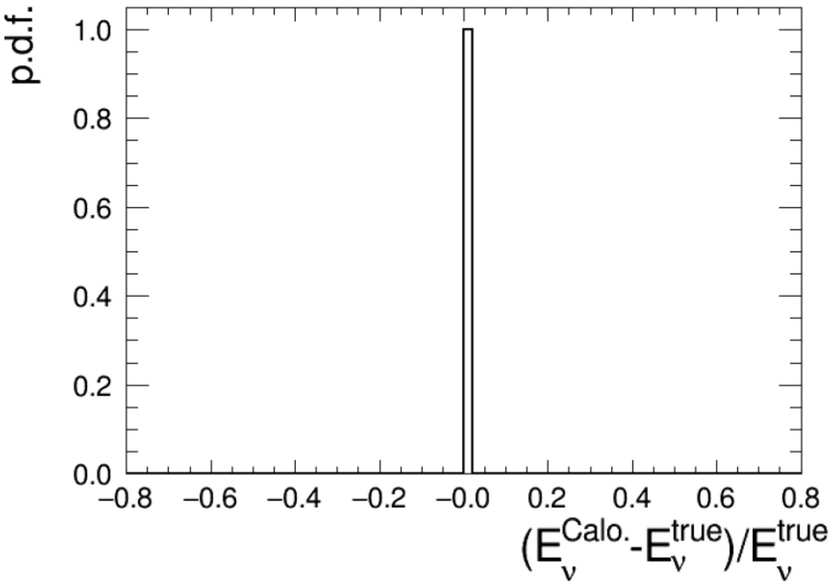
Nuclear removal energy effects introduce a **bias**

2p2h and pion abs. *FSI* cause further **bias**

Nuclear effects and E_ν (DUNE/NOvA)





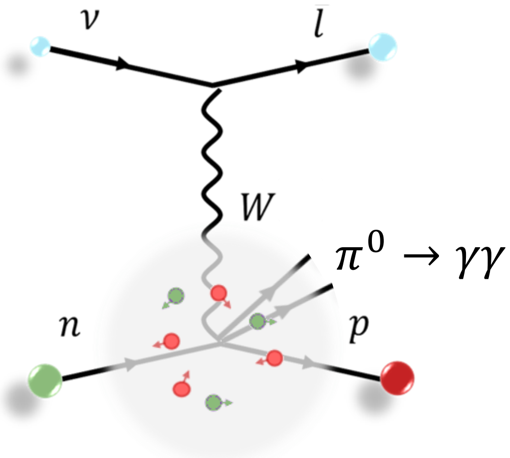


$$E_\nu = E_\ell + E_{had,vis} \approx E_\ell + \Sigma T_p + \Sigma T_{\pi^\pm} + \Sigma E_\gamma$$

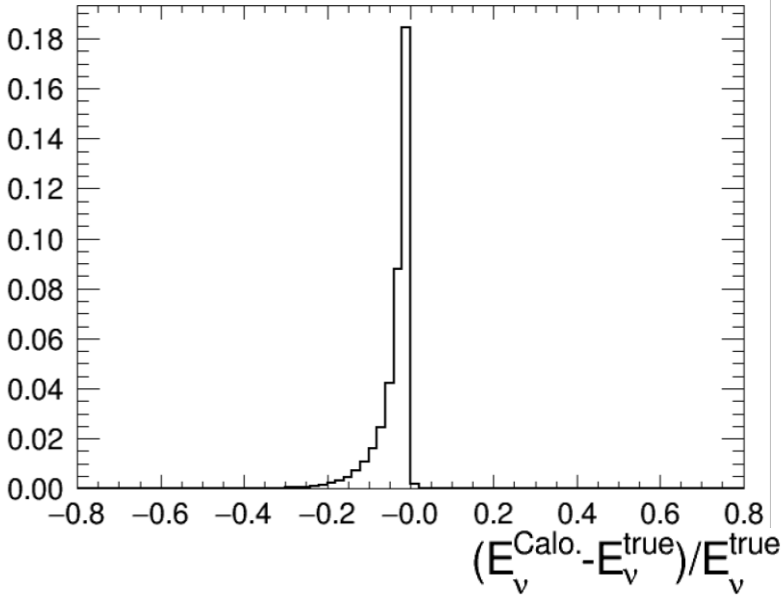
Calculation is exact only for interactions **without neutrons and charged pions** (ignore heavier mesons here) off a **free nucleon**

Usefulness is not restricted to QE-like interactions (no final state pions)

Nuclear effects and E_ν (DUNE/NOvA)



All events without n or $\pi^\pm$



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Impact of initial state effects (Fermi motion and removal energy) smaller than in QE approach

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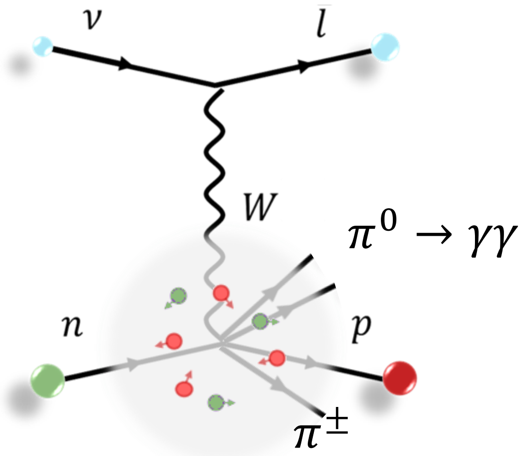

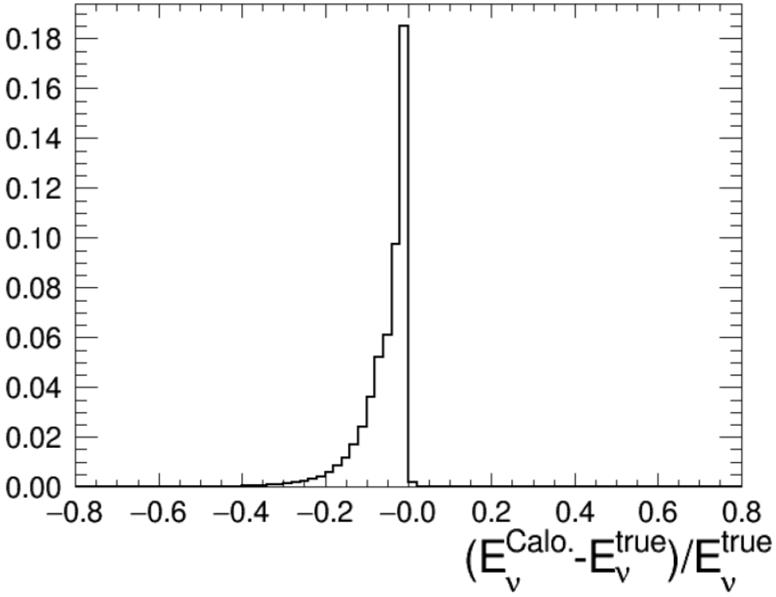
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Missed charged pion mass energy causes a bias

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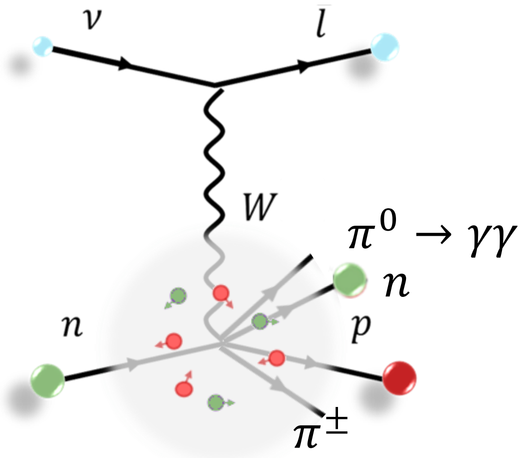

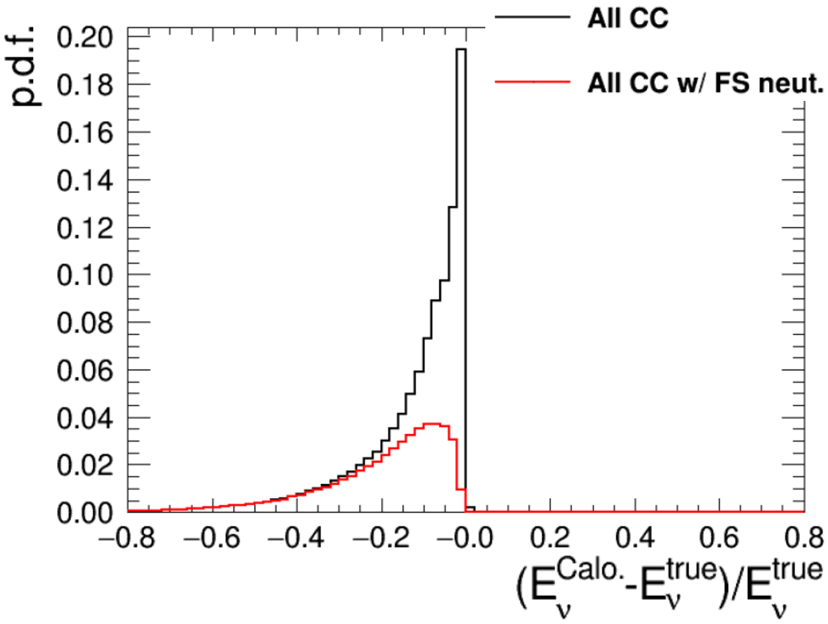

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Fraction of E_ν in **Neutrons** is critical

What we need to know (a non exhaustive list!)

T2K/HK

("kinematic" E_ν proxy)

Critical

- Nuclear ground state: **Fermi motion** and "**binding energy**"
- **2p2h** and **pion absorption FSI** contributions to 0π final states

DUNE/NOvA/SBN

("calorimetric" E_ν proxy)

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- Fraction of energy found in **neutrons**
- Charged **pion multiplicity**

- Subtle nuclear physics processes are crucial in order to understand how we can translate from what our detectors see to true neutrino energy

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DUNE/NOvA/SBN

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- Charged **pion multiplicity**

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Neutrino interaction modelling is crucial for all upcoming experiments, but different experiments have different priorities: **complementary approaches**!

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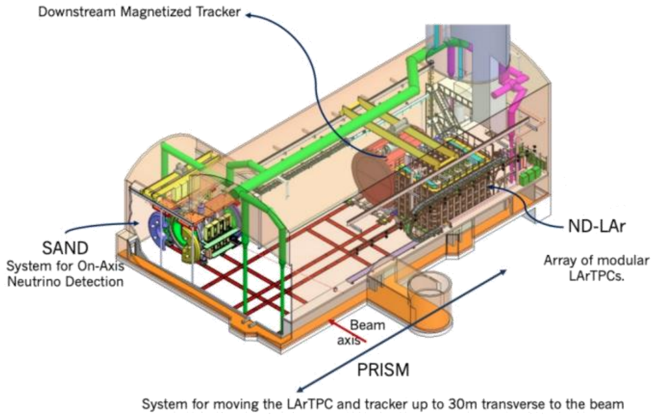
Path to Precision Measurements

Improved near detector capabilities



Engagement with the nuclear theory community

GiBUU

Genie

UNIVERSAL NEUTRINO GENERATOR & GLOBAL FIT

NuWro

Wroclaw Neutrino Event Generator

Dedicated lepton-nucleus cross-section measurement programs

SBN Program

MINERvA

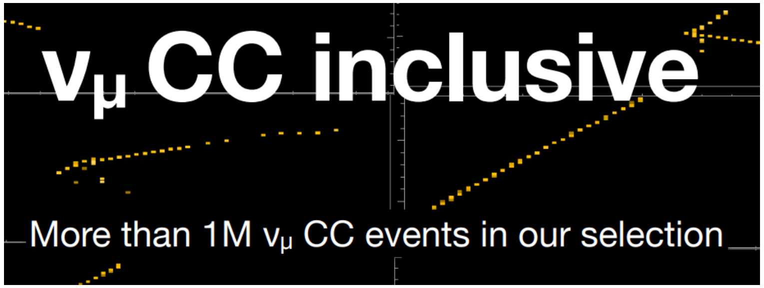
e4ν

Undetectable, you say?

"I have done something very bad today by proposing a particle that cannot be detected; it is something no theorist should ever do." *Wolfgang Pauli, 1930*

Well, have I got *νs* for you!

L. Cremonesi



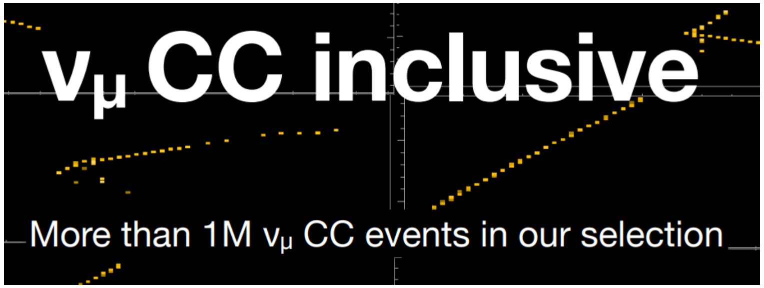
ν_μ CC inclusive

More than 1M ν_μ CC events in our selection

“I have done something very bad today by proposing a particle that cannot be detected; it is something no theorist should ever do.” *Wolfgang Pauli, 1930*

Well, have I got νs for you!

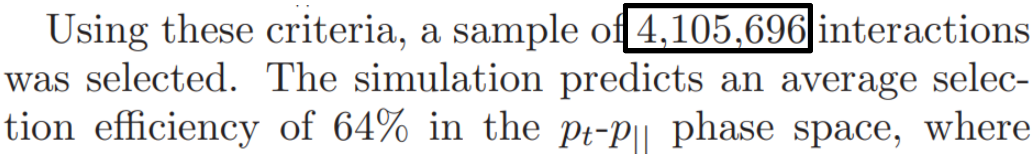
L. Cremonesi NEUTRINO 2020



ν_μ CC inclusive

More than 1M ν_μ CC events in our selection

Phys. Rev. D **104**, 092007



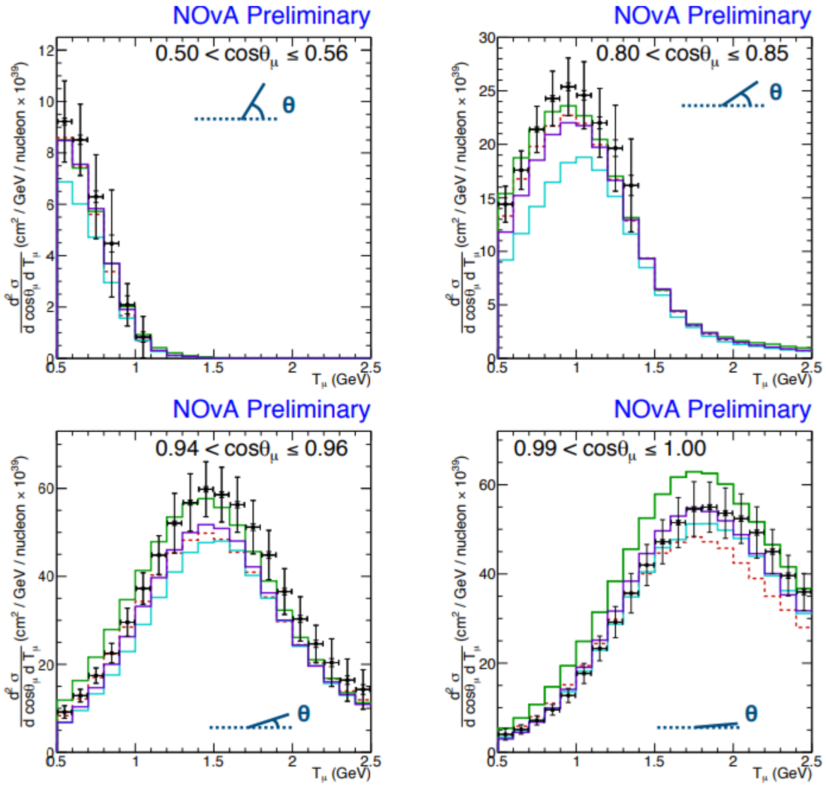
Using these criteria, a sample of $\boxed{4{,}105{,}696}$ interactions was selected. The simulation predicts an average selection efficiency of 64% in the p_t-$p_{||}$ phase space, where

> "I have done something very bad today by proposing a particle that cannot be detected; it is something no theorist should ever do." *Wolfgang Pauli, 1930*

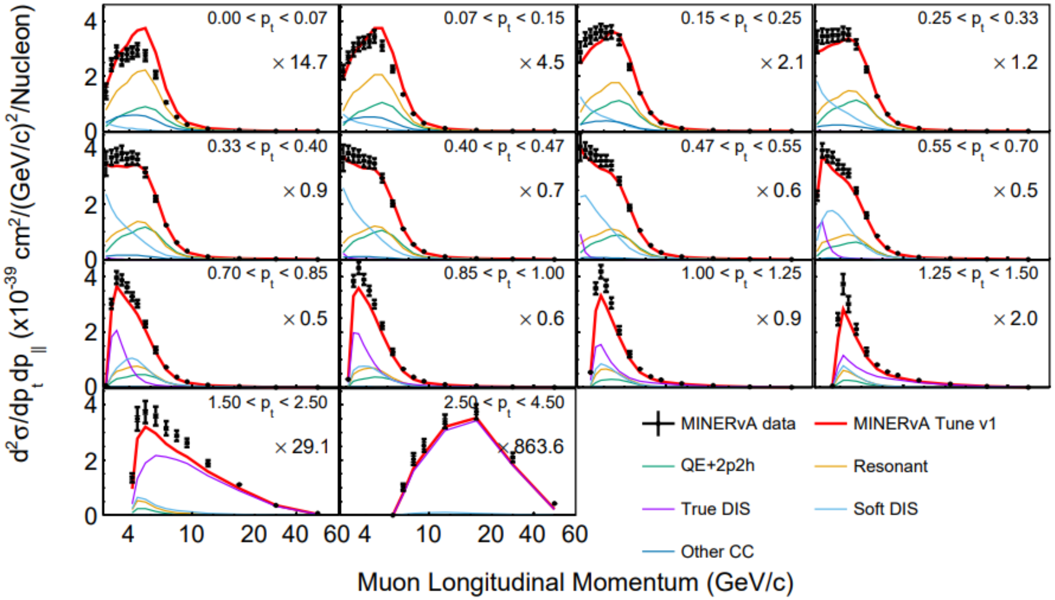
Well, have I got *νs* for you!

L. Cremonesi NEUTRINO 2020

Data (Stat. + Syst.)
GENIE 3.00.06*
GiBUU 2019
NEUT 5.4.0
NuWro 2019

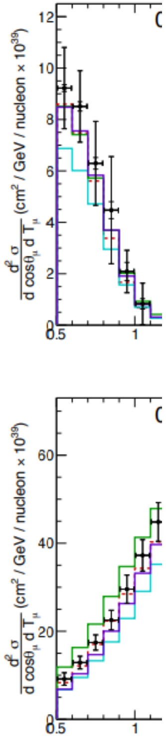


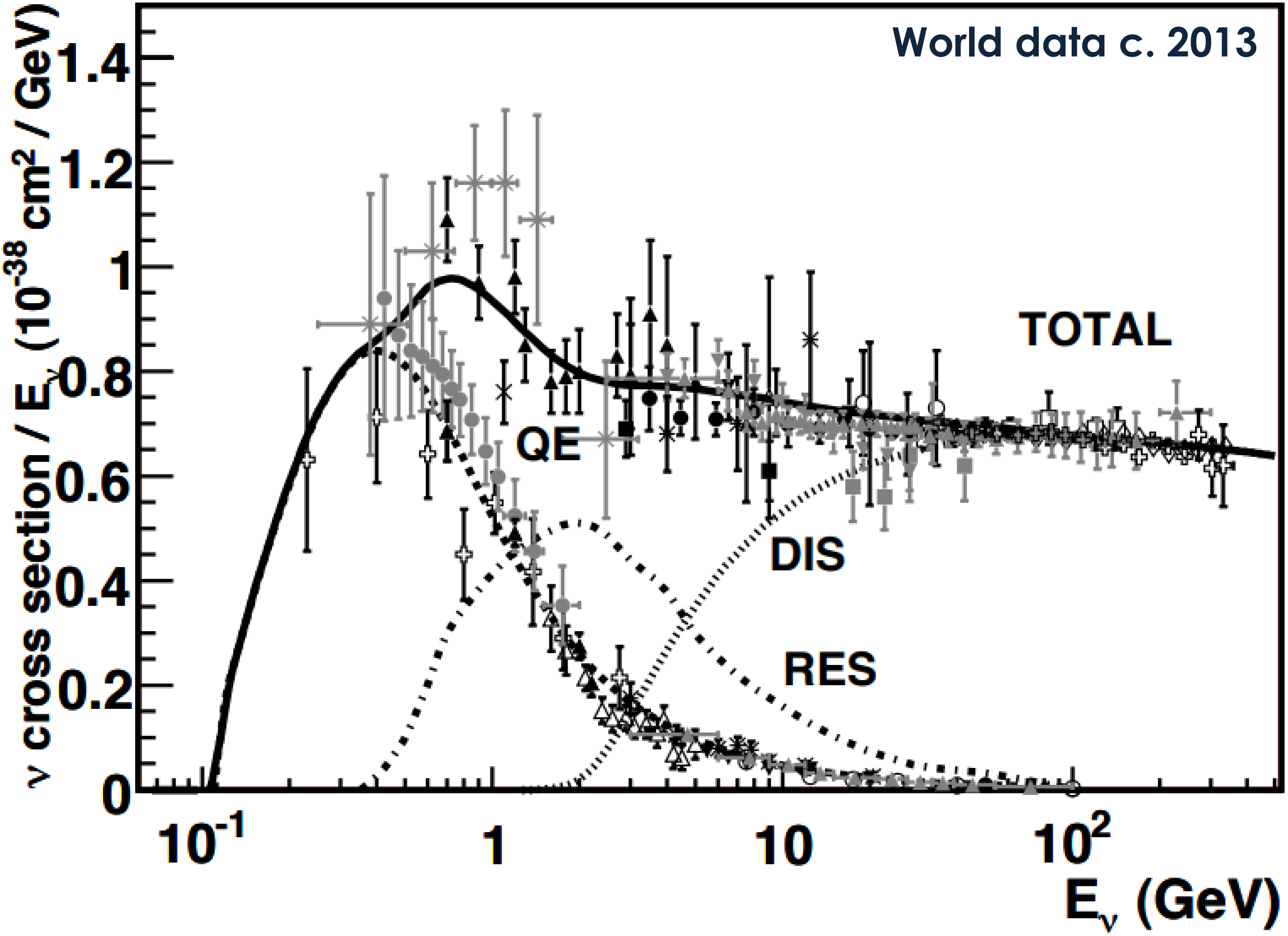
Phys. Rev. D **104**, 092007



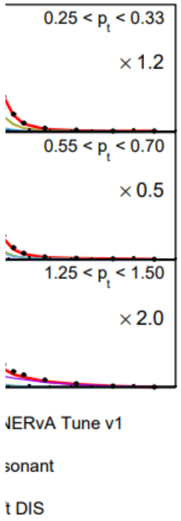
Well, have I got *νs* for you!

L. Crer



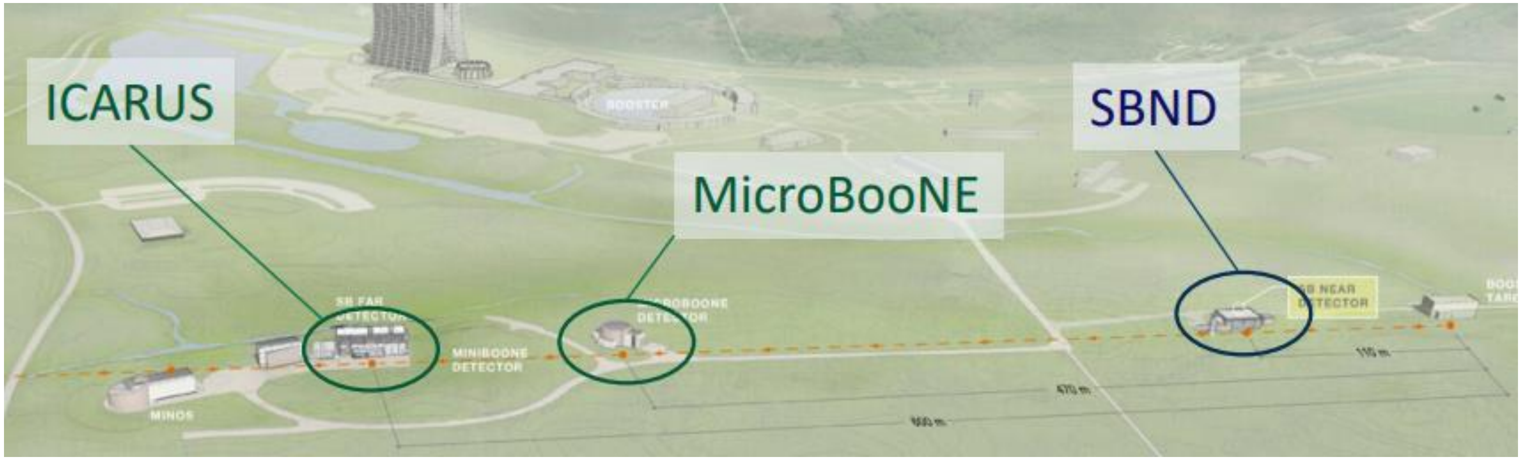


092007



A bright future for Argon

Short Baseline Program: Fermilab liquid Argon detectors in “Booster” beam (~0.8 GeV)



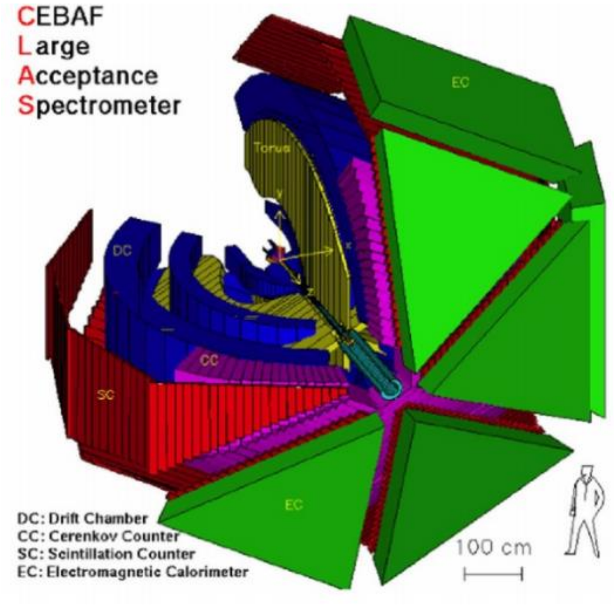
- **MicroBooNE**: already producing interesting results
- **ICARUS**: taking physics data
- **SBND:** enormous event rates coming soon (1M ν/y)

Beyond SBN:

- **DUNE “2x2” prototype:** measurements at DUNE energies

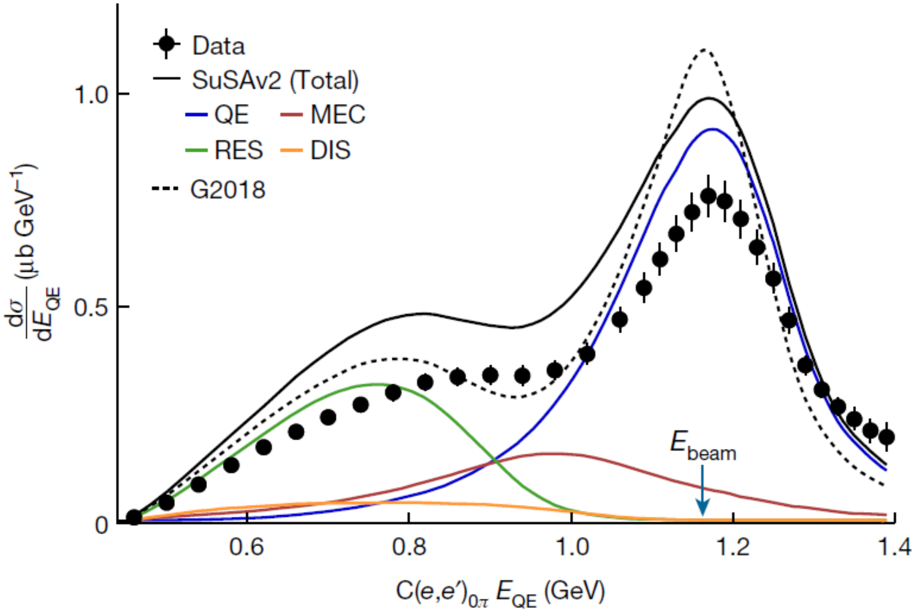
Tailored electron scattering

Nature **volume 599**, pages565–570 (2021)

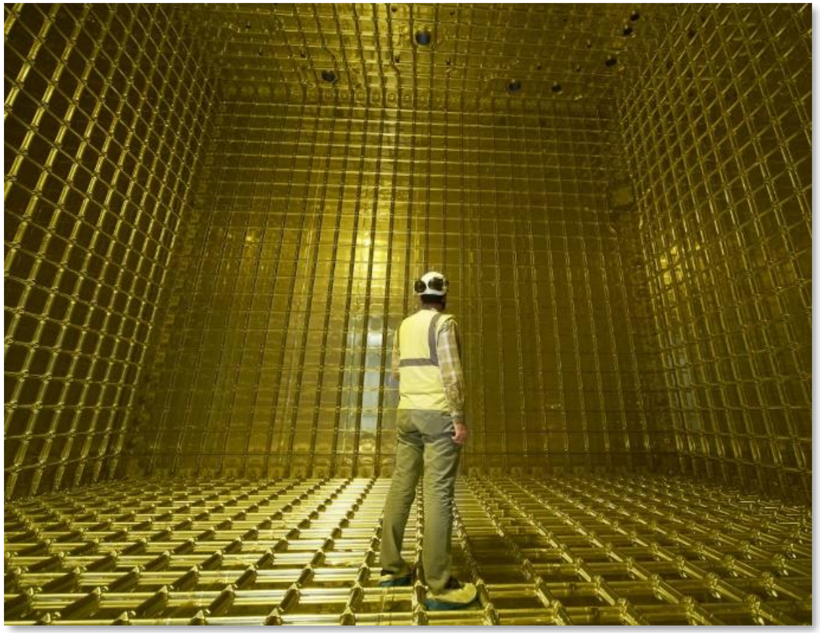


- Our models are becoming more able to make neutrino and electron scattering predictions in the same framework

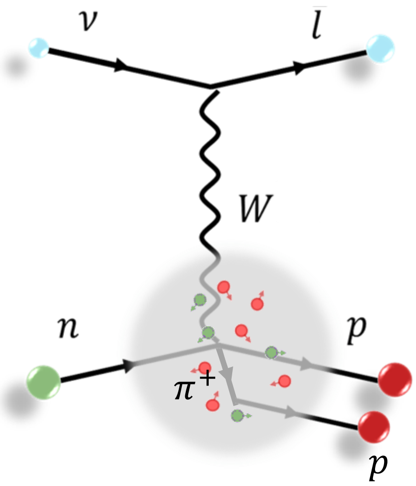
- New data from CLAS (e-scatting): specifically to help better understand neutrino scattering



Tailored hadron scattering

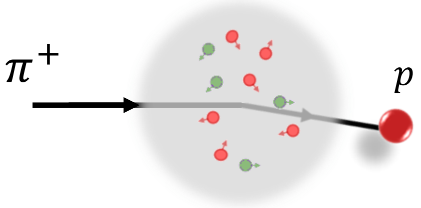


FSI in neutrino interactions

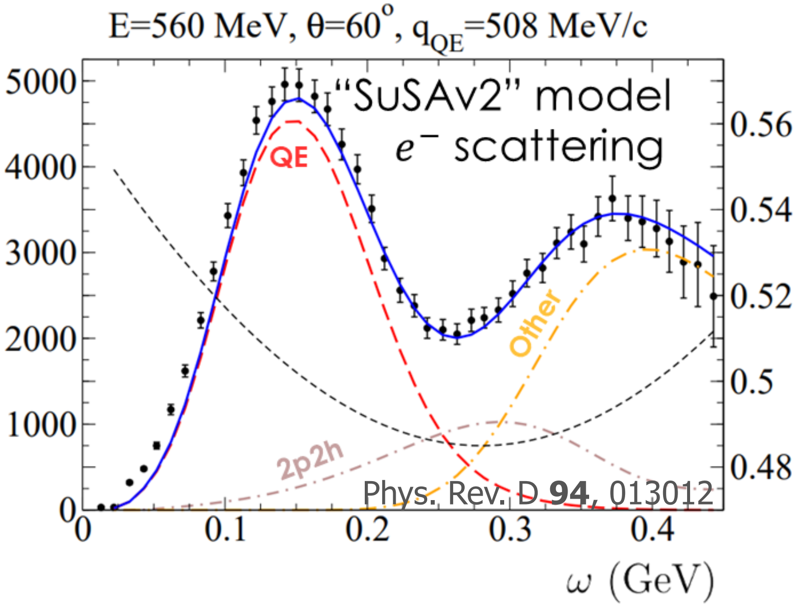


- Similarly, new hadron scattering data from ProtoDUNE and beyond can help constrain FSI processes

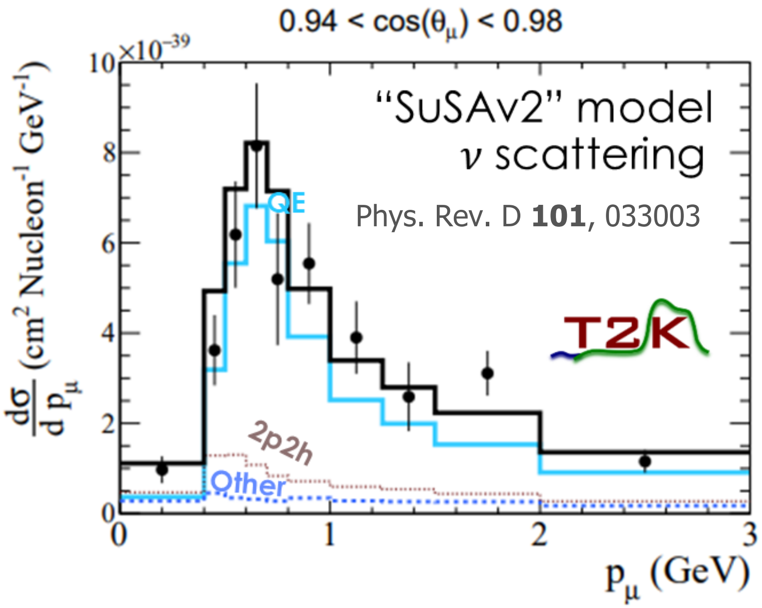
Hadron-nucleus scattering



New models, new constraints



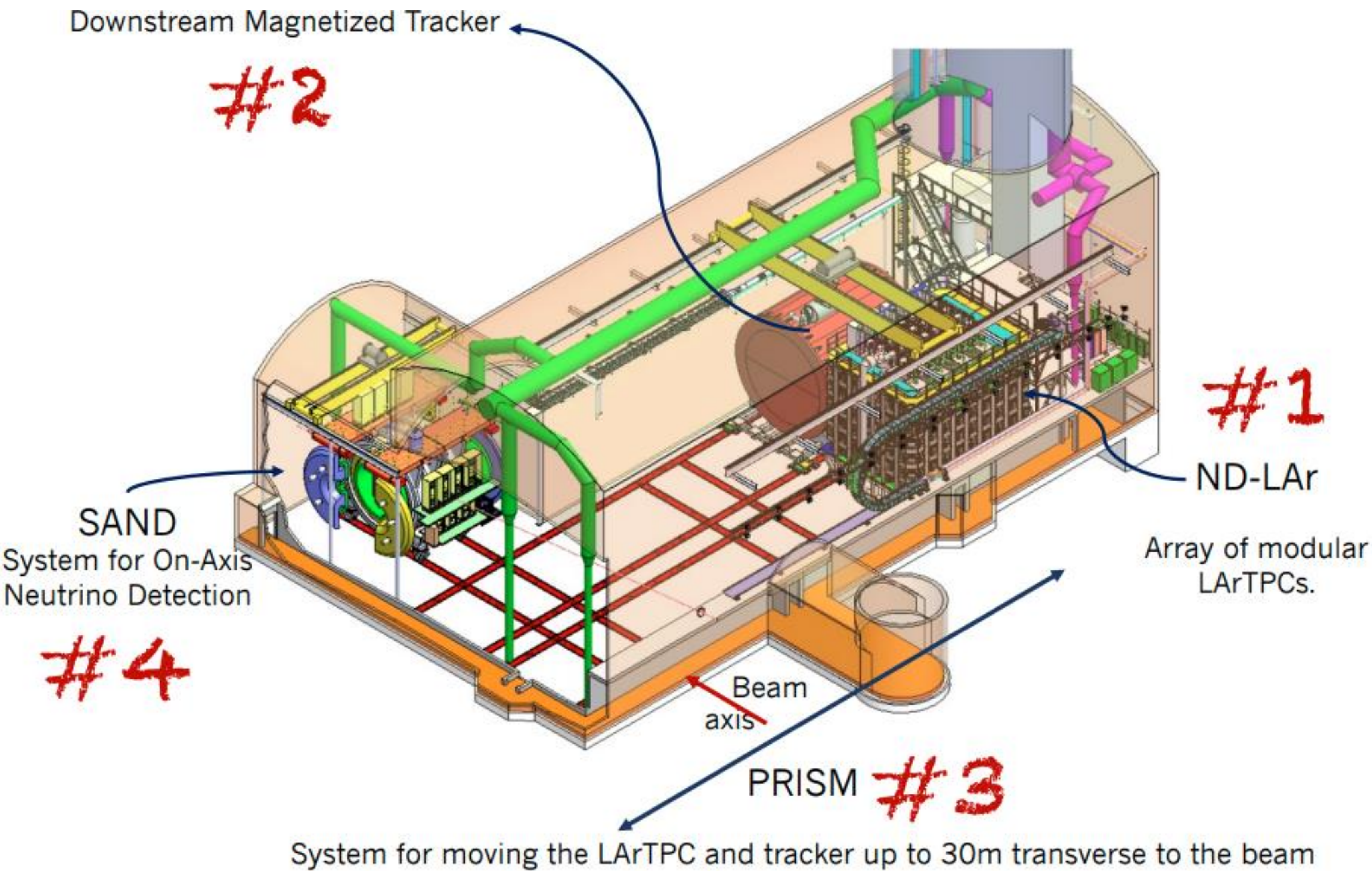
Phys. Rev. D **94**, 013012



Phys. Rev. D **101**, 033003

- New models, successful in describing electron scattering data, are now being implemented in neutrino interaction simulations
- Such models that describe e^- and ν interactions in the same framework can be directly constrained by precision e^- data
- New theoretical efforts are allowing models to be more predictive

Improved near detectors



Downstream Magnetized Tracker

#2

ND-LAr

#1

Array of modular LArTPCs.

SAND

System for On-Axis Neutrino Detection

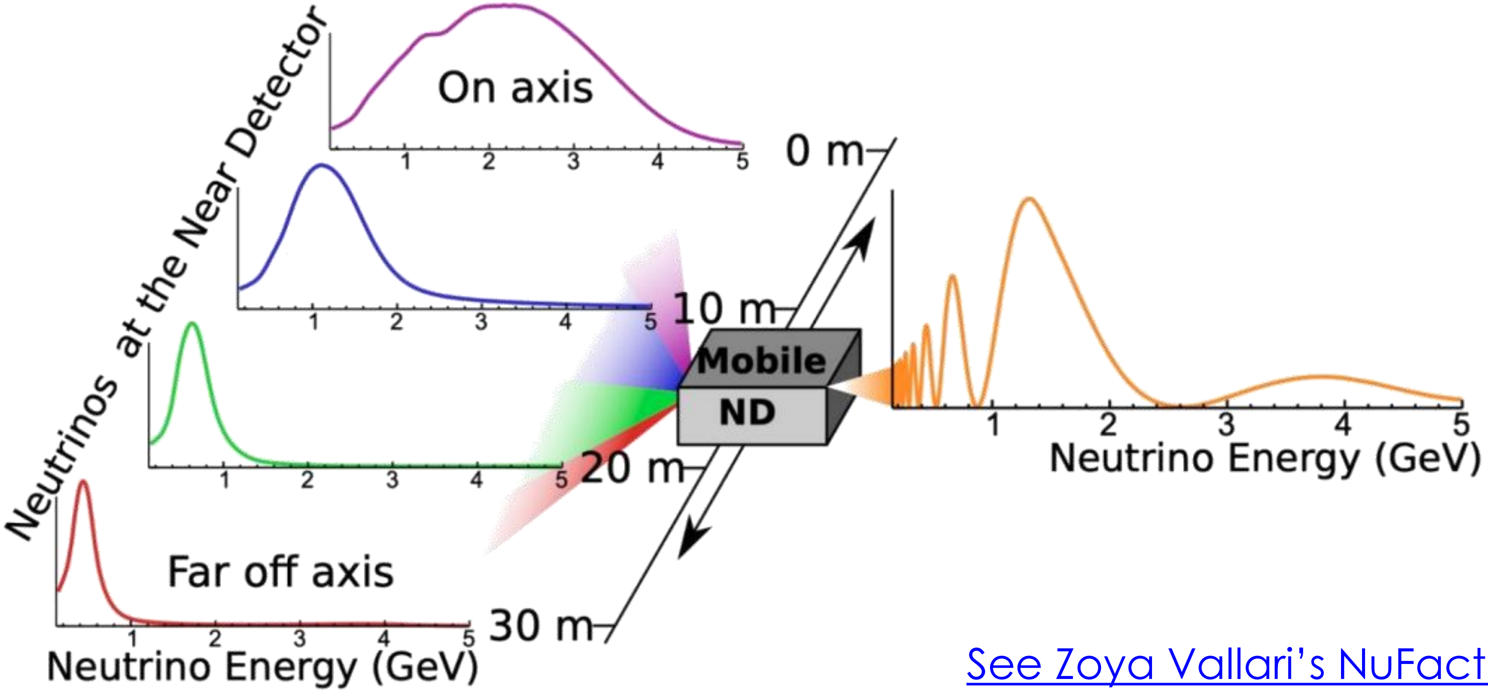
#4

PRISM #3

System for moving the LArTPC and tracker up to 30m transverse to the beam

DUNE PRISM

- A mobile 50 t liquid argon detector with a downstream spectrometer
  - ~59 M neutrino interactions per year!
- Moving the detector **changes the neutrino flux** in a predictable way, taking linear combinations of measurements at different positions allows a **construction of the oscillated spectrum at the near detector**
  - Better cancellation of uncertainties in oscillation measurements



See Zoya Vallari's NuFact 2022 talk

Summary

- A detailed understanding of neutrino-nucleus interactions is **crucial for current and future experiments** to realise their extraordinary goals (CP-violation, mass ordering, new BSM physics)

- This is a **challenging task**: neutrino interactions are complicated

- Mismodelling of subtle nuclear physics processes can cause **leading-order biases** on measurements of oscillations

- We've made **enormous progress** in modelling neutrino-nucleus interactions over the last 10 years, but **still have some way to go**

- Expect plenty of **exciting new results** and a continued exponential growth of the field in the run up to DUNE & Hyper-K.

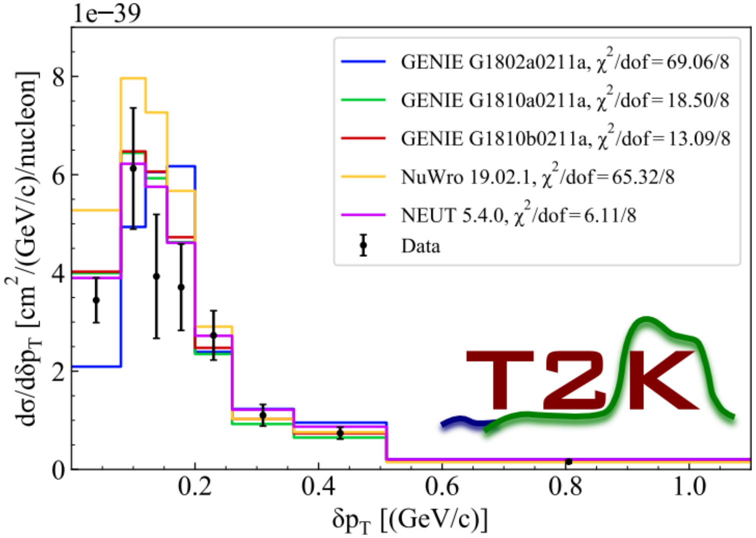
Backups

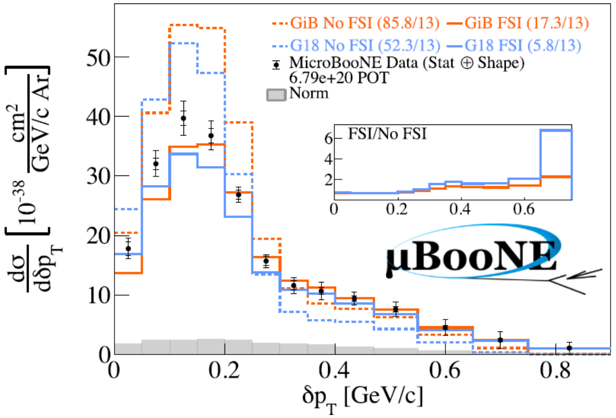
Generators vs data: a horror story

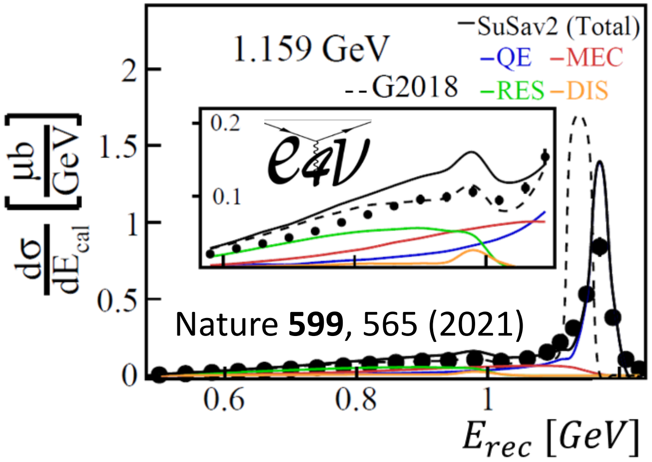
- No generator can come close to describing global lepton nucleus scattering data
  - All models are “wrong”, but they are each wrong in different ways

See many more informative generator comparisons in the TENSIONS 2019 report (arXiv:2112.09194)

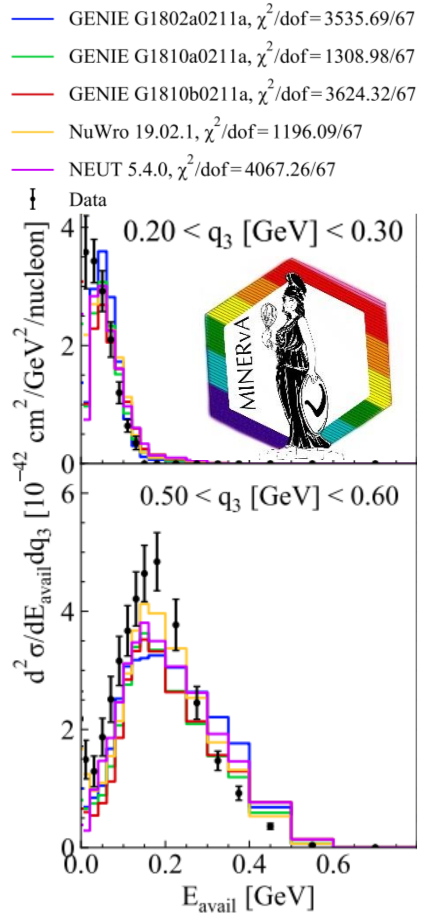
$$\chi^2_{allModels} \gg N_{bins}$$







Nature **599**, 565 (2021)



The hadronic current

Equation shamelessly lifted from
G. Perdue's other 2012 INSS lecture

$$J_H^\beta = \bar{u}_p \left[f_{1V}\gamma^\beta + i\frac{\xi f_{2V}}{2M}\sigma^{\beta\delta}q_\delta + \frac{f_{3V}}{M}q^\beta + f_A\gamma^\beta\gamma_5 + \frac{f_p}{M}q^\beta\gamma_5 + \frac{f_{3A}}{M}\left(P_p^\beta + P_n^\beta\right)\gamma_5 \right] u_n$$

$$M = \left(M_p + M_n\right)/2 \qquad q = p_\nu - p_\mu = P_p - P_n \qquad \xi = \mu_p - \mu_n \qquad \sigma^{\mu\nu} = \frac{i}{2}\left[\gamma^\mu, \gamma^\nu\right]$$

f_{3V}, f_{3A} are "second class currents", typically set to 0 for cross-section calculations, ξ is the difference between proton and neutron anomalous magnetic moments

- The other f factors are the "form factors" (read "fudge factors")
- These give us a way of parameterising the fact that the nucleon we interact with in an extended object.
- It turns out that the Fourier transform of form factors are what represents a physical distribution
- A dipole form factor represents an exponential distribution

$$f_A\left(q^2\right) = \frac{f_A(0)}{\left(1 - \frac{q^2}{M_A^2}\right)^2}$$

The hadronic current

Equation shamelessly lifted from
G. Perdue's other 2012 INSS lecture

$$J_H^\beta = \bar{u}_p \left[f_{1V}\gamma^\beta + i\frac{\xi f_{2V}}{2M}\sigma^{\beta\delta}q_\delta + \frac{f_{3V}}{M}q^\beta + f_A\gamma^\beta\gamma_5 + \frac{f_p}{M}q^\beta\gamma_5 + \frac{f_{3A}}{M}\left(P_p^\beta + P_n^\beta\right)\gamma_5 \right] u_n$$

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- f_{1V}, f_{2V} (vector form factors) can be extracted from electron scattering experiments. f_p can be related to f_A ("Partially Conserved Axial Current Hypothesis")
- f_A, we guess the form of! Usually we take a dipole with one free parameter: the infamous nucleon axial mass (M_A)
- We constrain the axial form factor with bubble chamber neutrino-nucleon (or light nucleus) data.

$$f_A\left(q^2\right) = \frac{f_A(0)}{\left(1 - \frac{q^2}{M_A^2}\right)^2}$$

Llewellyn-Smith CCQE

- Putting this all together gets us to the cross section

$$\frac{d\sigma}{d|q^2|}\begin{pmatrix}\nu n\to \ell^- p\\ \bar{\nu} p\to \ell^+ n\end{pmatrix}=\frac{M^2G^2\cos^2\theta_c}{8\pi E_\nu^2}\left[A(q^2)\mp B(q^2)\frac{(s-u)}{M^2}+\frac{C(q^2)(s-u)^2}{M^4}\right]$$

$$(s-u=4ME_\nu+q^2-m^2)\,.$$

Neutrino reactions at accelerator energies, Llewellyn Smith, 1972

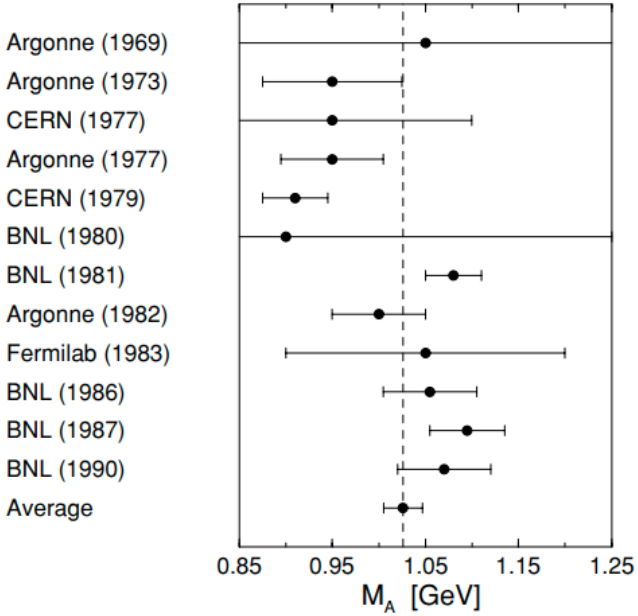
$$A \simeq \frac{t}{M^2}\left(|f_{1V}|^2-|f_A|^2\right)+\frac{t^2}{4M^2}\left(|f_{1V}|^2+\xi^2|f_{2V}|^2+|f_A|^2+4\xi\mathrm{Re}\left(f_{1V}f_{2V}^*\right)\right)$$
$$+\frac{t^3\xi^2}{16M^6}|f_{2V}|^2$$

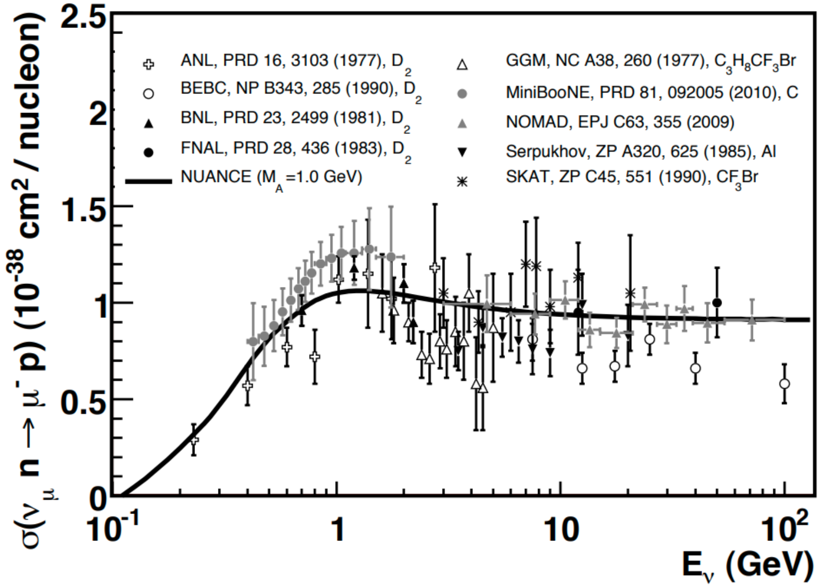
$$B \simeq \frac{1}{M^2}\left(\mathrm{Re}\left(f_{1V}f_A^*\right)+\xi\mathrm{Re}\left(f_{2V}f_A^*\right)\right)t \qquad C = \frac{1}{4}\left(|f_{1V}|^2+|f_A|^2-\frac{\xi^2|f_{2V}|^2}{4M^2}t\right)$$

The nucleon axial mass

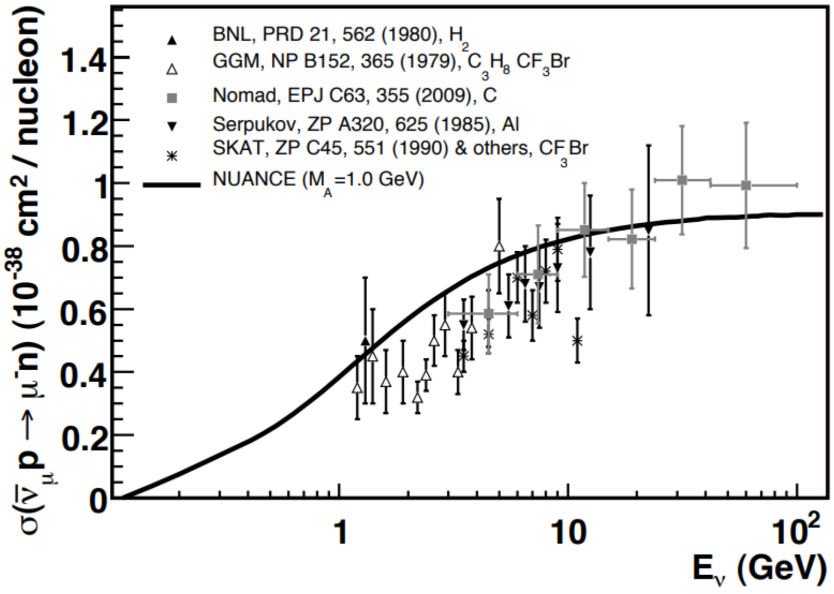
- We constrain the axial form factor with bubble chamber neutrino-nucleon (or light nucleus) data.
- The results seem pretty consistent with $M_A \sim 1\ GeV$

Journal of Physics G: Nuclear and Particle Physics, Volume 28, Number 1



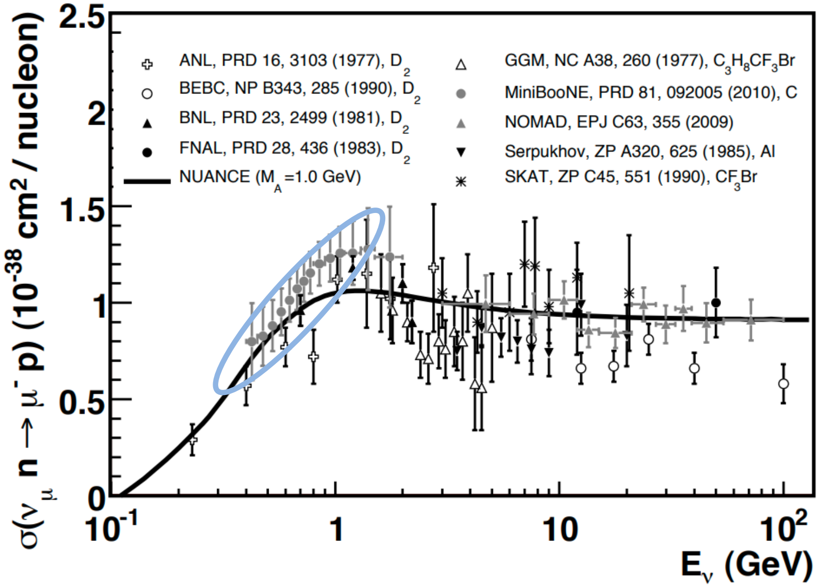


J. A. Formaggio and G. P. Zeller Rev. Mod. Phys. **84**, 1307



The nucleon axial mass “puzzle”

- Some heavier nuclear target experiments also try to measure M_A
- Now things don’t look so good
- We’ll come back to this …



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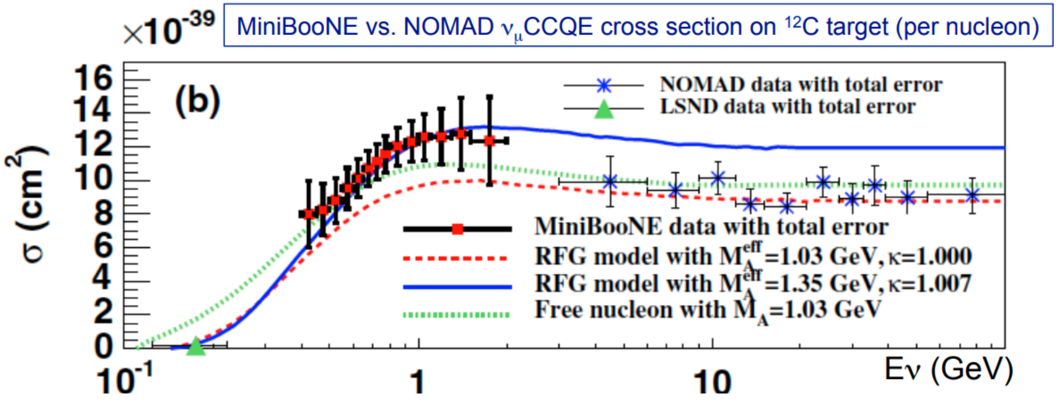
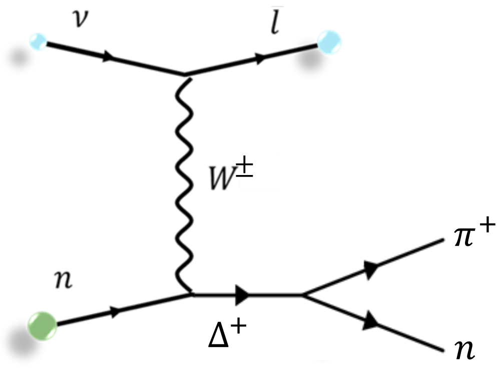


Resonant Pion Production

CCRES



CC Single Pion Production (SPP) final states

$$\nu_\mu\, p \to \mu^-\, p\, \pi^+, \qquad \overline{\nu}_\mu\, p \to \mu^+\, p\, \pi^-$$
$$\nu_\mu\, n \to \mu^-\, p\, \pi^0, \qquad \overline{\nu}_\mu\, p \to \mu^+\, n\, \pi^0$$
$$\nu_\mu\, n \to \mu^-\, n\, \pi^+, \qquad \overline{\nu}_\mu\, n \to \mu^+\, n\, \pi^-$$

D. Rein and L. Sehgal, Ann. Phys. 133, 79 (1981)

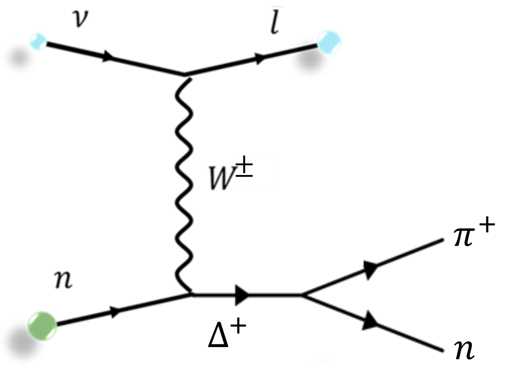
- Neutrinos can excite a nucleon into a resonance state, which then decays to give a nucleon + meson final state
- The dominant intermediate resonance is the $\Delta(1232)$ but others can contribute, as can non-resonant pion production
- And the contributions from each should have interference terms …
- Resonance models are complicated!
- Whilst CCQE scattering on the nucleon can described fully with one variable the multi-particle final state for SPP requires 4:

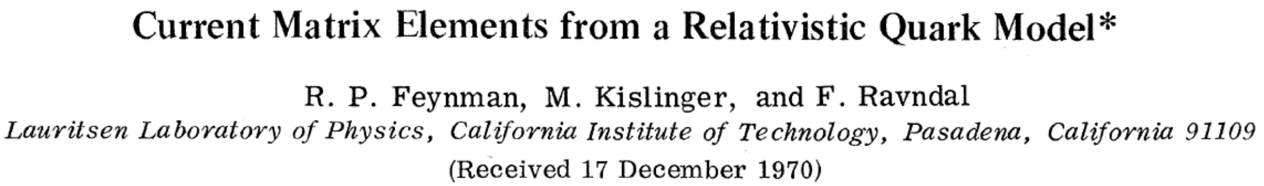
$$\frac{d\sigma}{dW\,dQ^2\,d\Omega_\pi}$$

Contains polar and azimuthal angle

Resonant Pion Production

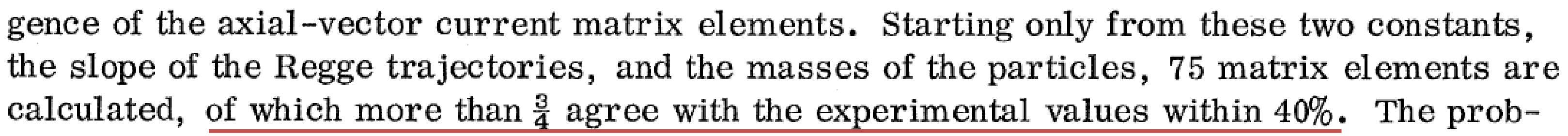
CCRES

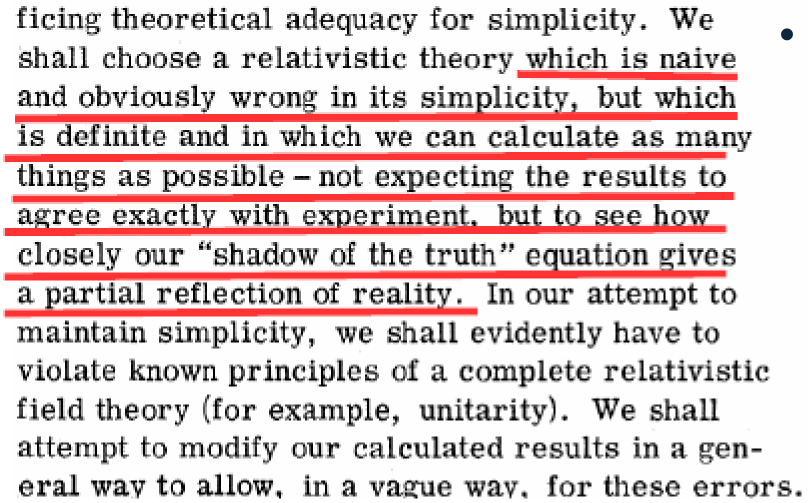



Current Matrix Elements from a Relativistic Quark Model*

R. P. Feynman, M. Kislinger, and F. Ravndal
Lauritsen Laboratory of Physics, California Institute of Technology, Pasadena, California 91109
(Received 17 December 1970)

- The model's used in today's neutrino experiments are based on an approximate model from the 1970s


gence of the axial-vector current matrix elements. Starting only from these two constants, the slope of the Regge trajectories, and the masses of the particles, 75 matrix elements are calculated, of which more than $\frac{3}{4}$ agree with the experimental values within 40%. The prob-


ficing theoretical adequacy for simplicity. We shall choose a relativistic theory which is naive and obviously wrong in its simplicity, but which is definite and in which we can calculate as many things as possible – not expecting the results to agree exactly with experiment, but to see how closely our "shadow of the truth" equation gives a partial reflection of reality. In our attempt to maintain simplicity, we shall evidently have to violate known principles of a complete relativistic field theory (for example, unitarity). We shall attempt to modify our calculated results in a general way to allow, in a vague way, for these errors.

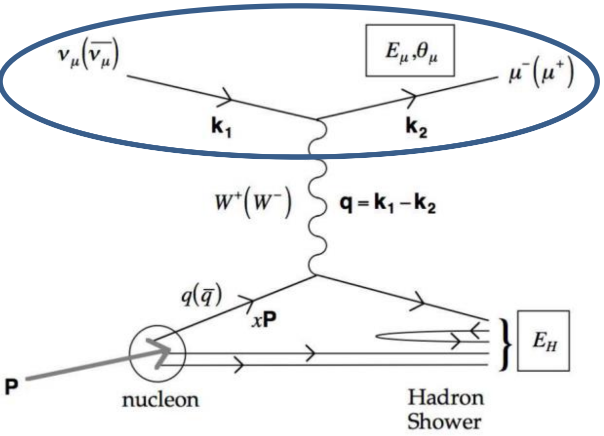
- The model includes its own form factors, including an axial part with an analogous M_A (and an additional uncertainty in the form factor numerator)

$$f_A\left(q^2\right) = \frac{f_A(0)}{\left(1 - \frac{q^2}{M_A^2}\right)^2}$$

- Theoretical developments are underway but it's safe to say CCRES is less well understood than CCQE!

Deep inelastic scattering

CCDIS



- Given enough energy, neutrinos can resolve the quarks within a nucleon. This is deep inelastic scattering.
- At high energies, the *inclusive* (i.e. integrating over possible hadronic final states) cross-section is fairly well understood (perturbative QCD):

$$\frac{d^2\sigma^{\nu,\overline{\nu}}}{dx\,dy} = \frac{G_F^2 M E_\nu}{\pi\left(1+Q^2/M_{W,Z}^2\right)^2}\left[\begin{array}{c}\frac{y^2}{2}2xF_1(x,Q^2)+\left(1-y-\frac{Mxy}{2E}\right)F_2(x,Q^2)\\ \pm y\left(1-\frac{y}{2}\right)xF_3(x,Q^2)\end{array}\right]$$

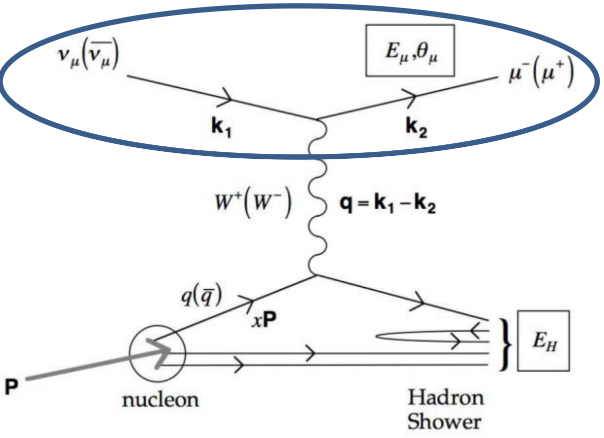
$$\begin{aligned} x &= \frac{Q^2}{2M\nu} = \frac{Q^2}{2ME_\nu y} \\ y &= E_{had}/E_\nu \\ Q^2 &= -m_\mu^2 + 2E_\nu(E_\mu - p_\mu \cos\theta_\mu) \end{aligned}$$

- The $F_i(x, Q^2)$ are nuclear structure functions, which are dimensionless and encompass the quark structure of nucleons
- The first two can be measured with e-scattering, the last one is from the weak VA interference term: only accessible with neutrinos!

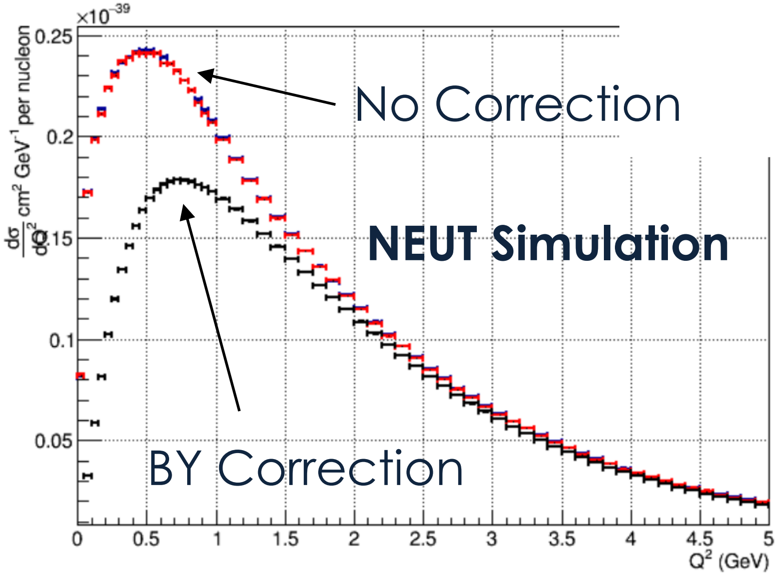
Deep inelastic scattering

https://arxiv.org/abs/1012.0261

CCDIS



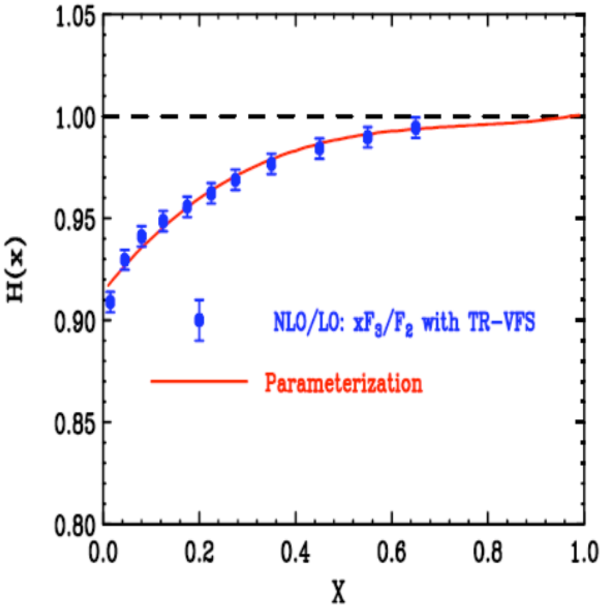
- At low energies (or actually low Q^2) QCD becomes non- perturbative.
- Bodek-Yang: extrapolate down to low Q^2 assuming some parametrised scaling. Fix the details with e-scatting, apply to ν- scattering
- But this is an empirical treatment that comes with uncertainties



No Correction

NEUT Simulation

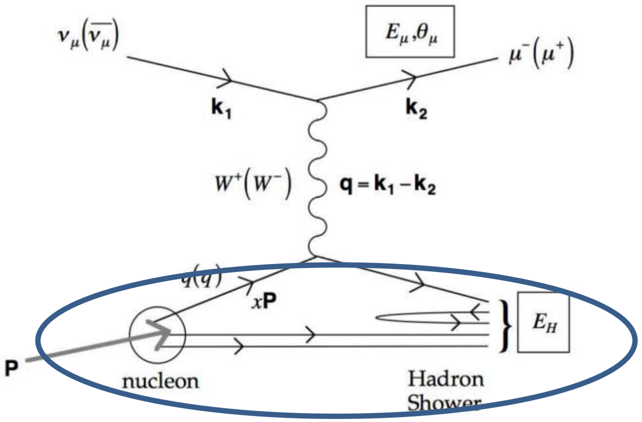
BY Correction



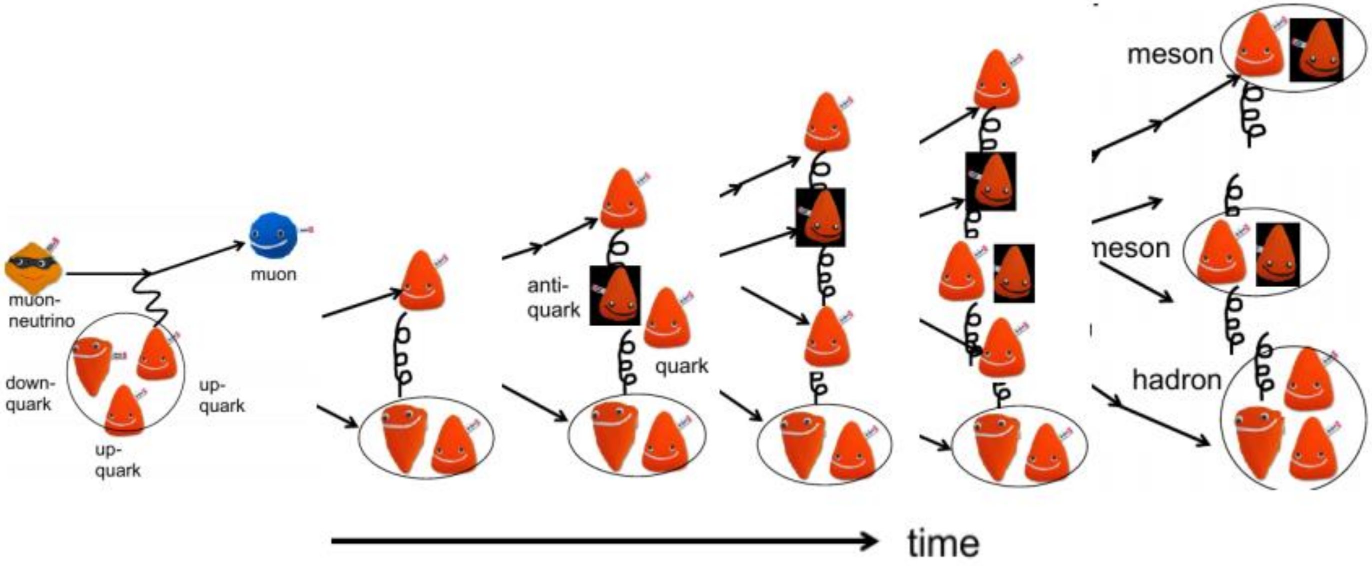
Deep inelastic scattering

https://arxiv.org/abs/1012.0261

CCDIS



- The hadronic side of DIS interactions requires more empirical treatments
- Often the PYTHIA generator is used, but this is really built for much higher energies than used in most neutrino experiments



T. Katori

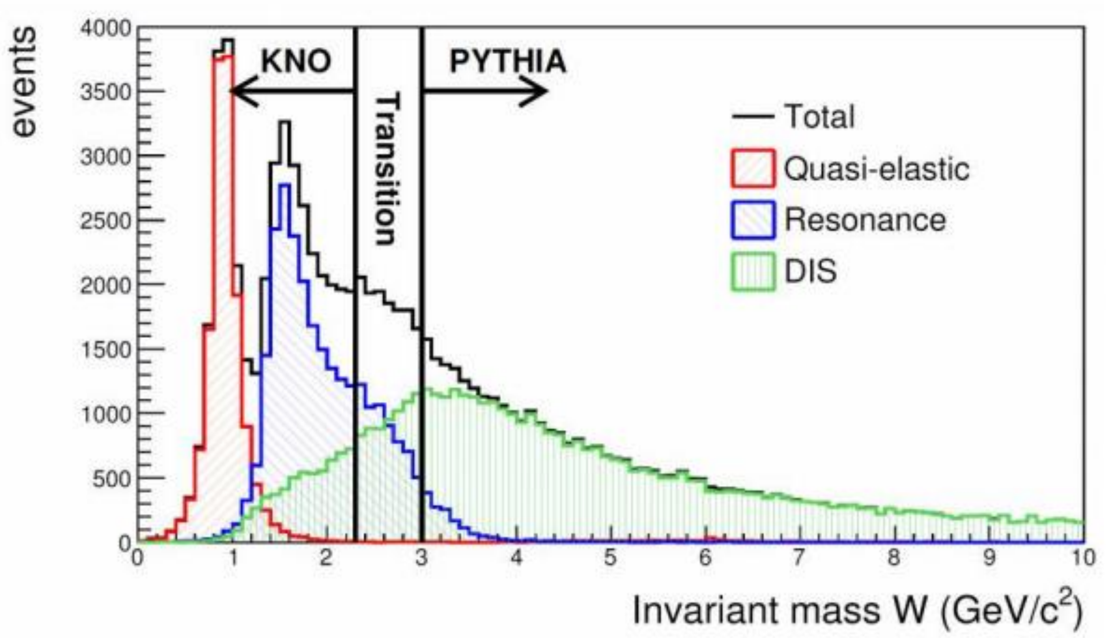
DIS-RES Transition Region

- There is no cut off where we better describe interactions in a DIS framework compared to In a RES framework

- In general we use models that extrapolate between regions which are definitely DIS (e.g. W>5 GeV) and that are definitively RES (e.g. W<2 GeV)

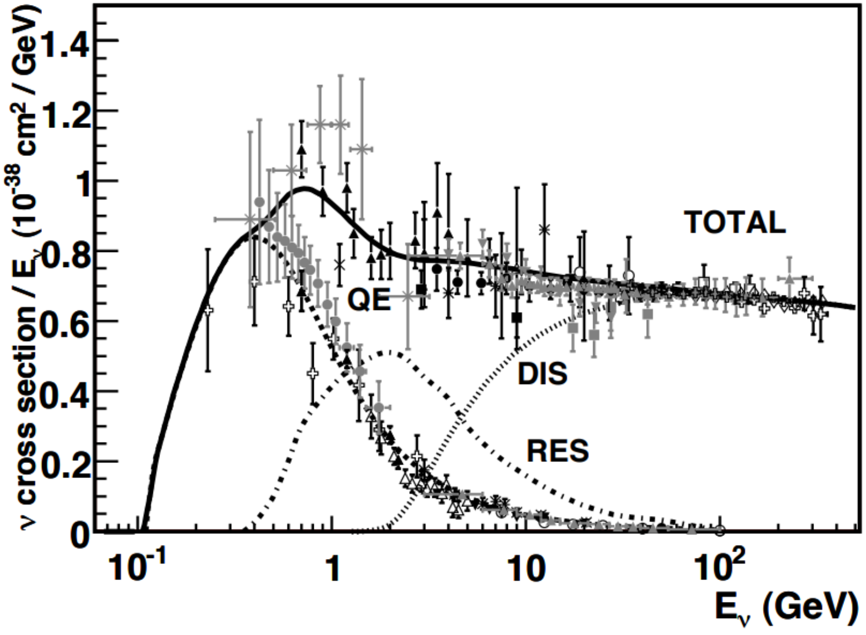
W = interaction invariant mass

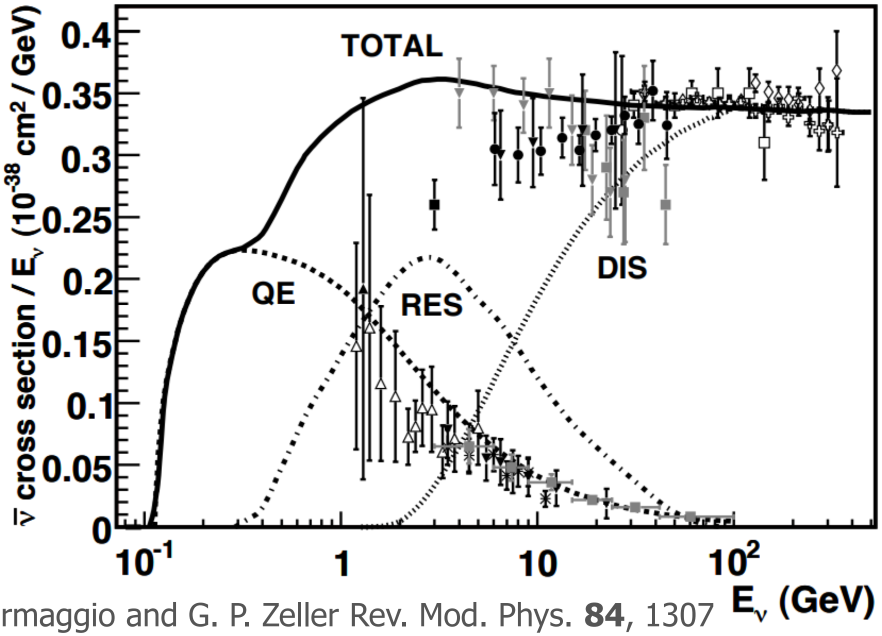
- But this is an imprecise method applied to a region that will be important for DUNE



Neutrino-nucleon cross sections

- Discussed neutrino-nucleon interactions
- But it's been a long time since we've measured this process!
- Almost all modern experiments use nuclear targets



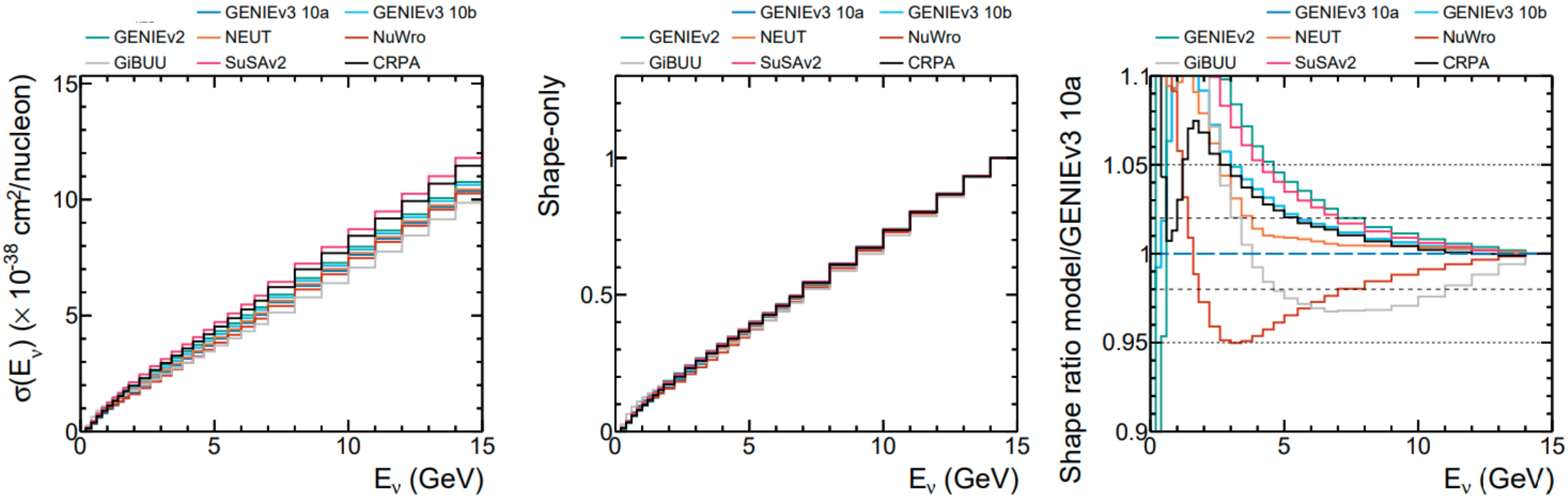


J. A. Formaggio and G. P. Zeller Rev. Mod. Phys. **84**, 1307

Energy dependence

Plots from
Wilkinson, Dolan, Pickering, Wret,
A substandard candle: the low-ν method at few-GeV neutrino energies
arXiv 2203.11821, accepted by EPJC

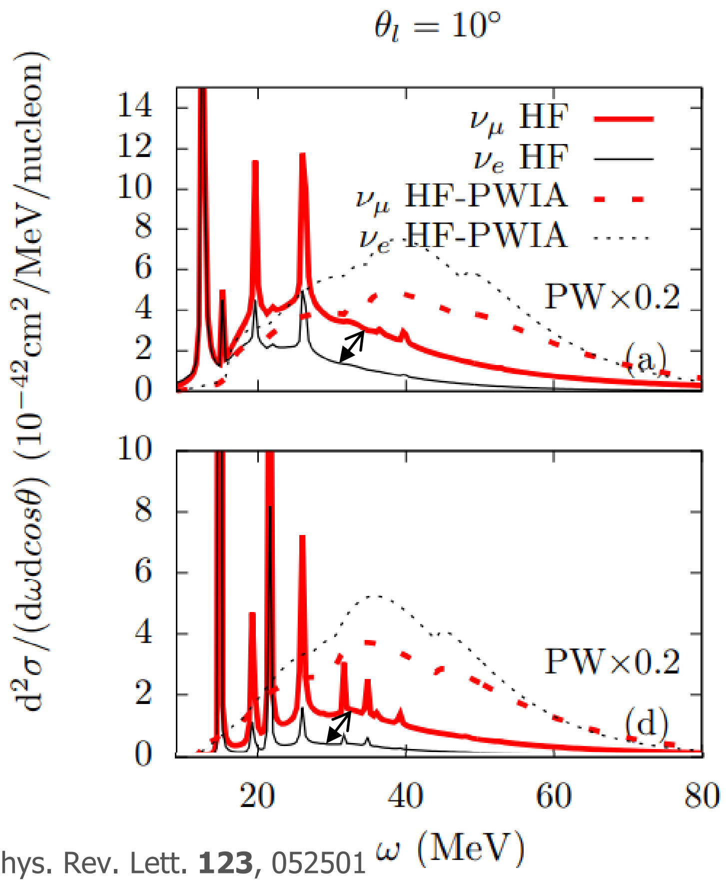
- What matters ND→FD extrapolation is the shape of total cross section as a function of E_ν
- Models differ by 5-10% in the region of interest for DUNE and Hyper-K
- Given expected statistics (~1000 ν_e, ~6000 ν_μ), this may be concerning
- Mitigation by direct measurements of cross section energy dependence (e.g. via multiple off-axis samples) is likely to be crucial



Nuclear effects and $\sigma(\nu_e)/\sigma(\nu_\mu)$

- Ratio of ν_e to ν_μ critical for future oscillation analyses
  - Measure ν_μ at ND but need to know about ν_e to measure δ_{CP}

- This is also subject to subtleties in the nuclear physics…



Phys. Rev. Lett. **123**, 052501

- If the outgoing nucleon exits the nucleus as a "plane wave" (no FSI): $\sigma(\nu_e) > \sigma(\nu_\mu)$

- If the outgoing nucleon is distorted by the nuclear potential (FSI): $\sigma(\nu_e) < \sigma(\nu_\mu)$

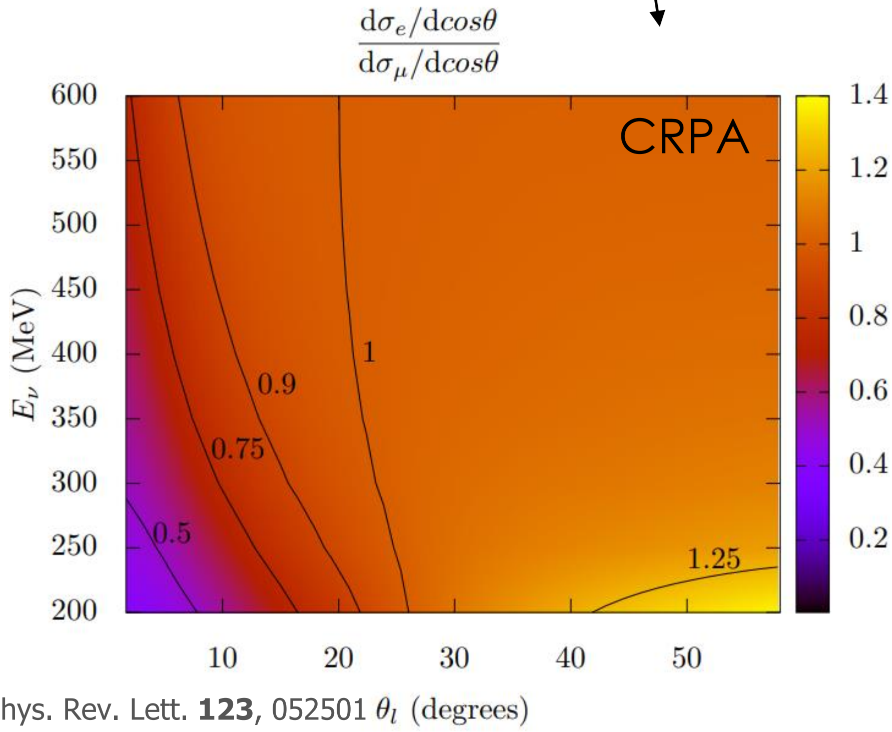
Nuclear effects and $\sigma(\nu_e)/\sigma(\nu_\mu)$

- Different models can predict quite different cross section ratios!
- Important for T2K/HK?

$$\frac{d\sigma_e/d\cos\theta}{d\sigma_\mu/d\cos\theta}$$

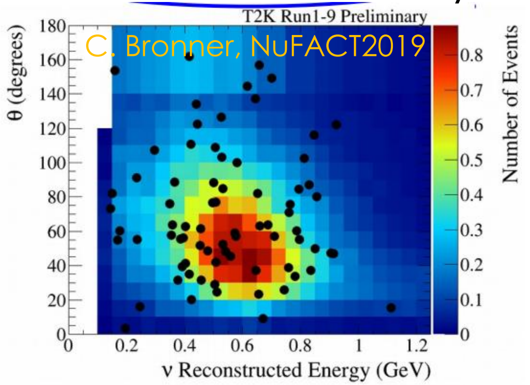
	$E_\nu = 200\ MeV$		$E_\nu = 600\ MeV$	
Model	5°	60°	5°	60°
RFG (w/PB)	0.64	1.61	0.97	1.03
SF (full)	1.41	1.92	1.04	1.03
CRPA	~0.5	~1.4	~0.9	~1.0

Tabulated from Phys. Rev. C **96**, 035501 and the left figure



Phys. Rev. Lett. **123**, 052501

These differences are predicted in regions that are relevant to T2K/HK oscillation analyses



Nuclear effects and $\sigma(\nu_e)/\sigma(\nu_\mu)$

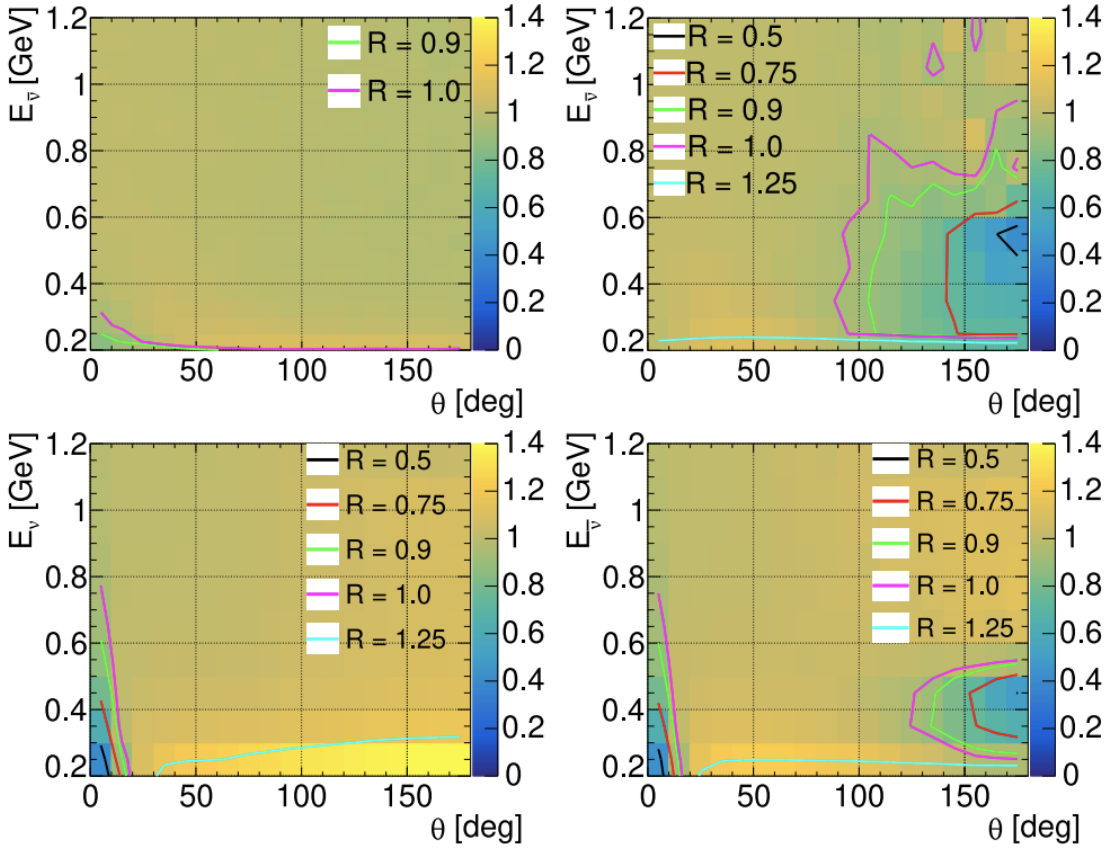
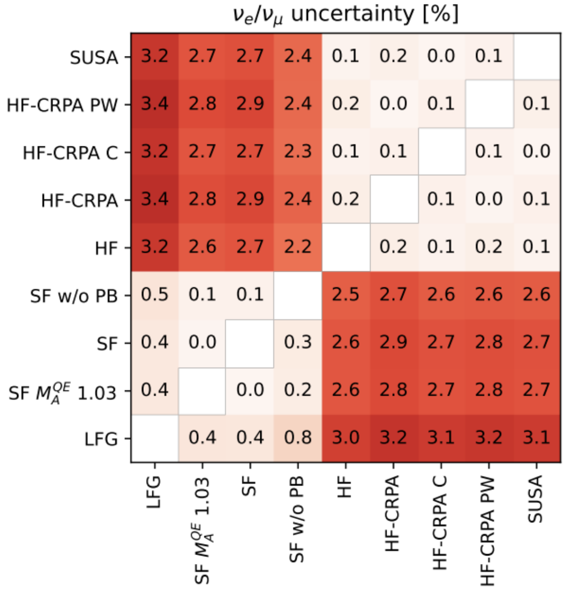


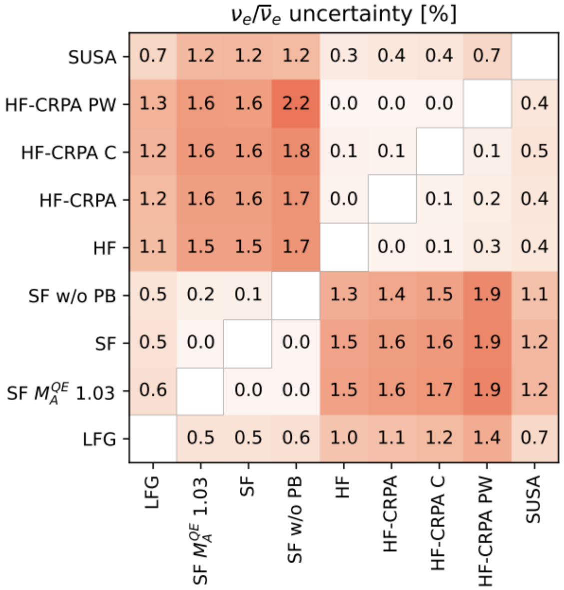
FIG. 1: $R^{\rm SF}_{\nu_e/\nu_\mu}$ (top-left), $R^{\rm HF\text{-}CRPA}_{\nu_e/\nu_\mu}$ (bottom left), $R^{\rm SF}_{\bar{\nu}_e/\bar{\nu}_\mu}$ (top-right) and $R^{\rm HF\text{-}CRPA}_{\bar{\nu}_e/\bar{\nu}_\mu}$ (bottom-right) are shown as a function of outgoing lepton angle and the neutrino energy. The contour lines highlight the regions where the ratio significantly deviates from unity.

arXiv:2301.08065



ν_e/ν_μ uncertainty [%]

	LFG	SF M_A^{QE} 1.03	SF	SF w/o PB	HF	HF-CRPA	HF-CRPA C	HF-CRPA PW	SUSA
SUSA	3.2	2.7	2.7	2.4	0.1	0.2	0.0	0.1	
HF-CRPA PW	3.4	2.8	2.9	2.4	0.2	0.0	0.1		0.1
HF-CRPA C	3.2	2.7	2.7	2.3	0.1	0.1		0.1	0.0
HF-CRPA	3.4	2.8	2.9	2.4	0.2		0.1	0.0	0.1
HF	3.2	2.6	2.7	2.2		0.2	0.1	0.2	0.1
SF w/o PB	0.5	0.1	0.1		2.5	2.7	2.6	2.6	2.6
SF	0.4	0.0		0.3	2.6	2.9	2.7	2.8	2.7
SF M_A^{QE} 1.03	0.4		0.0	0.2	2.6	2.8	2.7	2.8	2.7
LFG		0.4	0.4	0.8	3.0	3.2	3.1	3.2	3.1

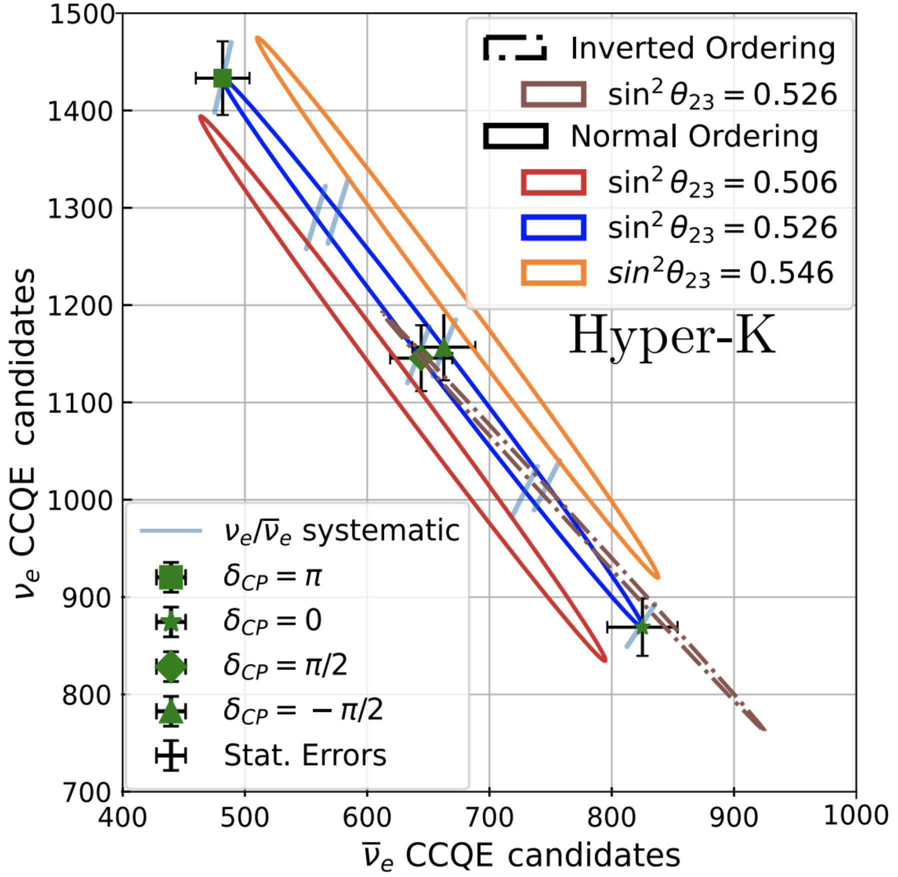


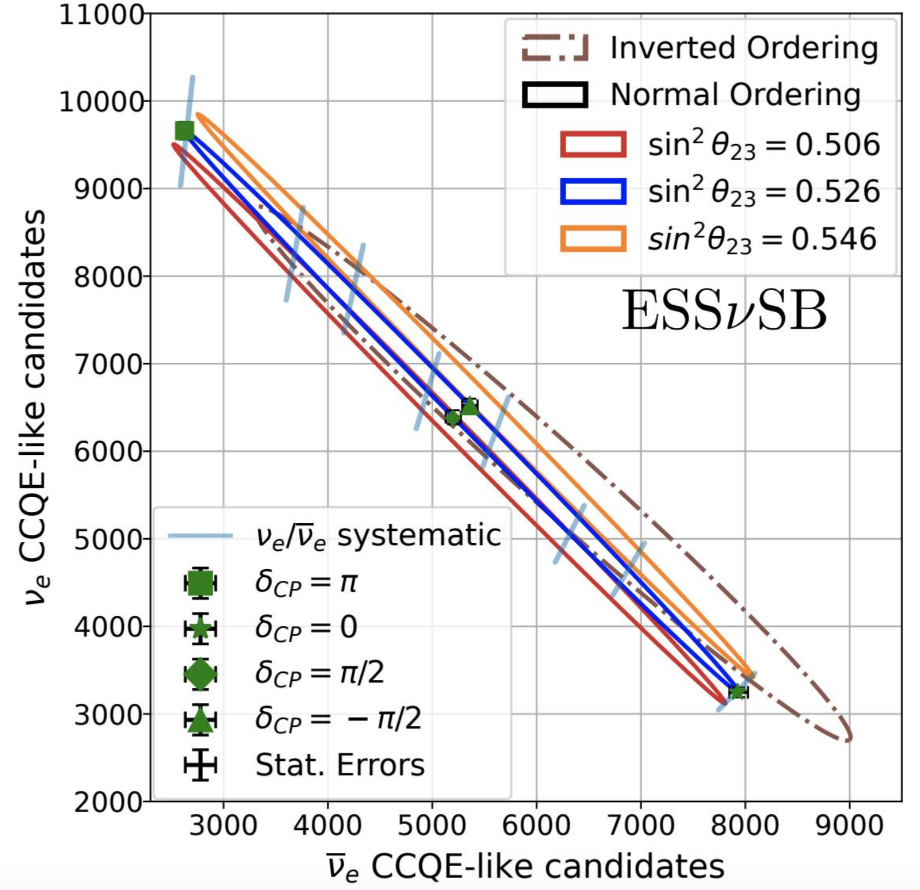
$\nu_e/\bar{\nu}_e$ uncertainty [%]

	LFG	SF M_A^{QE} 1.03	SF	SF w/o PB	HF	HF-CRPA	HF-CRPA C	HF-CRPA PW	SUSA
SUSA	0.7	1.2	1.2	1.2	0.3	0.4	0.4	0.7	
HF-CRPA PW	1.3	1.6	1.6	2.2	0.0	0.0	0.0		0.4
HF-CRPA C	1.2	1.6	1.6	1.8	0.1	0.1		0.1	0.5
HF-CRPA	1.2	1.6	1.6	1.7	0.0		0.1	0.2	0.4
HF	1.1	1.5	1.5	1.7		0.0	0.1	0.3	0.4
SF w/o PB	0.5	0.2	0.1		1.3	1.4	1.5	1.9	1.1
SF	0.5	0.0		0.0	1.5	1.6	1.6	1.9	1.2
SF M_A^{QE} 1.03	0.6		0.0	0.0	1.5	1.6	1.7	1.9	1.2
LFG		0.5	0.5	0.6	1.0	1.1	1.2	1.4	0.7

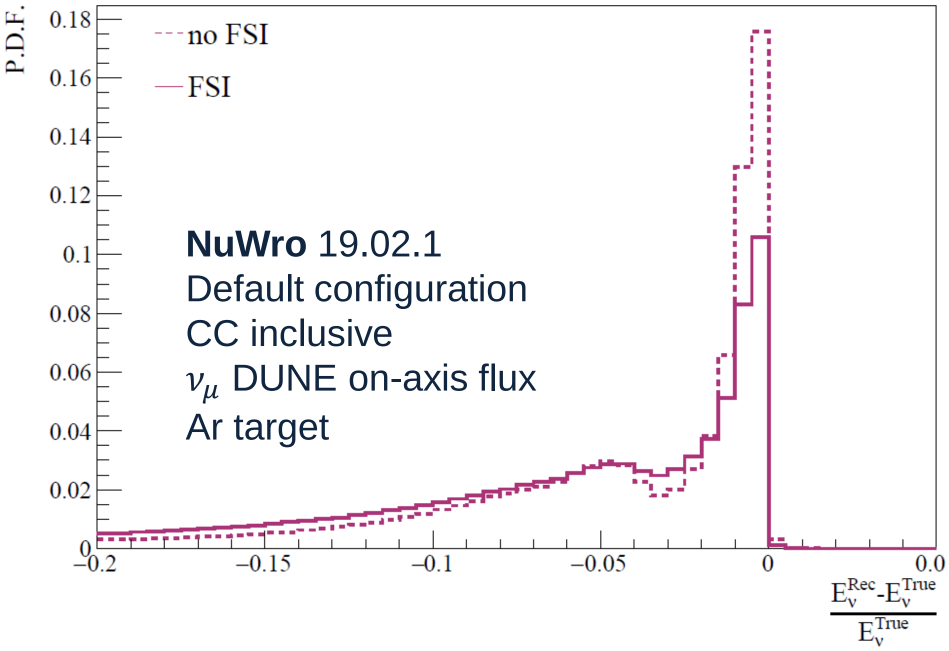
Nuclear effects and $\sigma(\nu_e)/\sigma(\nu_\mu)$

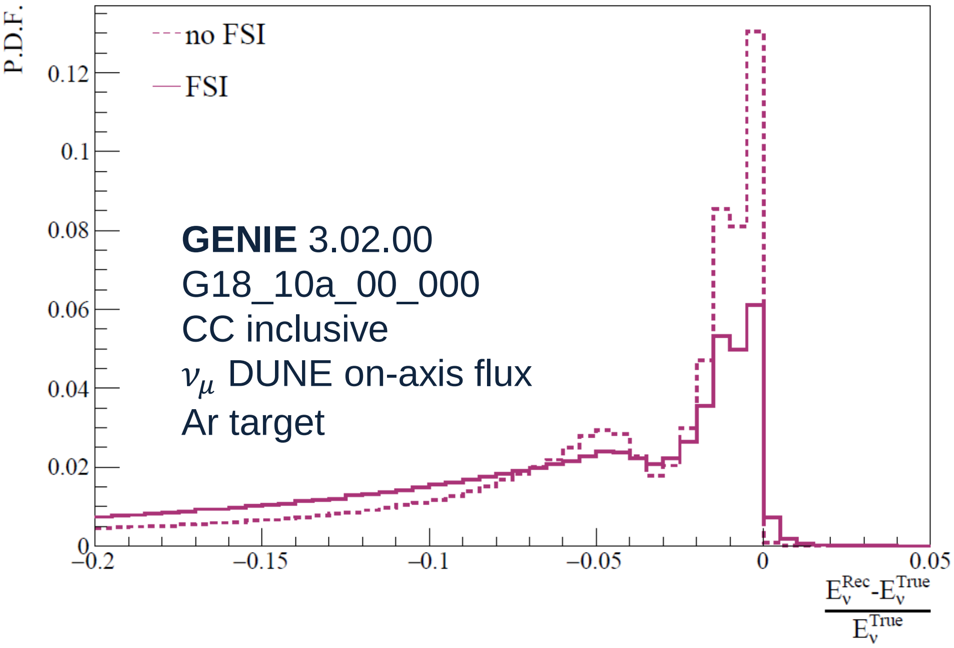
arXiv:2301.08065





FSI and neutrino energy reconstruction

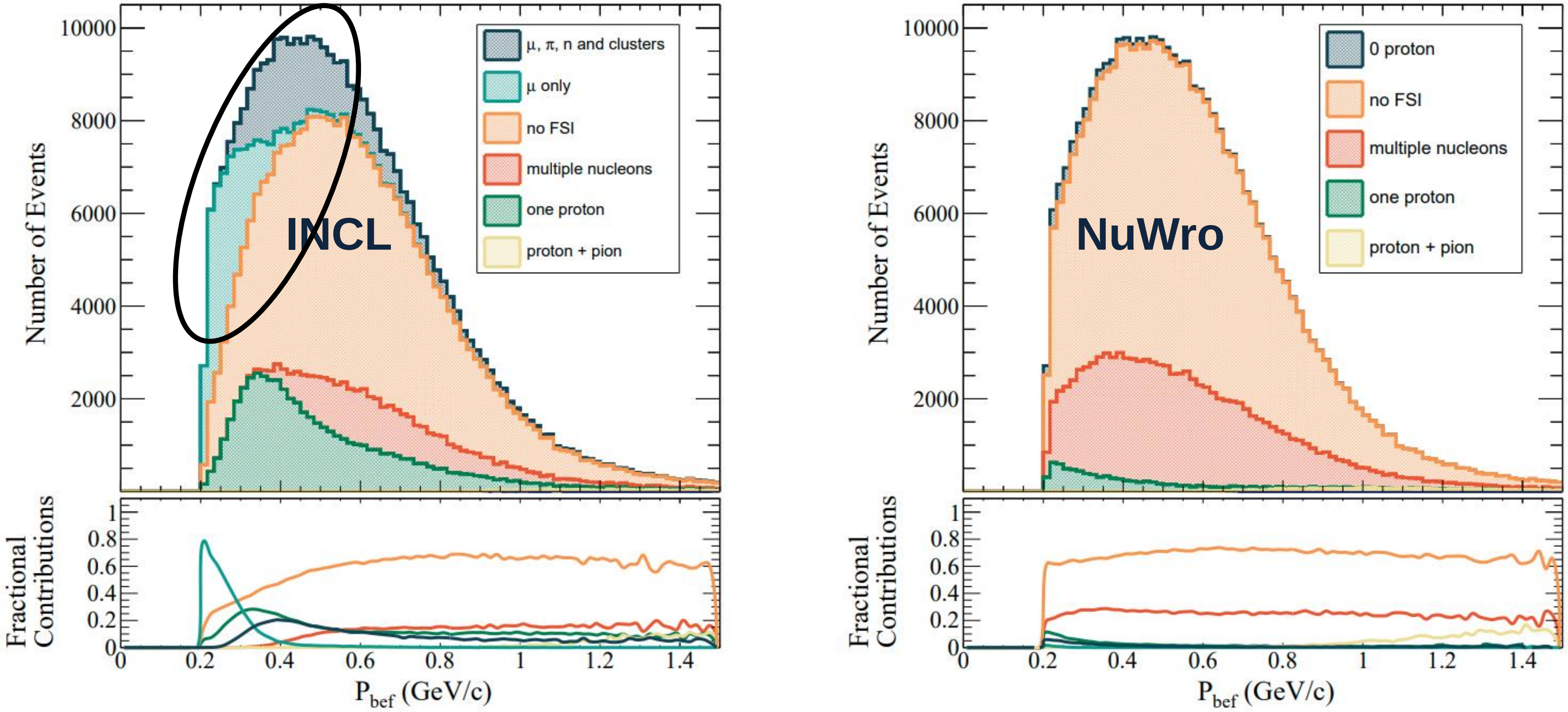




Advanced FSI cascades

Plots from
Ershova et al.,
Study of FSI of protons ith INCL and NuWro cascade models
Phys. Rev. D **106**, 032009

- More advanced treatment of FSIs is available via the INCL model (Phys. Rev. C **87** 014606)
- INCL's treatment of nucleon absorption and nuclear cluster production gives a different distribution of energy among outgoing hadrons
- Might expect a significant impact on neutrino energy smearing



FSI beyond the cascade

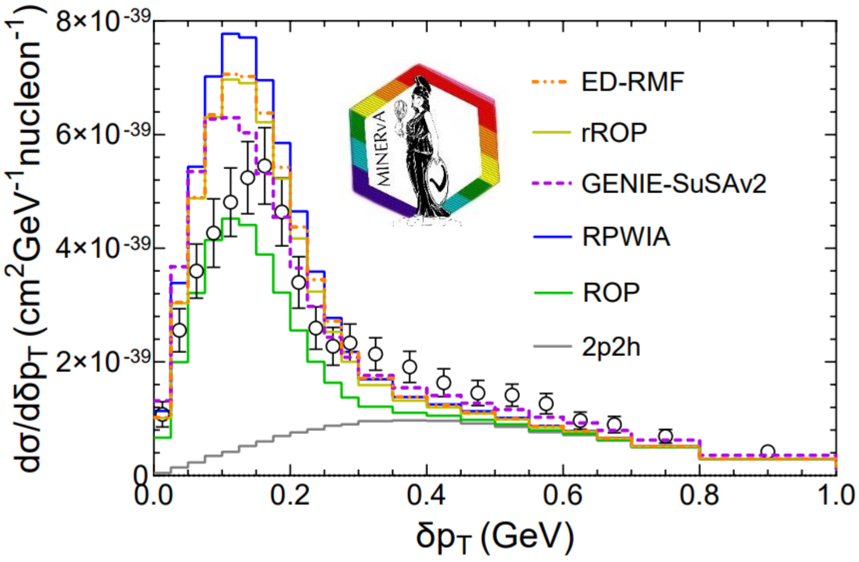
Plots from:
Franco-Patino et al.,
arXiv:2207.02086

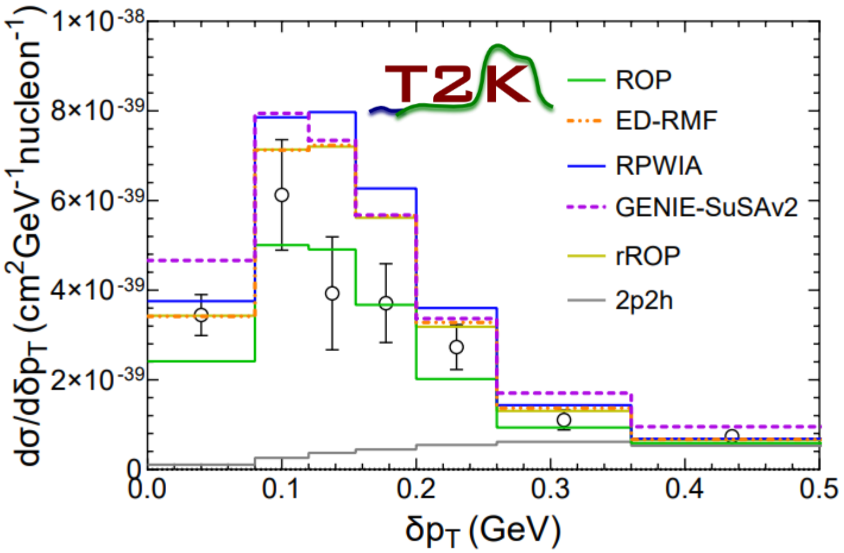
See also:
Nikolakopoulos et al.,
Phys. Rev. C **105**, 054603

- Instead of cascades, FSI can be modelled via a distortion of the outgoing nucleon wave function by a nuclear potential

- Recent theory effort has allowed a calculation of exclusive observables with such treatments
    - Example below: missing transverse momentum
    - In general: high δp_T → more missing hadronic energy → larger E_ν reconstruction bias

- Key conclusions
    - Significant differences in predictions for different nuclear potentials
    - Sometimes all of these deviate strongly from the cascade approach





Impact on analyses

Plot from:
DUNE physics TDR,
arXiv:2002.03005

- DUNE runs a study where it fits as data a model where **20% of final state proton energy** in its nominal model instead **goes into neutrons**
  - A **plausible consequence of alternative FSI models**
- At the same time, the cross section is altered to leave the proton momentum distribution unchanged
  - Another plausible change to the cross section model
- The result: **a large bias in oscillation parameters**
- Possible mitigation by creative use of the near detector
  - Off-axis samples
  - Additional nuclear targets