



University of
Zurich^{UZH}

ELUCIDATING THE NATURE OF NEUTRINOS: THE STATE-OF-THE ART IN SEARCHES FOR NEUTRINOLESS DOUBLE BETA DECAY

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FERMILAB COLLOQUIUM
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FIRST, SOME HISTORY

- ▶ Zürich, December 4, 1930: Wolfgang Pauli, a 30 years old professor at the ETH (since 1928), writes perhaps one of the most famous letters in modern physics: "Dear radioactive ladies and gentlemen..."
- ▶ The letter was addressed mainly to Lise Meitner*, who had been working on radioactivity since 1907 and was attending a meeting in Tübingen (Pauli could not attend, because "a ball which takes place in Zürich the night of the sixth to sevenths of December makes my presence here indispensable")

Original - Photocopy of PLC 039
Abschrift/15.12.36

Offener Brief an die Gruppe der Radioaktiven bei der Gauvereins-Tagung zu Tübingen.

Abschrift



Physikalisches Institut
der Eidg. Technischen Hochschule
Zürich

Zürich, 4. Des. 1930
Cloriastrasse

- Pauli was suggesting "a desperate way out" of some paradox that had arisen in the nascent field of nuclear physics
- He was proposing "a terrible thing" - a new subatomic particle, the neutrino, a particle "which can not be detected"
- In 1930, only the electron, the proton and the photon were known, and Pauli's idea was quite radical

Liebe Radioaktive Damen und Herren,

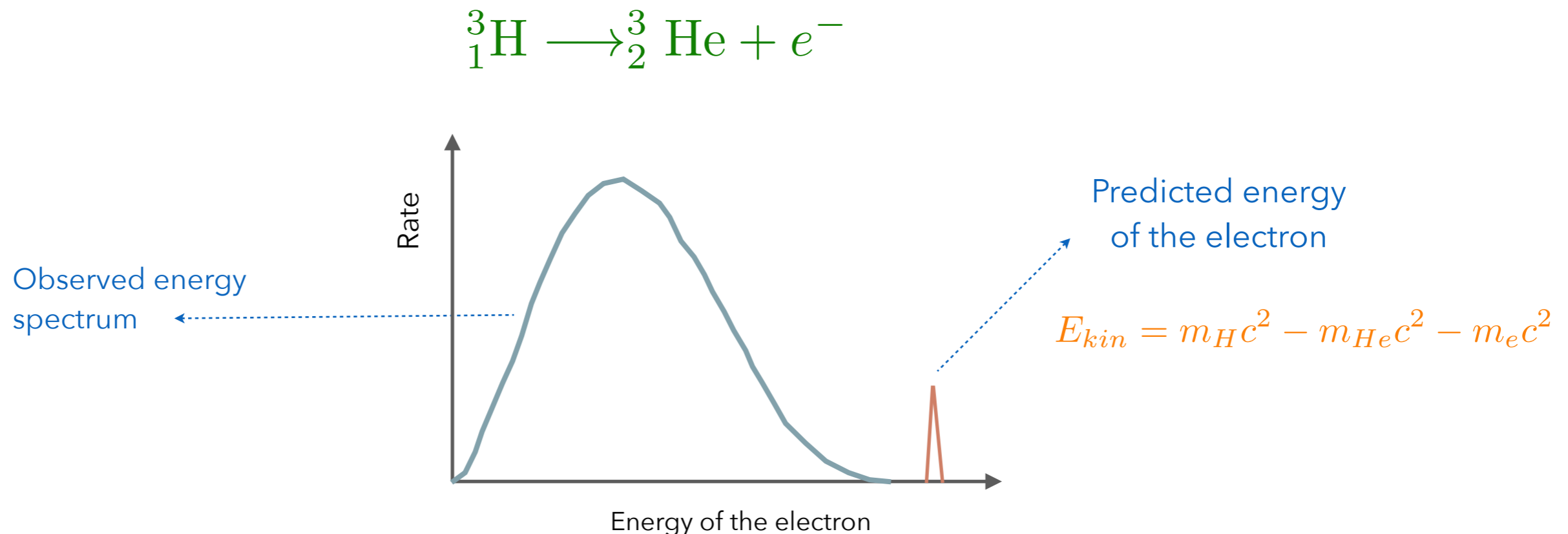
Wie der Ueberbringer dieser Zeilen, den ich kuldvollst anhören bitte, Ihnen des näheren auseinandersetzen wird, bin ich angesichts der "falschen" Statistik der β - und α -Kerne, sowie des kontinuierlichen β -Spektrums auf einen verweifelten Ausweg "selbst" (1) der Statistik und den Energiesatz hin, die Möglichkeit, es könnten elektrisch neutrale Teilchen nennen will, in den Kernen existieren, die den β -Strahlen und das Ausschliessungsprinzip befolgen und ausserdem noch dadurch unterscheiden, dass sie eine Masse haben, die grösser ist als die Elektronenmasse. Die Masse der Neutronen ist in der Ordnung wie die Elektronenmasse klein und nur ein wenig grösser als die Protonenmasse. Das kontinuierliche β -Spektrum ist verständlich unter der Annahme, dass beim β -Zerfall jedesmal ein Neutron in ein Elektron und ein Neutrino zerfällt. Die Summe der Energien von Neutron und Elektron

*Lise Meitner had made Pauli aware of the β -decay problem

THE PARADOX WAS... “THE ENERGY CRISIS”

- ▶ It had been observed by experimental physicists that some nuclei are not stable, but decay under the emission of “beta rays” (electrons)
- ▶ The energy of the emitted electrons could be measured - **the spectrum was continuous**
- ▶ This seemed to violate a respected laws in physics: the conservation of energy and momentum



ONLY ONE REASONABLE WAY OUT...

- ▶ A new particle: the neutrino (Pauli: "my foolish child"). It would share the energy with the electron, but would not be observed because of its incredible weak interaction with matter



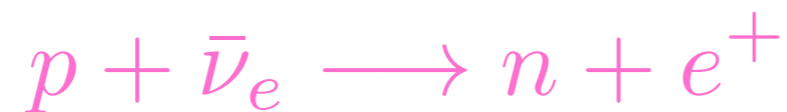
- ▶ Niels Bohr, 1934: *"I must confess that I don't really feel fully convinced of the physical existence of the neutrino"*
- ▶ Arthur Eddington, 1939: *"I am not much impressed by the neutrino theory.... Dare I say that physicists will not have sufficient ingenuity to make neutrinos?"*
- ▶ *Thus, while the idea was considered by many as a very useful hypothesis, few* believed it is a real particle (or that it can ever be detected**), until...*

***Enrico Fermi** did take the idea seriously and formulated a theoretical basis for the interaction between a neutrino, an electron, a proton and a neutron (1934, Z. Phys. 88)

** **Hans Bethe**: "there is a considerable evidence for the neutrino hypothesis. Unfortunately, all this evidence is indirect; and more unfortunately, there seems at present to be no way of getting any direct evidence."

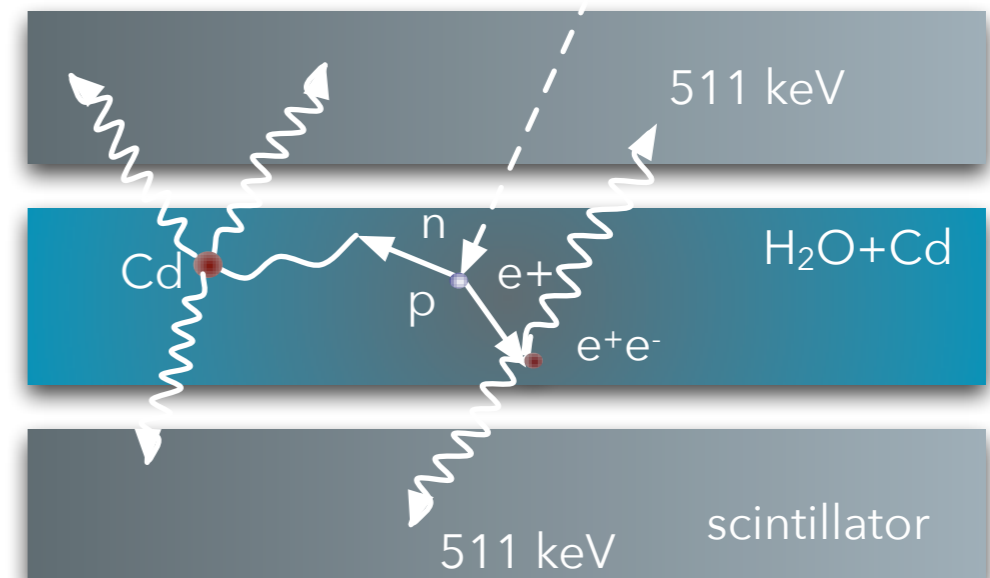
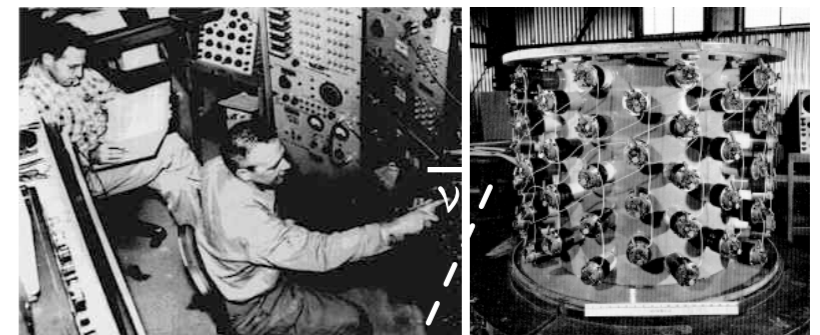
NEUTRINO DETECTION

- ▶ ... some 30 years later in 1956, when Clyde Cowan and Frederick Reines started the "Project Poltergeist" and finally detected (anti)neutrinos at the Savannah River Reactor in South Carolina



- ▶ Detector: 400 l water + CdCl₂ seen by 90 photodetectors

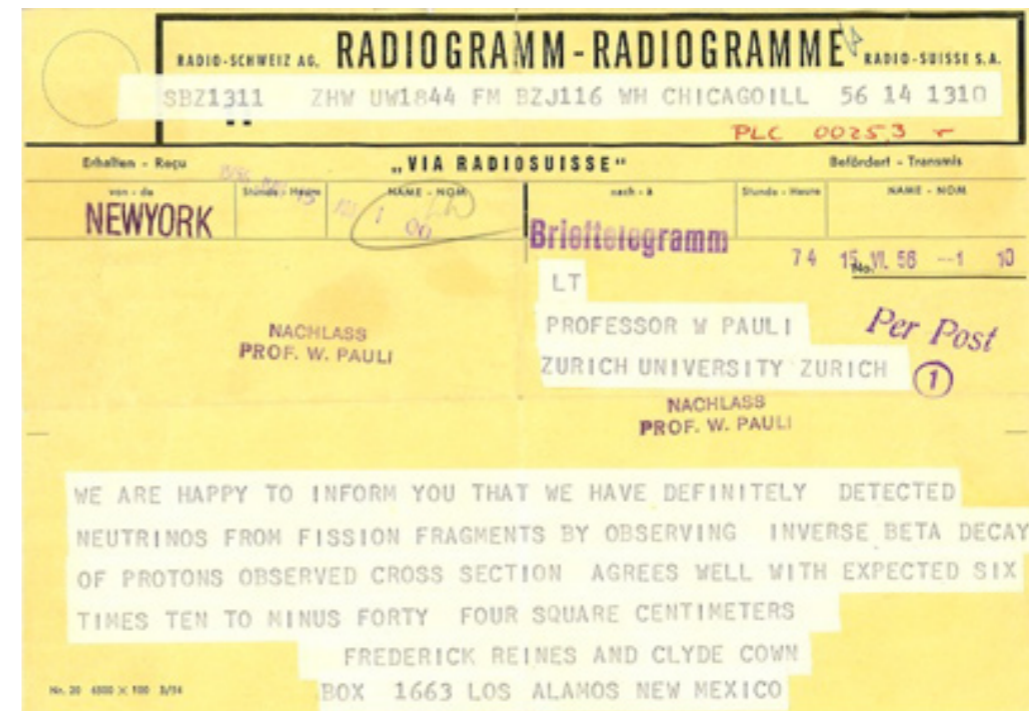
Detection via delayed (a few μs) coincidence reaction:



Nobel Prize 1995: (1/2) to Frederick Reines "for the detection of the neutrino"

A RADIOGRAMME TO PAULI, A SHORT ANSWER...

- ▶ June 1956: Pauli was at a CERN Symposium, and announced the most exciting news of the meeting* - he had just received a telegram from Cowan & Reines
 - ◉ "We are happy to inform you that we have definitely detected neutrinos..."
- ▶ Pauli's reply: "Thanks for message. Everything comes to him who knows how to wait."

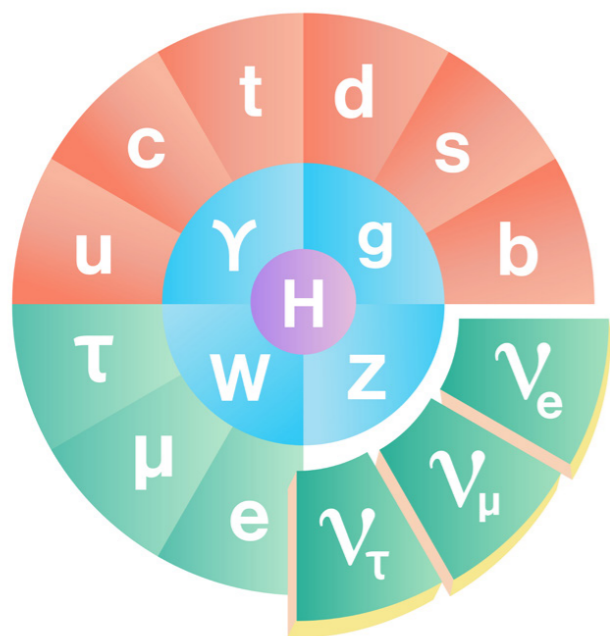


Frederick REINES and Clyde COWAN
 Box 1663, LOS ALAMOS, New Mexico
 Thanks for message. Everything comes to
 him who knows how to wait.
 Pauli

*See: Cecilia Jarlskog, "Birth of the neutrinos, from Pauli to the Reines-Cowan experiment", 2019 - International Conference of the History of the Neutrino

WHAT ARE NEUTRINOS?

- ▶ Elementary particles in the Standard Model which only interact via the weak interaction (they participate in charged current interactions other with the corresponding charged lepton)
 - The interactions are of "V-A" type: neutrinos are left-handed, anti-neutrinos are right-handed
- ▶ In the SM: flavour lepton number is conserved and neutrinos are exactly massless
 - Today many known sources of neutrinos



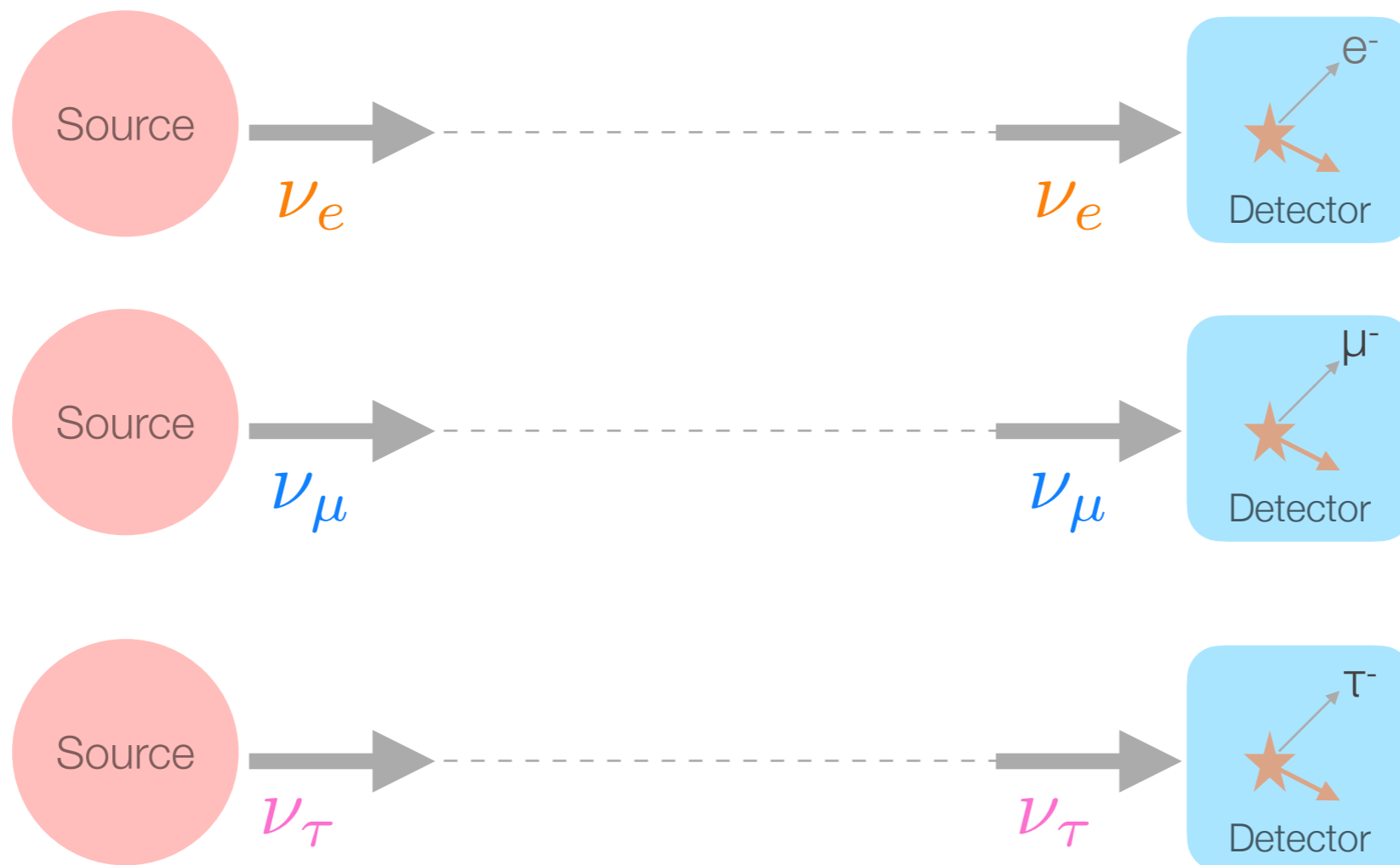
$$\begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} \nu_e \\ e^- \end{pmatrix}$$



WHAT DO WE KNOW ABOUT NEUTRINOS?

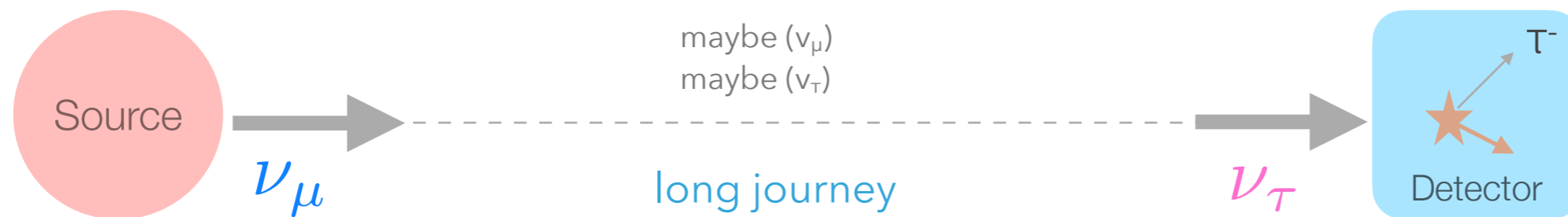
- ▶ They come in 3 flavours

ν_e electron ν_μ muon ν_τ tau



WHAT DO WE KNOW ABOUT NEUTRINOS?

- ▶ However when they propagate over macroscopic distances, they oscillate between flavours



- ▶ This is a well-studied effect in quantum mechanics
- ▶ It means that flavour is not conserved over macroscopic distances (ν states with different flavours ν_α mix with ν states with different masses ν_i)

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sum_{i,j=1}^3 U_{\alpha i} U_{\beta i}^* U_{\alpha j}^* U_{\beta j} \exp\left(-i \frac{m_{\nu_i}^2 - m_{\nu_j}^2}{2E} x\right)$$

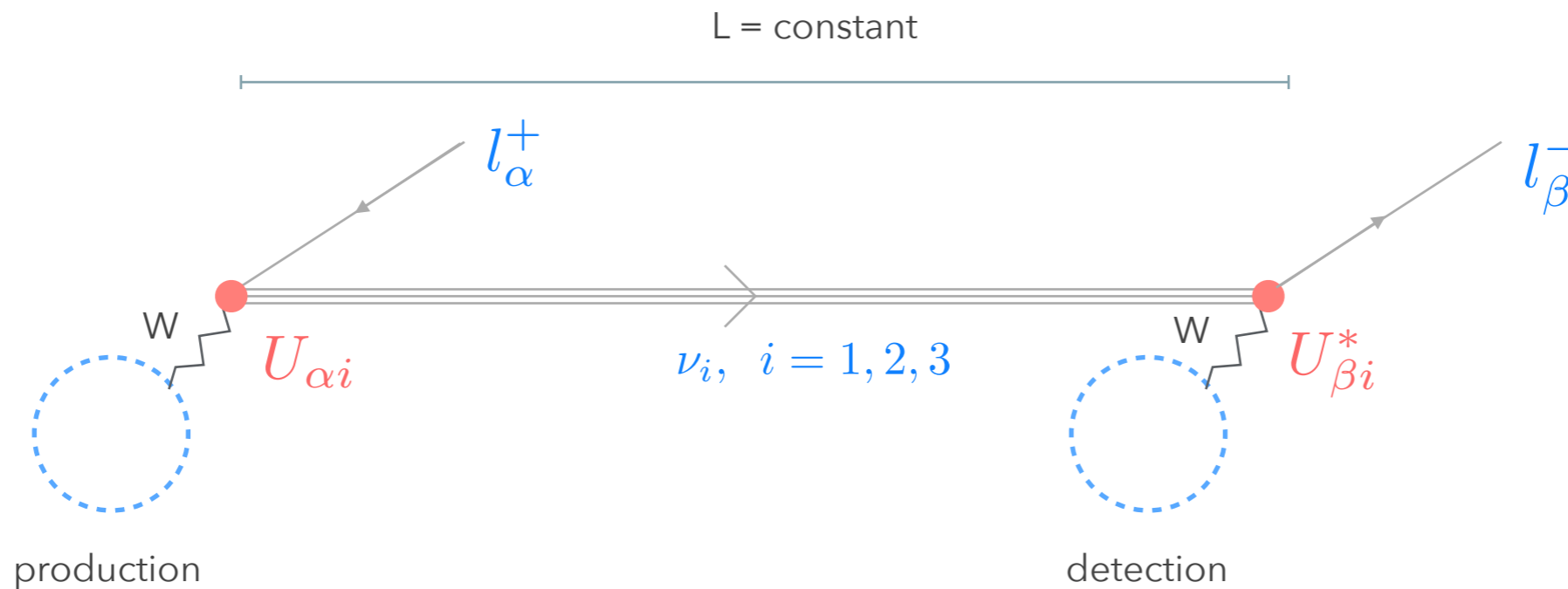
Unitary neutrino mixing matrix (PMNS matrix)

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

-> the effect of the mass is to generate flavour oscillations as a function of distance

WHAT DO WE KNOW ABOUT NEUTRINOS?

- ▶ From oscillation experiments: non-zero masses and non-trivial mixing



Nobel Prize 2015: to Takaaki Kajita and Arthur McDonald “for the discovery of neutrino oscillations, which shows that neutrinos have mass”

WHAT DO WE KNOW ABOUT NEUTRINOS?

- ▶ In general: 3 mixing angles, 1 CP violating phase, 2 independent Δm^2

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Data from
atmospheric ν 's
and accelerators
 $\theta_{23} \approx 48$ deg

Data from
reactors and
accelerators
 $\theta_{13} \approx 8.6$ deg

Data from solar
and reactor
neutrinos
 $\theta_{12} \approx 34$ deg

$$c_{ij} = \cos\theta_{ij}$$

$$s_{ij} = \sin\theta_{ij}$$

$$0 \leq \delta < 2\pi \quad \delta \simeq 3\pi/2$$

*Very different than the CKM mixing angles:

$$\theta_{12} \approx 13^\circ, \quad \theta_{23} \approx 2.4^\circ, \quad \theta_{13} \approx 0.2^\circ$$

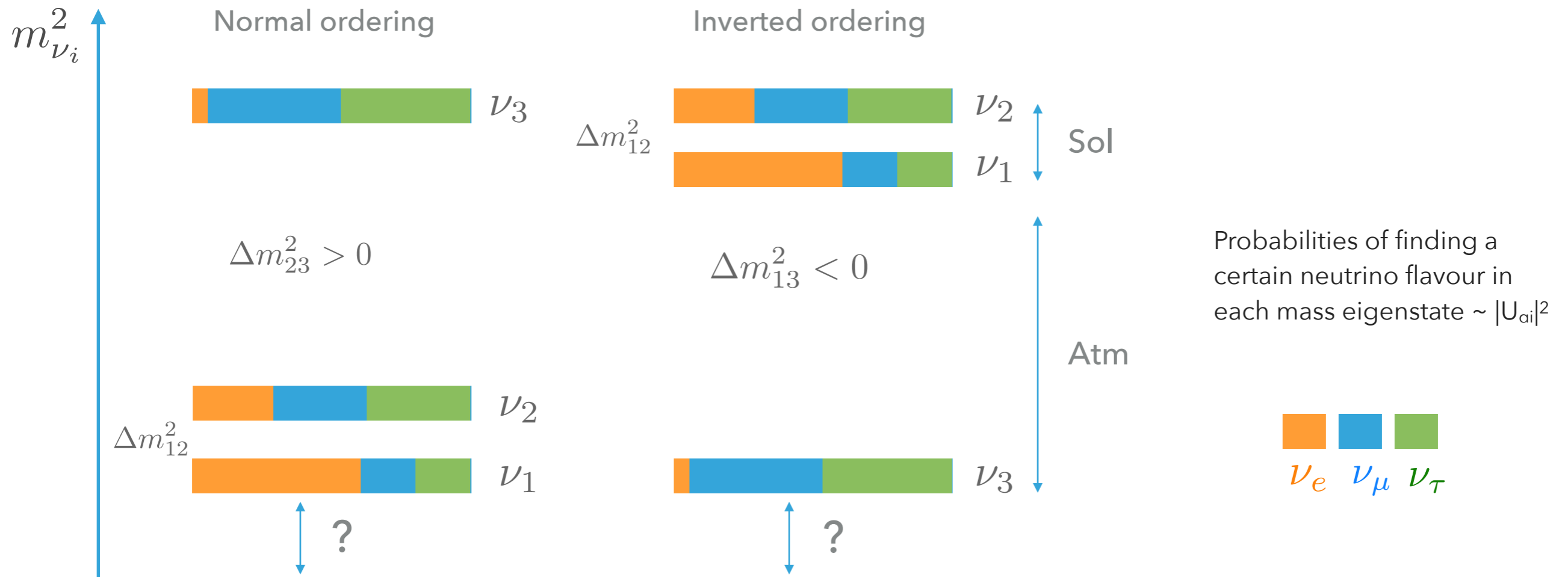
OPEN QUESTIONS IN NEUTRINO PHYSICS

- ▶ From oscillation experiments: we know the mixing angles (or the $U_{\alpha i}$) and the Δm^2

$$\Delta m_{atm}^2 \approx 2.5 \times 10^{-3} \text{ eV}^2$$

$$\Delta m_{sol}^2 \approx 8 \times 10^{-5} \text{ eV}^2$$

- ▶ However: 2 possible mass orderings and no information on the mass scale



OPEN QUESTIONS IN NEUTRINO PHYSICS

▶ Many questions remain open:

- ⦿ What are the absolute values of neutrino masses, and the mass ordering?
- ⦿ What is the nature of neutrinos? Are they Dirac or Majorana particles?
- ⦿ What is the origin of small neutrino masses?
- ⦿ What are the precise values of the mixing angles, and the origin of the large ν mixing?
- ⦿ Is the standard three-neutrino picture correct, or do other, sterile neutrinos exist?
- ⦿ What is the precise value of the CP violating phase δ ?
- ⦿ ...

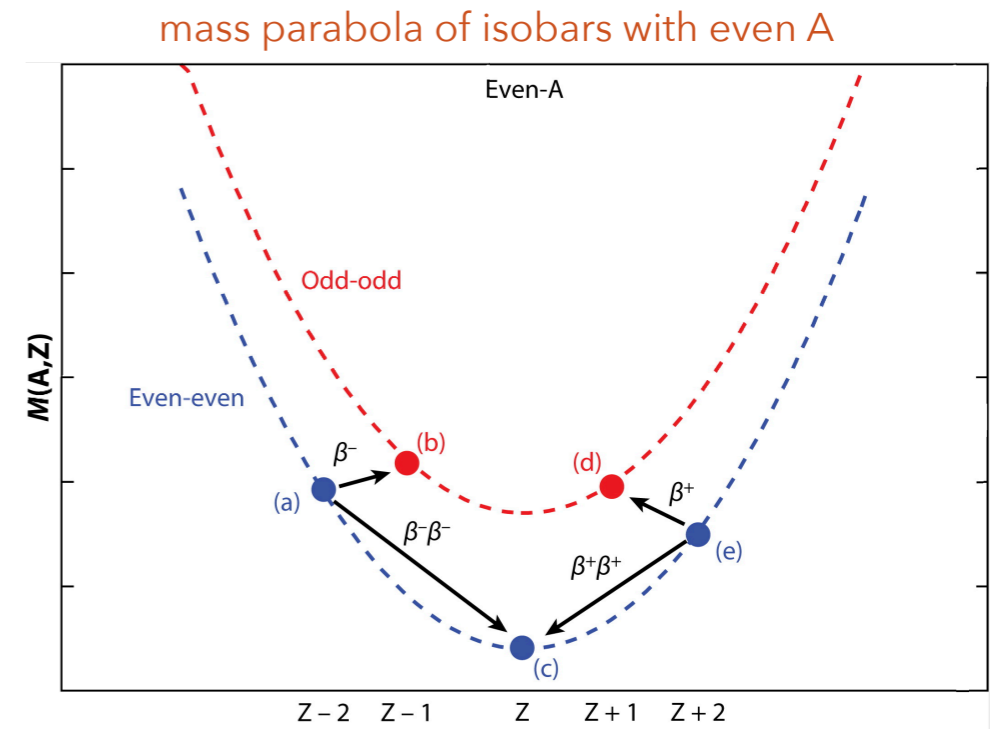
THE DOUBLE BETA DECAY

- ▶ Some of these open questions can be addressed with an extremely rare nuclear decay process
 - What are the absolute values of neutrino masses, and the mass ordering?
 - What is the nature of neutrinos? Are they Dirac or Majorana particles?
 - What is the origin of small neutrino masses?



THE DOUBLE BETA DECAY

- ▶ If simple β^- or β^+ -decay is forbidden on energetic grounds
- ▶ Predicted by Maria-Goeppert Mayer in 1935
- ▶ The probability for a decay is very small, the mean lifetime of a nucleus is much larger than the age of the universe ($\tau_U \sim 1.4 \times 10^{10}$ a)



Ruben Saakyan, Annu. Rev. Nucl. Part. Sci. 63 (2013)

$$\tau_{2\nu} \approx 10^{20} y$$

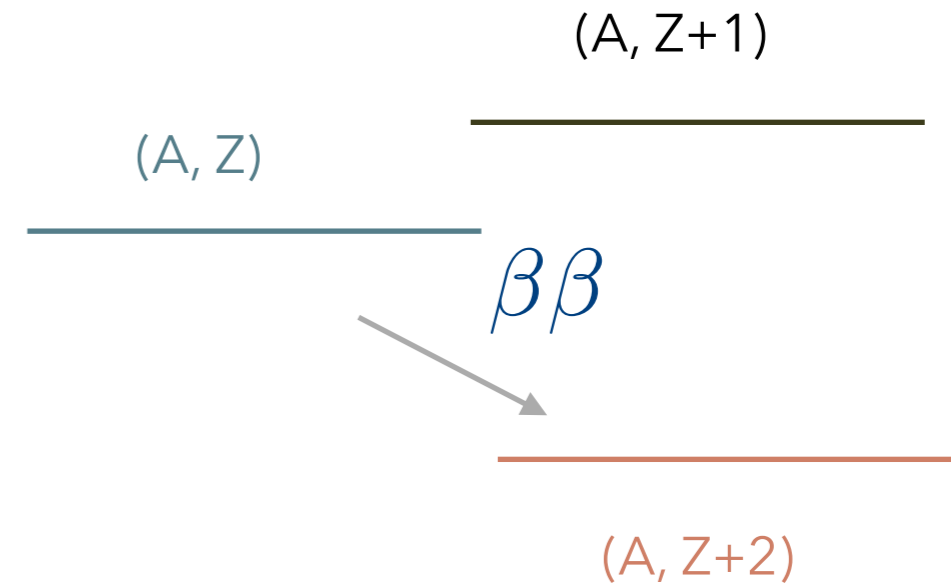
- Thus: a very rare process
- However, if a large amount of nuclei is used, the process can be observed experimentally



Nobel Prize in physics, 1963 for her discoveries concerning the nuclear shell structure

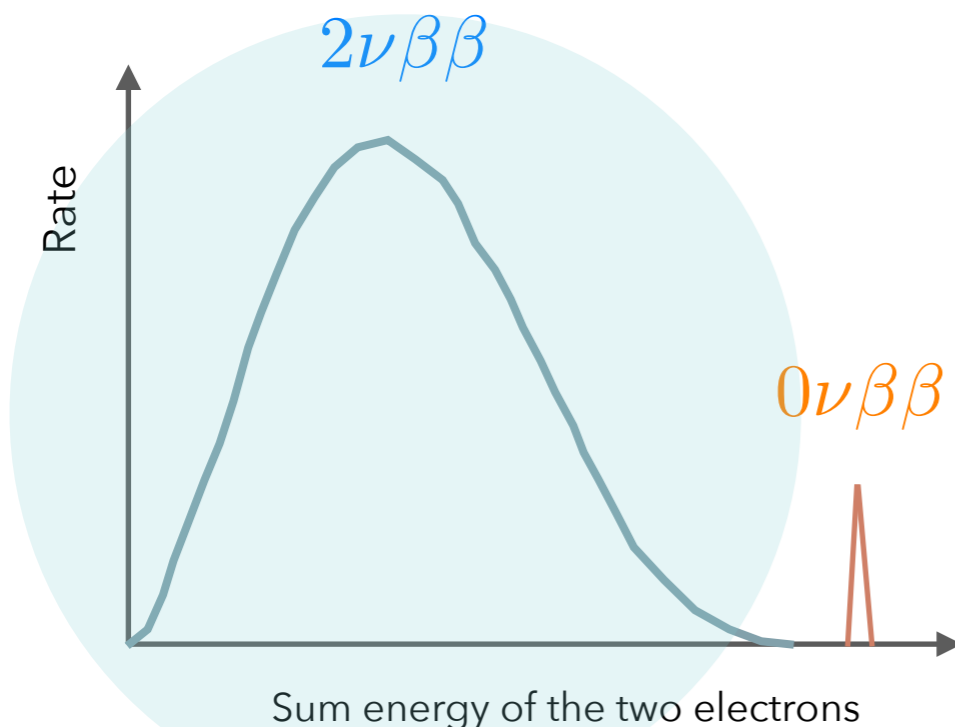
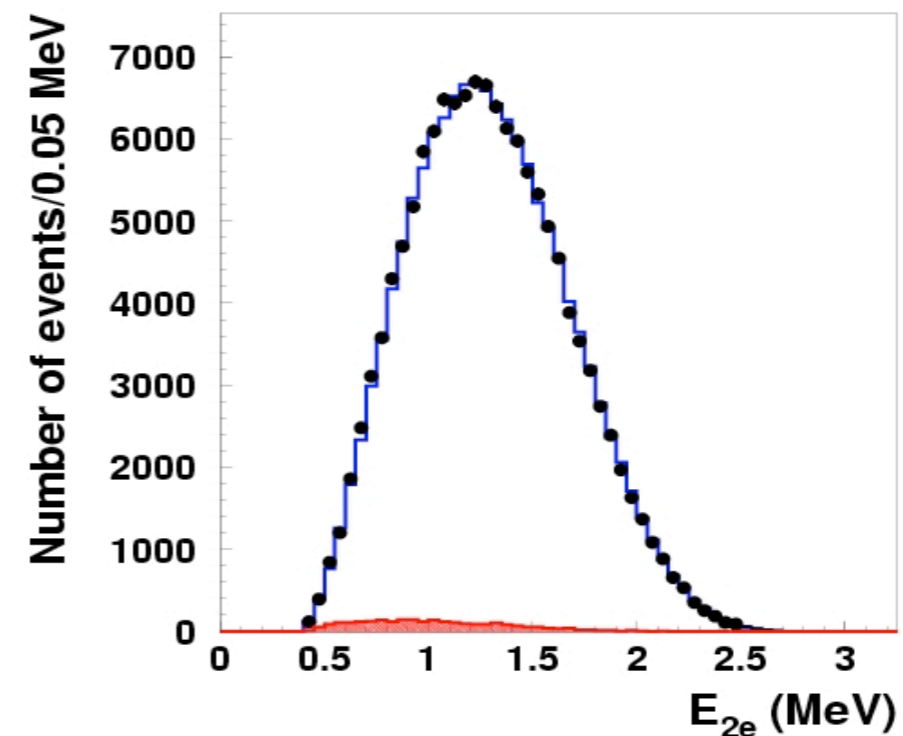
THE DOUBLE BETA DECAY

- ▶ The Standard Model decay, with 2 neutrinos, was observed in 14 nuclei
- ▶ $T_{1/2} > 10^{18}$ y: ^{48}Ca , ^{76}Ge , ^{82}Se , ^{96}Zr , ^{100}Mo , ^{116}Cd , ^{128}Te , ^{130}Te , ^{136}Xe , ^{150}Nd , ^{238}U



^{100}Mo : $T_{1/2} = 7.15 \times 10^{18}$ a

NEMO Experiment in Modane/Frejus



THE DOUBLE BETA DECAY

- ▶ The decay rate $\Gamma^{2\nu}$ depends on the matrix element $M^{2\nu}$ and on the phase space factor $G^{2\nu}$ (which determines the energy spectrum):

$$\Gamma^{2\nu} = \frac{\ln 2}{T_{1/2}^{2\nu}} = G^{2\nu}(Q, Z) |M^{2\nu}|^2$$

- ▶ The phase space factor (Z = charge of daughter nucleus) from the leptonic degrees of freedom:

$$G^{2\nu} \propto (G_F \cos \theta_C)^4 Q^7 \cdot \left(\frac{Q^4}{1980} + \frac{Q^3}{90} + \frac{Q^2}{9} + \frac{Q}{2} + 1 \right) \propto (G_F \cos \theta_C)^4 \cdot Q^{11}$$

- The decay rate scales with $Q^{11} \times (G_F)^4 \Rightarrow$ we expect indeed very long $T_{1/2}$ of $\sim 10^{20}$ y

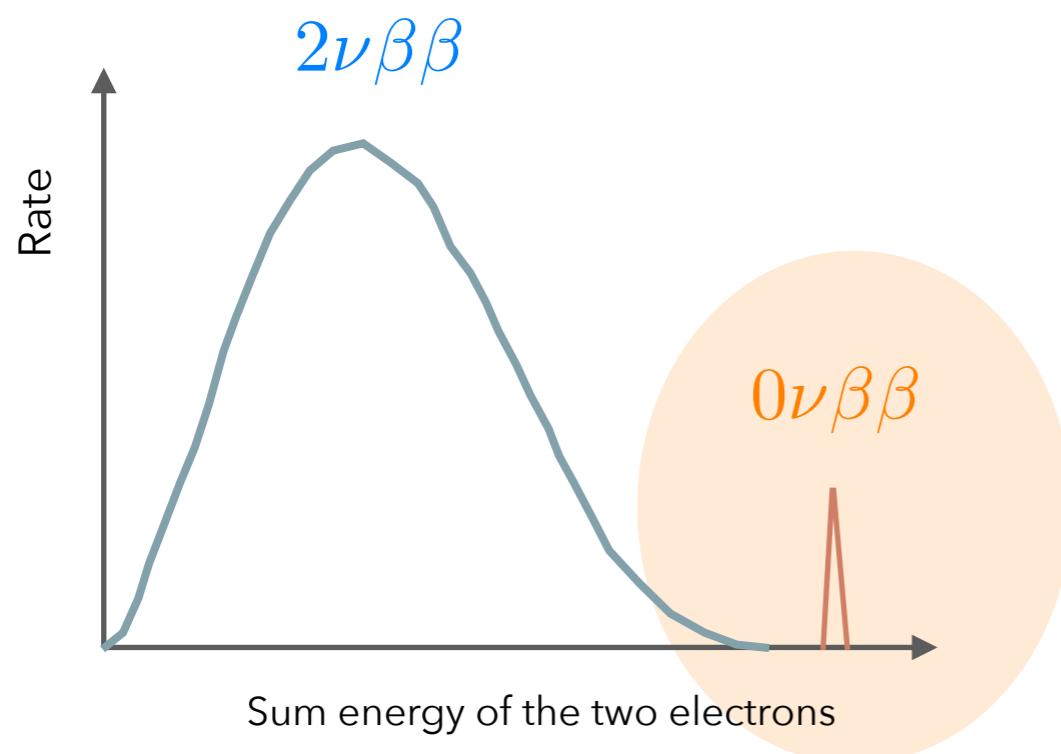
THE NEUTRINOLESS DOUBLE BETA DECAY

- ▶ More interesting: the decay *without* emission of neutrinos $\Rightarrow \Delta L = 2$

$$T_{1/2}^{0\nu\beta\beta} > 10^{24} \text{ y}$$

- ▶ Expected signature: *sharp peak at the Q-value of the decay*

$$Q = E_{e1} + E_{e2} - 2m_e$$

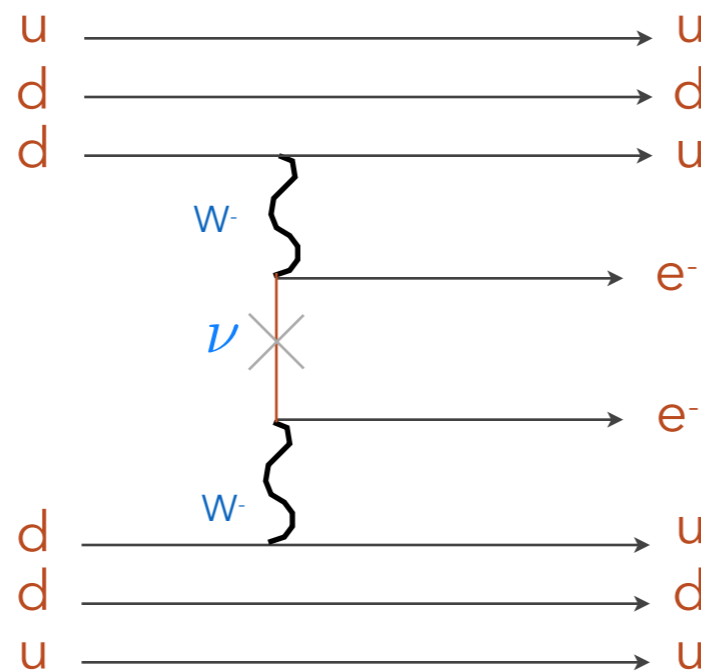


The double beta decay without neutrinos: first discussed by Wendell H. Furry in 1939

Ettore Majorana had proposed in 1937 that neutrinos could be their own antiparticles

THE NEUTRINOLESS DOUBLE BETA DECAY

- ▶ In this decay, a light virtual neutrino could be exchanged



Charge conjugate spinor

$$\psi^c = C\bar{\psi}^T$$

A Majorana field

$$\psi = \psi^c$$

$$\psi = \psi_L + \psi_L^c$$

has 2 spin d.o.f.

- ▶ The neutron decays under emission of a right handed 'anti-neutrino' ν_L^c
 - the ν_L^c has to be absorbed at the second vertex as left handed 'neutrino' ν_L
 - for the decay to happen: neutrinos and anti-neutrinos must be identical, thus Majorana particles
 - & the helicity must change

MAJORANA AND DIRAC NEUTRINOS

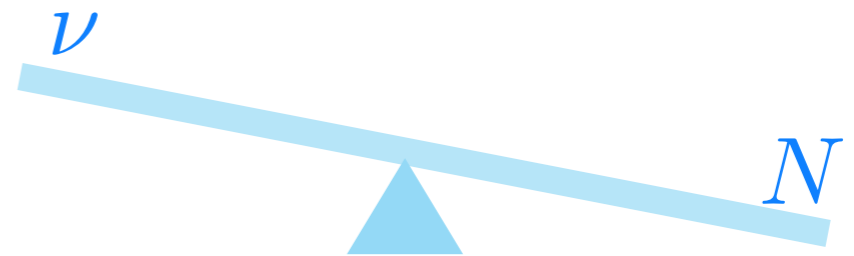


- ▶ Most general Lagrangian: both type of neutrinos masses

$$\mathcal{L}_{\mathcal{M}_\nu} = -\frac{1}{2} [m_D(\bar{\psi}_R^c \psi_L^c + \bar{\psi}_R \psi_L) + M\bar{\psi}_L^c \psi_L] + h.c.$$

- ▶ **Dirac term:** generated after SSB from Yukawa interactions; **Majorana term:** singlet of the SM gauge group and can appear as bare mass term
- ▶ **Masses of physical neutrinos:** from the eigenvalues of the mass matrix. In the “see saw” mechanism: $M \gg m_D \Rightarrow$ a very light neutrinos state ν and a heavy state N with masses:

$$m_\nu \approx \frac{m_D^2}{M} \quad m_N \approx M$$



- ▶ If Dirac mass term m_D : of similar size as of other fermions & M at the GUT scale ($\sim 10^{14}$ GeV) \Rightarrow **explanation of the smallness of neutrino masses**

THE NEUTRINOLESS DOUBLE BETA DECAY

- ▶ The expected rate can be calculated as:

$$\Gamma^{0\nu} = \frac{\ln 2}{T_{1/2}^{0\nu}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \frac{|m_{\beta\beta}|^2}{m_e^2}$$

← from the leptonic part of the matrix element

↑
the matrix element of the nuclear transition

- ▶ with the phase space integral (now spanned only by 2 electrons):

$$G^{0\nu} \propto (G_F \cos \theta_C)^4 \cdot \left(\frac{Q^5}{30} - \frac{2Q^2}{3} + Q - \frac{2}{5} \right) \propto (G_F \cos \theta_C)^4 \cdot Q^5$$

THE EFFECTIVE MAJORANA NEUTRINO MASS

- ▶ The effective Majorana neutrino mass parameter: embeds all the dependence on neutrino quantities

$$|m_{\beta\beta}| = |m_1 U_{e1}^2 + m_2 U_{e2}^2 e^{2i\phi_1} + m_3 U_{e3}^2 e^{2i(\phi_2 - \delta)}|$$

- ▶ A mixture of $m_1, m_2, m_3 \sim$ to the U_{ei}^2 (the complex entries in the PMNS matrix)

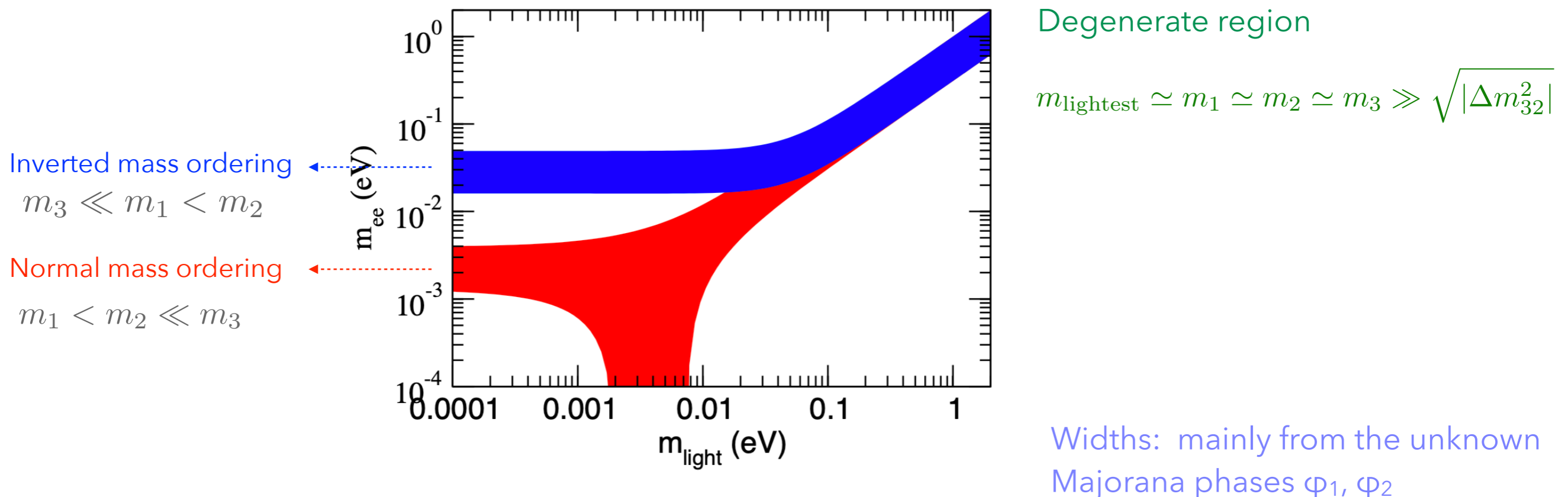
$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \times \begin{pmatrix} e^{i\phi_1} & 0 & 0 \\ 0 & e^{i\phi_2} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

- ▶ $\phi_1, \phi_2 =$ Majorana phases and $|U_{e1}|^2$ is for instance the probability that ν_e has the mass m_1

◉ fewer phases can be removed by redefining the fields

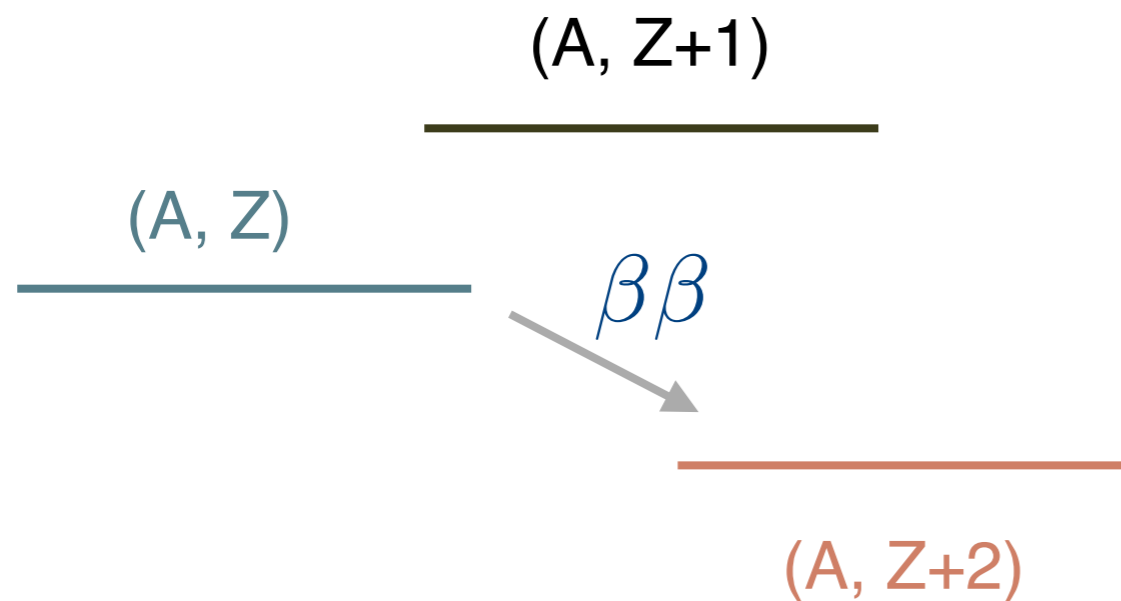
THE EFFECTIVE MAJORANA NEUTRINO MASS

- ▶ The values depend critically on the neutrino mass spectrum and on the values of the two Majorana phases in the PMNS matrix
- ▶ One can express $m_{\beta\beta}$ as a function of the lightest (m_{lightest}) mass state for the two mass orderings and obtain the allowed ranges



EMPLOYED NUCLEI

- Even-even nuclei
- Natural abundance is low (except ^{130}Te)
- Must use enriched material



Candidate*	Q [MeV]	Abund [%]
$^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$	4.271	0.187
$^{76}\text{Ge} \rightarrow ^{76}\text{Se}$	2.039	7.8
$^{82}\text{Se} \rightarrow ^{82}\text{Kr}$	2.995	9.2
$^{96}\text{Zr} \rightarrow ^{96}\text{Mo}$	3.350	2.8
$^{100}\text{Mo} \rightarrow ^{100}\text{Ru}$	3.034	9.6
$^{110}\text{Pd} \rightarrow ^{110}\text{Cd}$	2.013	11.8
$^{116}\text{Cd} \rightarrow ^{116}\text{Sn}$	2.802	7.5
$^{124}\text{Sn} \rightarrow ^{124}\text{Te}$	2.228	5.64
$^{130}\text{Te} \rightarrow ^{130}\text{Xe}$	2.530	34.5
$^{136}\text{Xe} \rightarrow ^{136}\text{Ba}$	2.479	8.9
$^{150}\text{Nd} \rightarrow ^{150}\text{Sm}$	3.367	5.6

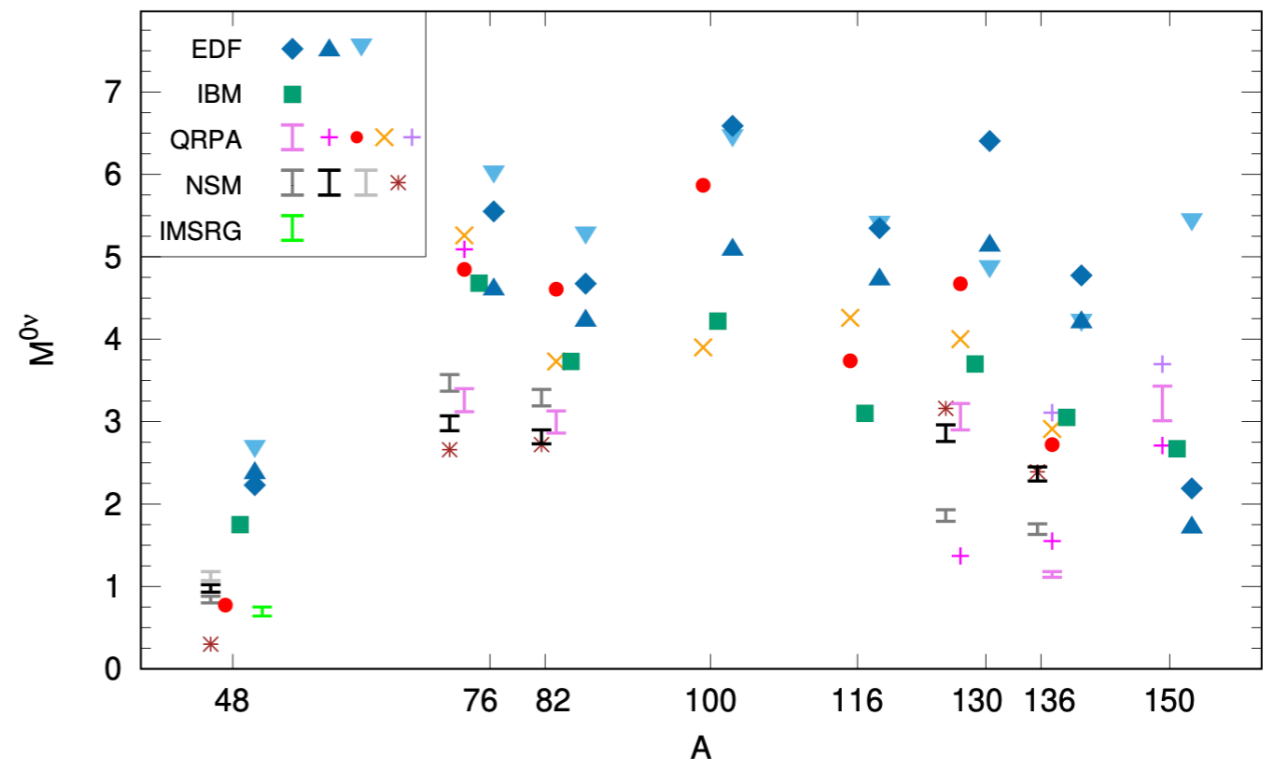
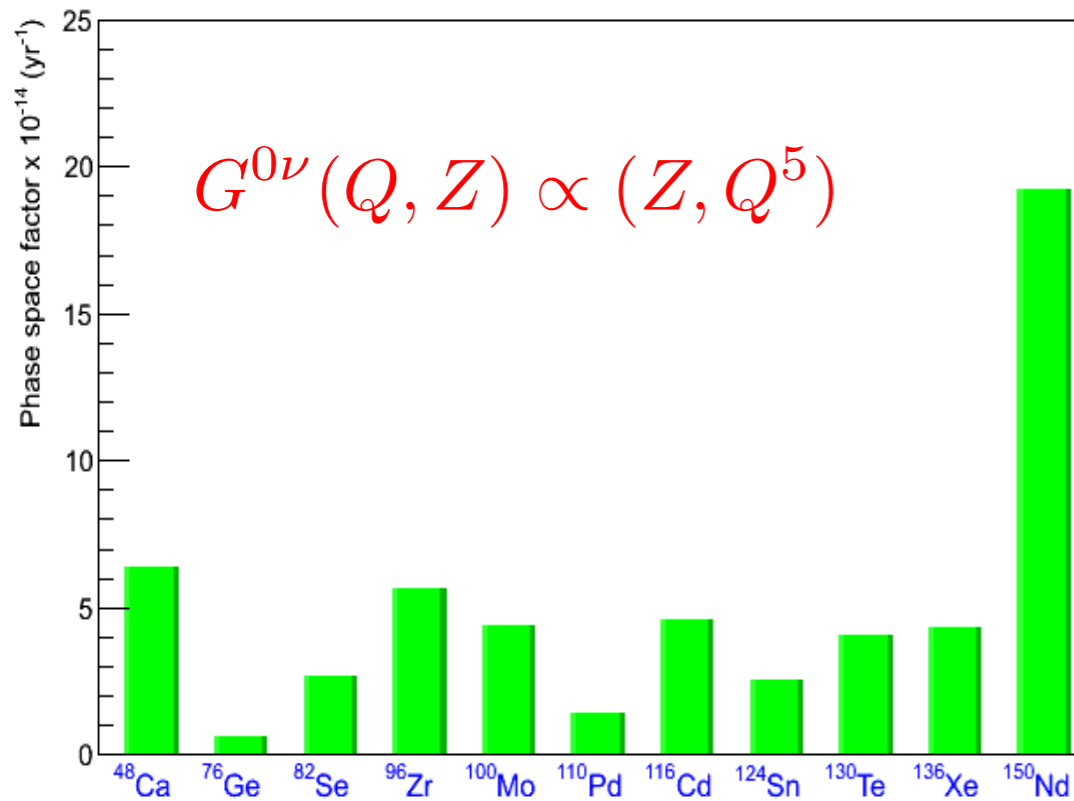
* Q-value > 2 MeV

PHASE SPACE AND MATRIX ELEMENTS

Matrix elements: vary by a factor of 2- 3 for a given A

$$\Gamma^{0\nu} = \frac{\ln 2}{T_{1/2}^{0\nu}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \frac{|m_{\beta\beta}|^2}{m_e^2}$$

$$\Gamma^{0\nu} = \frac{\ln 2}{T_{1/2}^{0\nu}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \frac{|m_{\beta\beta}|^2}{m_e^2}$$



J. Menendez Neutrino 2020

Jonathan Engel and Javier Menéndez 2017 *Rep. Prog. Phys.* **80** 046301

*See also Vergados, Ejiri, Simkovoc, *Int. Journal of Modern Physics E*, Vol 25 (2016)

EXPERIMENTAL REQUIREMENTS

- ▶ Experiments measure the half-life, with a sensitivity (for non-zero background)

$$T_{1/2}^{0\nu} \propto a \cdot \epsilon \cdot \sqrt{\frac{M \cdot t}{B \cdot \Delta E}}$$



Minimal requirements:

- high isotopic abundance (a)
- high efficiency (ε)
- large detector masses (M)
- ultra-low background noise (B)
- good energy resolution (ΔE)

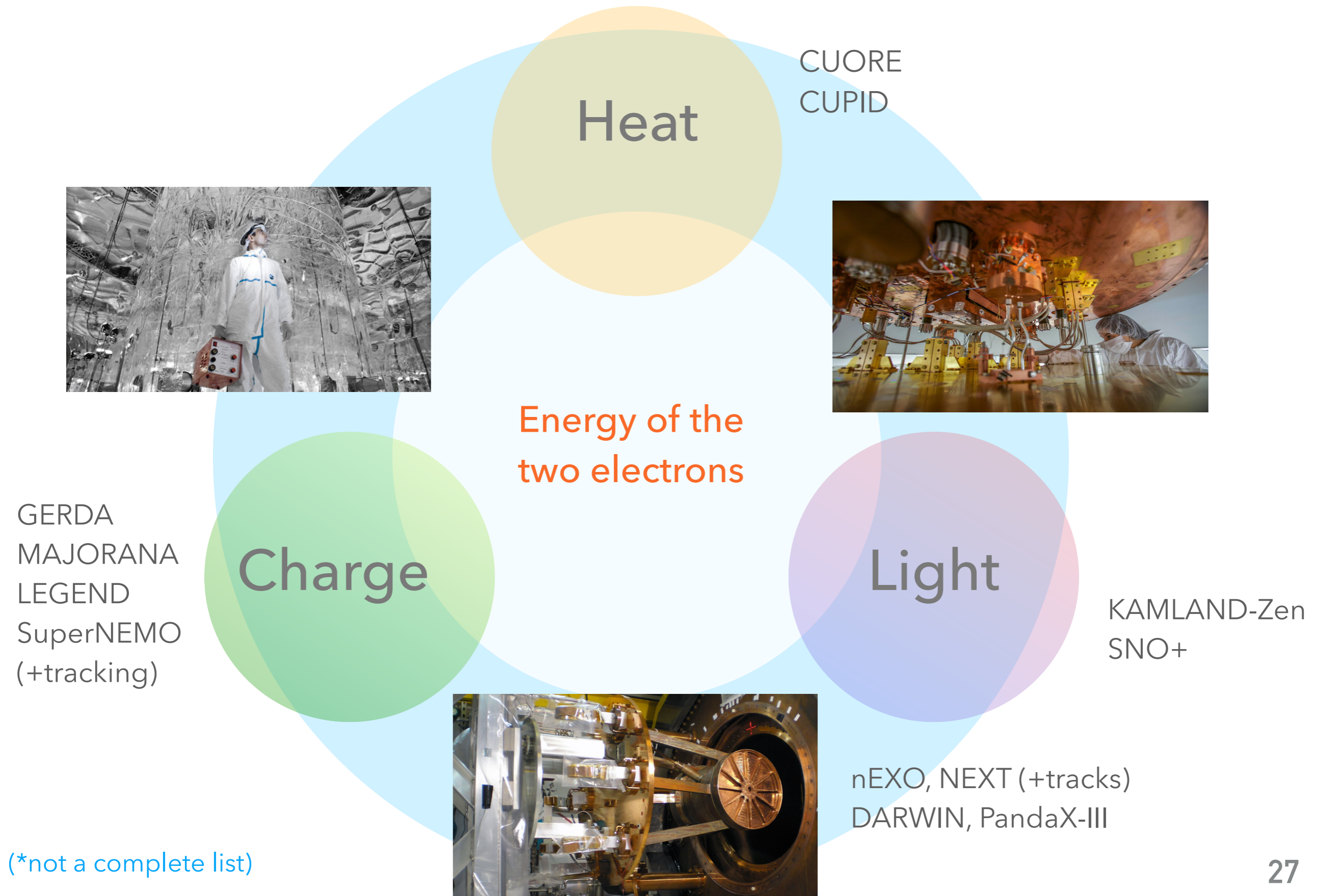


$$\langle m_{\beta\beta} \rangle \propto \frac{1}{\sqrt{T_{1/2}^{0\nu}}}$$

Additional tools to distinguish signal from background:

- event topology
- pulse shape discrimination
- particle identification

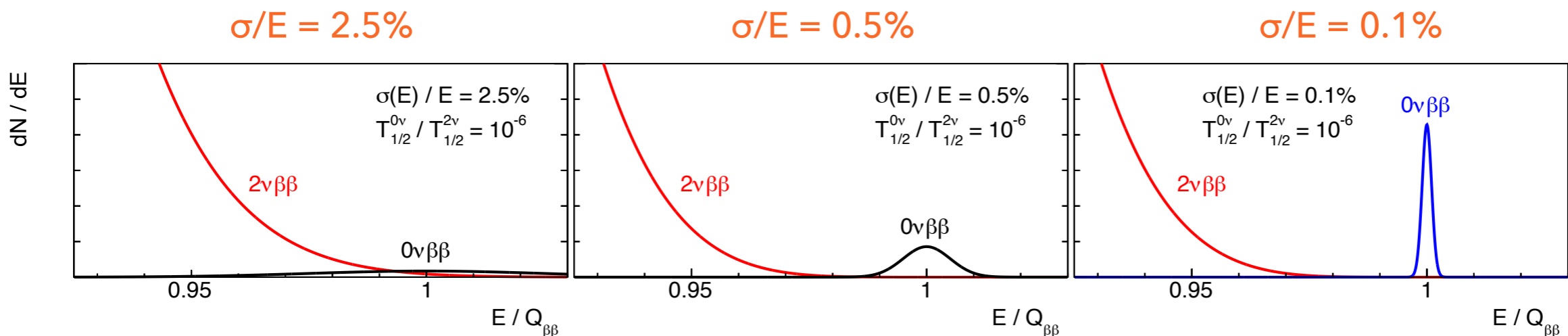
DOUBLE BETA DECAY: EXPERIMENTAL TECHNIQUES*



(*not a complete list)

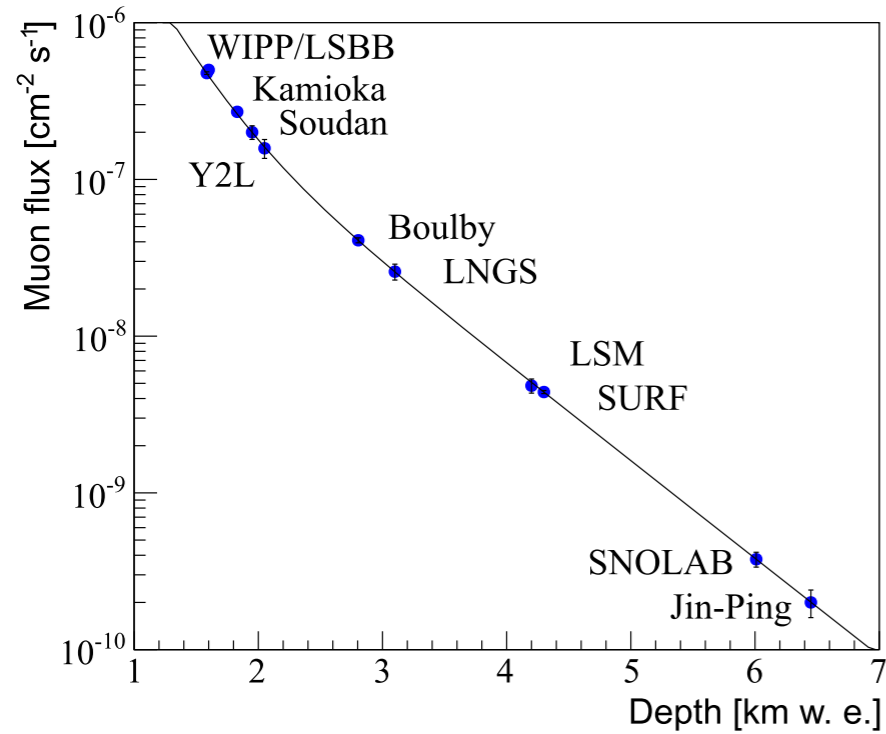
MAIN CHALLENGES

- ▶ Energy resolution (ultimate background from $2\nu\beta\beta$ -decay)
- ▶ Backgrounds
 - cosmic rays & cosmogenic activation (including in situ, e.g., ^{77}Ge , ^{137}Xe)
 - radioactivity of detector materials (^{238}U , ^{232}Th , ^{40}K , ^{60}Co , etc: α , β , γ -radiation)
 - anthropogenic (e.g., ^{137}Cs , $^{110\text{m}}\text{Ag}$)
 - neutrinos (e.g., ^8B from the Sun): $\nu + e^- \rightarrow \nu + e^-$



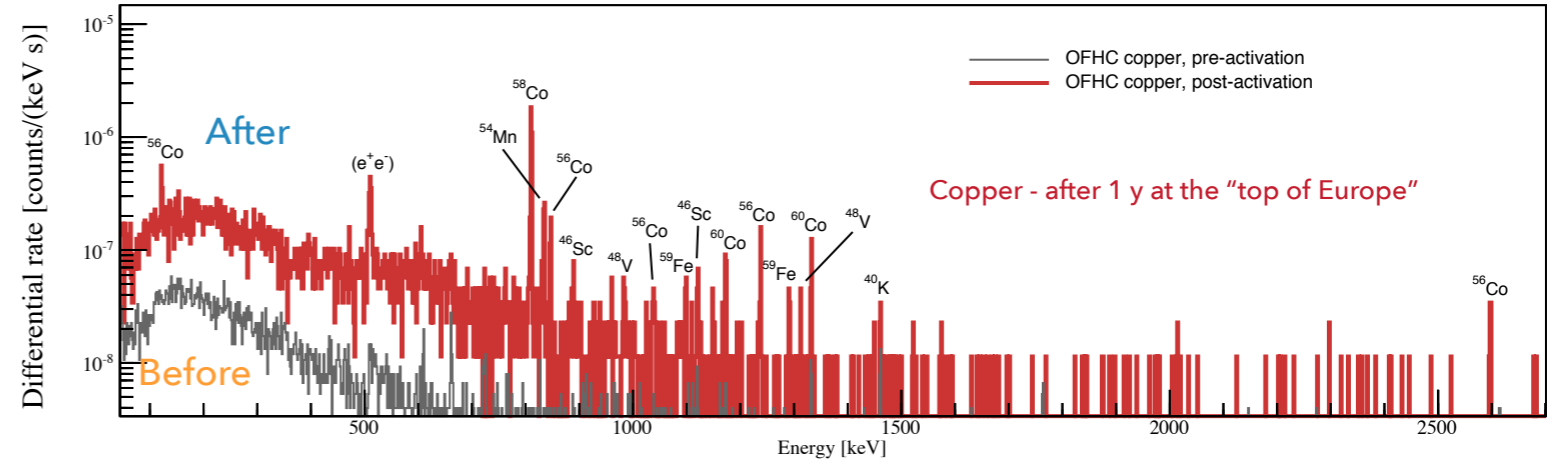
BACKGROUND REDUCTION

Go deep underground

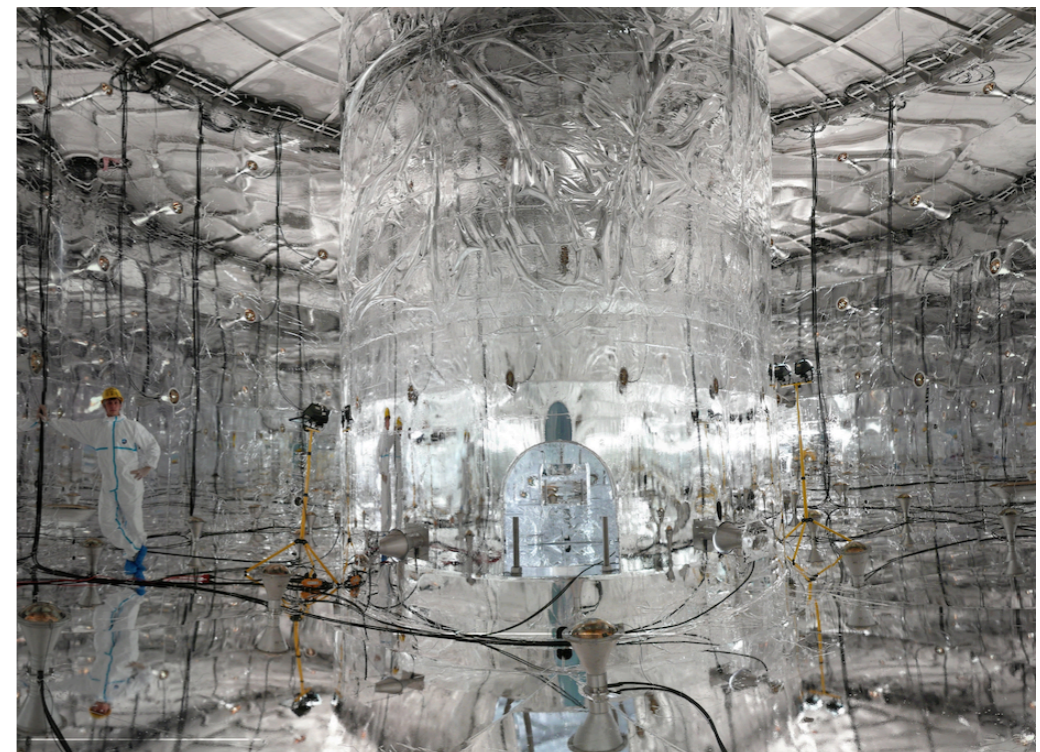


Avoid cosmic activation

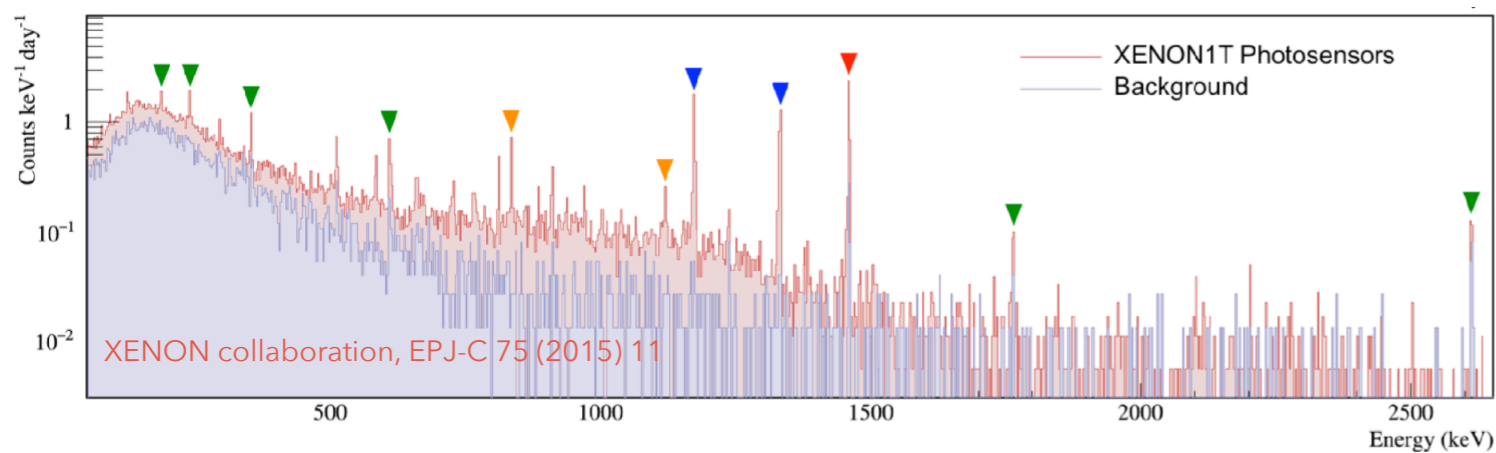
LB et al., Eur. Phys. J. C75 2015



Use active shields



Select low-radioactivity materials



VERY BRIEF CURRENT STATUS OF THE FIELD

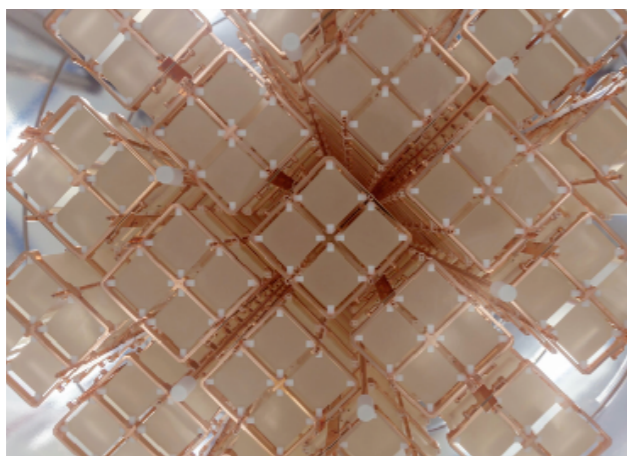
- ▶ No observation of this extremely rare nuclear decay (so far)
- ▶ Best *lower limits* on $T_{1/2}$: 1.07×10^{26} y (^{136}Xe), 1.8×10^{26} y (^{76}Ge), 1.5×10^{25} y (^{130}Te)

$$m_{\beta\beta} < (0.08 - 0.18) \text{ eV (90\% C.L.)}$$

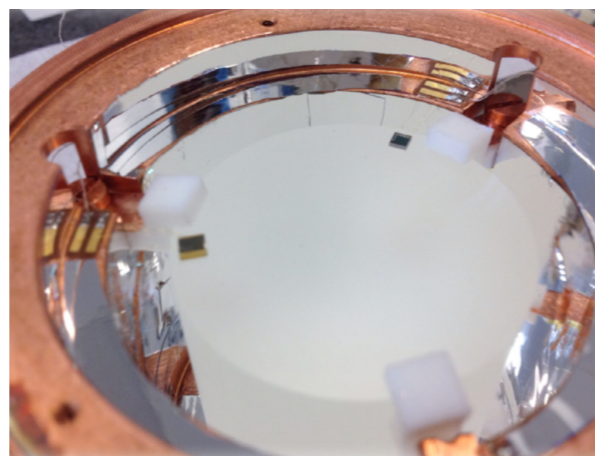
- ▶ Running and upcoming experiments (a selection)
 - ^{130}Te : CUORE, SNO+
 - ^{136}Xe : KAMLAND-Zen, KAMLAND2-Zen, EXO-200, nEXO, NEXT, DARWIN, PandaX-III
 - ^{76}Ge : GERDA Phase-II, Majorana, LEGEND (GERDA & Majorana + new groups)
 - ^{82}Se : CUPID (= CUORE with light read-out)
 - ^{82}Se (^{150}Nd , ^{48}Ca): SuperNEMO
 - ^{100}Mo : NEMO-3, AMoRE

CUORE AND CUPID

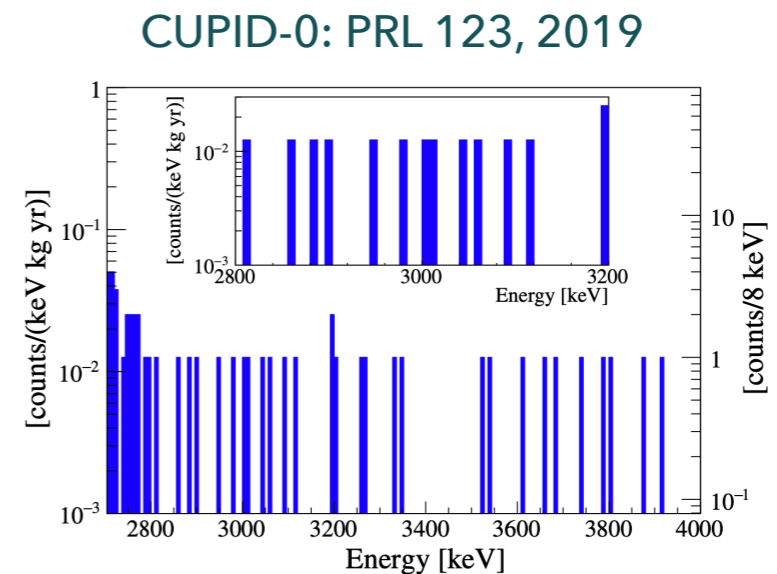
- ▶ **CUORE**: 988 crystals (206 kg ^{130}Te assembled in towers) at LNGS
- ▶ Background level: 14 events/(keV t y); energy resolution: 0.3% FWHM (7.7 keV in ROI)
 - ⦿ Results: $T_{1/2} > 1.5 \times 10^{25}$ y for ^{130}Te
- ▶ **CUPID**: R&D for ton-scale detector using $\text{Li}_2^{100}\text{MoO}_4$ and Zn^{82}Se crystals as scintillating bolometers (to identify major α -particle background)
- ▶ **CUPID-0**: pilot project at LNGS, 24 Zn^{82}Se crystals, best limit on $T_{1/2}$ of ^{82}Se



CUORE: PRL 120, 2018



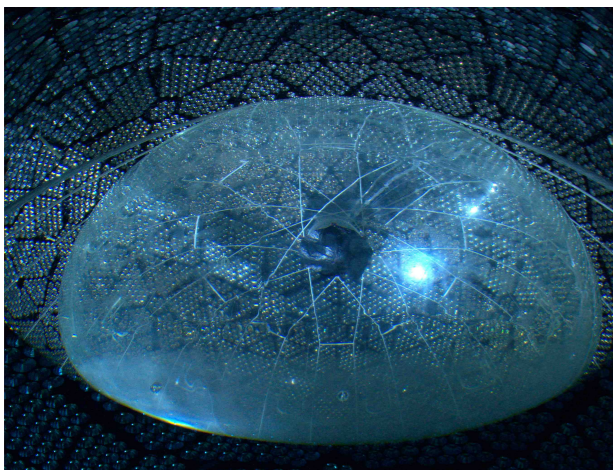
Q value of ^{82}Se (2997.9 ± 0.3 keV)



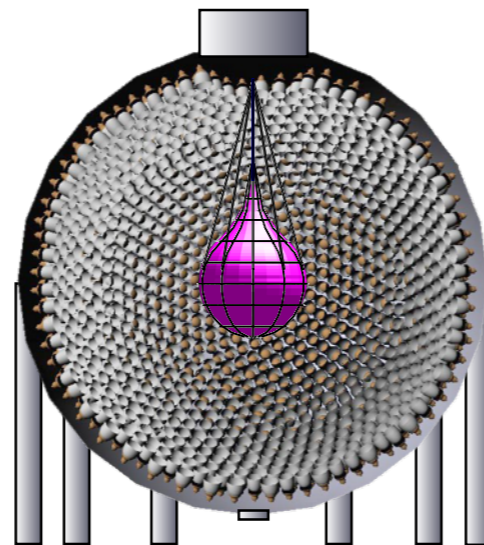
$T_{1/2} > 3.5 \times 10^{24}$ y

SNO+ AND KAMLAND-ZEN

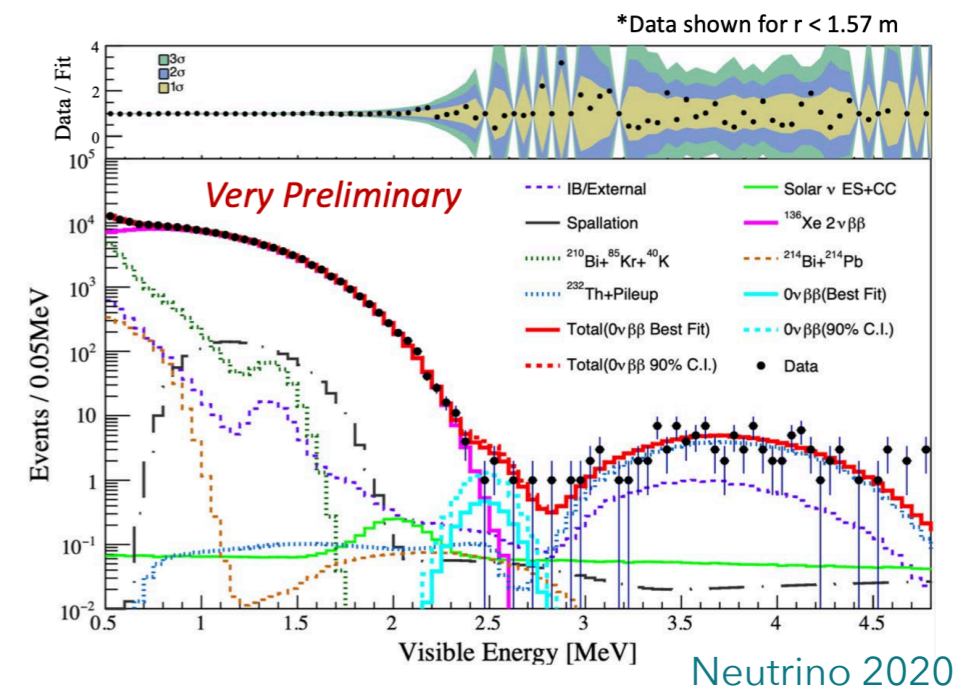
- ▶ **SNO+:** 0.5% $^{\text{nat}}\text{Te}$ ~ 1330 kg ^{130}Te in liquid scintillator at SNOLAB
 - ⦿ Scintillator fill in 2020, preparing for Te loading and $\beta\beta$ -decay phase to start
- ▶ **KamLAND-Zen:** 745 kg ^{136}Xe in liquid scintillator at Kamioka, ongoing since Jan 2019
 - ⦿ Previous results (phase I + II): $T_{1/2} > 1.07 \times 10^{26}$ y (5.6×10^{25} y sensitivity)
- ▶ **KamLAND2-Zen:** ~ 1 t ^{136}Xe , higher LCE: $\sigma/E(2.6 \text{ MeV}) = 4\% \longrightarrow < 2.5\%$



SNO+ J.Phys.Conf.Ser. 1137 (2019)
 $T_{1/2} > 1.9 \times 10^{26}$ y, 5 y of data

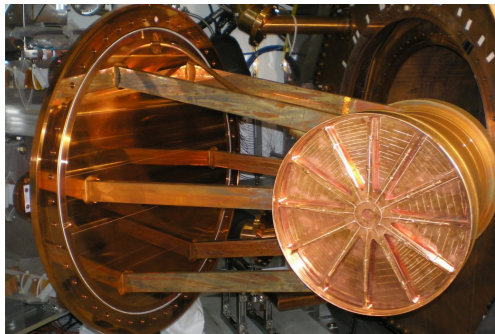


KamLAND-Zen: PRL 117, 2016



EXO-200, NEXO, NEXT, PANDAX-III

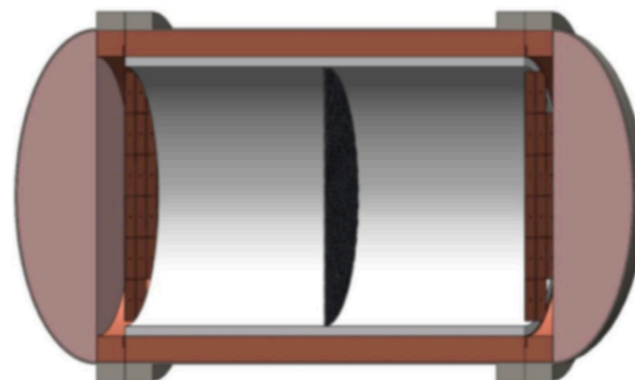
- ▶ **EXO-200**: TPC with 75 kg ^{136}Xe in fiducial region, $\sigma/E = 1.1\%$; $T_{1/2} > 3.5 \times 10^{25}$ y
- ▶ **nEXO**: TPC with 5 t of LXe enriched in ^{136}Xe , goal $T_{1/2} \sim 9.2 \times 10^{27}$ y after 10 y
- ▶ **NEXT**: high-pressure (15 bar) ^{136}Xe gas TPC: e⁻ track reconstruction
 - ◉ Demonstrated: $\sigma/E = 0.43\%$; NEXT-100: operation in 2021, $T_{1/2} \sim 6 \times 10^{25}$ y after 3 y
 - ◉ R&D on Ba ion tagging ongoing (e.g., NEXT-BOLD, for ton-scale detector)
- ▶ **PandaX-III**: high-pressure (10 bar) ^{136}Xe gas TPC, multiple modules with 200 kg each



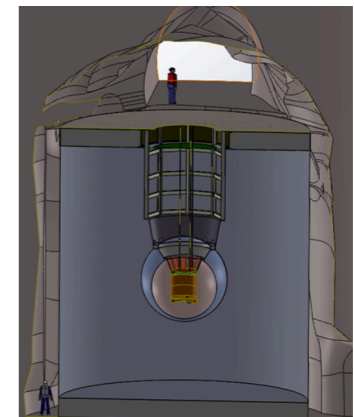
EXO-200: PRL 123, 2019



NEXT arXiv:1910.07314



PandaX-III NIM-A 958, 2020

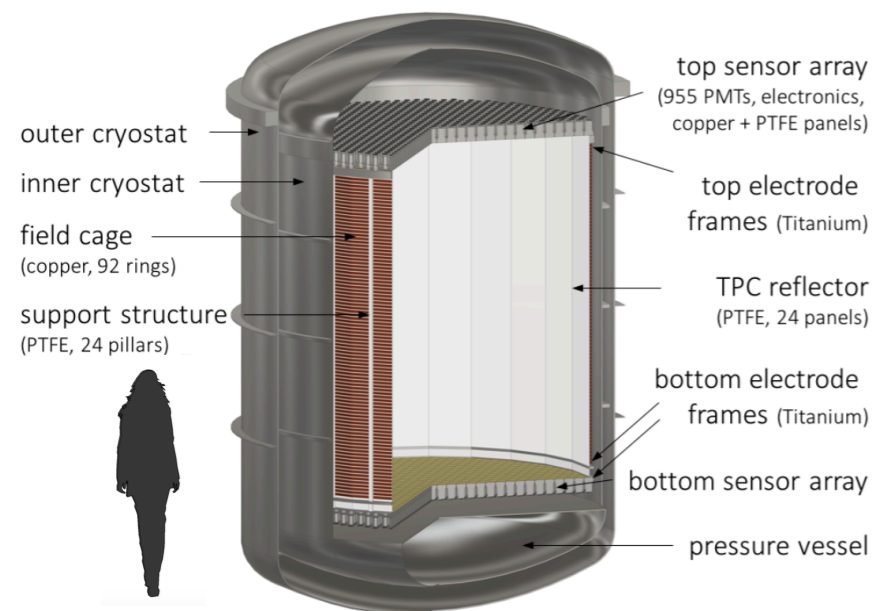


nEXO arXiv:1805.11142

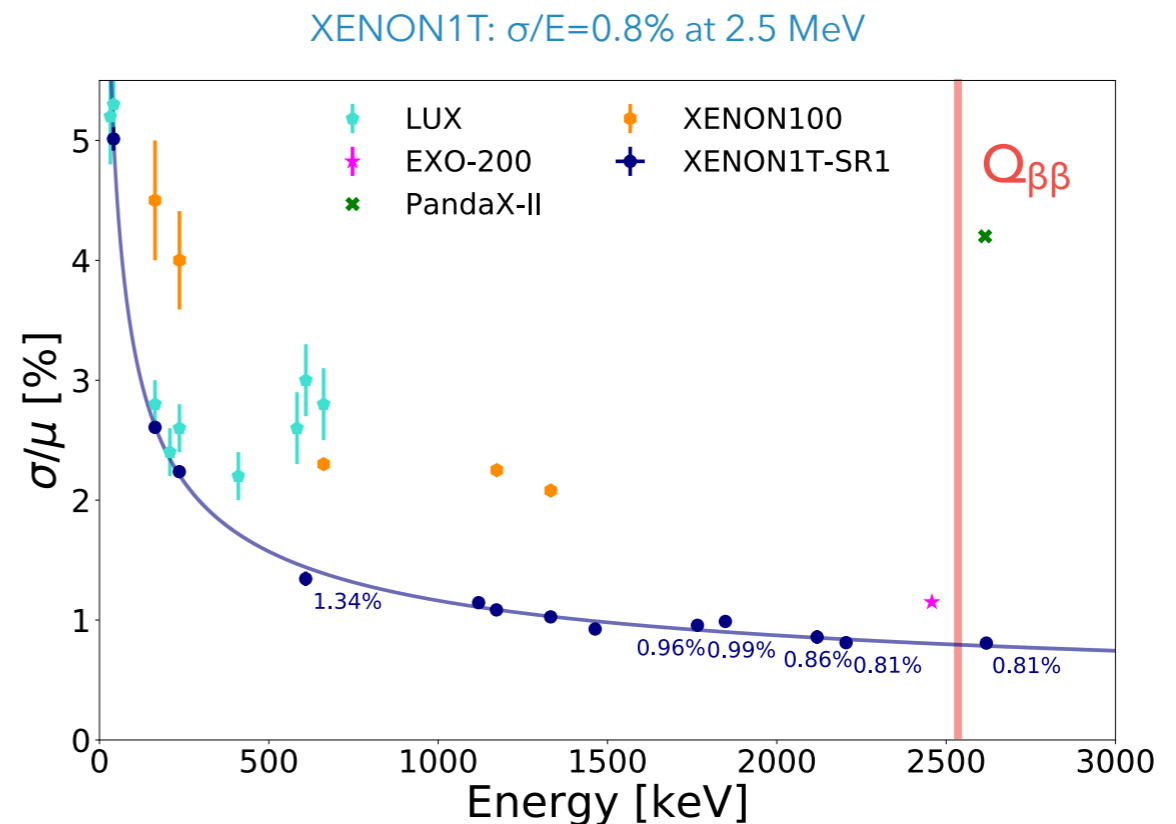
DARWIN



- ▶ TPC with 40 t ^{nat}Xe (50 t in total) for DM searches; **8.9% ¹³⁶Xe ≈ 3.6 t of ¹³⁶Xe**
- ▶ **Goal: $T_{1/2} \sim \text{few} \times 10^{27} \text{ y}$, with background rate $< 0.2 \text{ events}/(\text{t y})$ in ROI**
- ▶ Energy resolution: $\sigma/E = 0.8\%$ (achieved in XENON1T)
- ▶ Detailed $\beta\beta$ -sensitivity study: [arXiv:2003.13407](https://arxiv.org/abs/2003.13407) (EPJ-C 80, 2020, 9, DARWIN collaboration)



DARWIN Collaboration, JCAP 1611 (2016) 017

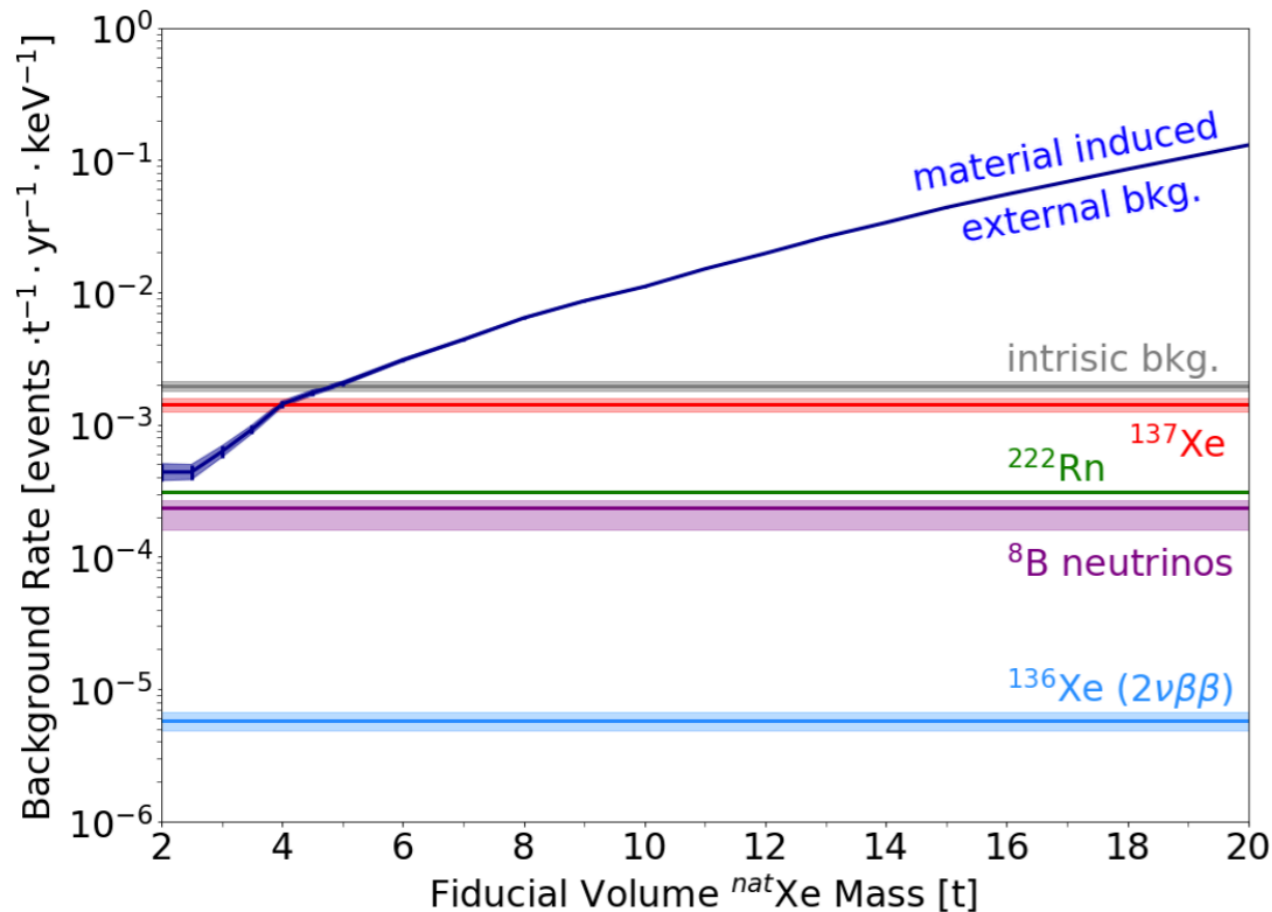


XENON Collaboration, arXiv:2003.03825, EPJ-C 80., 2020, 8

DARWIN BACKGROUNDS

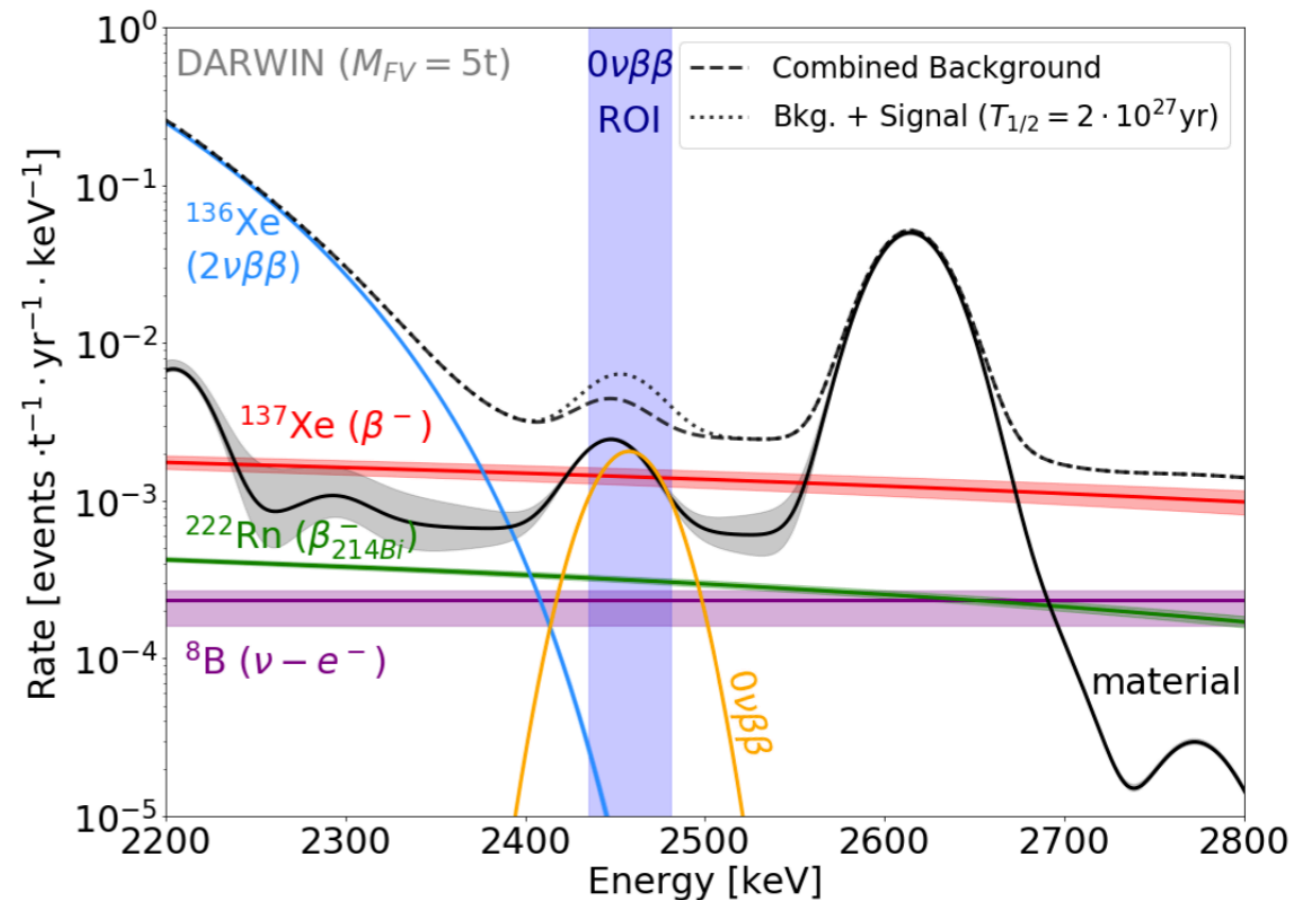
- ▶ ROI: [2435-2481] keV = FWHM around $Q_{\beta\beta}$
- ▶ ^{137}Xe : β -decay with $Q=4173$ keV, $T_{1/2}=3.82$ min (via n-capture on ^{136}Xe)

DARWIN collaboration, arXiv:2003.13407



Rate versus fiducial mass

Signal: $T_{1/2} = 2 \times 10^{27}$ y

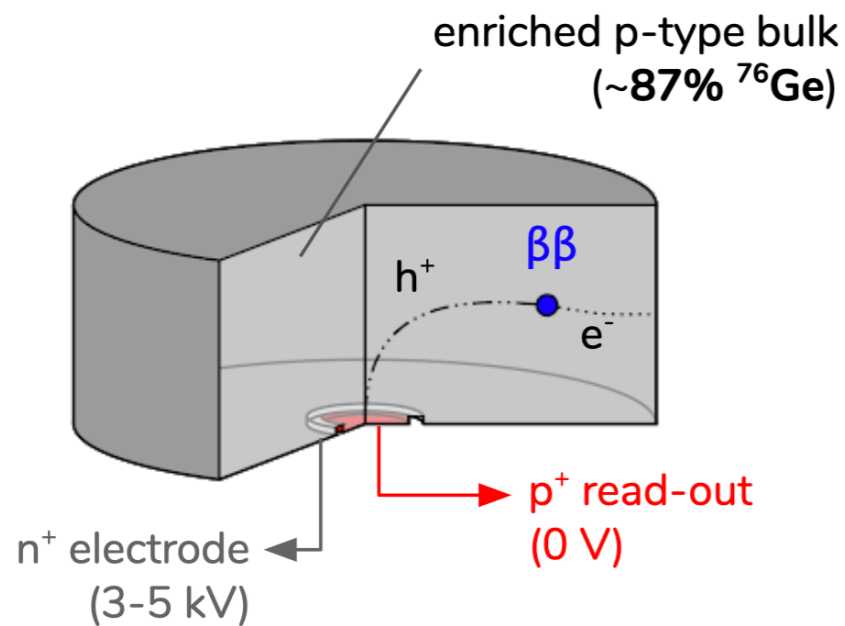


Rate in 5 tonnes fiducial region (0.45 t ^{136}Xe)

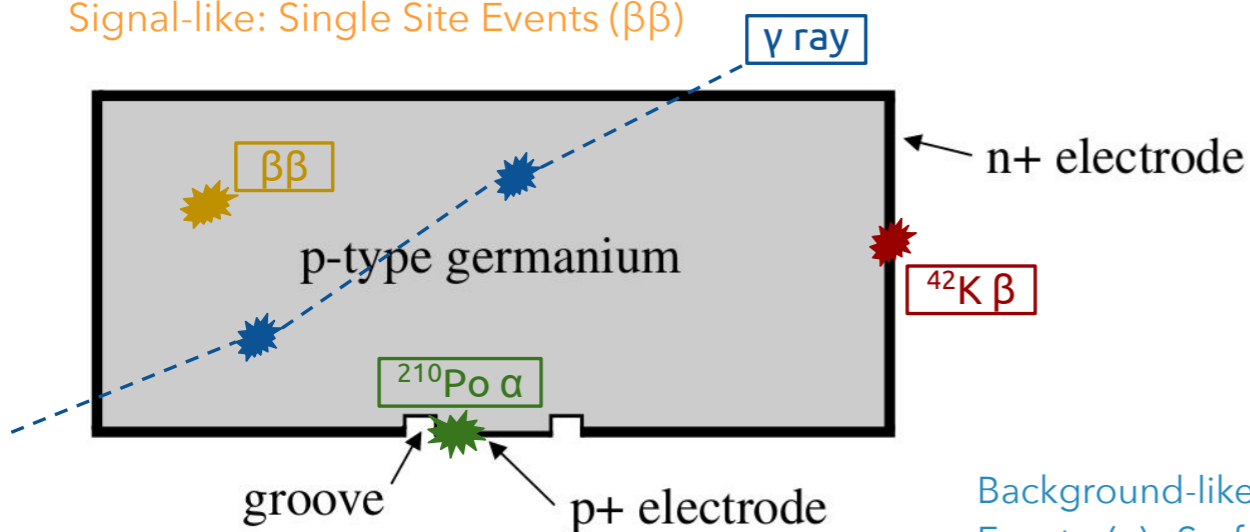
GERMANIUM IONISATION DETECTORS

▶ HPGe detectors enriched in ^{76}Ge

- Source = detector: high detection efficiency
- High-purity material: no intrinsic backgrounds
- Semiconductor: $\sigma/E < 0.1\%$ at $Q_{\beta\beta} = 2039.061 \text{ keV}$
- High stopping power: β absorbed within $O(1) \text{ mm}$



Signal-like: Single Site Events ($\beta\beta$)



Background-like: Multiple Site Events (γ), Surface events (α, β)

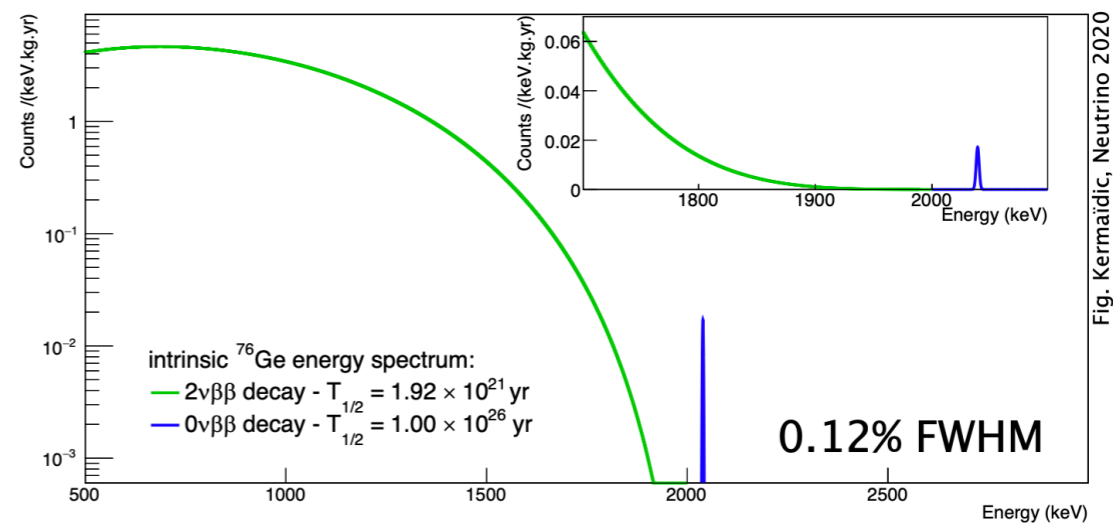
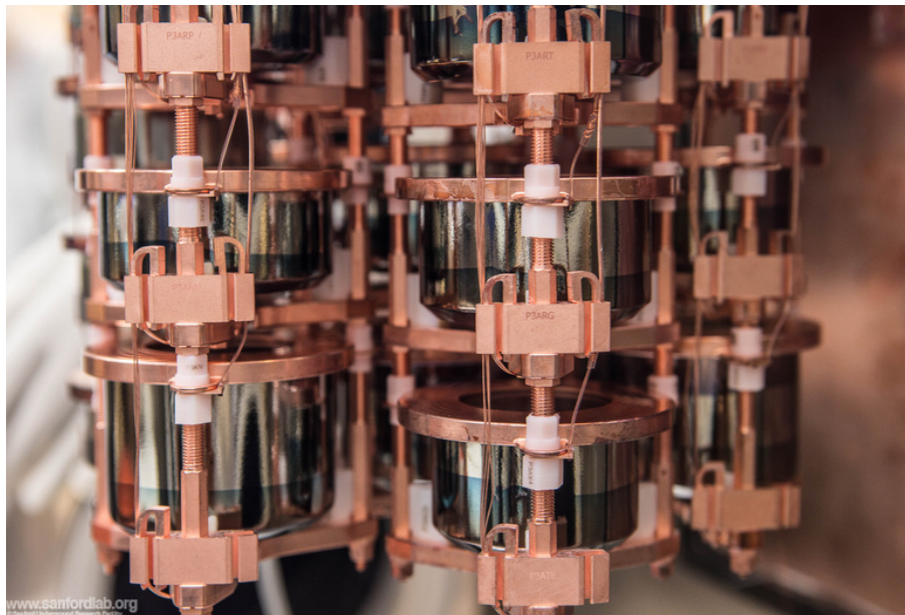


Fig. Kermaidic, Neutrino 2020

Neutrino 2020

RECENT GERMANIUM EXPERIMENTS



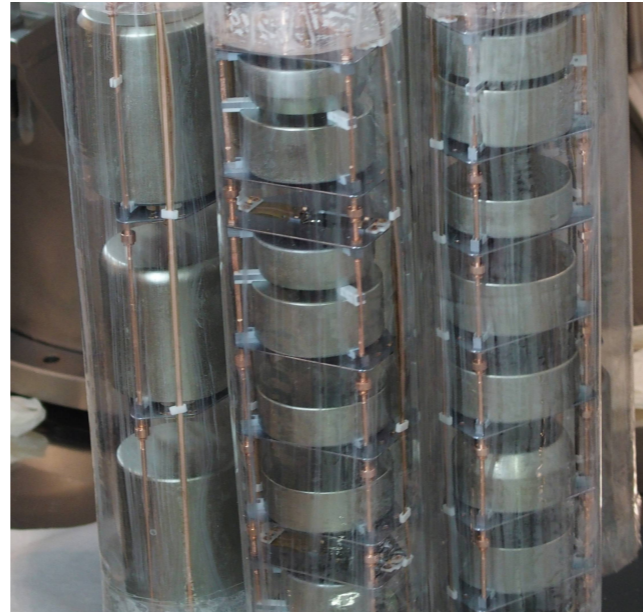
MAJORANA at SURF

29.7 kg of 88% enriched ^{76}Ge crystals

2.5 keV FWHM at 2039 keV

26 kg y exposure; PRL 120 (2018)

$T_{1/2} > 2.7 \times 10^{25} \text{ y}$ (90% CL)



GERDA at LNGS

35.6 kg of 86% enriched ^{76}Ge crystals in LAr

3.0 keV FWHM at 2039 keV

58.9 kg y exposure; Science 365 (2019), 127.2 kg y exposure: Neutrino 2020 & sub to PRL

$T_{1/2} > 1.8 \times 10^{26} \text{ y}$ (90% CL)

$Q_{\beta\beta} = 2039.061 \pm 0.007 \text{ keV}$

THE HEIDELBERG-MOSCOW EXPERIMENT

- ▶ Detectors in conventional shield: five ^{76}Ge detectors, mass 10.96 kg
- ▶ Concept to operate directly in cryogenic liquid:
 - ◉ Genius -> now GERDA->upcoming LEGEND

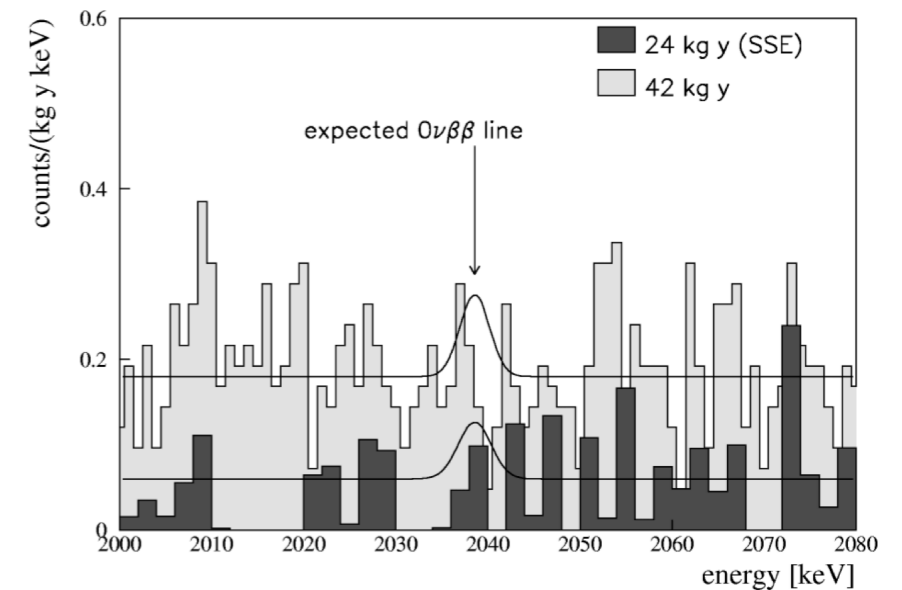


A first "bare" HPGe detector

GENIUS background and technical studies:
L. Baudis et al, NIM A 426 (1999)



Heidelberg-Moscow HPGe detector in conventional shield



Limits on the Majorana neutrino mass in the 0.1 eV range, L. Baudis et al., Phys. Rev. Lett. 83, 1999

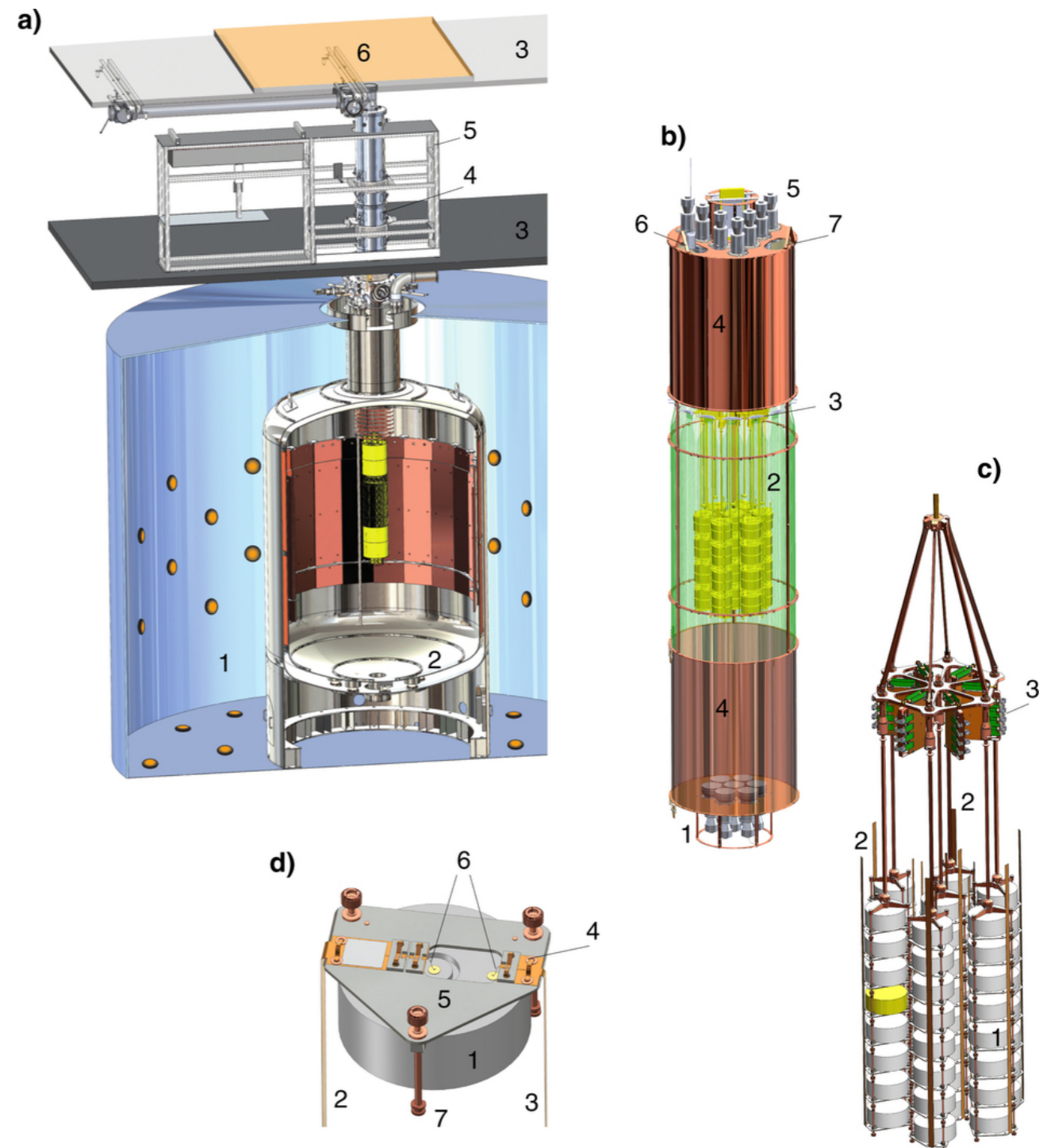
$$T_{1/2} > 1.6 \times 10^{25} \text{ y } 90\% \text{ C.L.}$$

Sensitivity

THE GERDA EXPERIMENT

GERDA collaboration, EPJ-C 78 (2018) 5

- ▶ Liquid Ar (64 m³) as cooling medium and shielding
- ▶ Surrounded by 590 m³ of ultra-pure water as muon Cherenkov veto
- ▶ U/Th in LAr < 7×10^{-4} $\mu\text{Bq/kg}$
- ▶ A minimal amount of surrounding material
- ▶ Two phases
 - Phase I: 2011-2014
 - Phase II: 2015-2019

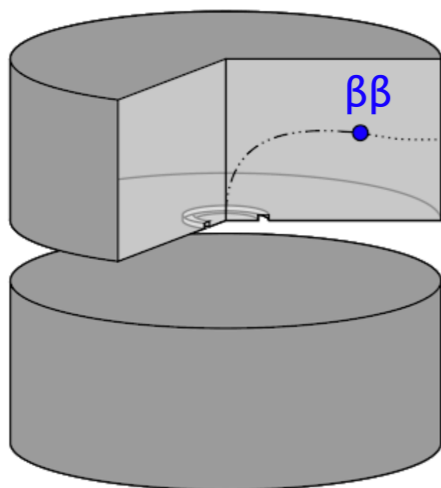


BACKGROUND SUPPRESSION

► Several handles:

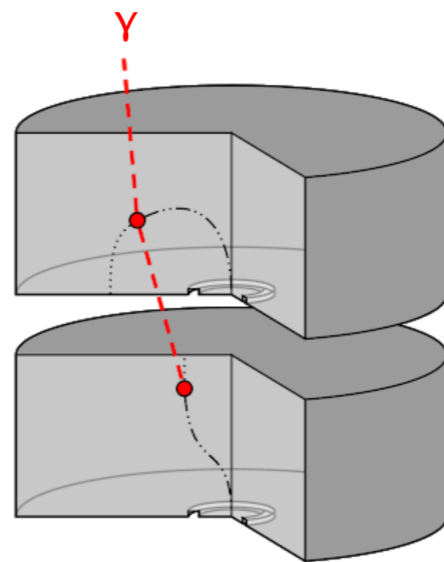
- Event topology + anti-coincidence between HPGe detectors + pulse shape discrimination + liquid argon veto

event topology



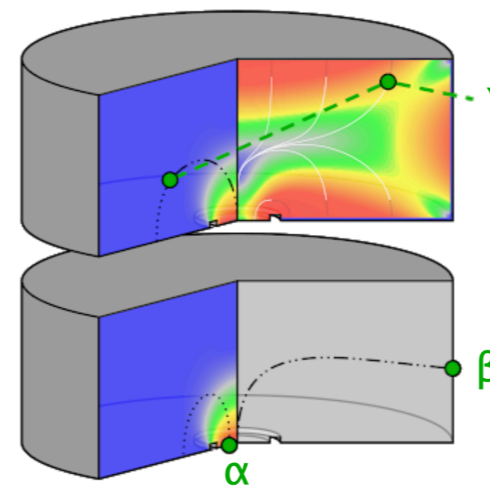
differentiate **point-like**
(single-detector, single-site)
 $\beta\beta$ topology from:

detector anti-coincidence



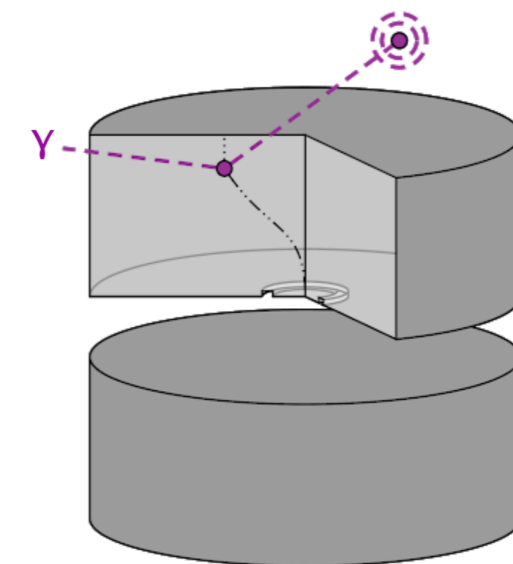
multi-detector
interactions

pulse shape discrimination (PSD)



multi-site/surface
interactions

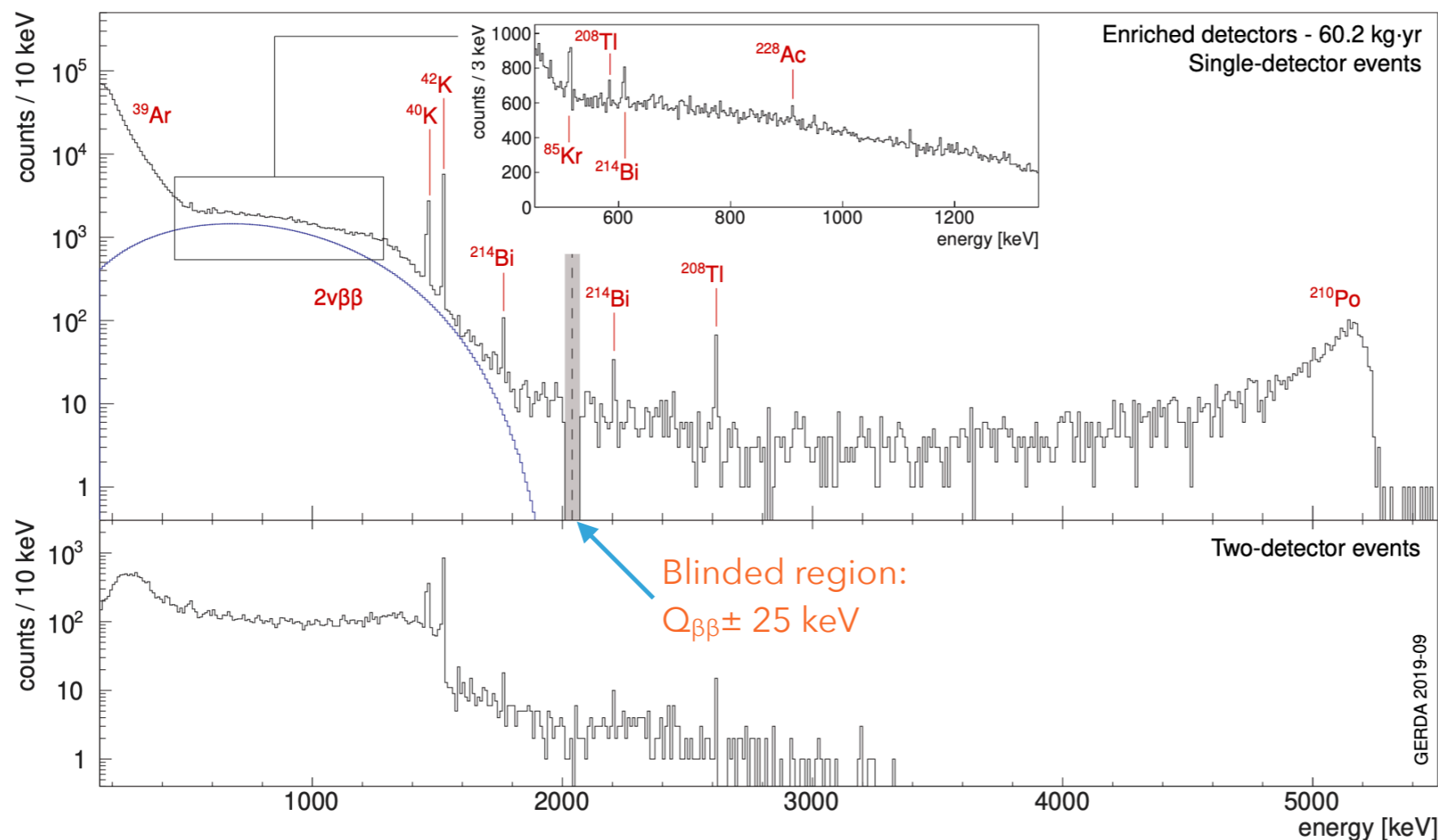
detector-LAr anti-coincidence (LAr veto)



interactions with **coincident**
energy deposition in
surroundings

BACKGROUND MODEL IN GERDA

- ▶ Intrinsic $2\nu\beta\beta$ -events, ^{39}Ar ($T_{1/2} = 269$ y), ^{42}Ar ($T_{1/2} = 33$ y) and ^{85}Kr ($T_{1/2} = 11$ y) in liquid argon
- ▶ ^{60}Co , ^{40}K , ^{232}Th , ^{238}U in materials, α -decays (^{210}Po) on the thin p^+ contact

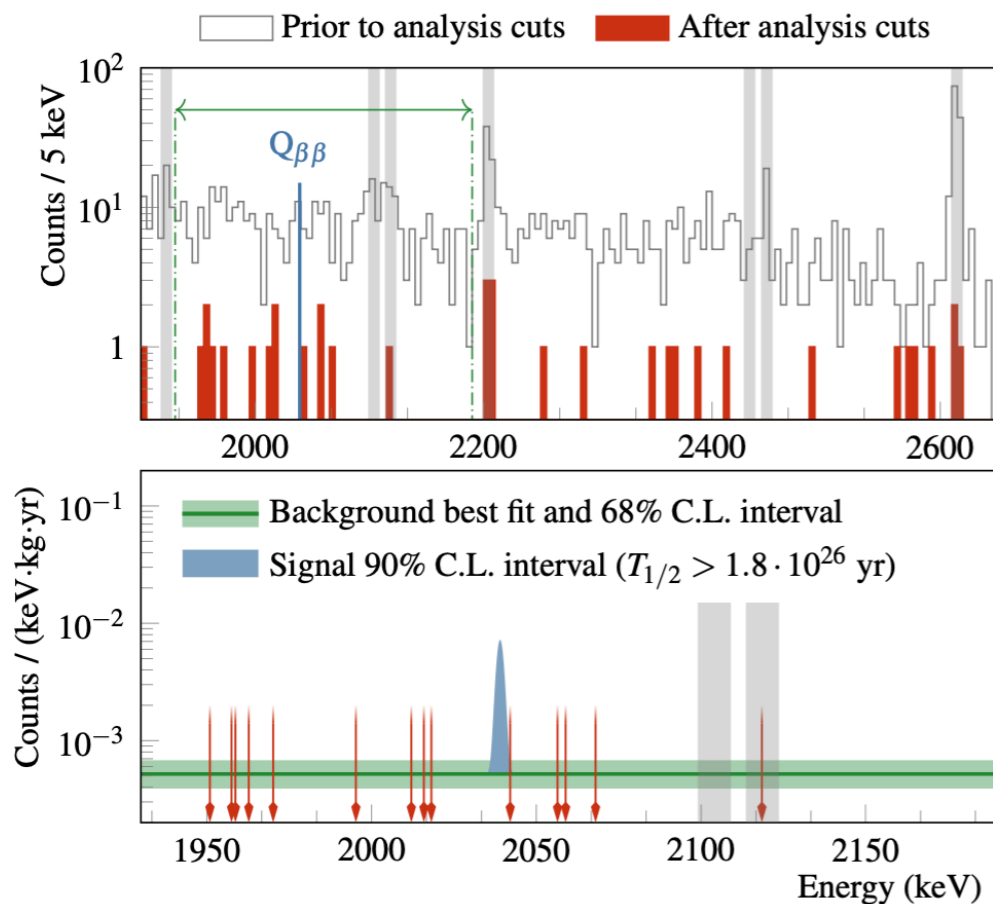


Sum spectrum,
single-detector
events

Sum spectrum,
two-detector
events

DOUBLE BETA DECAY FINAL RESULTS

- ▶ Measured $T_{1/2}$ of the $2\nu\beta\beta$ -decay: 1.92×10^{21} y
- ▶ Liquid argon veto: factor 5 background suppression at 1525 keV (^{42}K line)
- ▶ Background level: 5.2×10^{-4} events/(keV kg y) in 230 keV window around Q -value

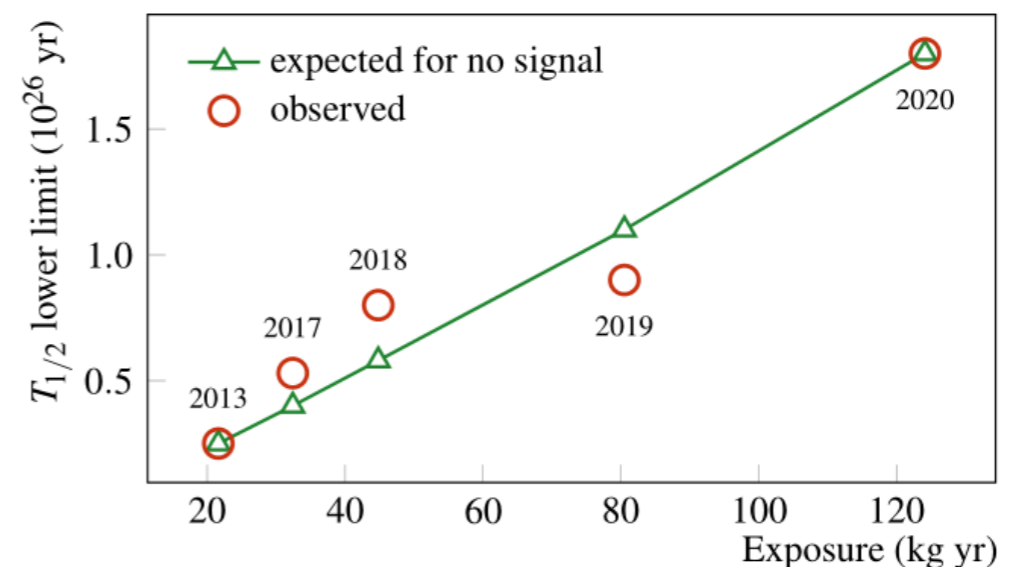


GERDA collaboration, Neutrino2020, submitted to PRL

Constraints on the ^{76}Ge $0\nu\beta\beta$ -decay

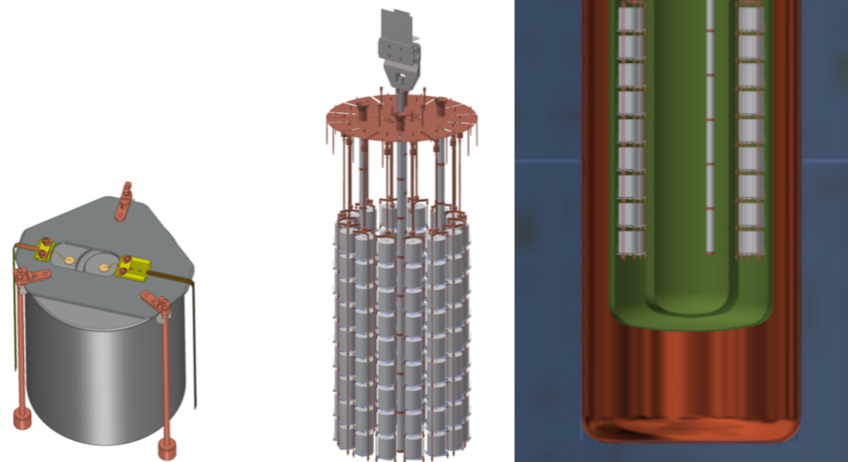
$$T_{1/2} > 1.8 \times 10^{26} \text{ y (90\% CL)}$$

$$m_{\beta\beta} < 80 - 182 \text{ meV}$$



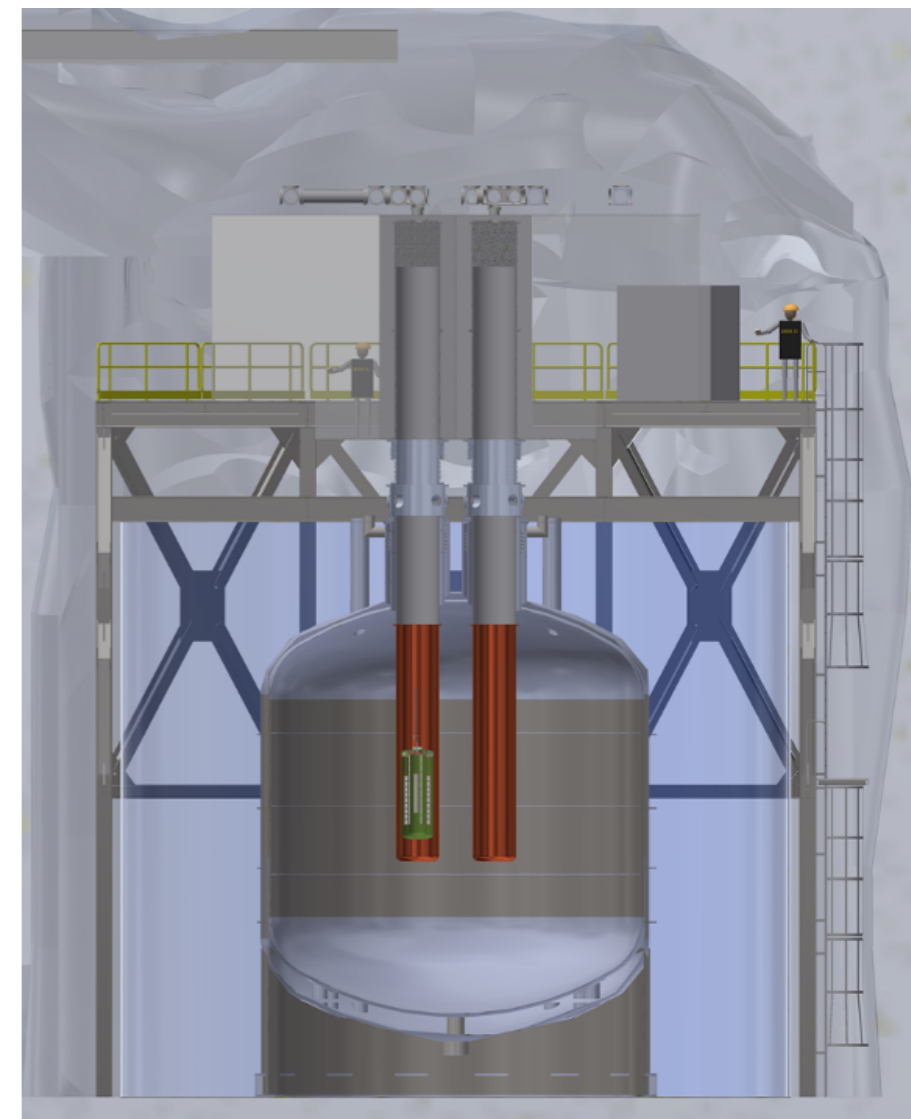
THE FUTURE: LEGEND

- ▶ Large enriched germanium experiment for $0\nu\beta\beta$ decay
- ▶ Collaboration formed in October 2016
- ▶ 219 members, 48 institutions, 16 countries
 - ▶ **LEGEND-200**: 200 kg in existing (upgraded) infrastructure at LNGS, to start in 2021
 - ▶ Background goal: 0.6 events/(FWHM t y)
 - ▶ **LEGEND-1000**: 1000 kg, staged, 4 modules
 - ▶ Background goal: 0.1 events/(FWHM t y)



LEGEND

Large Enriched
Germanium Experiment
for Neutrinoless $\beta\beta$ Decay

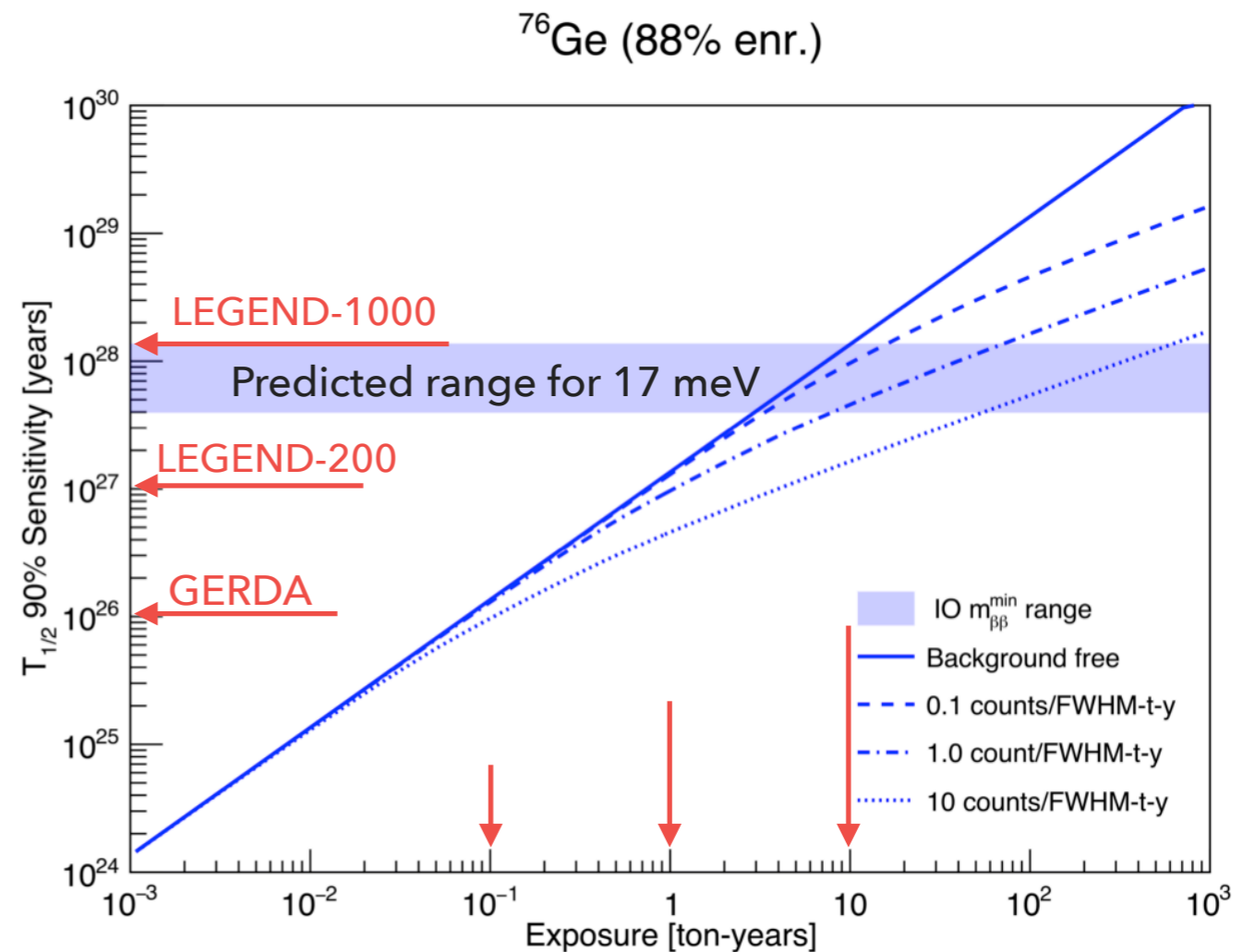


EXPECTED SENSITIVITY

- ▶ LEGEND-200: 10^{27} y
- ▶ LEGEND-1000: 10^{28} y
- ▶ $m_{\beta\beta} = 17$ meV (for worst case NME = 3.5)



Post GERDA tests with 20 Majorana, GERDA and new LEGEND detectors completed



Abgrall et al., AIP Conf. Proc. 1894(1), 020027 (2017)

Background

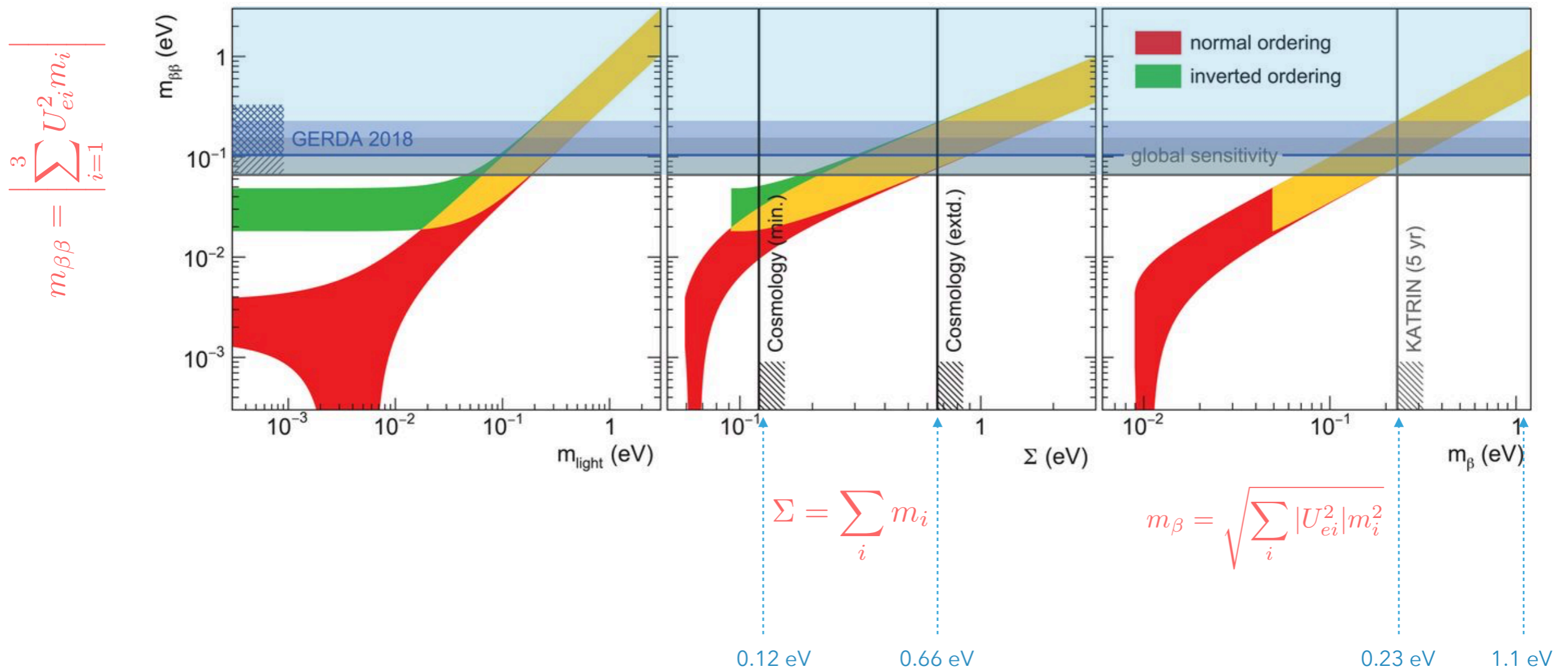
GERDA: 3 events/(ROI t y)
 LEGEND-200: 0.6 events/(ROI t y)
 LEGEND-1t: 0.1/(ROI t y)

LEADING RESULTS: OVERVIEW

Experiment	Isotope	FWHM [keV]	$T_{1/2}$ [10^{26} y]	$m_{\beta\beta}$ [meV]
CUORE	^{130}Te	7.4	0.15	162-757
CUPID-0	^{82}Se	23	0.024	394-810
EXO-200	^{136}Xe	71	0.18	93-287
KamLAND-Zen	^{136}Xe	270	1.1	76-234
GERDA	^{76}Ge	3.3	1.8	80-182
Majorana	^{76}Ge	2.5	0.27	157-346

MASS OBSERVABLES

- ▶ Constraints in the $m_{\beta\beta}$ parameters space in the 3 light ν scenario
- ▶ Global sensitivity from $0\nu\beta\beta$ -experiments & constraints from direct searches & cosmology



FUTURE PROJECTS: A SELECTION

$$|m_{\beta\beta}| \propto \left(\frac{B \cdot \Delta E}{M \cdot t} \right)^{\frac{1}{4}}$$

Experiment	Isotope	Iso mass [kg]	FWHM [keV]	$T_{1/2}$ [10^{27} y]	$m_{\beta\beta}$ [meV]
CUPID	^{130}Te	543	5	2.1	13-31
CUPID	^{82}Se	336	5	2.6	8-38
nEXO	^{136}Xe	4500	59	9	7-21
KamLAND2-Zen	^{136}Xe	1000	141	0.6	25-70
DARWIN	^{136}Xe	1068	20	2.4	11-46
PandaX-III	^{136}Xe	901	24	1.0	20-55
LEGEND-200	^{76}Ge	175	3	1	34-74
LEGEND-1t	^{76}Ge	873	3	6	11-28
SuperNEMO	^{82}Se	100	120	0.1	58-144

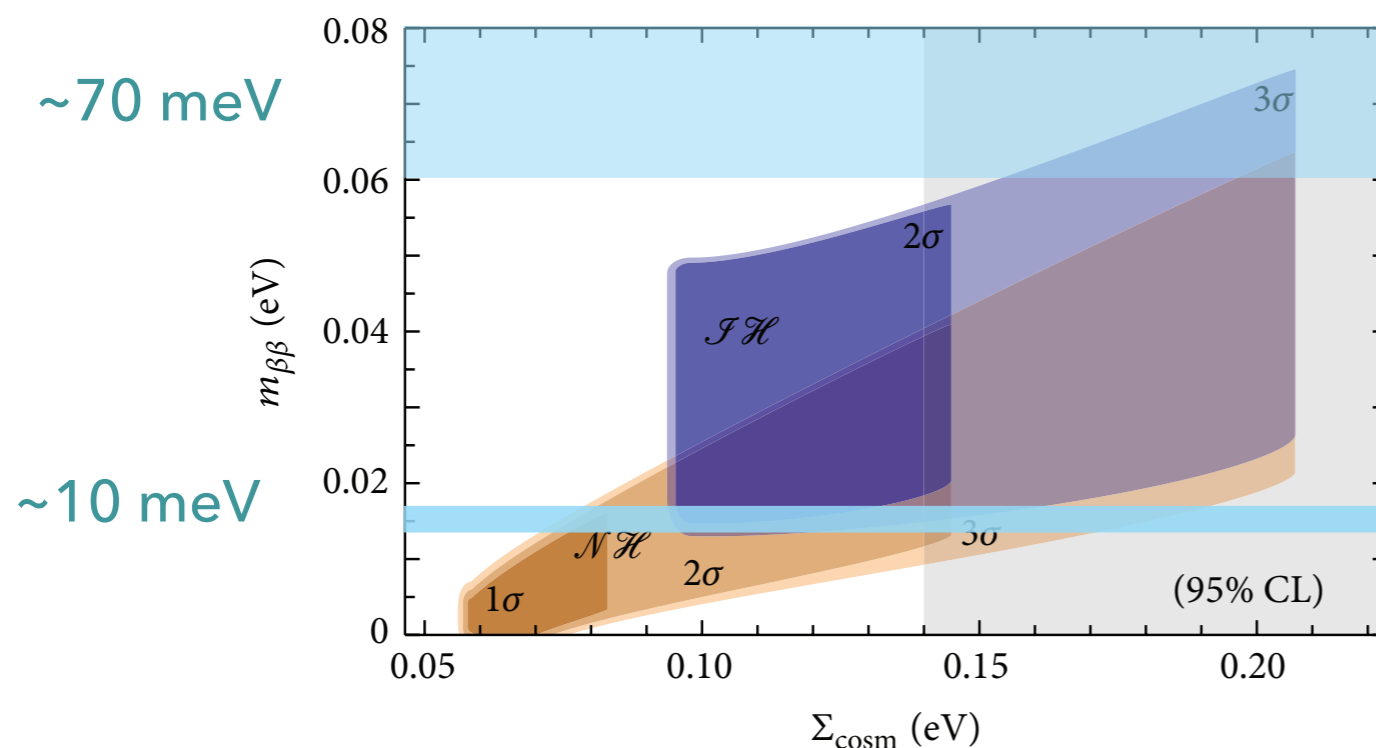
► Reminder

- Large exposures: 10 tonne x year, low background rates < 1 event/(FWHM tonne x year)
- Good energy resolution, large Q-value, high efficiency, demonstrated technology, etc

► Important to have multiple isotopes to make a convincing case for LNV

SUMMARY AND OUTLOOK

- ▶ Ninety years after Pauli postulated his “*silly child*”: many open questions in neutrino physics
- ▶ **Neutrinoless double beta decay: excellent tool to test LNV and the nature of neutrinos (Dirac vs Majorana)**
- ▶ Existing experiments probe $T_{1/2}$ up to $\sim 10^{26}$ years, with $T_{1/2} \sim (0.1 \text{ eV}/m_\nu)^2 \times 10^{26} \text{ y}$
- ▶ Ton-scale experiments are required to cover the *inverted mass ordering scenario*
 - ◉ Several technologies move into this direction
- ▶ Much larger experiments required to probe the *normal mass ordering*



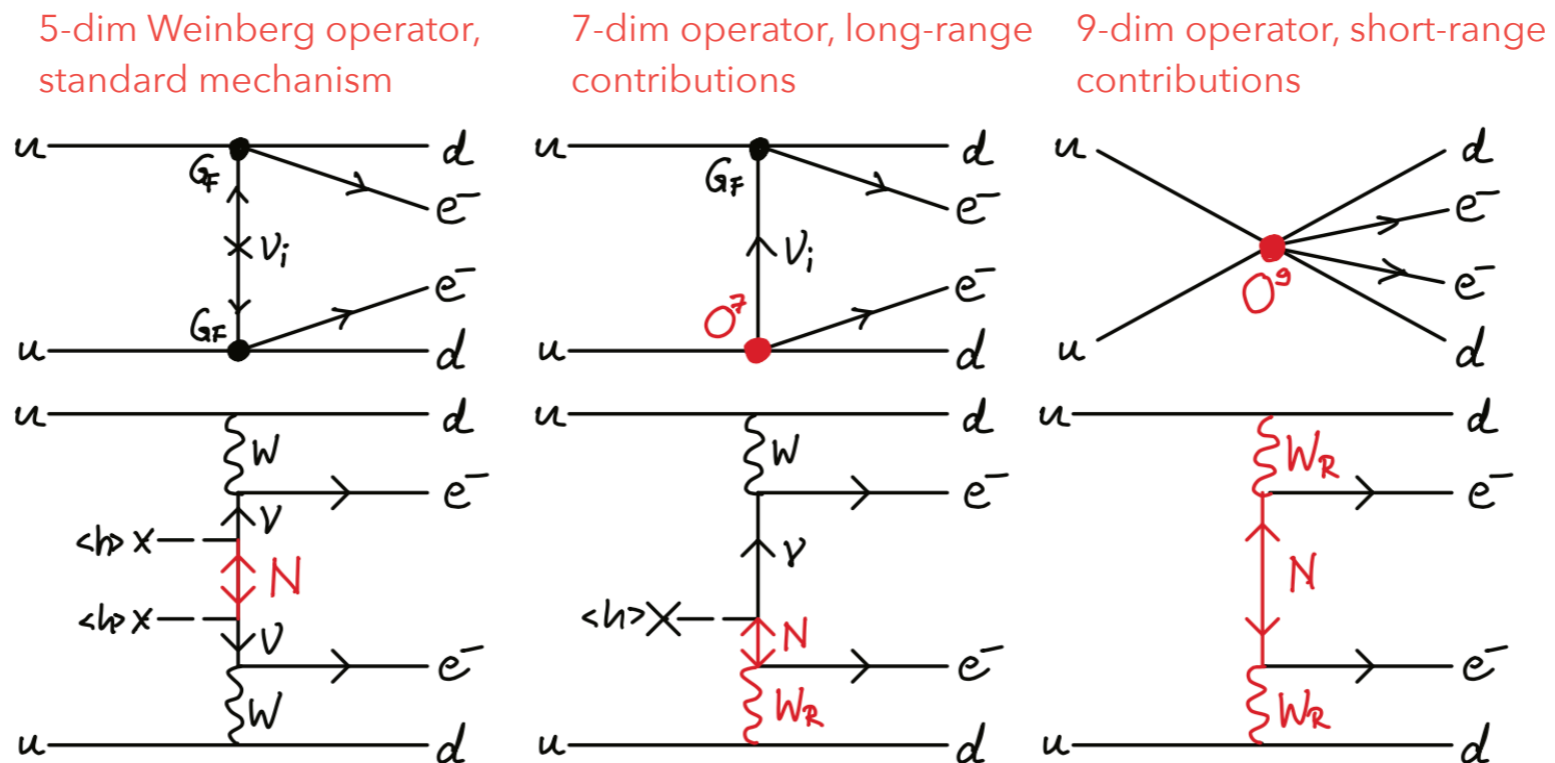
Current experiments

Future, ton-scale experiments

THANK YOU

OTHER MECHANISMS FOR DOUBLE BETA DECAY

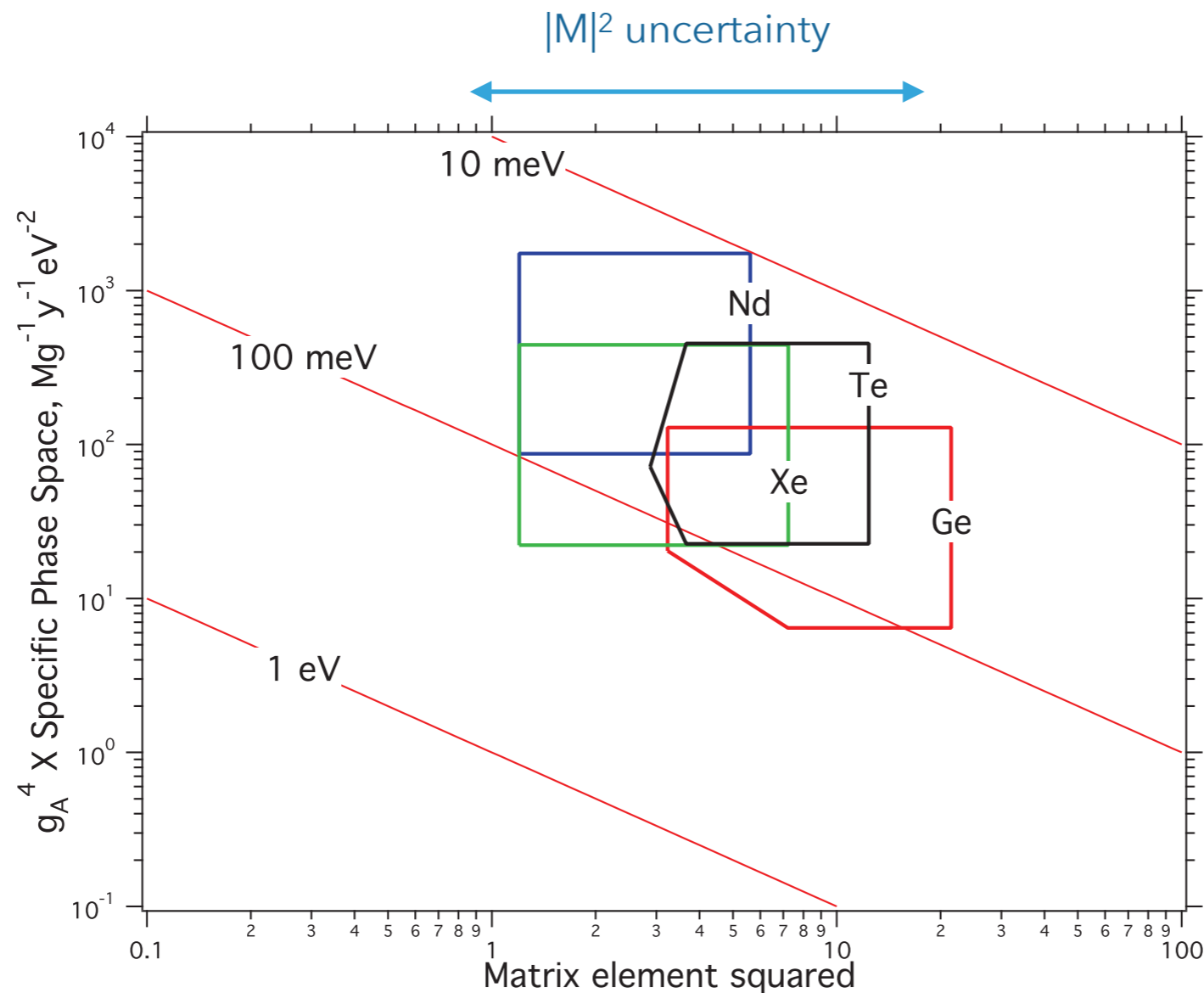
- ▶ LNV processes in extensions of the Standard Model generically contribute to $0\nu\beta\beta$ -decay (light or heavy sterile neutrinos, LR symmetric models, R-parity violating SUSY, leptoquarks, etc)
- ▶ Often classified as short- and long range processes, depending on the mass of the particles mediating the process (whether lighter or heavier than the momentum exchange scale $\sim O(100 \text{ MeV})$)
- ▶ In the effective Lagrangian picture, the effects at low energies can be summarised in terms of higher order operators, added to the SM Lagrangian



ISOTOPES AND SENSITIVITY TO THE DECAY

- ▶ Isotopes have comparable sensitivities in terms of rates per unit mass

$$g_A^4 \ln(2) \frac{N_A}{Am_e^2} G^{0\nu}$$

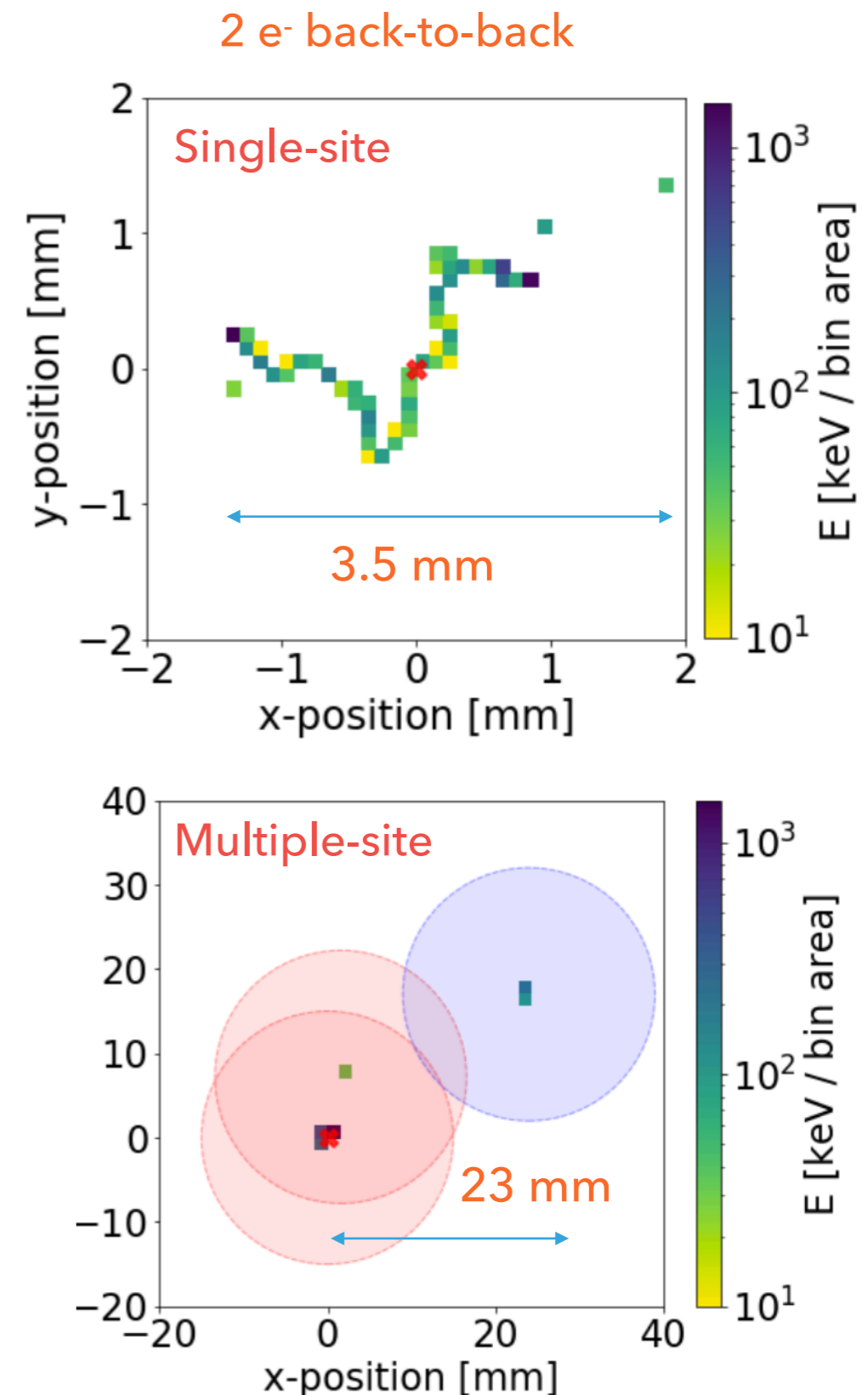


effective value for the axial vector coupling constant g_A : ~ 0.6 - 1.269 (free nucleon value)

SIGNAL EVENTS IN LIQUID XENON

- ▶ Electrons thermalise within $O(\text{mm}) \Rightarrow$ **single-site topology**
- ▶ Bremsstrahlung photons: may travel > 15 mm ($E > 300$ keV) \Rightarrow **multi-site event**
- ▶ Energy depositions: **spatially grouped using density-based spatial clustering algorithm**
 - ▶ New cluster, if distance to any previous $E_{\text{dep}} > \varepsilon$ (separation threshold)

Assumption: $\varepsilon = 15$ mm; 90% efficiency for $\beta\beta$ -events



MAIN BACKGROUND COMPONENTS

▶ Intrinsic:

- ▶ ^8B ν 's, ^{137}Xe , $2\nu\beta\beta$, ^{222}Rn

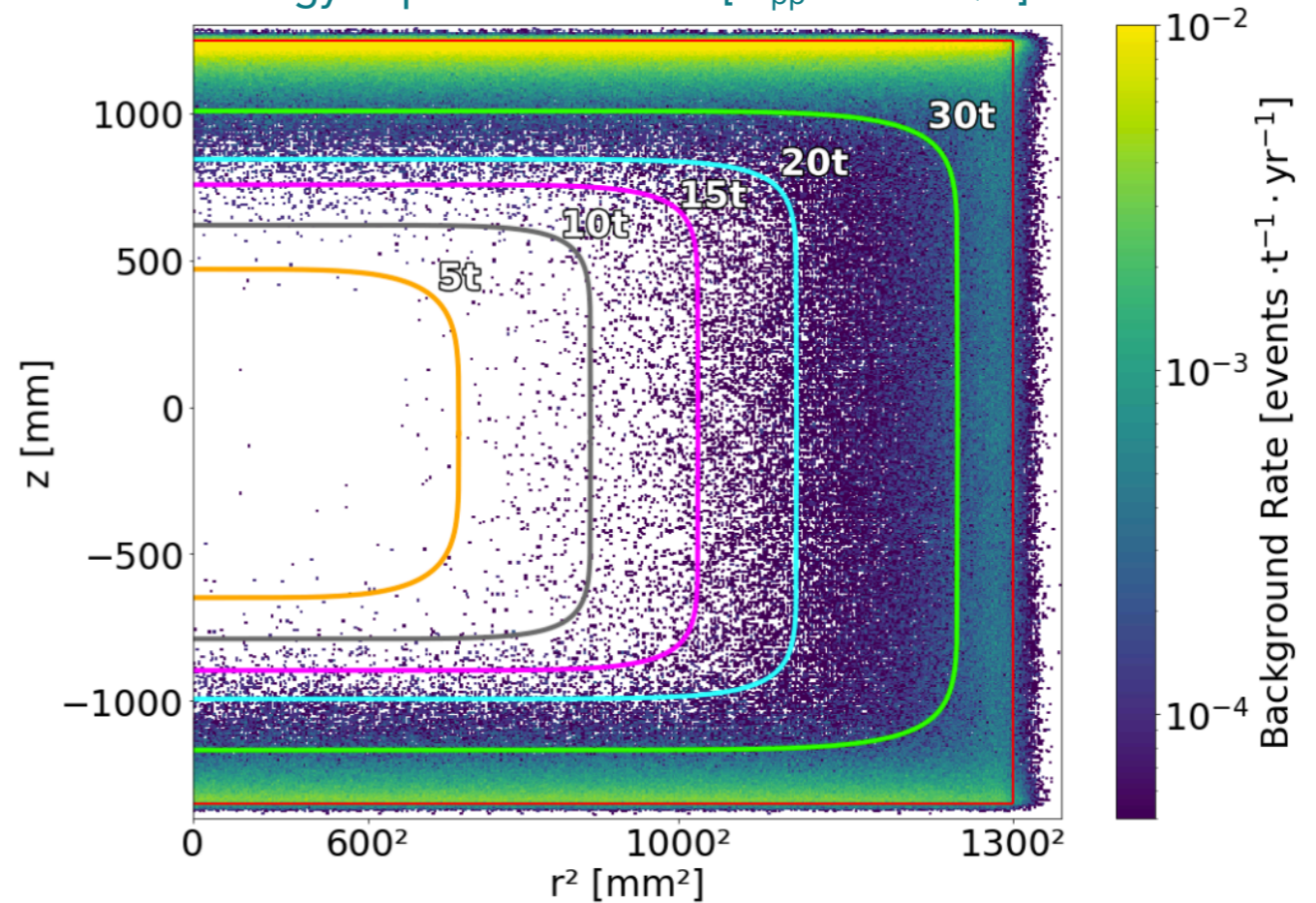
▶ Materials:

- ▶ ^{238}U , ^{232}Th , ^{60}Co , ^{44}Ti

▶ FV cut: super-ellipsoidal

$$\left(\frac{z + z_0}{z_{max}}\right)^t + \left(\frac{r}{r_{max}}\right)^t < 1$$

100 y of DARWIN run time, event with energy deposits in the ROI [$Q_{\beta\beta} \pm \text{FWHM}/2$]



Material	Unit	^{238}U	^{226}Ra	^{232}Th	^{228}Th	^{60}Co	^{44}Ti
Titanium	mBq/kg	<1.6	<0.09	0.28	0.25	<0.02	<1.16
PTFE	mBq/kg	<1.2	0.07	<0.07	0.06	0.027	-
Copper	mBq/kg	<1.0	<0.035	<0.033	<0.026	<0.019	-
PMT	mBq/unit	8.0	0.6	0.7	0.6	0.84	-
Electronics	mBq/unit	1.10	0.34	0.16	0.16	<0.008	-

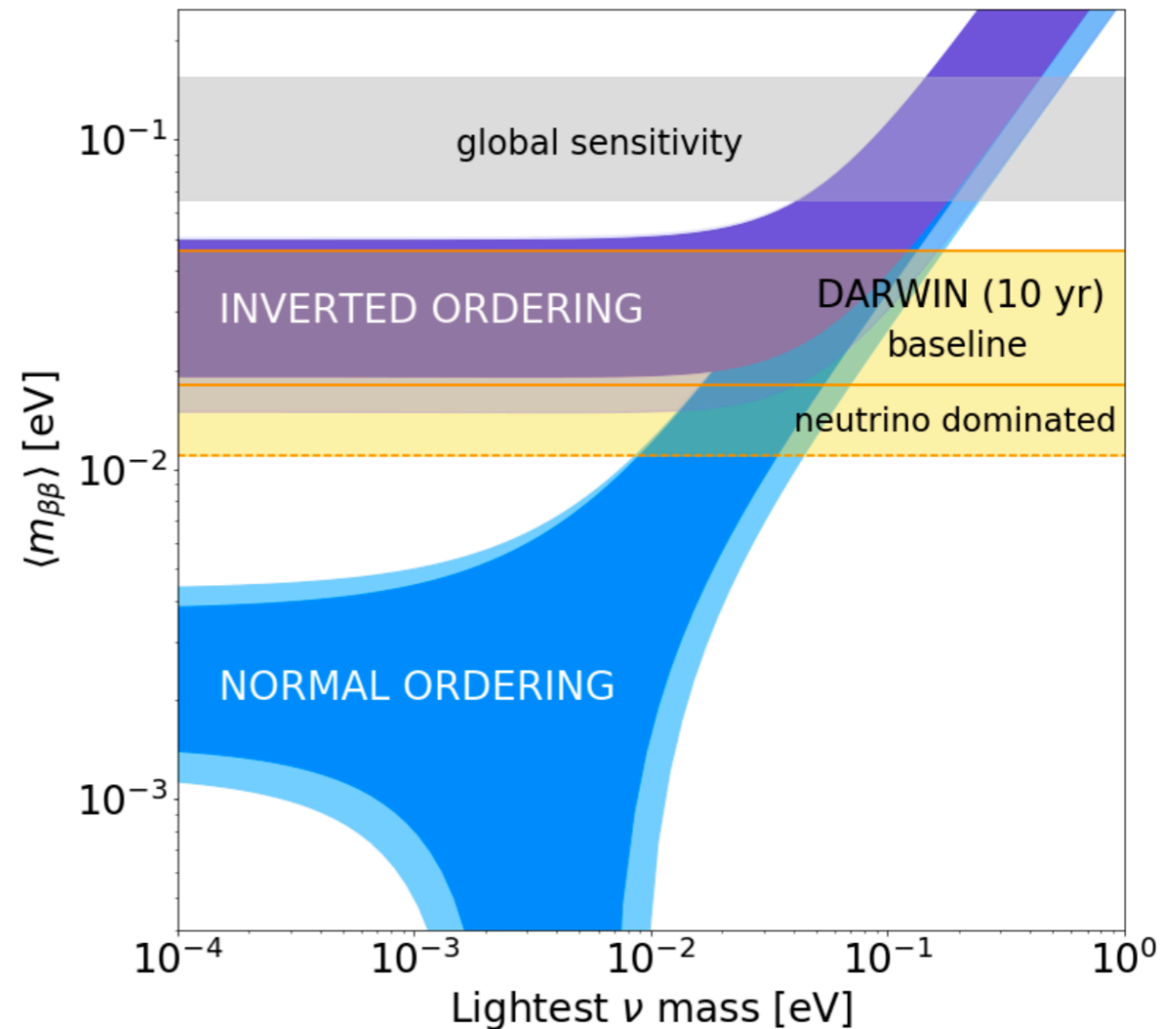
^{44}Ti : $T_{1/2} = 59$ y, cosmogenic

Ti: LZ, Astrop. Phys., 96 (2017)

Other: XENON, EPJ-C 77 (2017)

ROOM FOR IMPROVEMENT

- ▶ Reduce external backgrounds
 - ▶ SiPMs, cleaner materials & electronics
- ▶ Reduce internal background
 - ▶ Time veto for ^{137}Xe , deeper lab, BiPo tagging
- ▶ Improve signal/background discrimination; resolution...



Baseline: $m_{\beta\beta} = (18 - 46) \text{ meV}$

Neutrino dominated: $m_{\beta\beta} = (11 - 28) \text{ meV}$

NEUTRINO MASSES

- ▶ Three main methods: direct mass measurements, $0\nu\beta\beta$ -decay, cosmology
 - ▶ the observation of flavour oscillations imply a *lower bound on the mass of the heavier neutrino*
 - ▶ depending on the mass ordering, this lower bound is ≈ 0.05 eV

● The most direct probe: precision measurements of β -decays

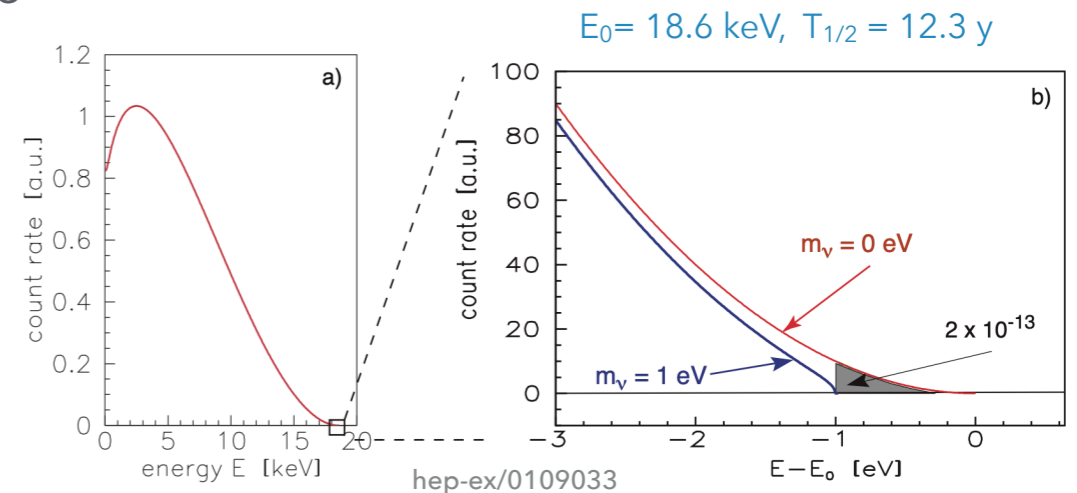


● The effect of a non-zero neutrino masses is observed kinematically: when a ν is produced, some of the energy exchanged in the process is spent by the non-zero neutrino mass

● The effects are however very small & difficult to observe

● KATRIN will probe the eff. ν_e mass down to 0.2 eV

$$m_{\nu_e}^2 = \sum_i |U_{ei}|^2 m_i^2$$



NEUTRINO MASSES

- ▶ Three main methods: direct mass measurements, $0\nu\beta\beta$ -decay, cosmology
 - ▶ the observation of flavour oscillations imply a *lower bound on the mass of the heavier neutrino*
 - ▶ depending on the mass ordering, this lower bound is ≈ 0.05 eV

- Cosmology: neutrinos influence the LSS and the CMB (with the ν density ratio):

$$\frac{\rho_\nu}{\rho_\gamma} = \frac{7}{8} N_{eff} \left(\frac{4}{11} \right)^{4/3} \quad N_{eff} = 3 \sim \text{number of active neutrinos}$$

- The constraints are on the sum of neutrino masses

$$\sum_i m_i$$

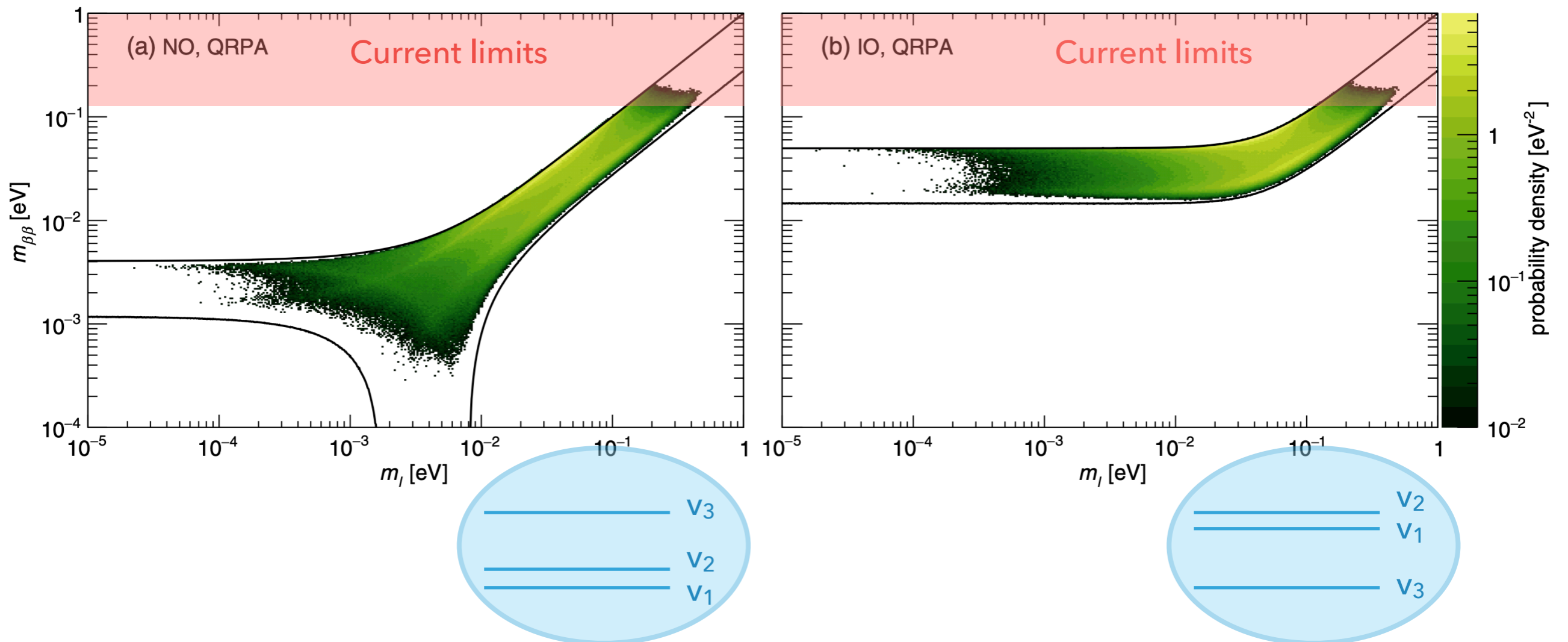
- Dependent on the parameters of the cosmological model (Λ CDM)
- In general, depending on which data is included (see e.g., review in PDG2020)

$$\sum_i m_i < (0.11 - 0.54) \text{ eV}$$

THE EFFECTIVE MAJORANA NEUTRINO MASS

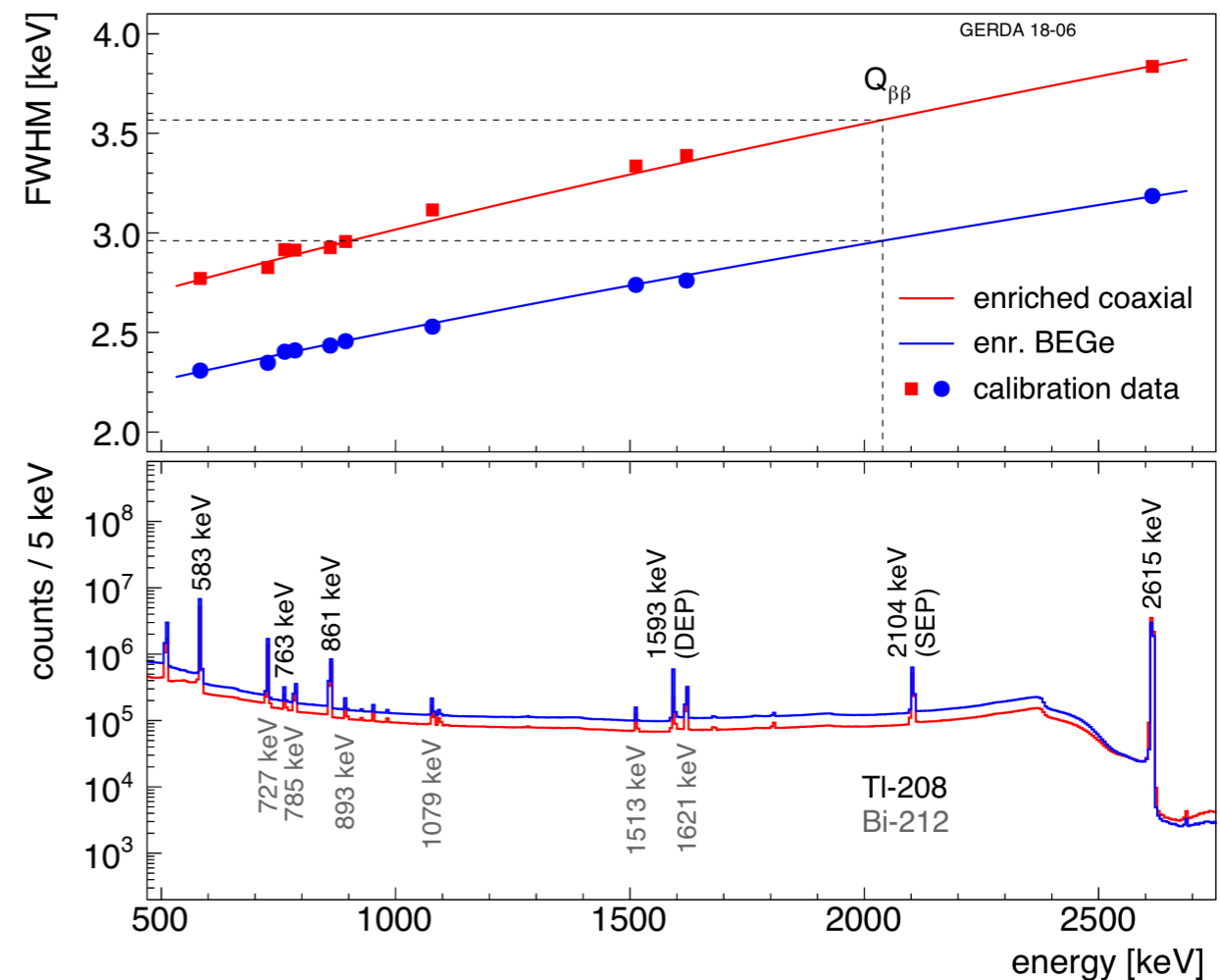
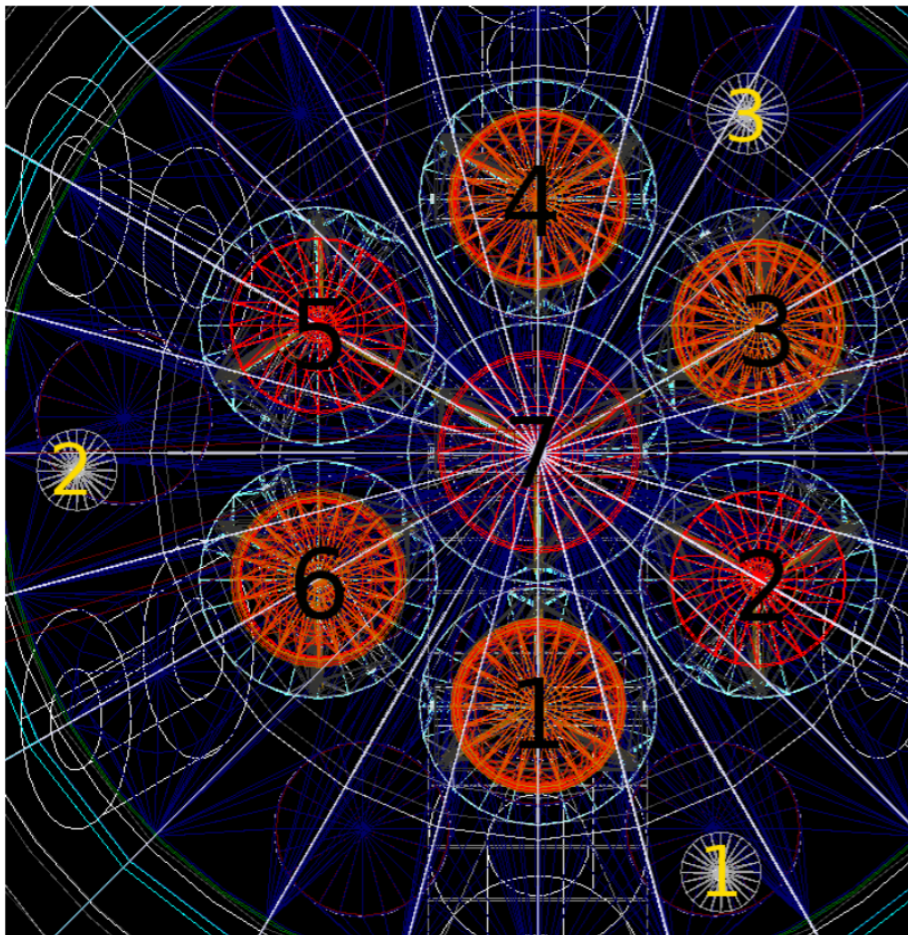
- ▶ Probability distribution of $m_{\beta\beta}$ via random sampling from the distributions of mixing angles and Δm^2
- ▶ Flat priors for the Majorana phases

Agostini, Benato, Detwiler, PRD 96, 2017



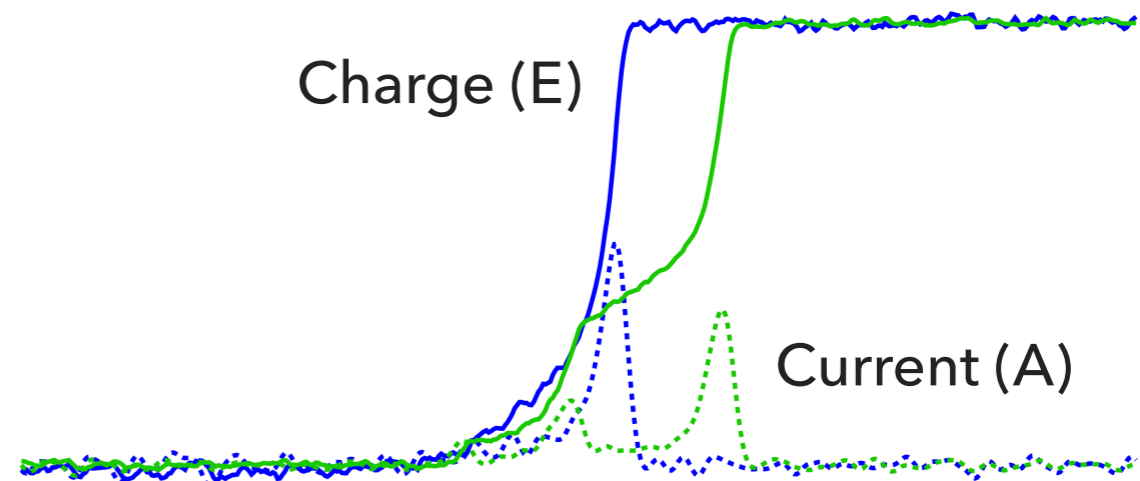
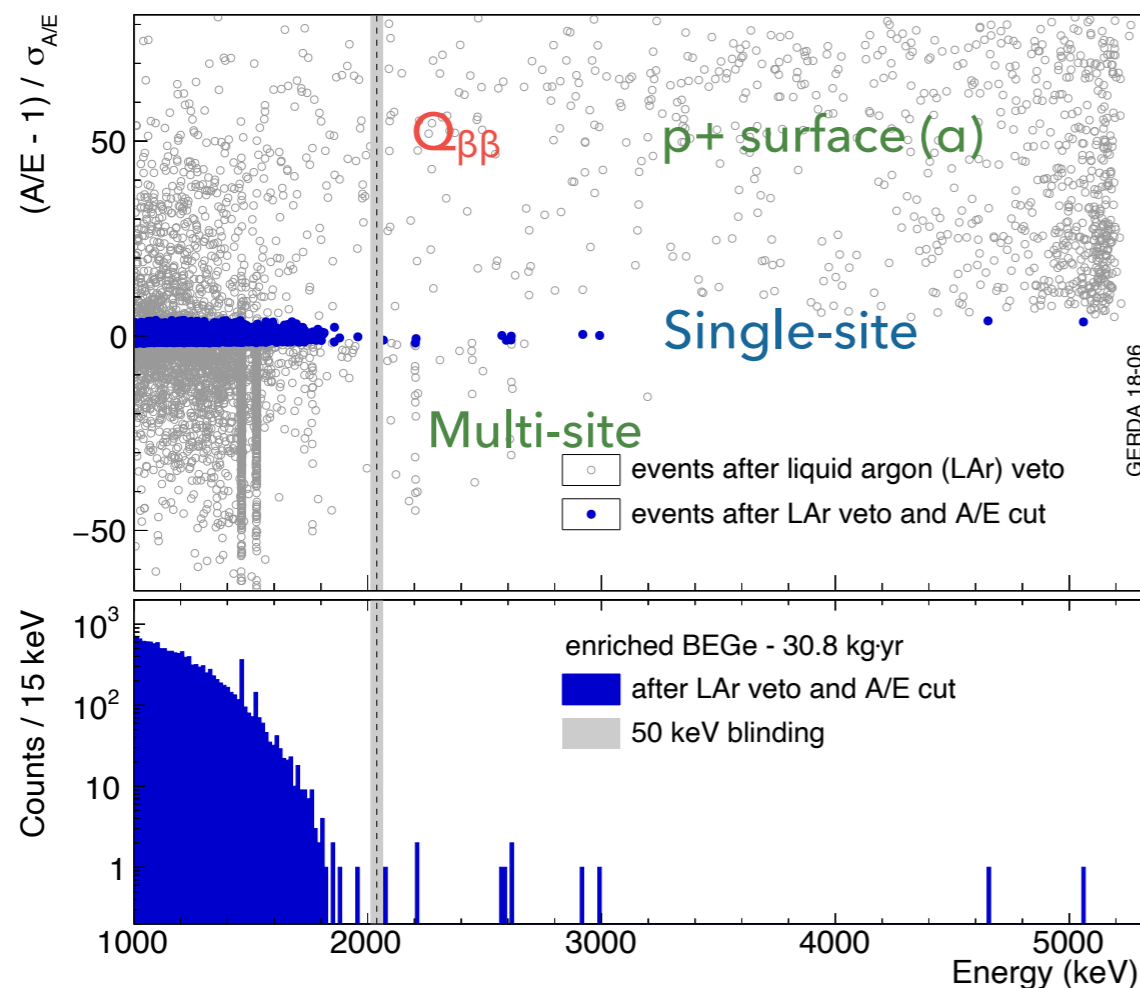
ENERGY CALIBRATION

- ▶ Three low neutron-emission ^{228}Th sources in SIS, deployed once every week
- ▶ FWHM at $Q_{\beta\beta}$: (3.0 ± 0.1) keV for BEGe, (3.6 ± 0.1) keV for coaxial detectors



PULSE SHAPE DISCRIMINATION

- ▶ Cut based on 1 parameter: max of current pulse (A) normalised to total energy (E) (BEGe)
- ▶ Tuned on calibration data (90% ^{208}Tl DEP acceptance)
- ▶ Acceptance at $0\nu\beta\beta$: $(87.6\pm 2.5)\%$



PSD parameter: $(A/E - 1) / \sigma_{A/E}$

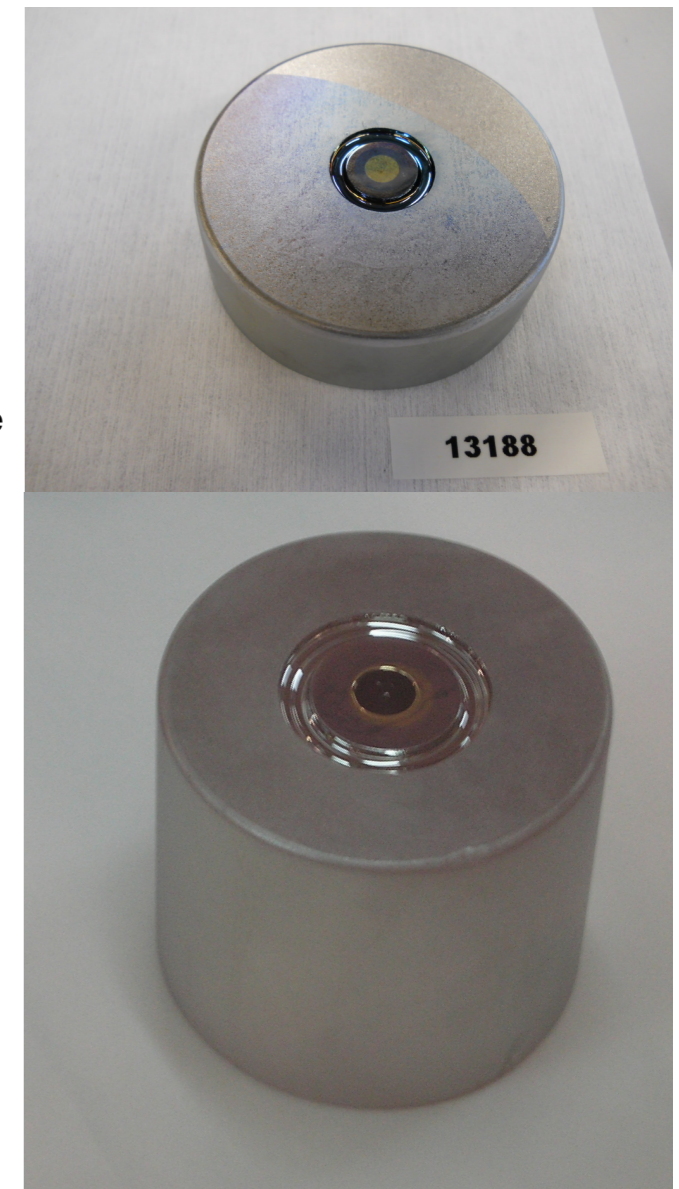
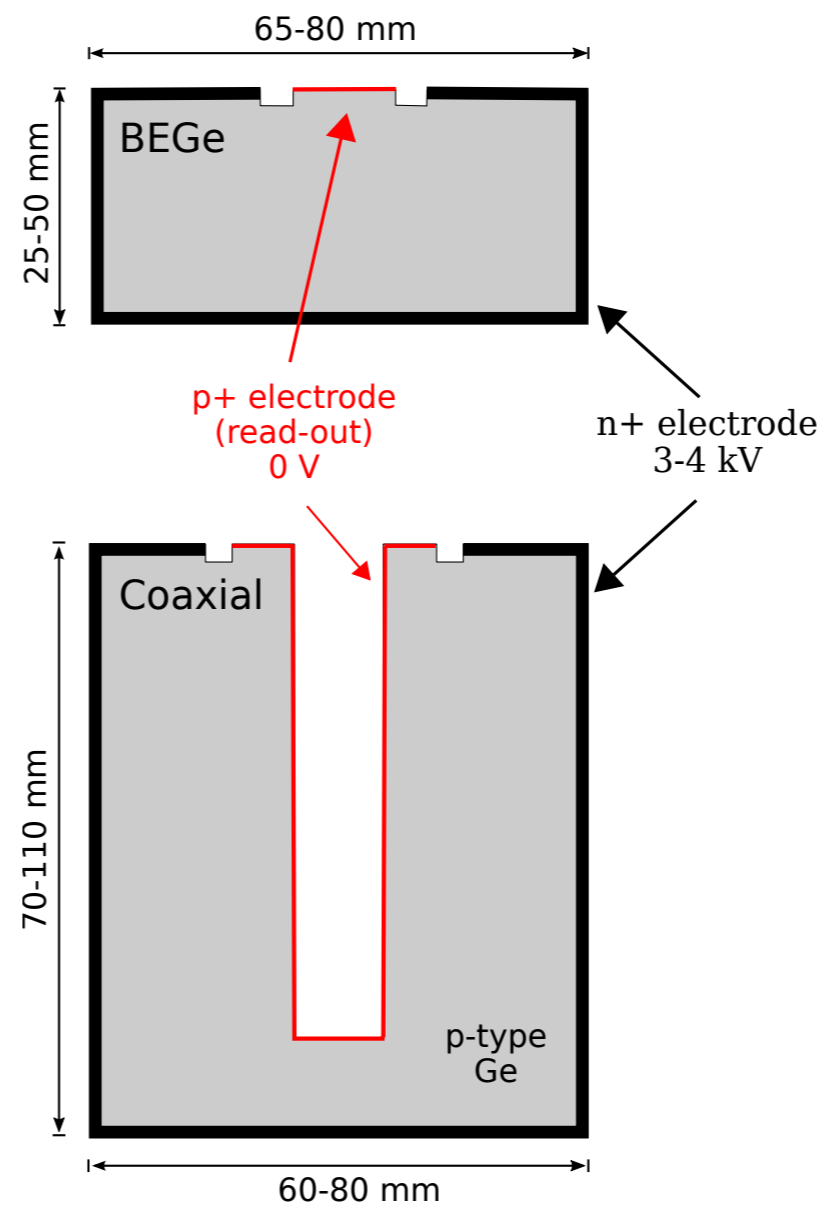
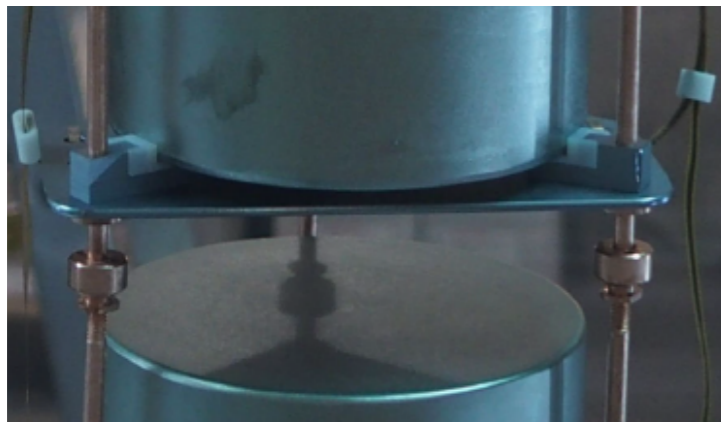
Mean and resolution corrected for E-dependance

A/E normalised to 1

Accept events around $(A/E - 1) / \sigma_{A/E} = 0$

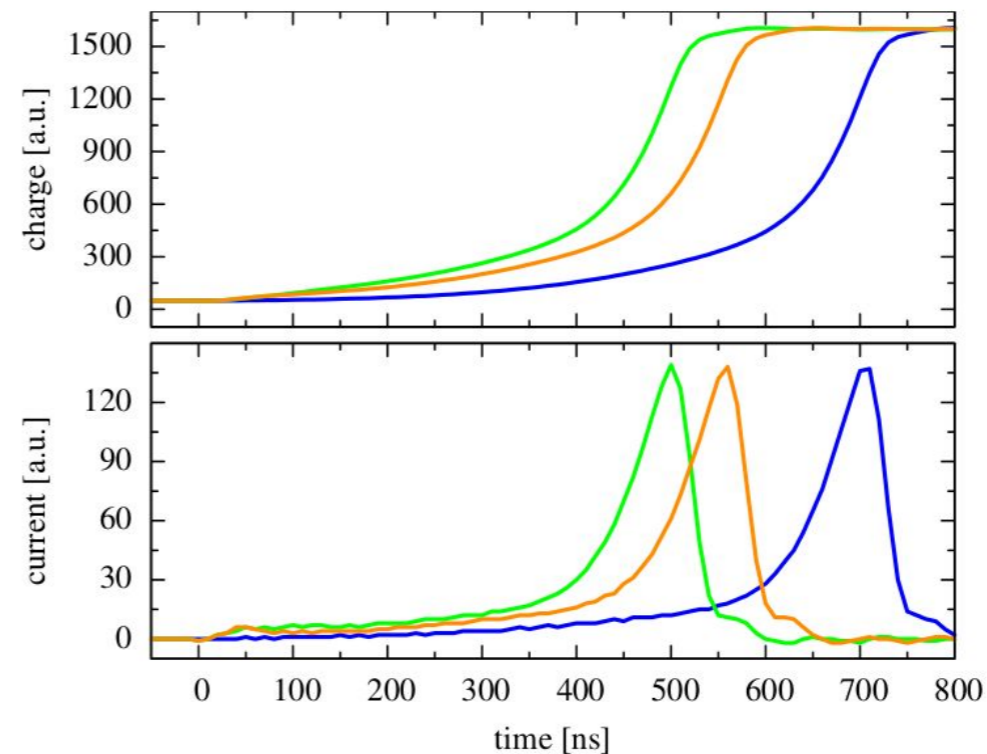
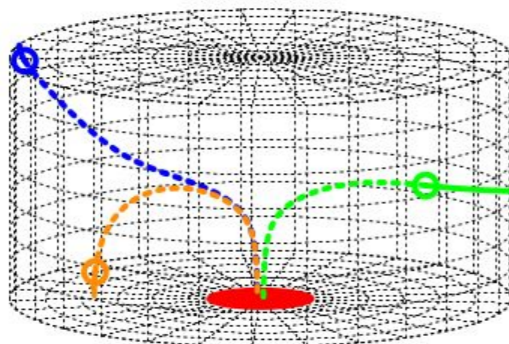
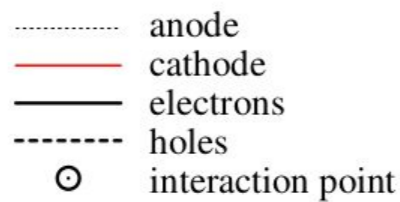
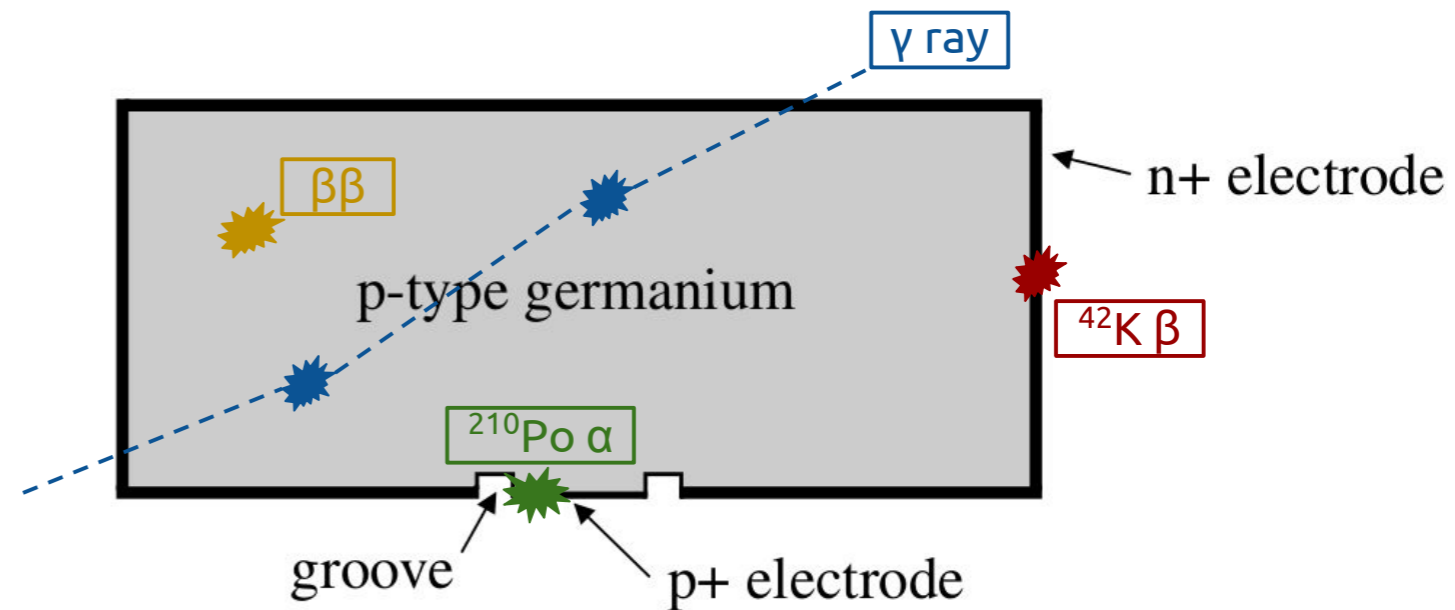
GERDA PHASE-II DETECTORS

- BEGe and coaxial
- p+ electrodes:
 - 0.3 μm boron implantation
- n+ electrodes:
 - 1-2 mm lithium layer (biased up to +4.5 kV)
- Low-mass detector holders (Si, Cu, PTFE)



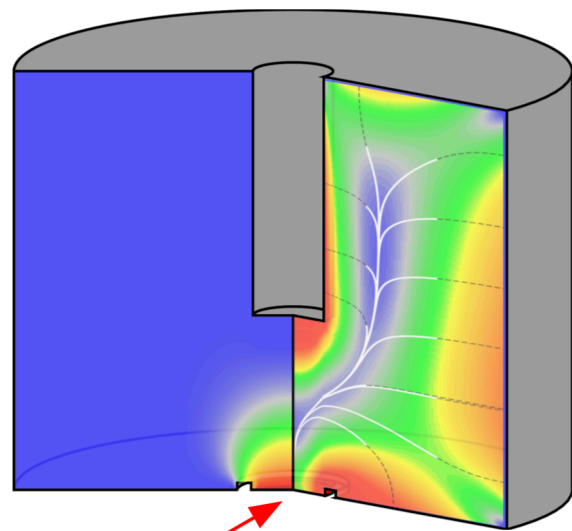
GERDA PULSE SHAPE DISCRIMINATION

- Signal-like: Single Site Events (SSE)
- Background-like: Multiple Site Events (MSE)
- BEGe detectors: E-field and weighting potential has special shape: pulse-height nearly independent of position



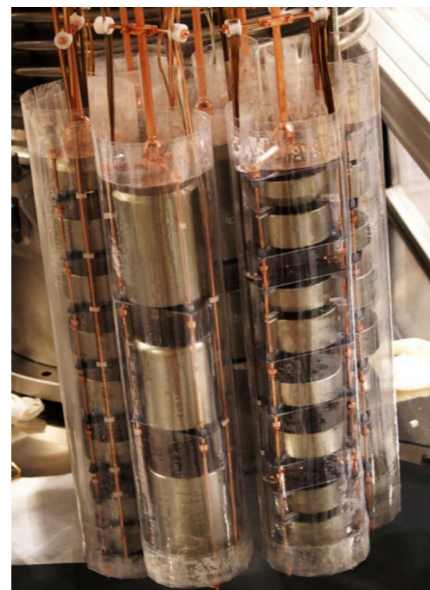
INVERTED COAXIAL DETECTORS

- ▶ Large point-contact detectors with ~ 3 kg mass, excellent PSD performance
- ▶ First 5 enriched IC detectors installed in spring 2018; baseline for LEGEND

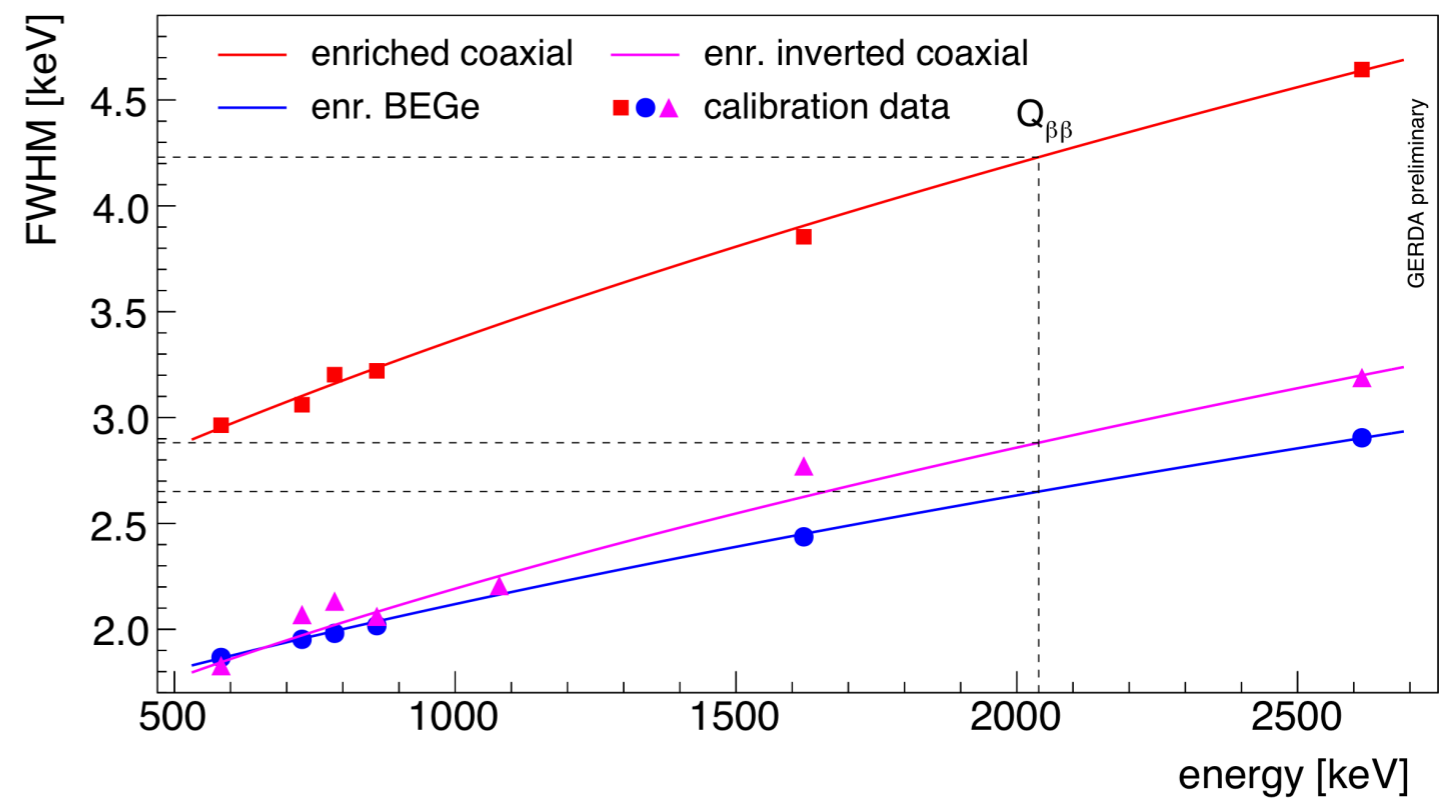


point contact

R.J Cooper et al.,
NIM A 665 (2011) 25



Detector mass
increase: 35.6 kg \rightarrow
44.2 kg



FWHM at $Q_{\beta\beta}$ [keV]: 4.2 ± 0.1 coax; 2.7 ± 0.1 BEGe; 2.9 ± 0.1 IC