

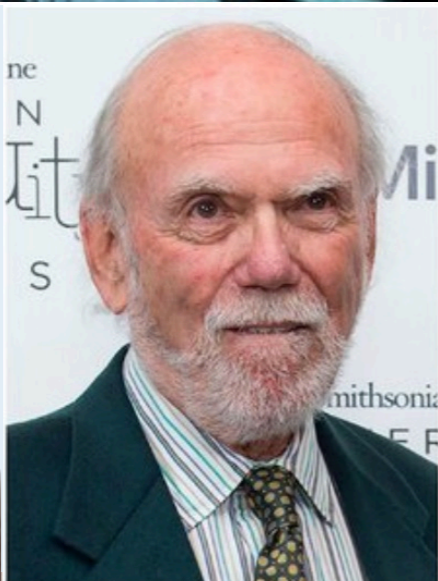
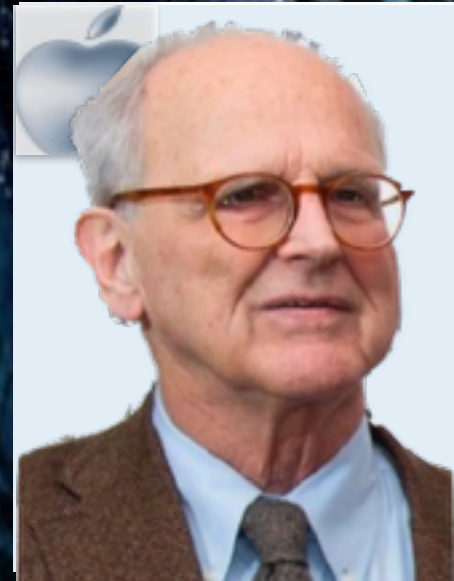
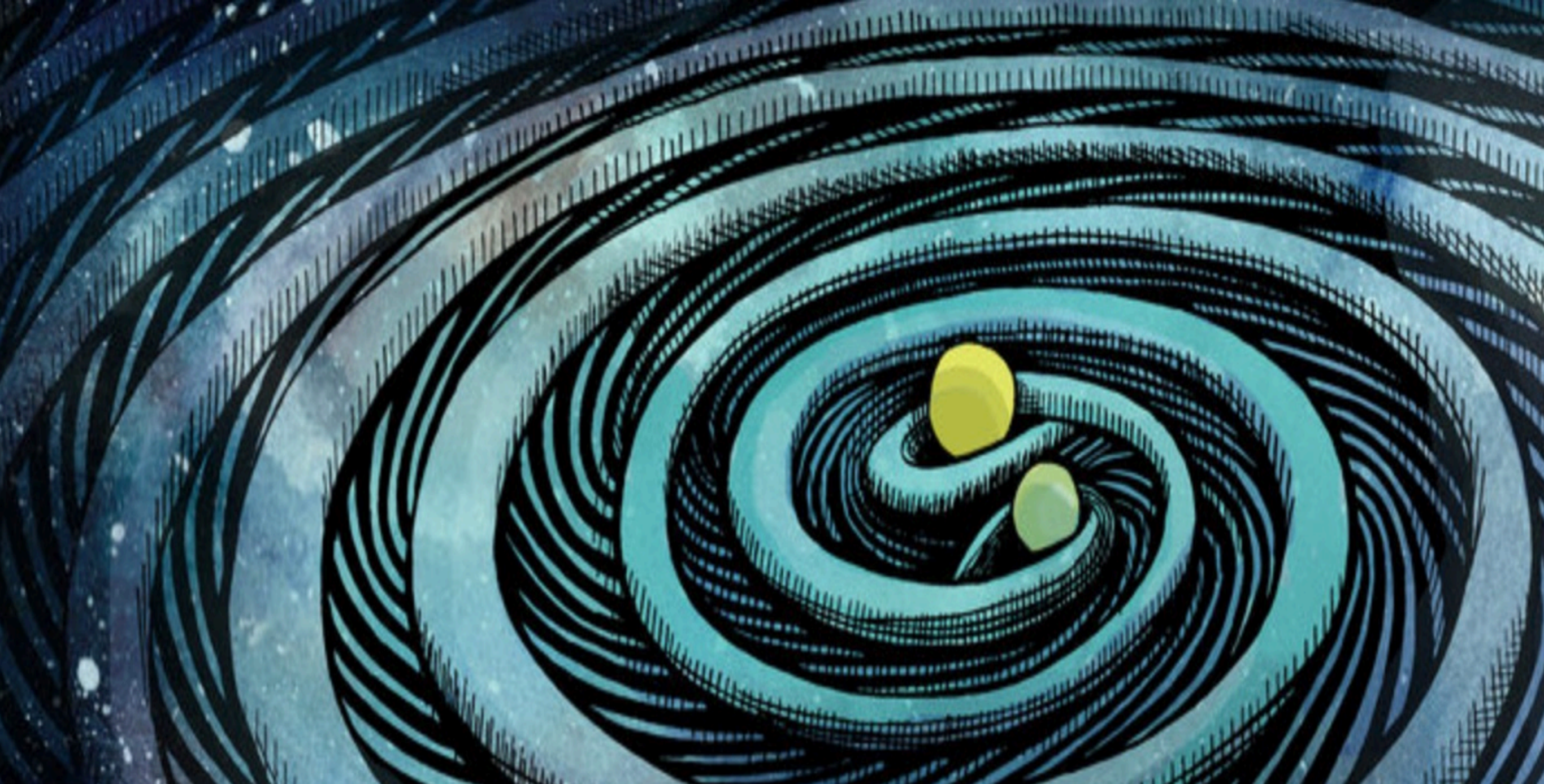
FNAL, Oct 4 2017

Physics at a 100 TeV pp collider



Michelangelo L. Mangano
michelangelo.mangano@cern.ch
Theoretical Physics Department
CERN





Rainer Weiss

Barry Barish

Kip Thorne

A “real” story from the past ...

Barcelona, 15 March 1493



Cristoforo Colombo:

Your Majesty, the fleet needs an **upgrade**, we need to go back to the Indies with **10 times** more ships

King Ferdinand and Queen Isabella:

You discovered the Indies, your theory is right, why do you need more?

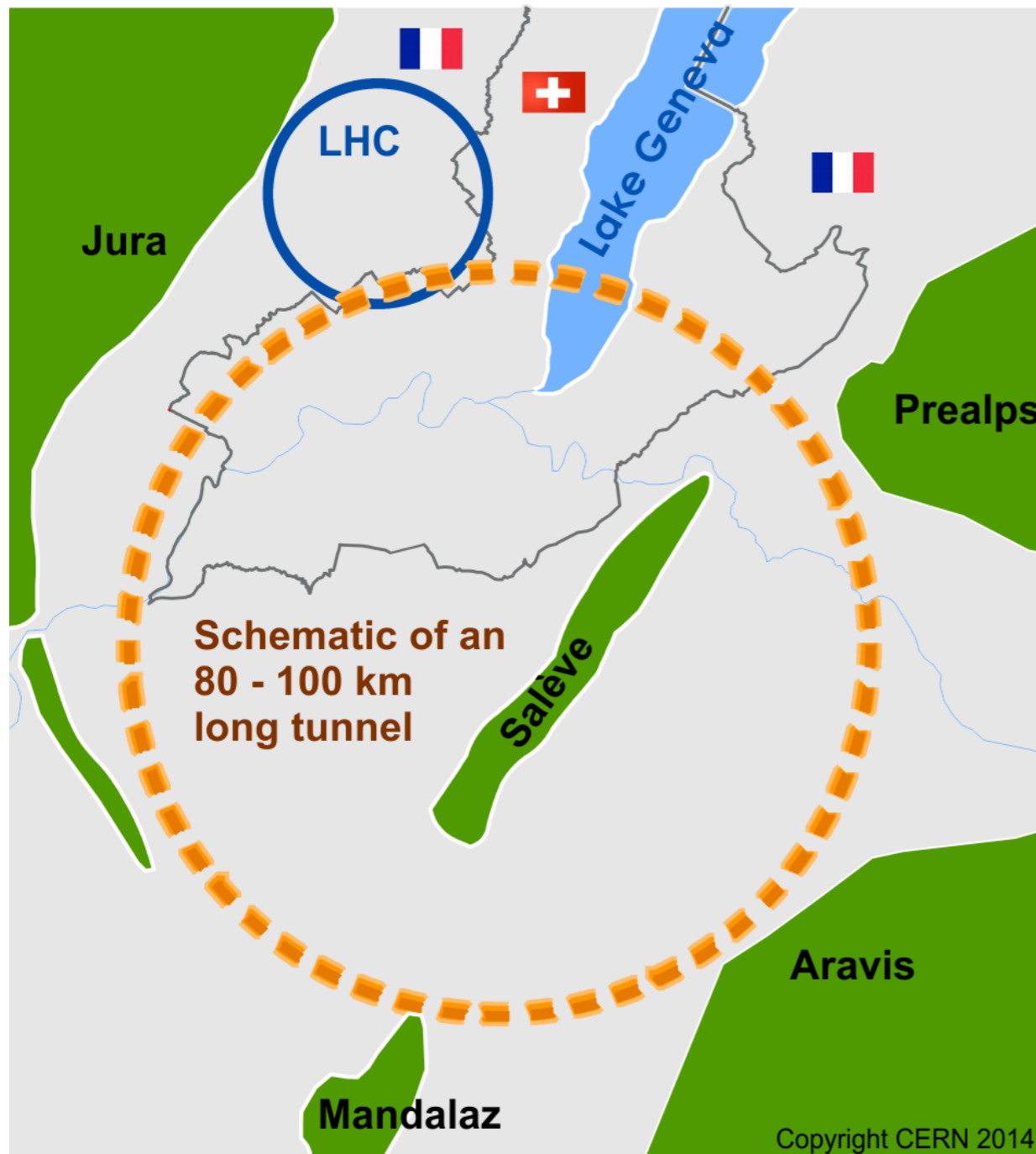
Cristoforo Colombo:

Theorists* say these may not be the **standard Indies**. They calculated the Earth radius, and the standard Indies cannot be so close: these are likely to be **beyond the standard Indies** (*moving eastward ...*)

** If the King had listened to theorists to start with, he would have never authorized the mission: everyone would have died of starvation well before reaching the “standard” Indies ...*

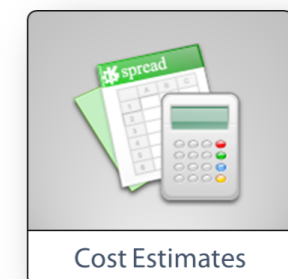
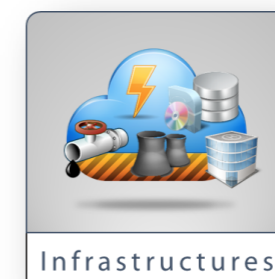
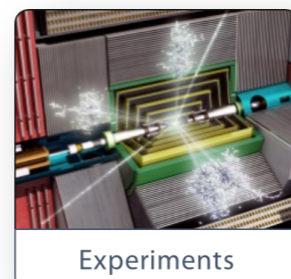
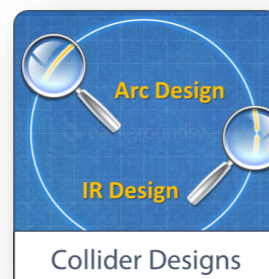
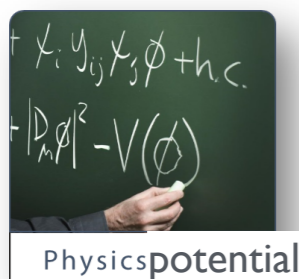
The context of this talk: Future Circular Colliders (FCC)

with emphasis on the pp facility, see *Blondel Wine&Cheese* for the e^+e^- facility

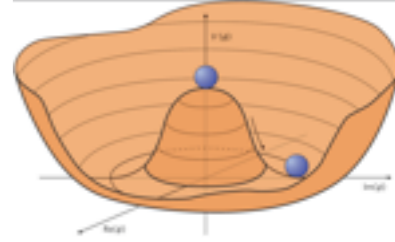


International FCC collaboration (CERN as host lab) to study:

- pp -collider (*FCC-hh*)
→ main emphasis, defining infrastructure requirements
~16 T \Rightarrow 100 TeV pp in 100 km
- ~100 km tunnel infrastructure in Geneva area, site specific
- e^+e^- collider (*FCC-ee*), as potential first step
- HE-LHC with *FCC-hh* technology
- $p-e$ (*FCC-he*) option, integration of one IP, e from ERL
- CDR for end 2018

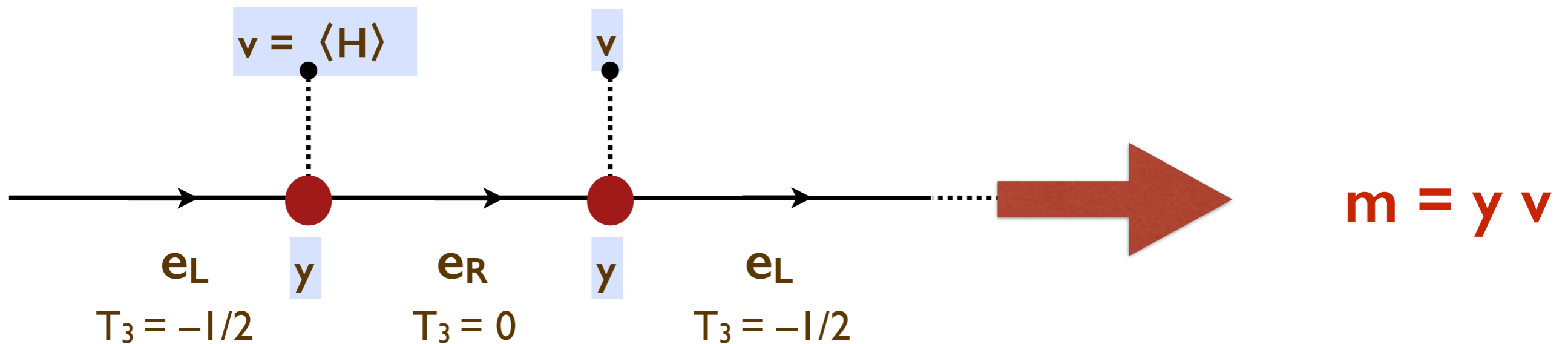
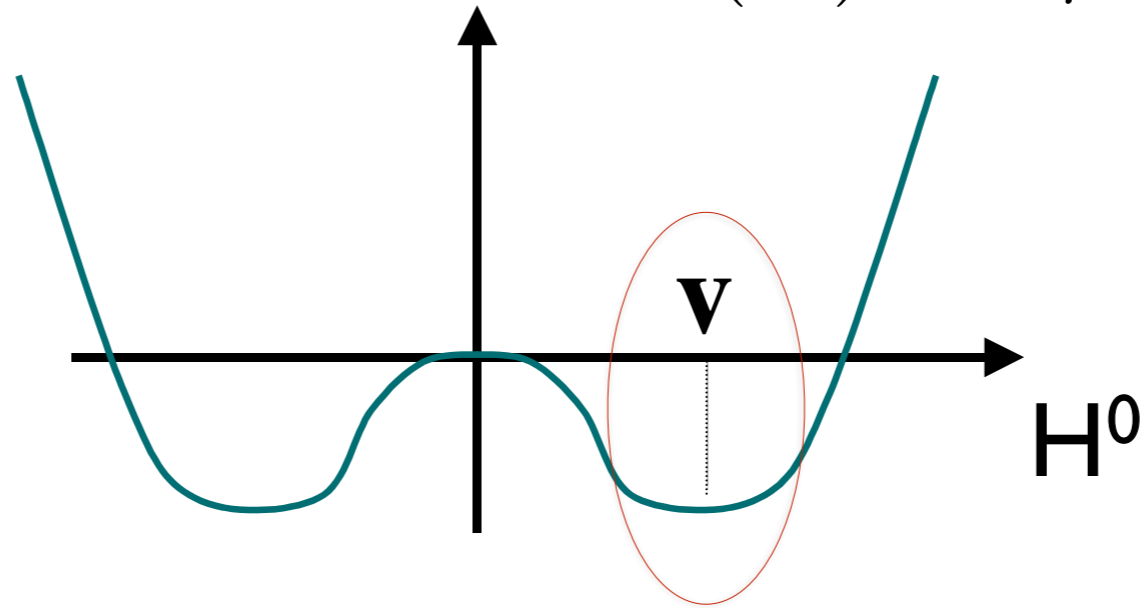


The Higgs mechanism in a nutshell ...



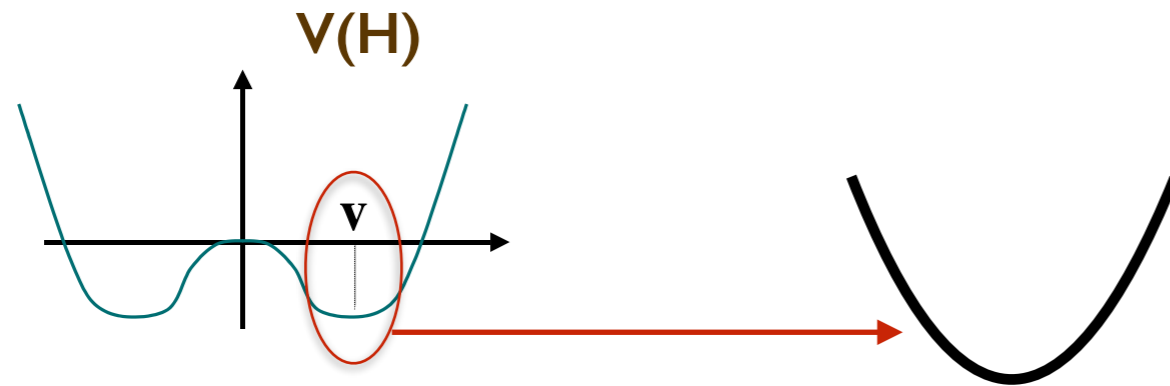
$$V_{SM}(H) = -\mu^2 |H|^2 + \lambda |H|^4$$

$$H = \begin{pmatrix} H^0 \\ H^- \end{pmatrix}$$



How far have we tested the Higgs mechanism?

parameters of the potential

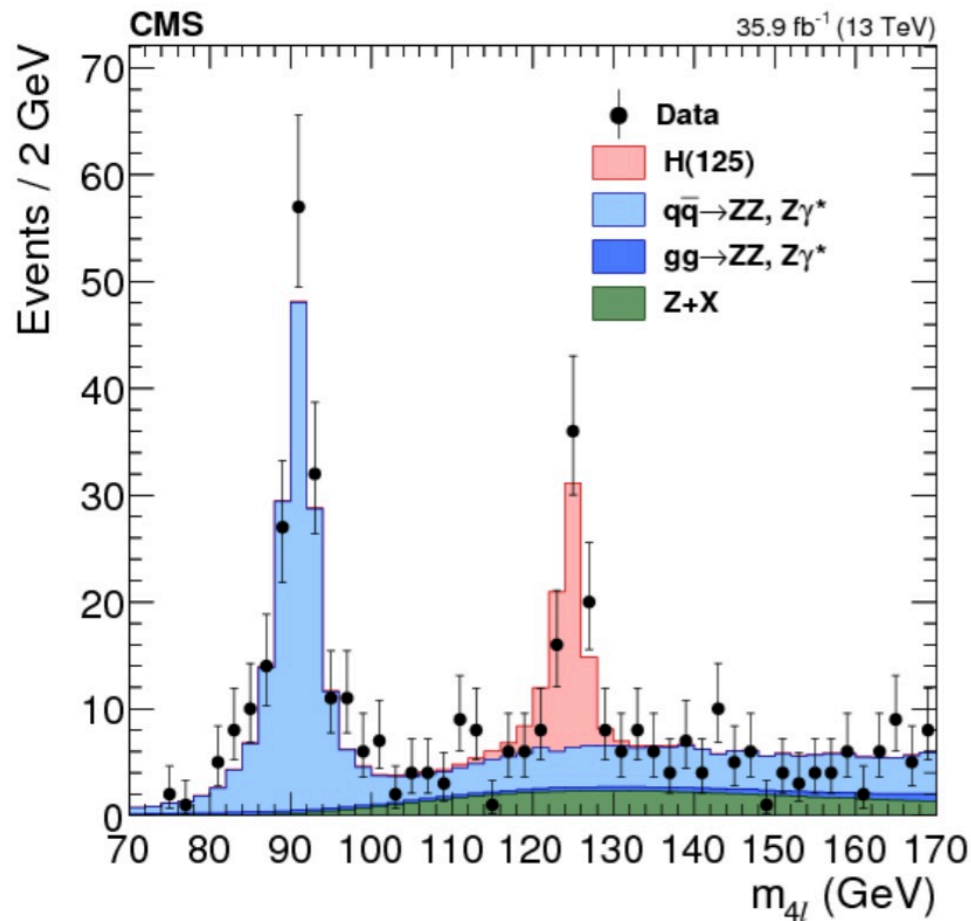


$$V(H) \sim m_H^2 (H - \mathbf{v})^2$$

$v=246$ GeV, from
weak decays

Higgs mass, 2017

CMS



[arXiv:1706.09936](https://arxiv.org/abs/1706.09936)

3D likelihood fit (m_{4l} , ZZ bg, δm) \Rightarrow

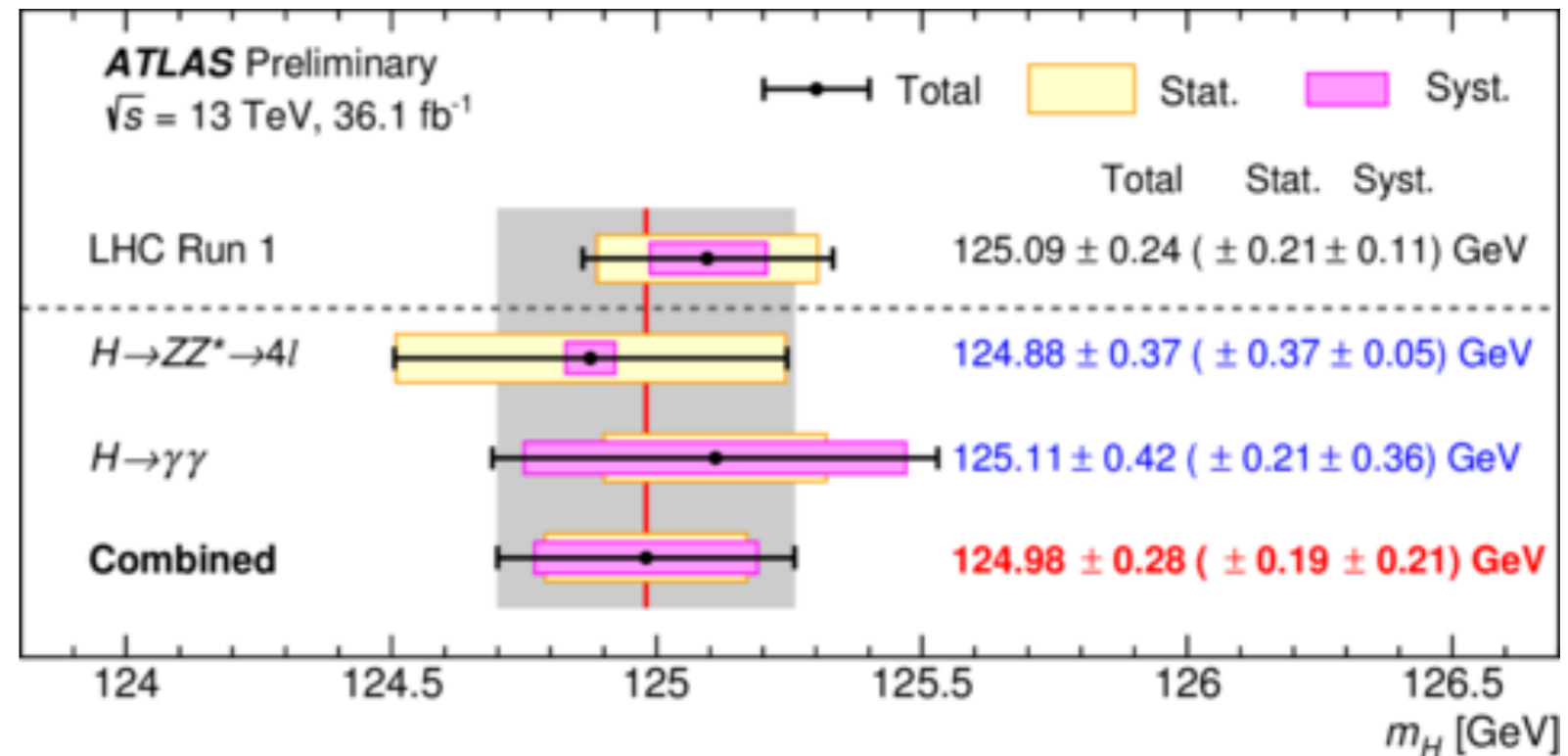
$$m_H = 125.26 \pm 0.20_{\text{stat}} \pm 0.08_{\text{syst}} \text{ GeV}$$

$$= 125.26 \pm 0.22 \text{ GeV}$$

$\Rightarrow 2 \times 10^{-3}$ precision

it took over 6 years from 1983 discovery to get below 5×10^{-3} on m_z (1989: CDF, SLC, LEP) 7

ATLAS



[ATLAS-CONF-2017-046](https://atlas.conf.cern.ch/2017/046)

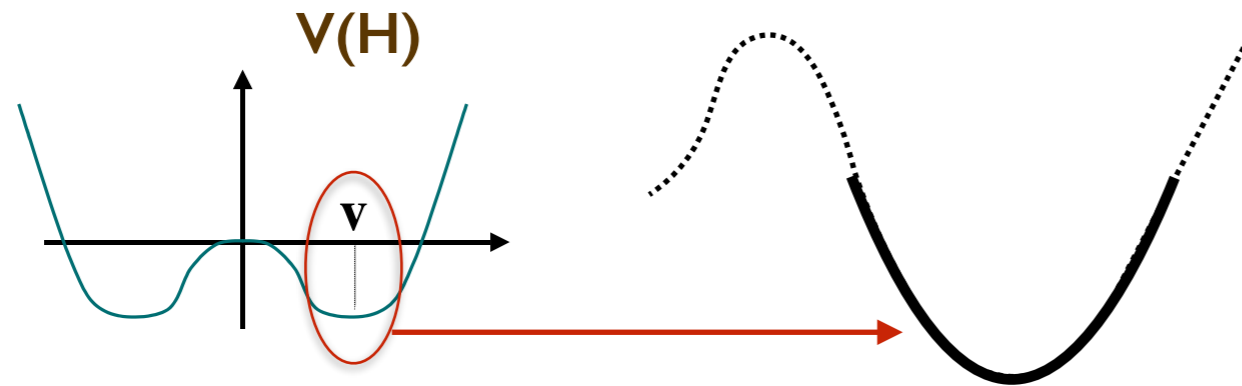
$\gamma\gamma$ and 4ℓ combination, run 1+2 \Rightarrow

$$m_H = 124.98 \pm 0.19_{\text{stat}} \pm 0.21_{\text{syst}} \text{ GeV}$$

$$= 124.98 \pm 0.26 \text{ GeV}$$

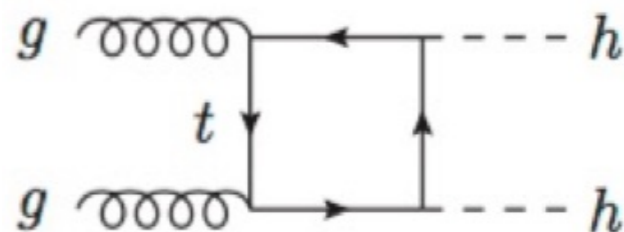
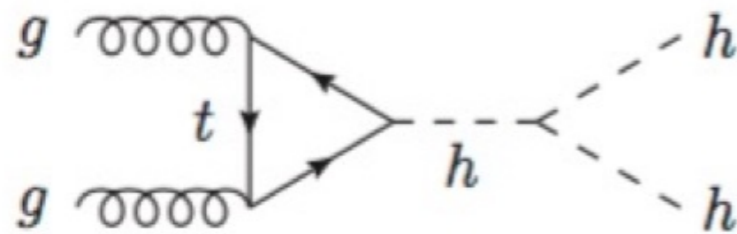
How far have we tested the Higgs mechanism?

parameters of the potential

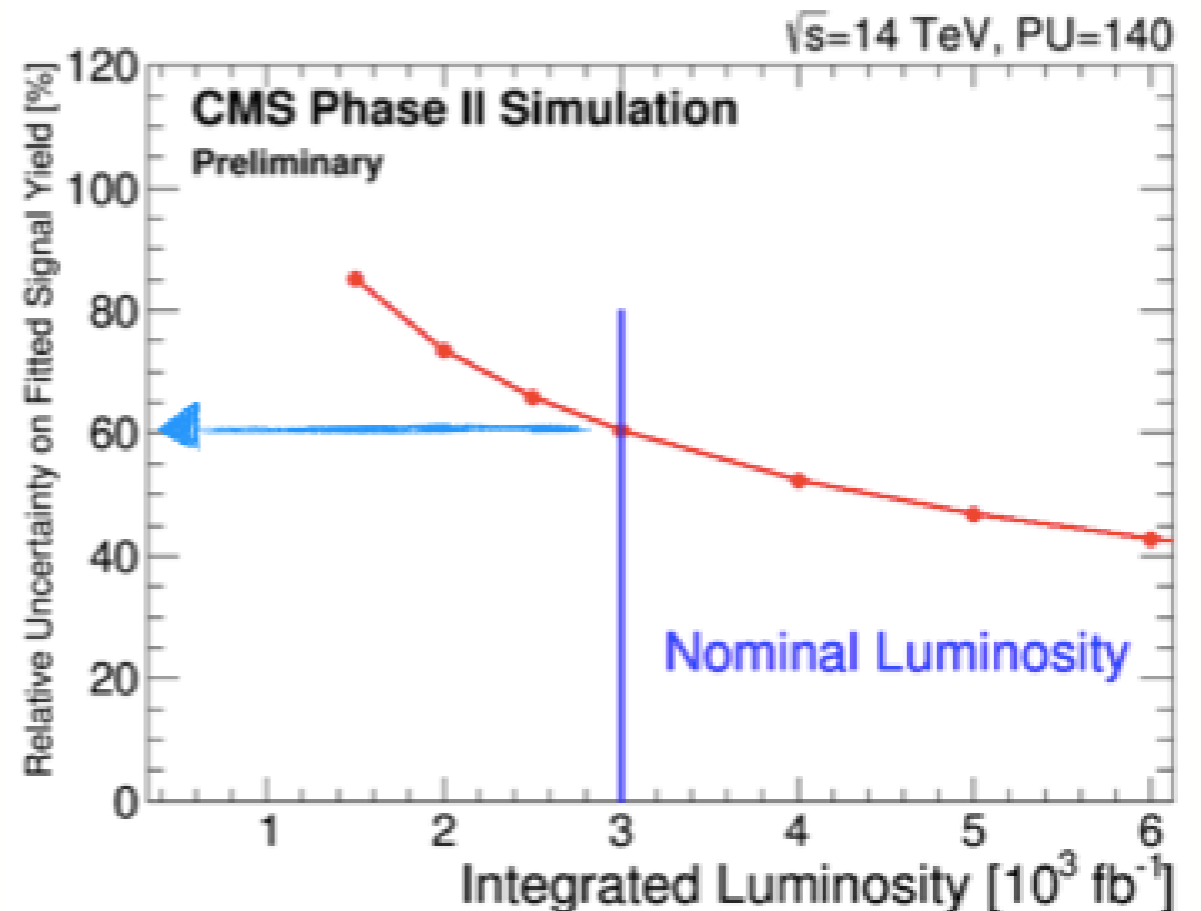


$$V(H) \sim m_H^2 (H-v)^2 + ???$$

Probing the cubic term of the Higgs potential will require at least 100x the current LHC statistics, and possibly more



Physics Performance for 2nd ECFA workshop

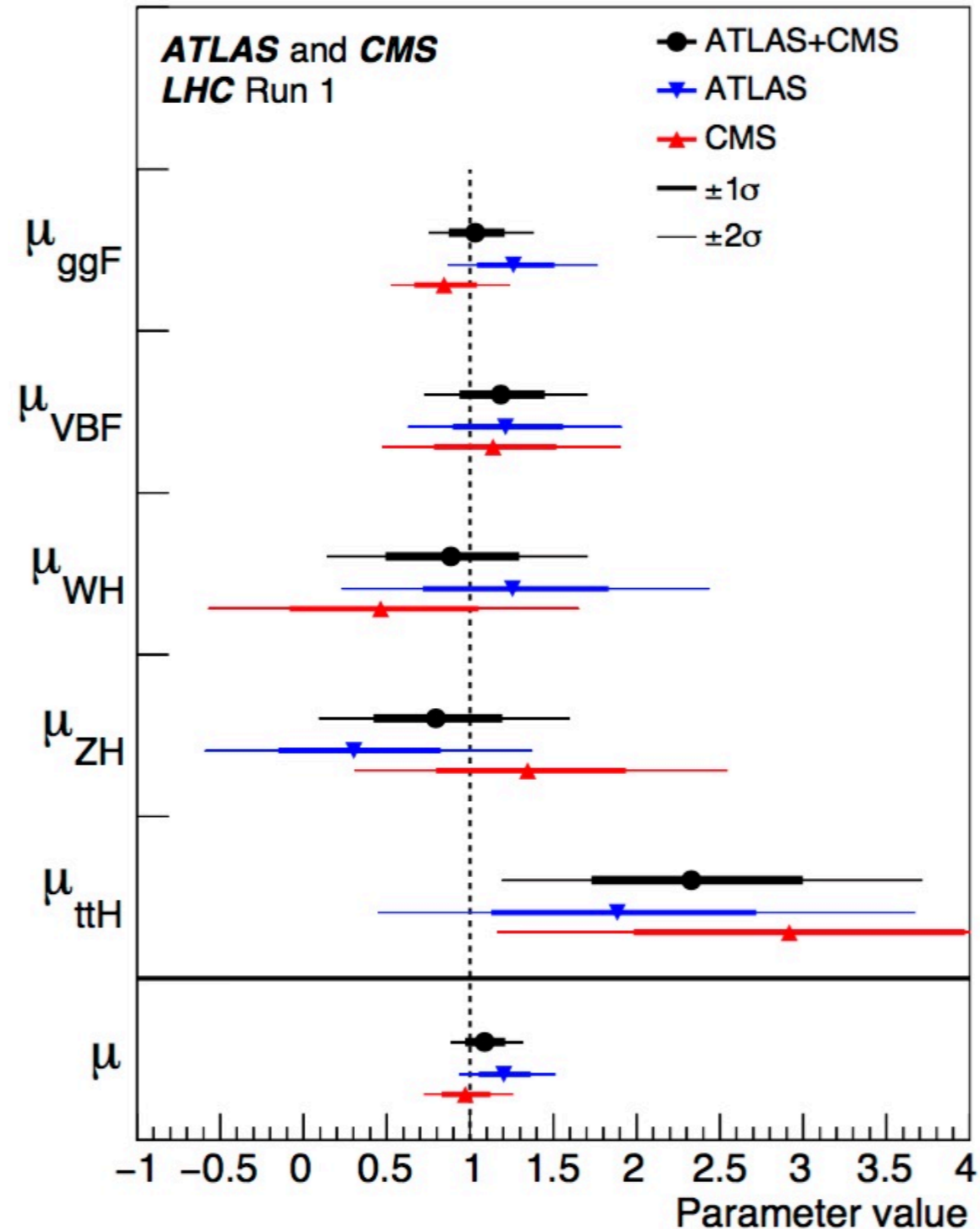


Higgs couplings: global fit of run I data

$\mu = \sigma \times \text{BR} / [\sigma \times \text{BR}]_{\text{SM}}$
assuming SM BR's in data

ATLAS+CMS
[JHEP 1608 \(2016\) 045](#)

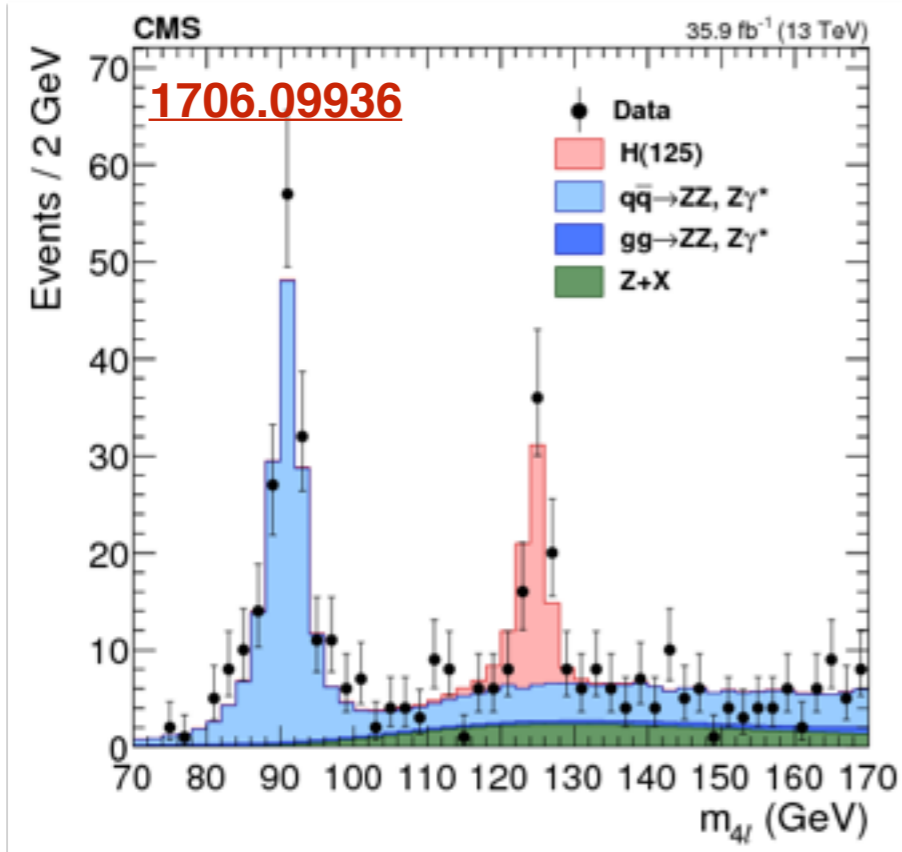
$\mu = 1.09 \pm 0.11$



- combination of different production and decay channels, explicit constraints on individual couplings are much less precise than 10% !!

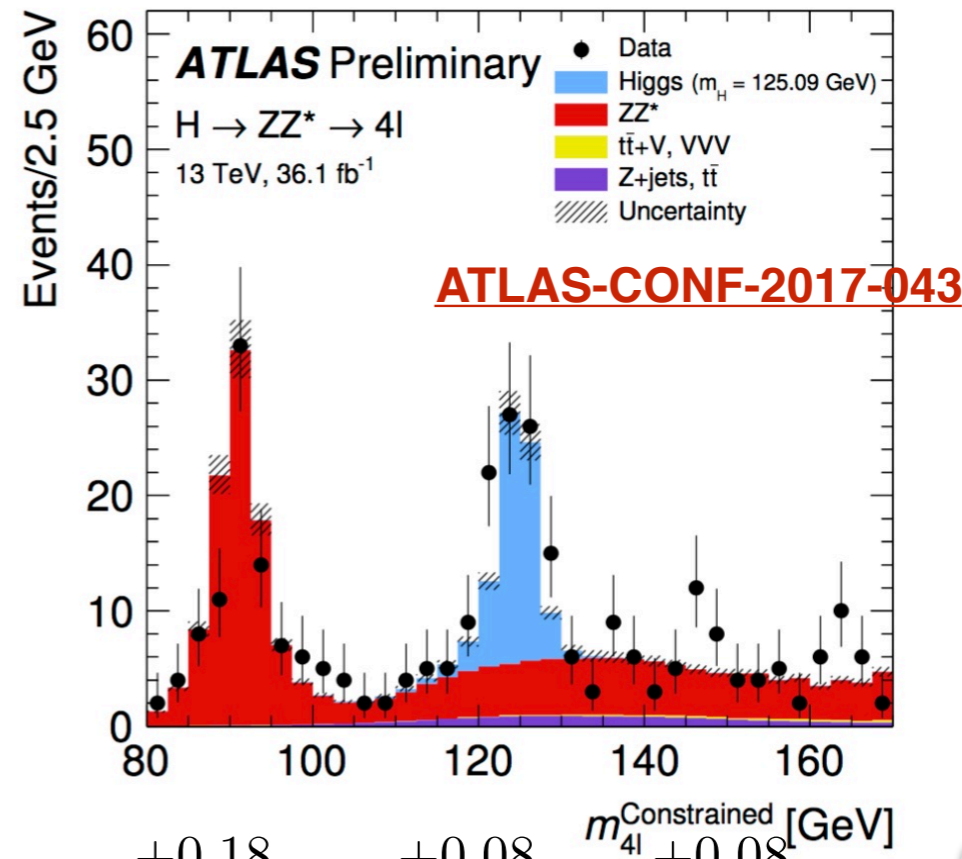
- essential to establish couplings individually, through combinations of different production and decay channels

$\mu = (\text{obs rate}) / (\text{SM rate}), 2017$

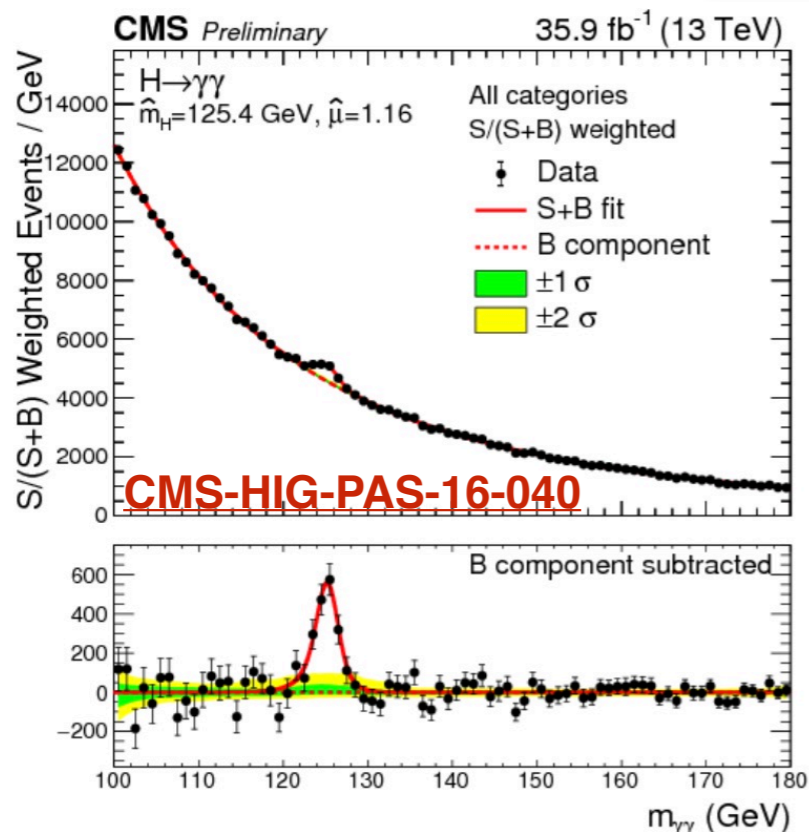


$pp \rightarrow H \rightarrow 4\ell$

$$\mu = 1.05^{+0.15}_{-0.14} (\text{stat})^{+0.11}_{-0.09} (\text{syst}) = 1.05^{+0.19}_{-0.17}$$

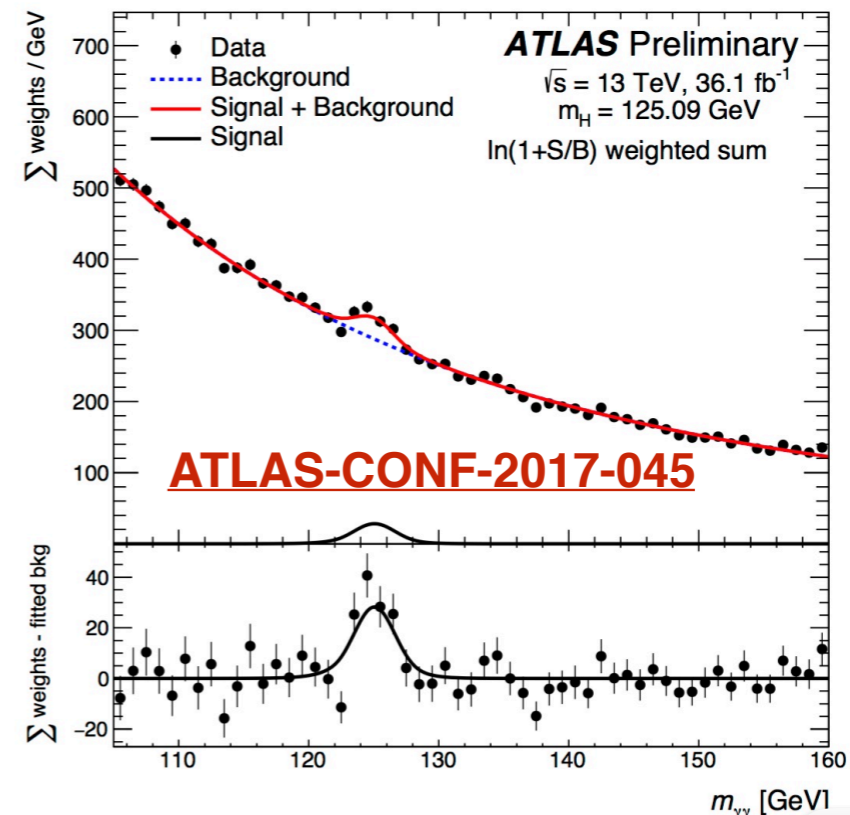


$$\mu = 1.28^{+0.18}_{-0.17} (\text{stat})^{+0.08}_{-0.06} (\text{exp})^{+0.08}_{-0.06} (TH) = 1.28^{+0.21}_{-0.19}$$



$pp \rightarrow H \rightarrow \gamma\gamma$

$$\mu = 1.16^{+0.11}_{-0.10} (\text{stat})^{+0.09}_{-0.08} (\text{exp})^{+0.06}_{-0.05} (TH) = 1.16^{+0.15}_{-0.14}$$



$$\mu = 0.99^{+0.12}_{-0.11} (\text{stat})^{+0.06}_{-0.05} (\text{exp})^{+0.06}_{-0.05} (TH) = 0.99 \pm 0.14$$

on the nature of EW symmetry breaking

- EW and strong interactions have free parameters (the symmetry groups, the strength of couplings, the charges of elementary particles). But at least we do have a deep understanding of their dynamical nature, namely the gauge principle. This allows us to speculate about an even deeper origin, e.g. from string theory or higher-dimensional Kaluza-Klein theories
- The Higgs mechanism relies on the quartic Higgs potential, in particular on the negative sign of its quadratic component. But we have no clue as to what is its dynamical origin, independently of whether we look at it with a SM or BSM perspective ...
- Understanding the origin of the Higgs potential and the nature of Higgs interactions is a paramount puzzle of modern physics, regardless of whether they eventually match the SM assumption or require new physics
- Having established the existence of the Higgs is similar to having established inflation, through cosmological observations. The real question (for both Higgs and inflation) is now **“where does it come from?”**

a historical example: superconductivity

- The relation between the Higgs phenomenon and the SM is similar to the relation between superconductivity and the Landau-Ginzburg theory of phase transitions: a quartic potential for a bosonic order parameter, with negative quadratic term, and the ensuing symmetry breaking. If superconductivity had been discovered after Landau-Ginzburg, we would be in a similar situations as we are in today: an experimentally proven phenomenological model. But we would still lack a deep understanding of the relevant dynamics.
- For superconductivity, this came later, with the identification of e^-e^- Cooper pairs as the underlying order parameter, and BCS theory. In particle physics, we still don't know whether the Higgs is built out of some sort of Cooper pairs (composite Higgs) or whether it is elementary, and in both cases we have no clue as to what is the dynamics that generates the Higgs potential. With Cooper pairs it turned out to be just EM and phonon interactions. With the Higgs, none of the SM interactions can do this, and **we must look beyond.**

The other *big* questions that press us to look *beyond* the Standard Model

- What's the real origin of EW symmetry breaking and particle's masses?
- What's the origin of Dark matter / energy ?
- What's the origin of matter/antimatter asymmetry in the universe?
- What's the origin of neutrino masses?
- What protects the smallness of $m_H / m_{\text{Plank,GUT}}$ (hierarchy problem)?
- ...

- The hierarchy problem, and the search for a *natural* explanation of the separation between the EW and Planck scales, provided so far an obvious setting for the exploration of the dynamics underlying the Higgs phenomenon.
- Lack of experimental evidence, so far, for a straightforward answer to naturalness (eg SUSY), forces us to review our biases, and to **take a closer look even at the most basic assumptions about Higgs properties**
- We often ask “is the Higgs like in SM?”The right way to set the issue is rather, more humbly, **“what is the Higgs?”** ...
 - in this perspective, even innocent questions like whether the Higgs gives mass also to 1st and 2nd generation fermions call for experimental verification.

Why do we care so much?

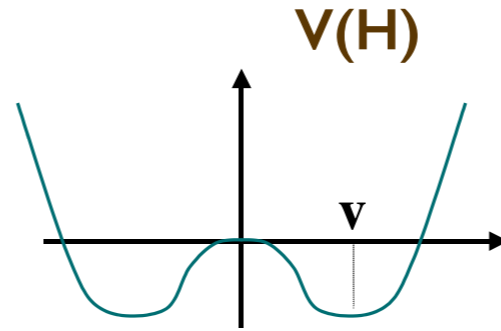
The Higgs boson is directly connected to several questions:

- Is the Higgs the only (fundamental?) scalar field, or are there other Higgs-like states (e.g. $H^\pm, A^0, H^{\pm\pm}, \dots$, EW-singlets,) ?
- What happens at the EW phase transition (PT) during the Big Bang?
 - what's the order of the phase transition?
 - are the conditions realized to allow EW baryogenesis?
 - does the PT wash out possible pre-existing baryon asymmetry?
- Is there a relation between any amongst Higgs/EWSB, baryogenesis, Dark Matter, inflation?
- Is there a deep reason for the apparent metastability of the Higgs vacuum?
- The hierarchy problem: what protects the smallness of $m_H / m_{\text{Planck, GUT, ...}}$?

Higgs selfcouplings

The Higgs sector is defined in the SM by two parameters, μ and λ :

$$V_{SM}(H) = -\mu^2 |H|^2 + \lambda |H|^4$$



$$\frac{\partial V_{SM}(H)}{\partial H} \Big|_{H=v} = 0 \quad \text{and} \quad m_H^2 = \frac{\partial^2 V_{SM}(H)}{\partial H \partial H^*} \Big|_{H=v} \Rightarrow$$

$$\begin{aligned} \mu &= m_H \\ \lambda &= \frac{m_H^2}{2v^2} \end{aligned}$$

These relations uniquely determine the strength of Higgs selfcouplings in terms of the two now-known parameters m_H and v

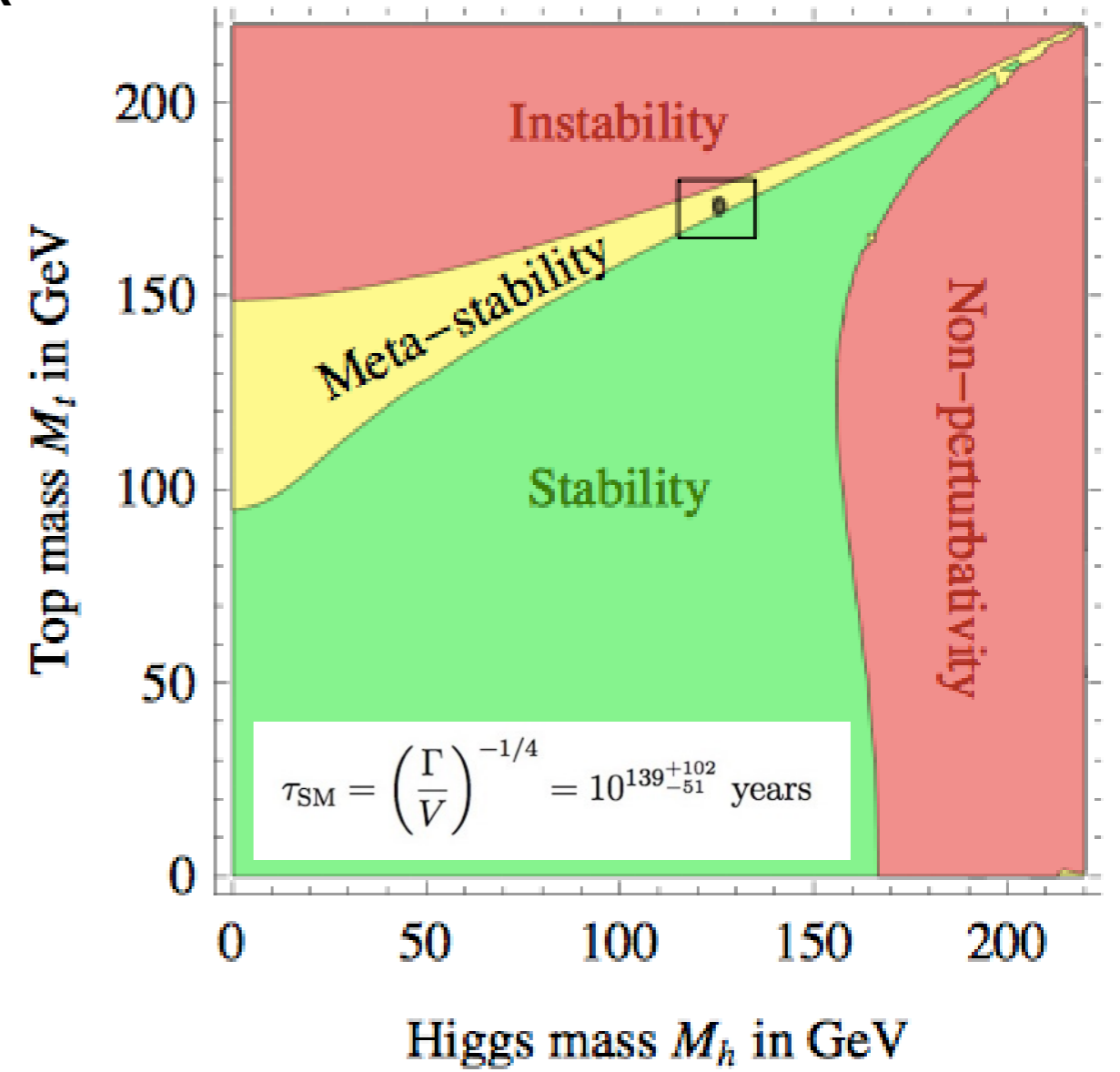
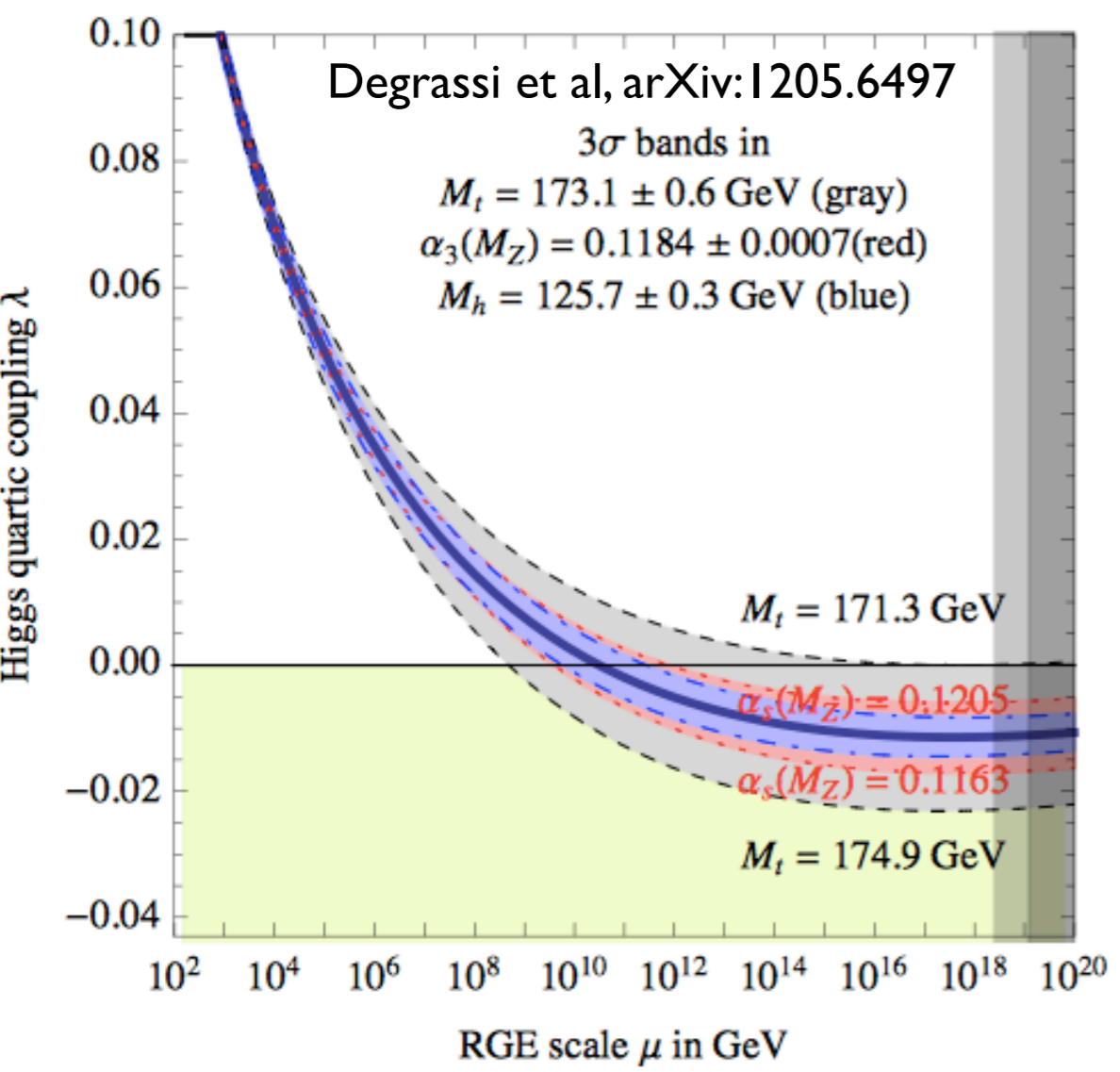
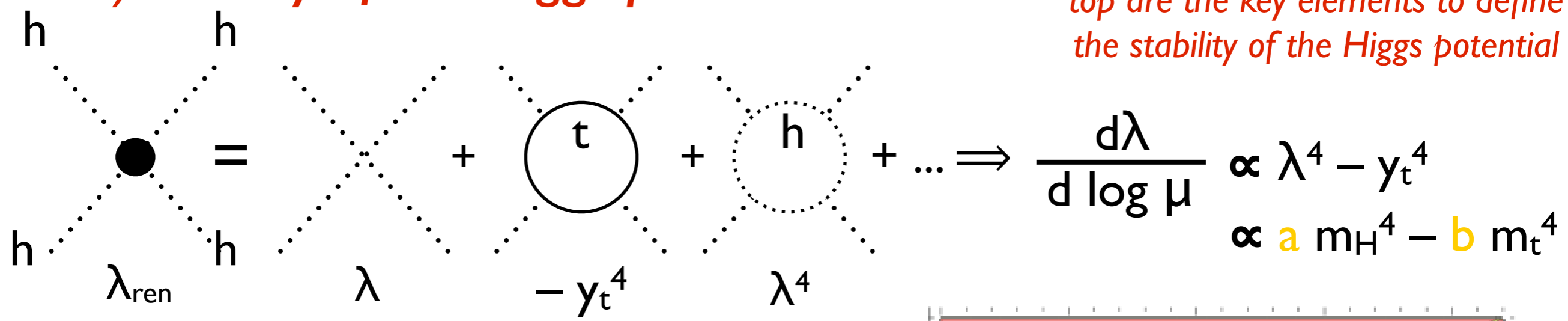
$$g_{3H} \Rightarrow 4\lambda v = \frac{2m_H^2}{v}$$

$$g_{4H} \Rightarrow \lambda = \frac{m_H^2}{2v^2}$$

These relations between Higgs self-couplings, m_H and v entirely depend on the functional form of the Higgs potential. Their measurement is therefore an important test of the SM nature of the Higgs mechanism

(meta)Stability of the Higgs potential

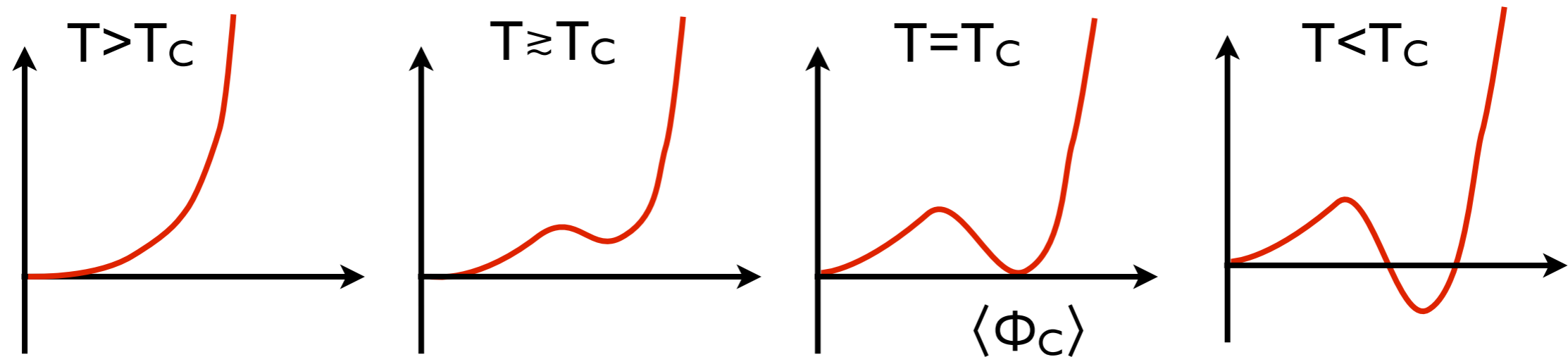
Higgs selfcoupling and coupling to the top are the key elements to define the stability of the Higgs potential



Not an issue of concern for the human race.... but the closeness of m_{top} to the critical value where the Higgs selfcoupling becomes 0 at M_{Planck} (namely 171.3 GeV) might be telling us something fundamental about the origin of EWSB ... incidentally, $y_{top}=1$ (!?)

The nature of the EW phase transition

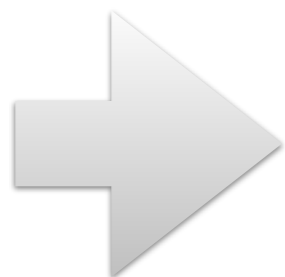
Strong 1st order phase transition is required to induce and sustain the out of equilibrium generation of a baryon asymmetry during EW symmetry breaking



Strong 1st order phase transition $\Rightarrow \langle \Phi_C \rangle > T_c$

In the SM this requires $m_H \approx 80$ GeV.

Since $m_H = 125$ GeV, **new physics**, coupling to the Higgs and effective at **scales $O(\text{TeV})$** , must modify the Higgs potential to make this possible



- Probe higher-order terms of the Higgs potential (selfcouplings)
- Probe the existence of other particles coupled to the Higgs

The basic motivation for Future Circular Colliders

- HEP has two priorities:
 - explore the physics of electroweak symmetry breaking:
 - experimentally, via the measurement of Higgs properties, Higgs interactions and selfinteractions, couplings of gauge bosons, flavour phenomena, etc
 - theoretically, to understand the nature of the hierarchy problem and identify possible natural solutions (to be subjected to exptl test)
 - explore the origin of known departures from the SM (DM, neutrino masses, baryon asymmetry of the universe)

The physics case of FCCs builds on the belief that these two directions are deeply intertwined, and equally worth investigating

Key question for the future developments of HEP:
Why don't we see the new physics we expected to be present around the TeV scale ?

- **Is the mass scale beyond the LHC reach ?**
- **Is the mass scale within LHC's reach, but final states are elusive to the direct search ?**

These two scenarios are a priori equally likely, but they impact in different ways the future of HEP, and thus the assessment of the physics potential of possible future facilities

Readiness to address both scenarios is the best hedge for the field:

- *precision*
- *sensitivity (to elusive signatures)*
- *extended energy/mass reach*

The potential of a Future Circular Collider

- Guaranteed deliverables:
 - study of Higgs and top quark properties, and exploration of EWSB phenomena, with unmatched **precision and sensitivity**
- Exploration potential:
 - **mass reach enhanced** by factor $\sim E / 14 \text{ TeV}$ (will be 5–7 at 100 TeV, depending on integrated luminosity)
 - *statistics enhanced by several orders of magnitude for BSM phenomena brought to light by the LHC*
 - benefit from both direct (large Q^2) and indirect (precision) probes
- Provide firm Yes/No answers to questions like:
 - is the SM dynamics all there is at the TeV scale?
 - is there a TeV-scale solution to the hierarchy problem?
 - is DM a thermal WIMP?
 - did baryogenesis take place during the EW phase transition?

Higgs physics

SM Higgs rates at 100 TeV

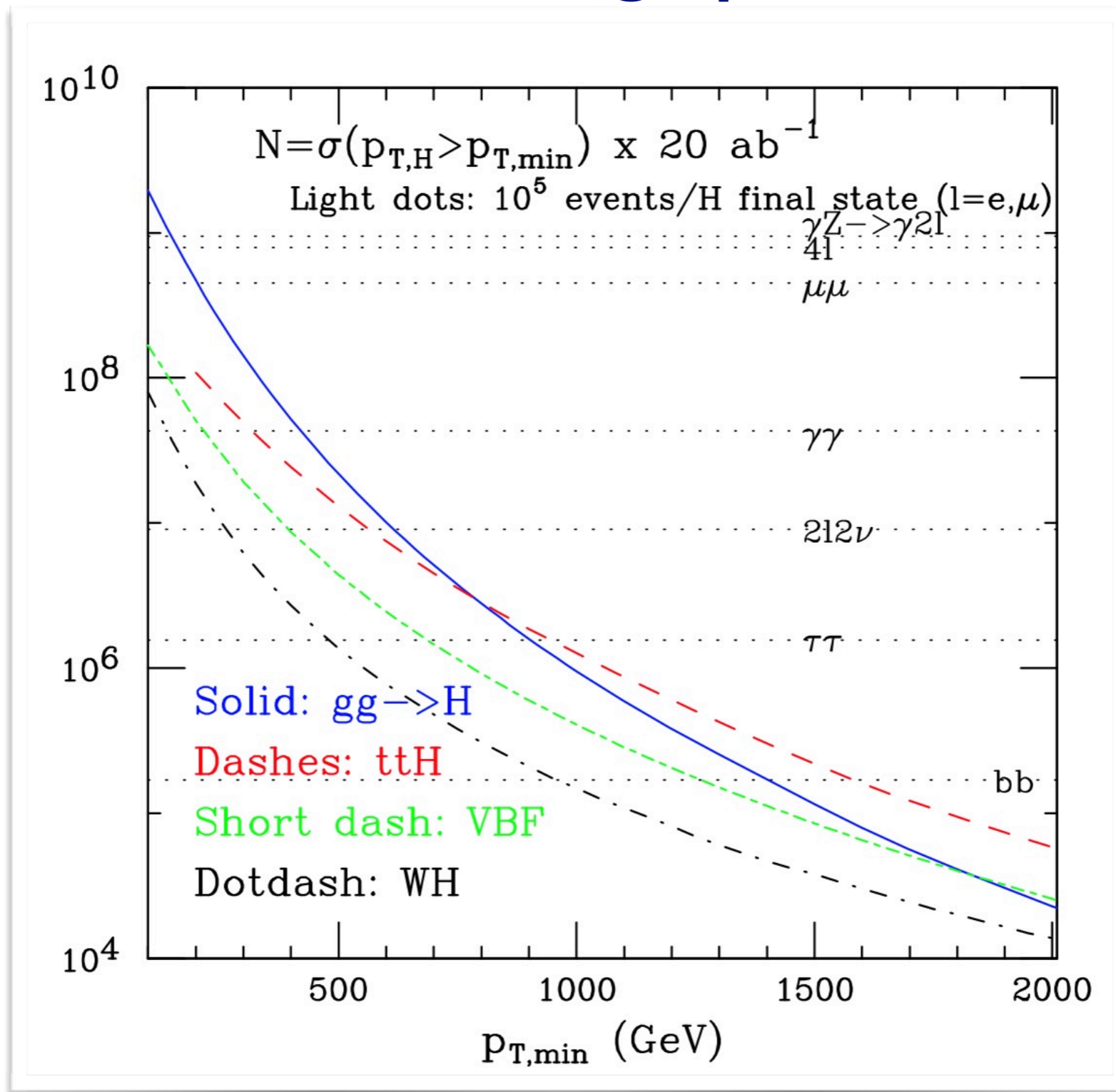
	N_{100}	N_{100}/N_8	N_{100}/N_{14}
$gg \rightarrow H$	16×10^9	4×10^4	110
VBF	1.6×10^9	5×10^4	120
WH	3.2×10^8	2×10^4	65
ZH	2.2×10^8	3×10^4	85
$t\bar{t}H$	7.6×10^8	3×10^5	420

$$N_{100} = \sigma_{100\text{TeV}} \times 20 \text{ ab}^{-1}$$

$$N_8 = \sigma_{8\text{TeV}} \times 20 \text{ fb}^{-1}$$

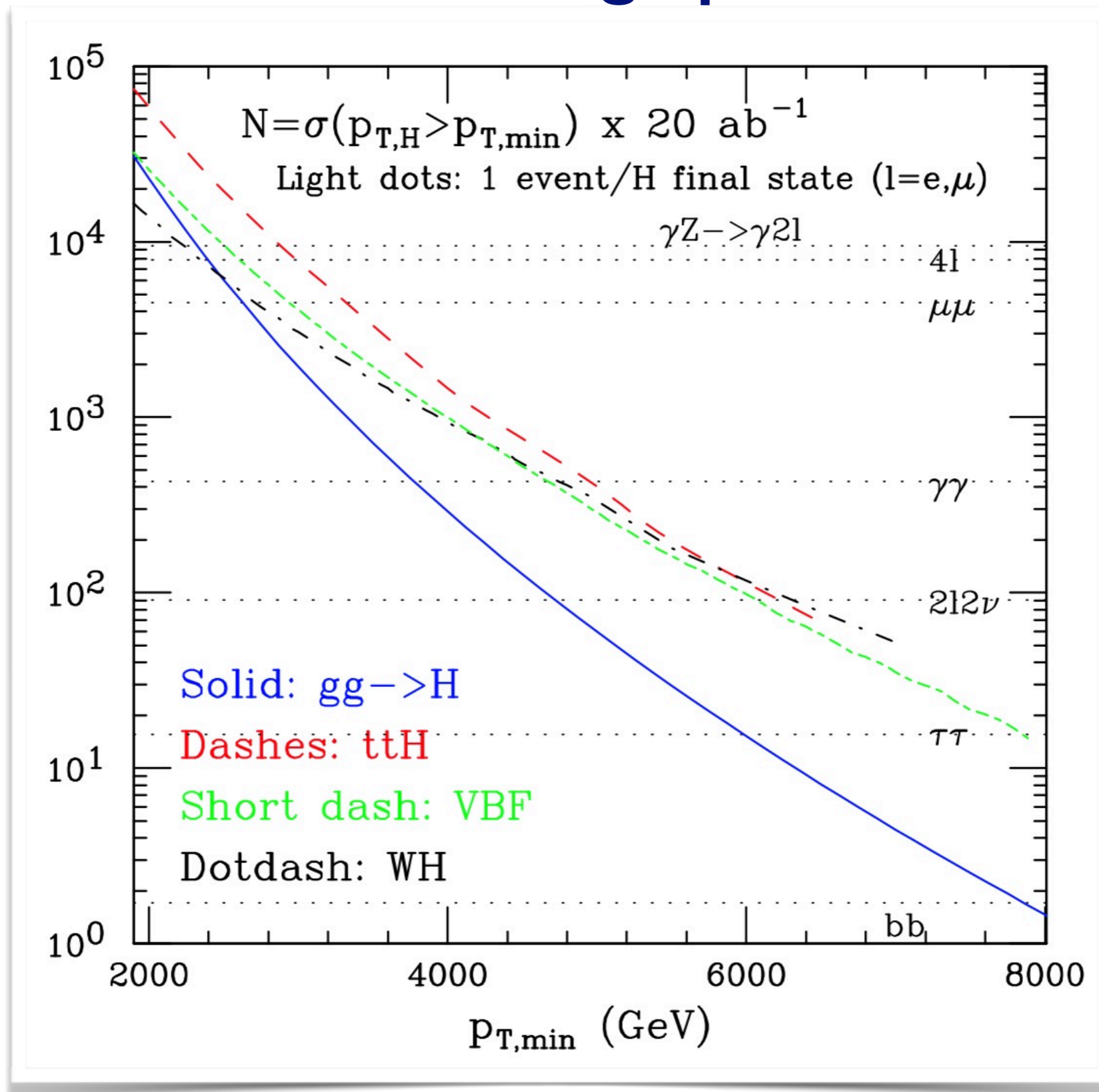
$$N_{14} = \sigma_{14\text{TeV}} \times 3 \text{ ab}^{-1}$$

H at large p_T



- Hierarchy of production channels changes at large $p_T(H)$:
 - $\sigma(ttH) > \sigma(gg \rightarrow H)$ above 800 GeV
 - $\sigma(VBF) > \sigma(gg \rightarrow H)$ above 1800 GeV

H at large p_T



- Statistics in potentially visible final states out to several TeV

Remarks

- Higher statistics shifts the balance between systematic and statistical uncertainties. It can be exploited to define different signal regions, with better S/B, better systematics, pushing the potential for better measurements beyond the “systematics wall” of low-stat measurements.
- We often talk about “**precise**” Higgs measurements. What we actually aim at, is “**sensitive**” tests of the Higgs properties, where *sensitive* refers to the ability to reveal BSM behaviours.
- ***Sensitivity*** may not require extreme precision
 - Going after “sensitivity”, rather than *just* precision, opens itself new opportunities ...

Higgs as a BSM probe: precision vs dynamic reach

$$L = L_{SM} + \frac{1}{\Lambda^2} \sum_k \mathcal{O}_k + \dots$$

$$O = | \langle f | L | i \rangle |^2 = O_{SM} [1 + O(\mu^2 / \Lambda^2) + \dots]$$

For H decays, or inclusive production, $\mu \sim O(v, m_H)$

$$\delta O \sim \left(\frac{v}{\Lambda}\right)^2 \sim 6\% \left(\frac{\text{TeV}}{\Lambda}\right)^2 \Rightarrow \text{precision probes large } \Lambda$$

$$\text{e.g. } \delta O = 1\% \Rightarrow \Lambda \sim 2.5 \text{ TeV}$$

For H production off-shell or with large momentum transfer Q , $\mu \sim O(Q)$

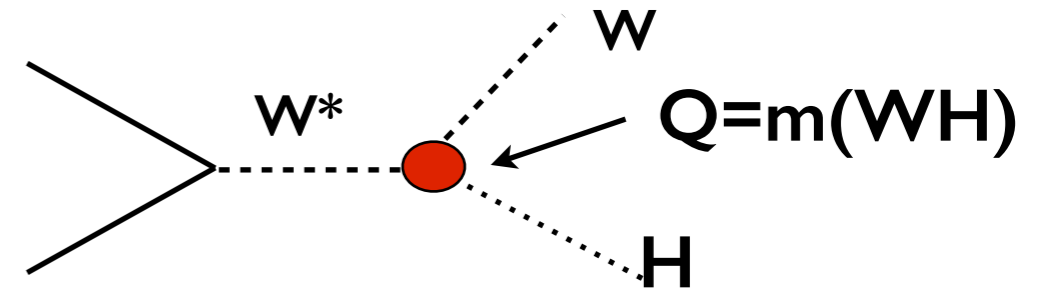
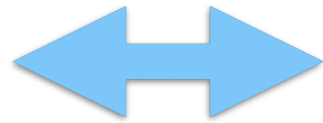
$$\delta O \sim \left(\frac{Q}{\Lambda}\right)^2 \Rightarrow \text{kinematic reach probes large } \Lambda$$

even if precision is low

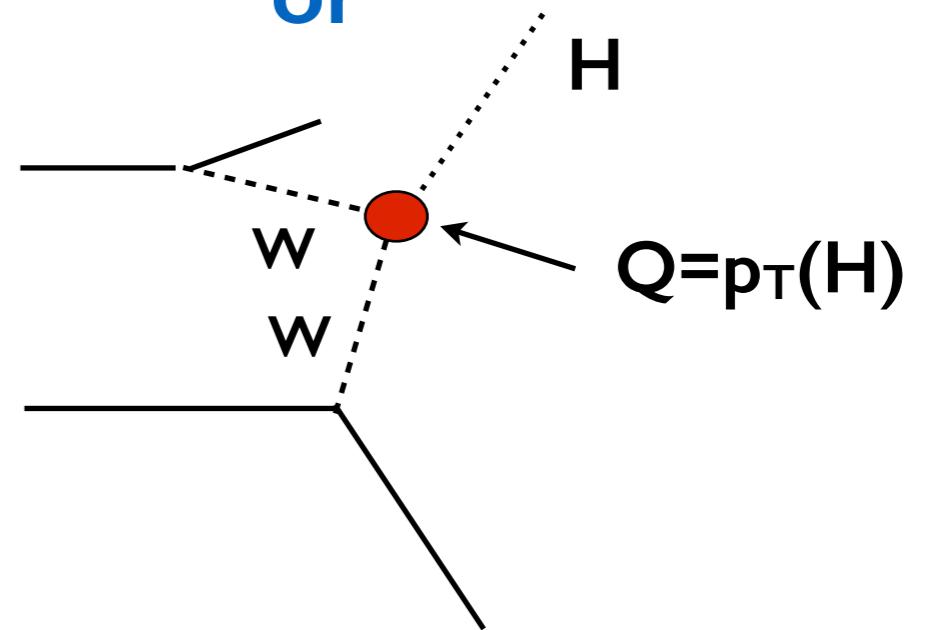
$$\text{e.g. } \delta O = 15\% \text{ at } Q = 1 \text{ TeV} \Rightarrow \Lambda \sim 2.5 \text{ TeV}$$

Examples

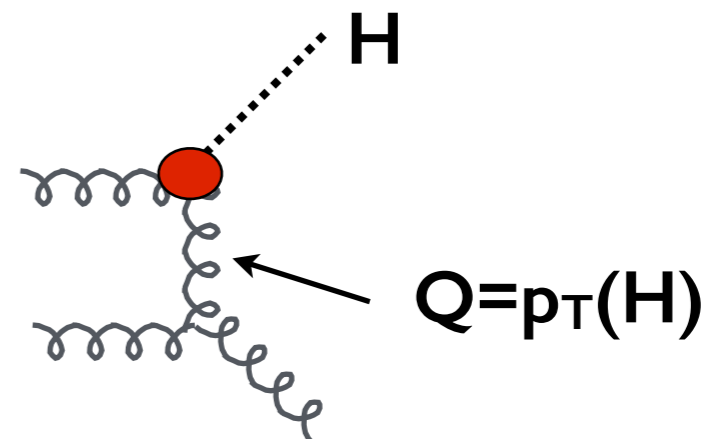
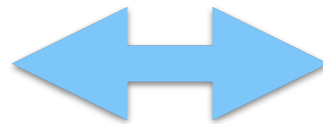
$\delta\text{BR}(H \rightarrow WW^*)$



or



$\delta\text{BR}(H \rightarrow gg)$

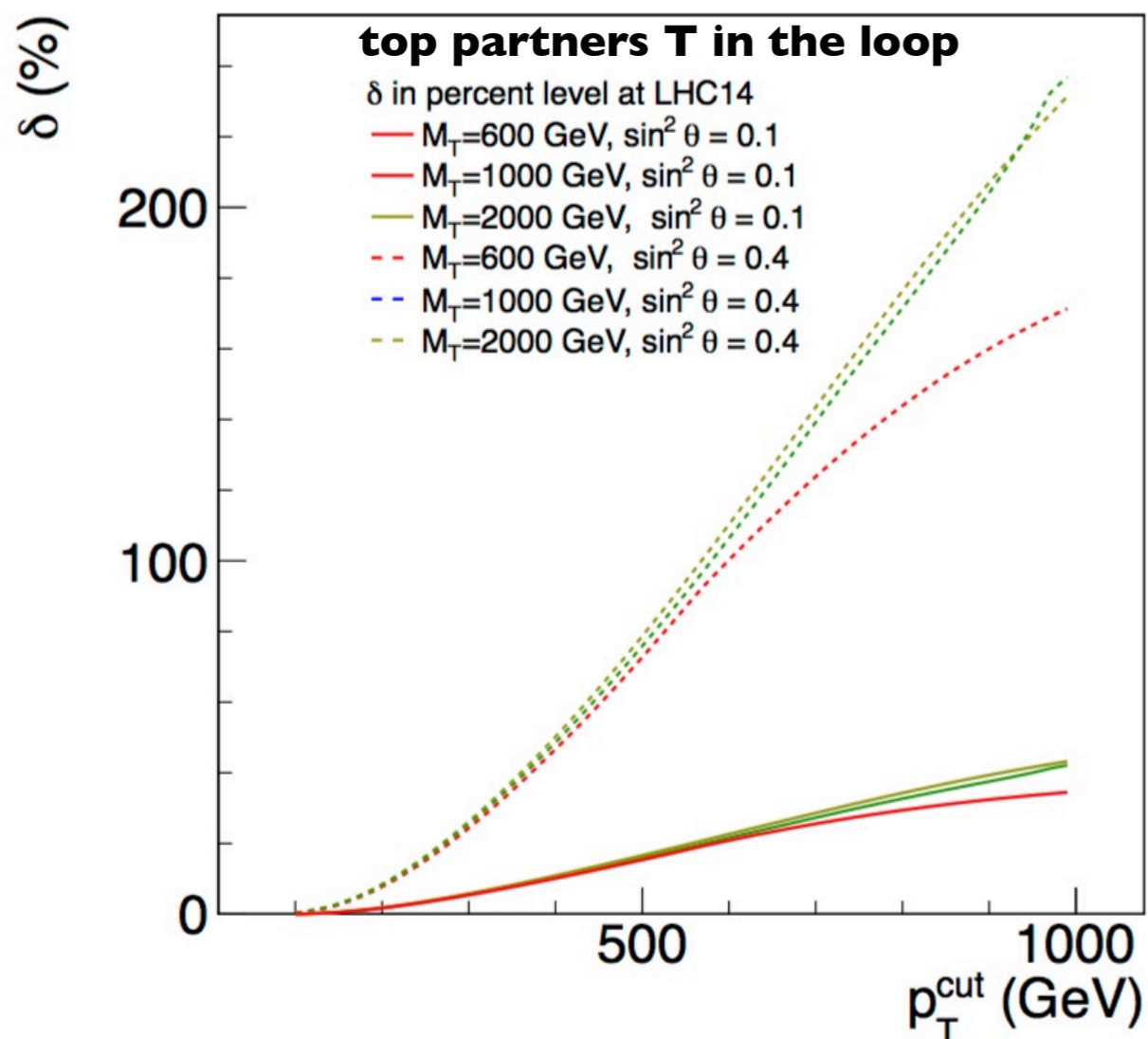


Examples of deviations of the Higgs p_T spectrum from SM, in presence of new particles in the ggH loop

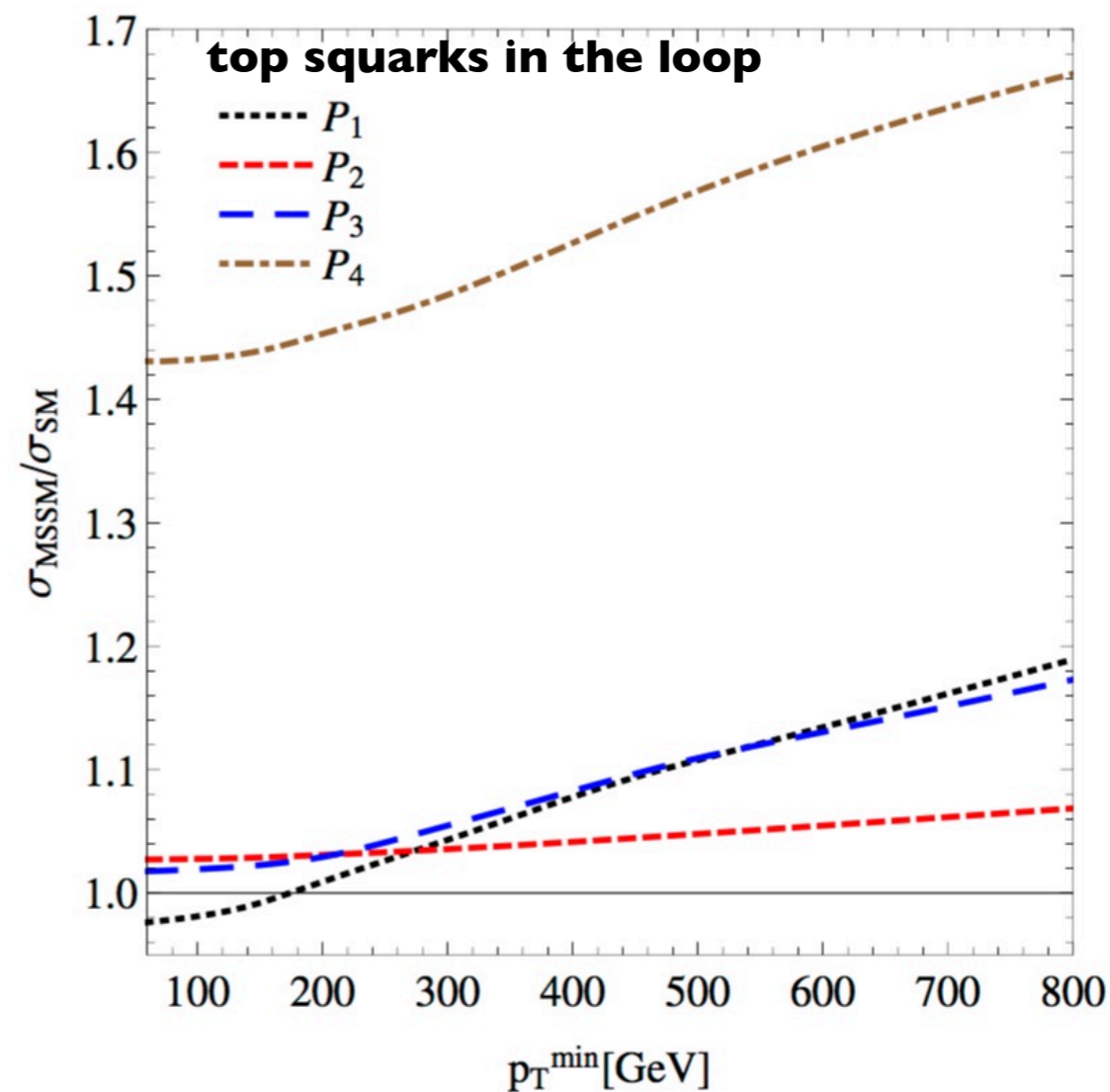
(See also
Azatov and Paul [arXiv:1309.5273v3](https://arxiv.org/abs/1309.5273v3))

Table 3: The benchmark points shown in Fig. 7. We set $\tan\beta = 10$, $M_{A^0} = 500$ GeV, $M_2 = 1000$ GeV, $\mu = 200$ GeV and all trilinear couplings to a common value A_t . The remaining sfermion masses were set to 1 TeV and the mass of the lightest CP -even Higgs was set to 125 GeV.

Point	$m_{\tilde{t}_1}$ [GeV]	$m_{\tilde{t}_2}$ [GeV]	A_t [GeV]	Δ_t
P_1	171	440	490	0.0026
P_2	192	1224	1220	0.013
P_3	226	484	532	0.015
P_4	226	484	0	0.18

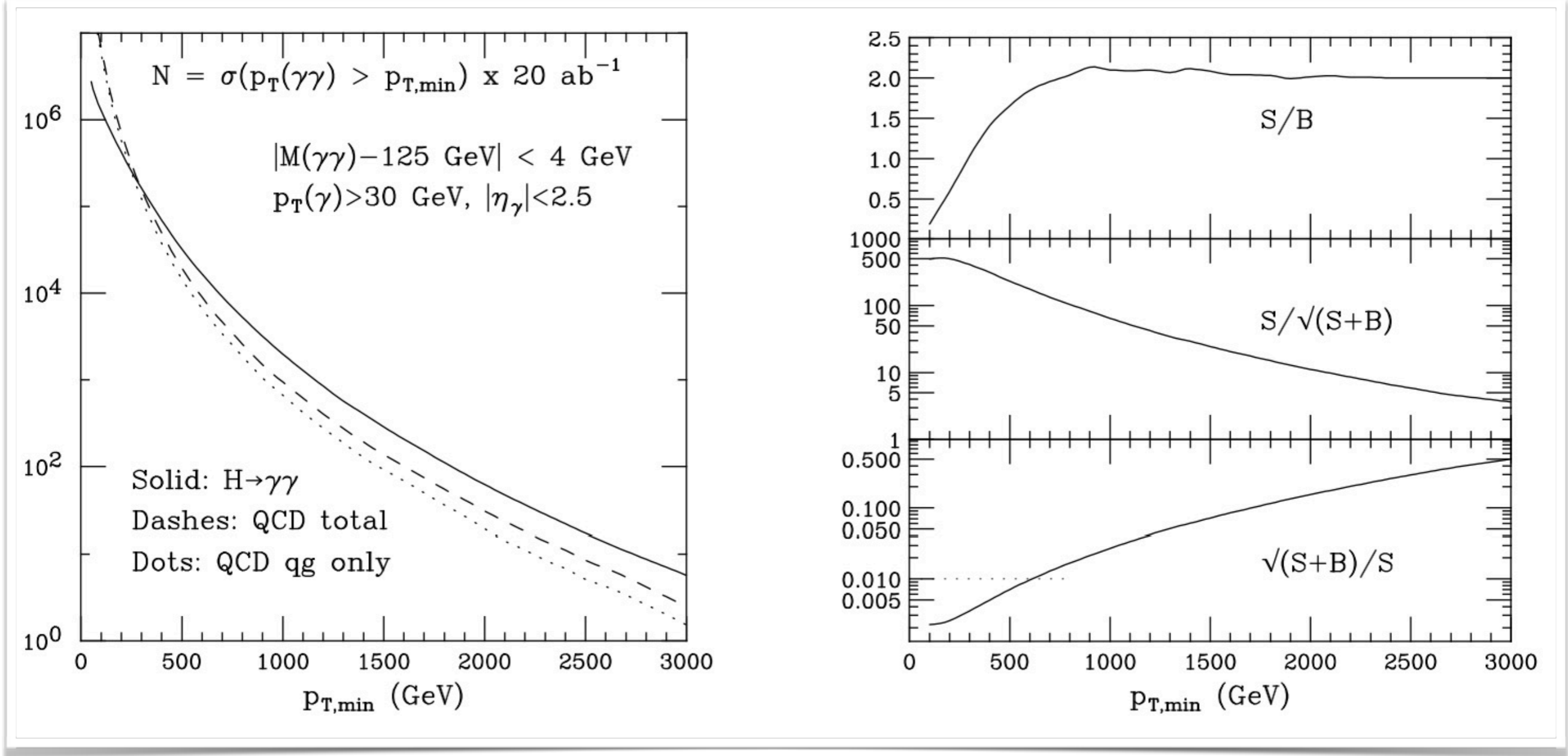


Banfi Martin Sanz, [arXiv:1308.4771](https://arxiv.org/abs/1308.4771)



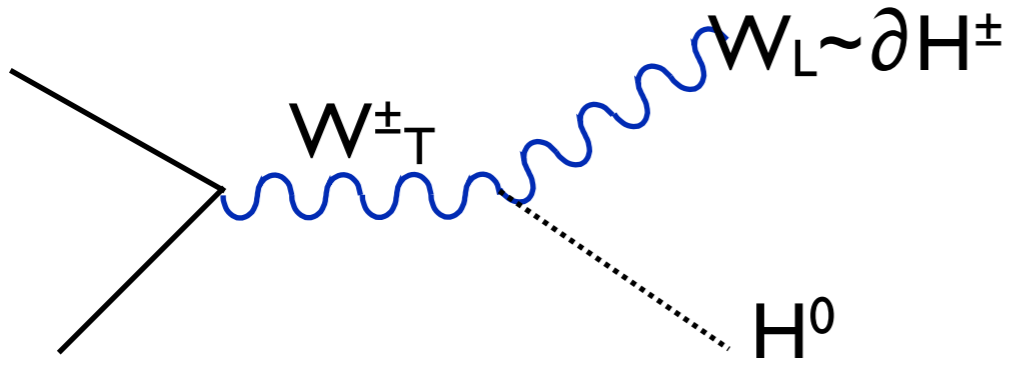
Grojean, Salvioni, Schläffer, Weiler [arXiv:1312.3317](https://arxiv.org/abs/1312.3317)

$gg \rightarrow H \rightarrow \gamma\gamma$ at large p_T at 100 TeV



- At 1 TeV, statistical sensitivity (accounting for bg) well below 10% !!
- What is a best BSM probe: $BR(\gamma\gamma)$ or shape of $p_T(H)$?
 - answer likely BSM-model dependent
 - \Rightarrow synergy/complementarity !!

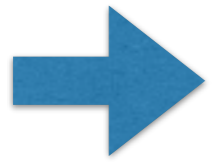
VH production at large $m(VH)$



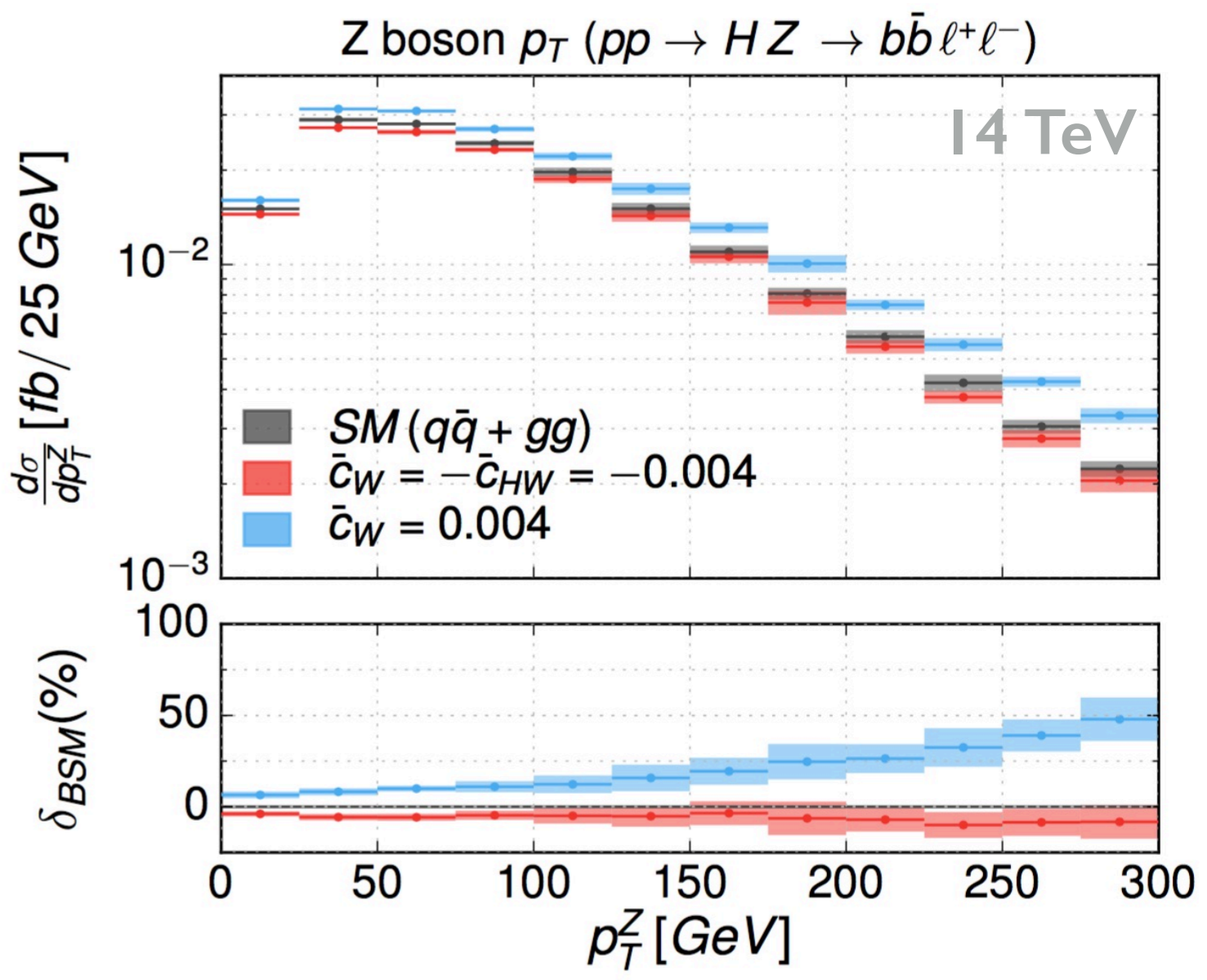
See e.g.
 Biekötter, Knochel, Krämer, Liu, Riva,
 arXiv:1406.7320

In presence of a higher-dim op
 such as:

$$L_{D=6} = \frac{ig}{2} \frac{c_W}{\Lambda^2} (H^\dagger \sigma^a D^\mu H) D^\nu V_{\mu\nu}^a$$



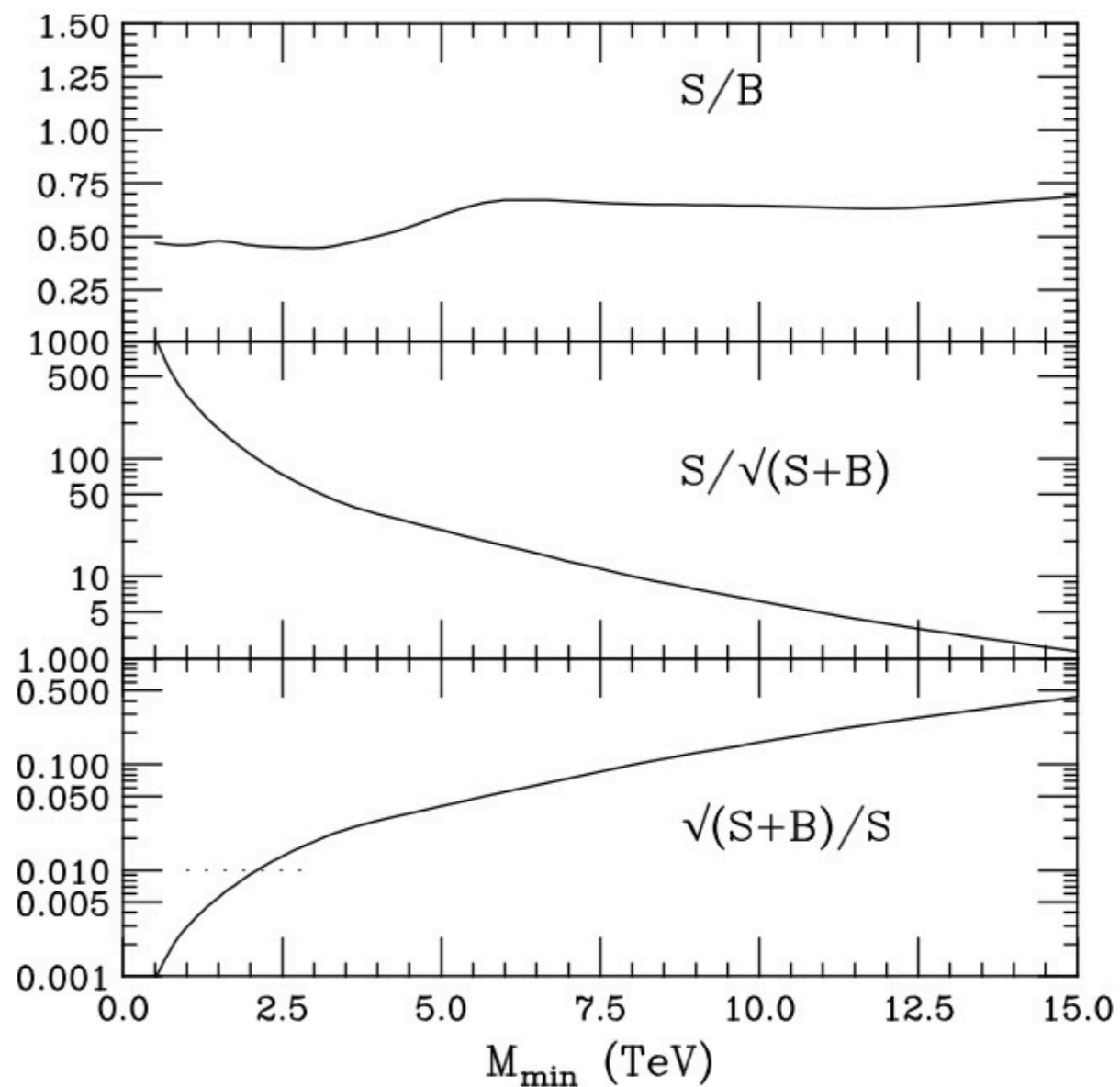
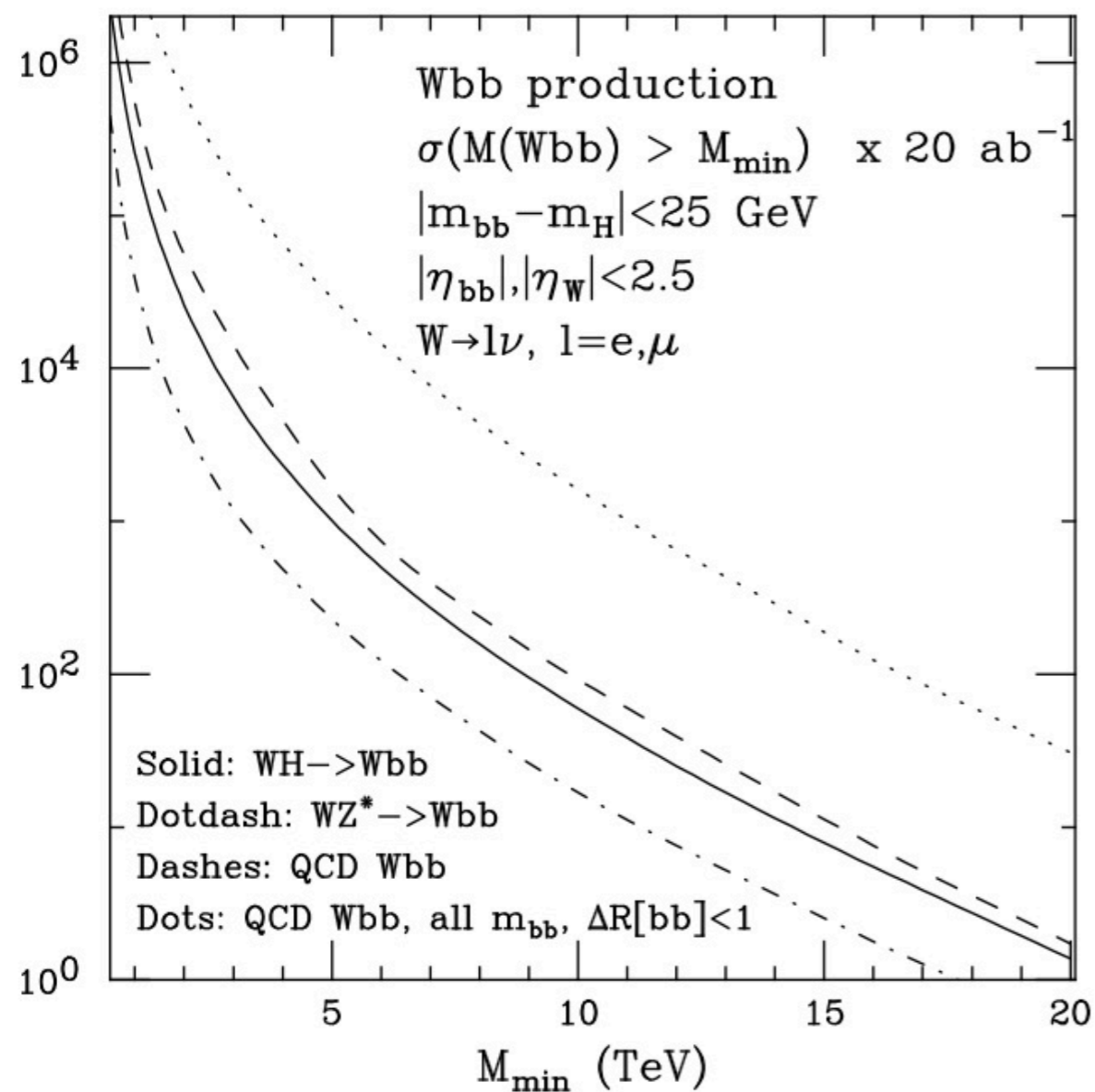
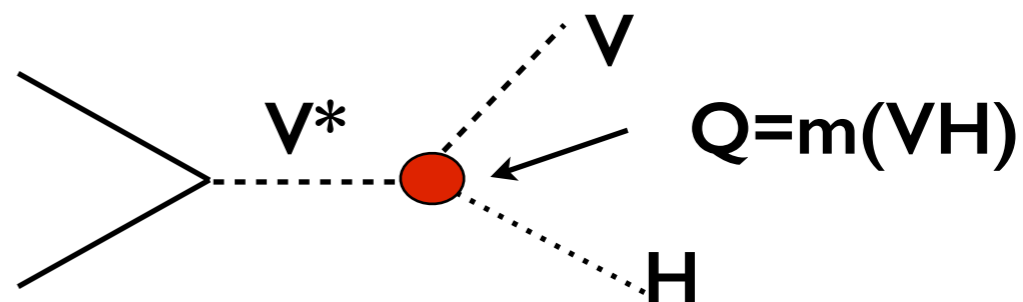
$$\frac{\sigma}{\sigma_{SM}} \sim \left(1 + c_W \frac{\hat{s}}{\Lambda^2} \right)^2$$



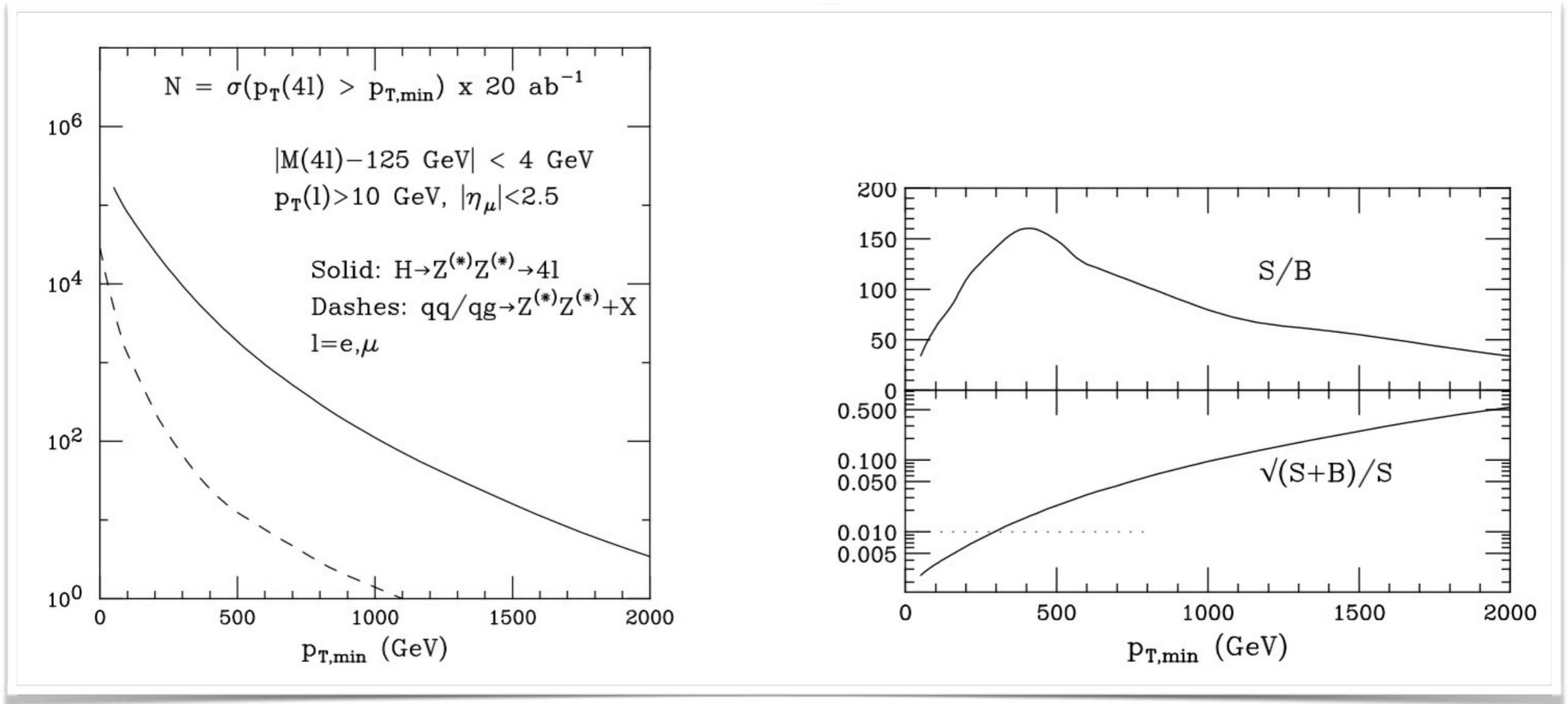
Mimasu, Sanz, Williams, arXiv:1512.02572v

WH → Wbb at large M_{WH}

100 TeV



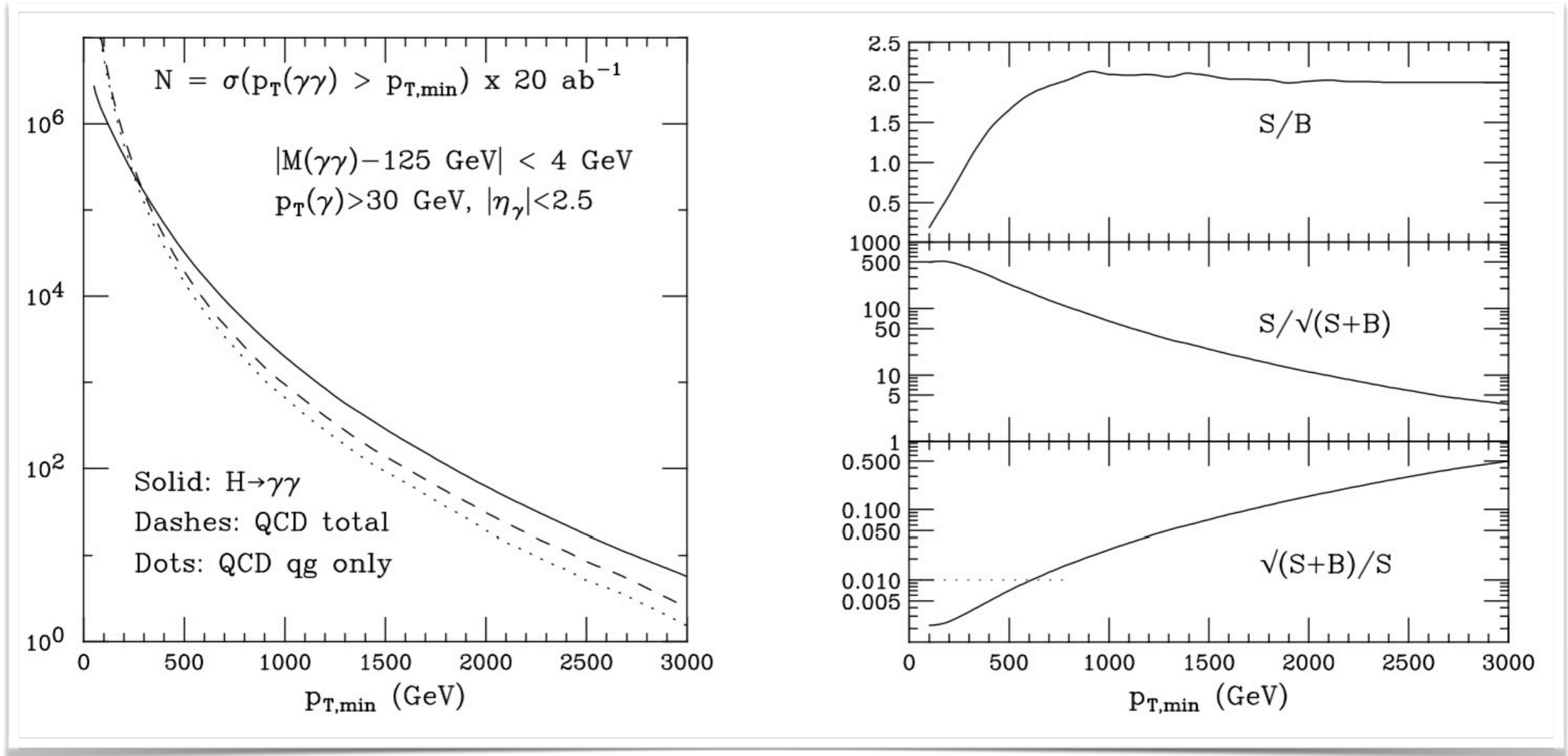
$gg \rightarrow H \rightarrow ZZ^* \rightarrow 4l$ at large p_T



- $S/B \sim 1$ for inclusive production at LHC
- Practically bg-free at large p_T at 100 TeV, maintaining large rates

$p_{T,min}$ (GeV)	δ_{stat}
100	0.3%
300	1%
1000	10%

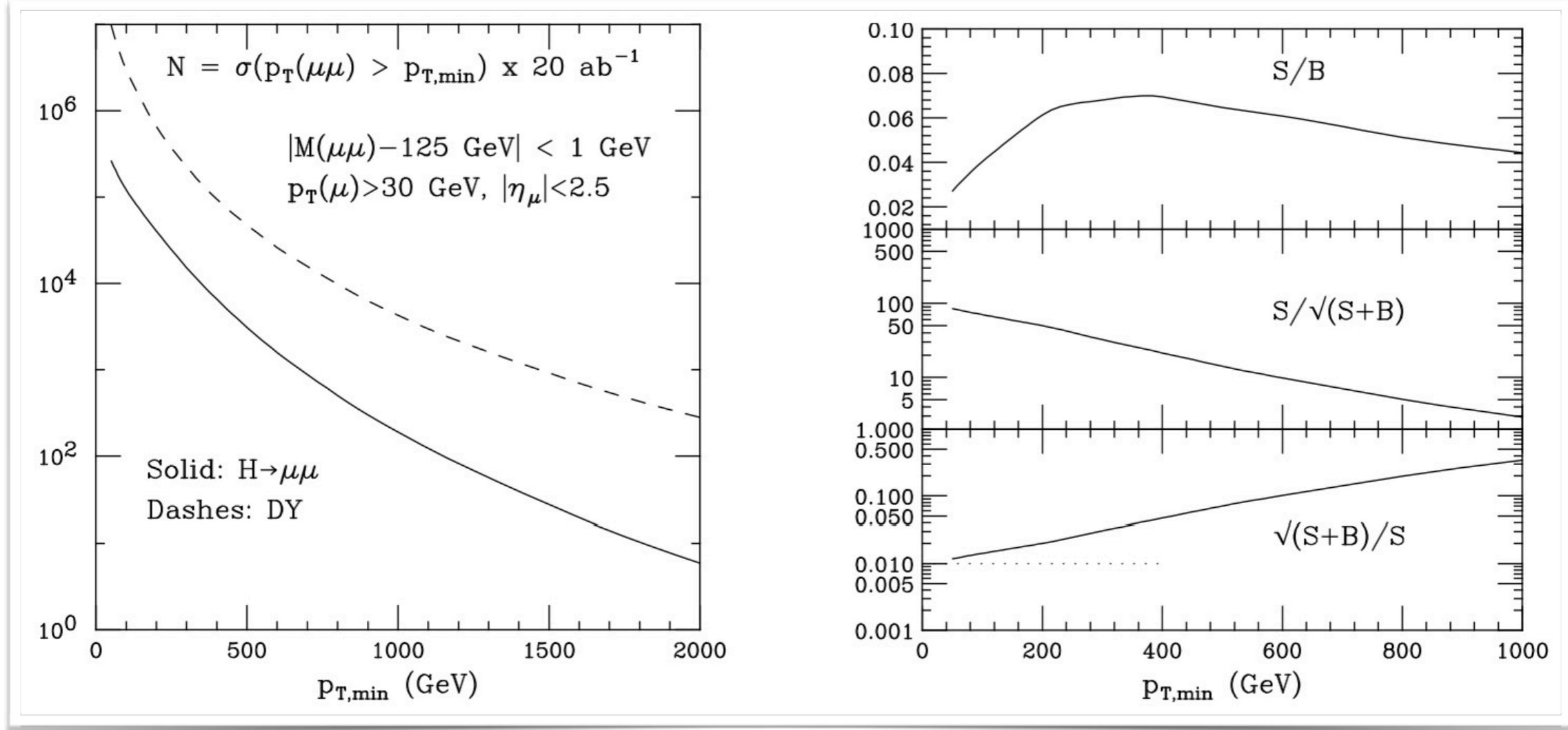
$gg \rightarrow H \rightarrow \gamma\gamma$ at large p_T



- At LHC, S/B in the $H \rightarrow \gamma\gamma$ channel is $O(\text{few } \%)$
- At FCC, for $p_T(H) > 300 \text{ GeV}$, $S/B \sim 1$
- Exptl systematics on $BR(\mu\mu)/BR(\gamma\gamma)$? (use same fiducial selection to remove H modeling syst's)
- Exptl mass resolution at large $p_T(H)$?
- Potentially accurate probe of the H p_T spectrum up to large p_T

$p_{T,\min}$	δ_{stat}
100	0.2%
400	0.5%
600	1%
1600	10%

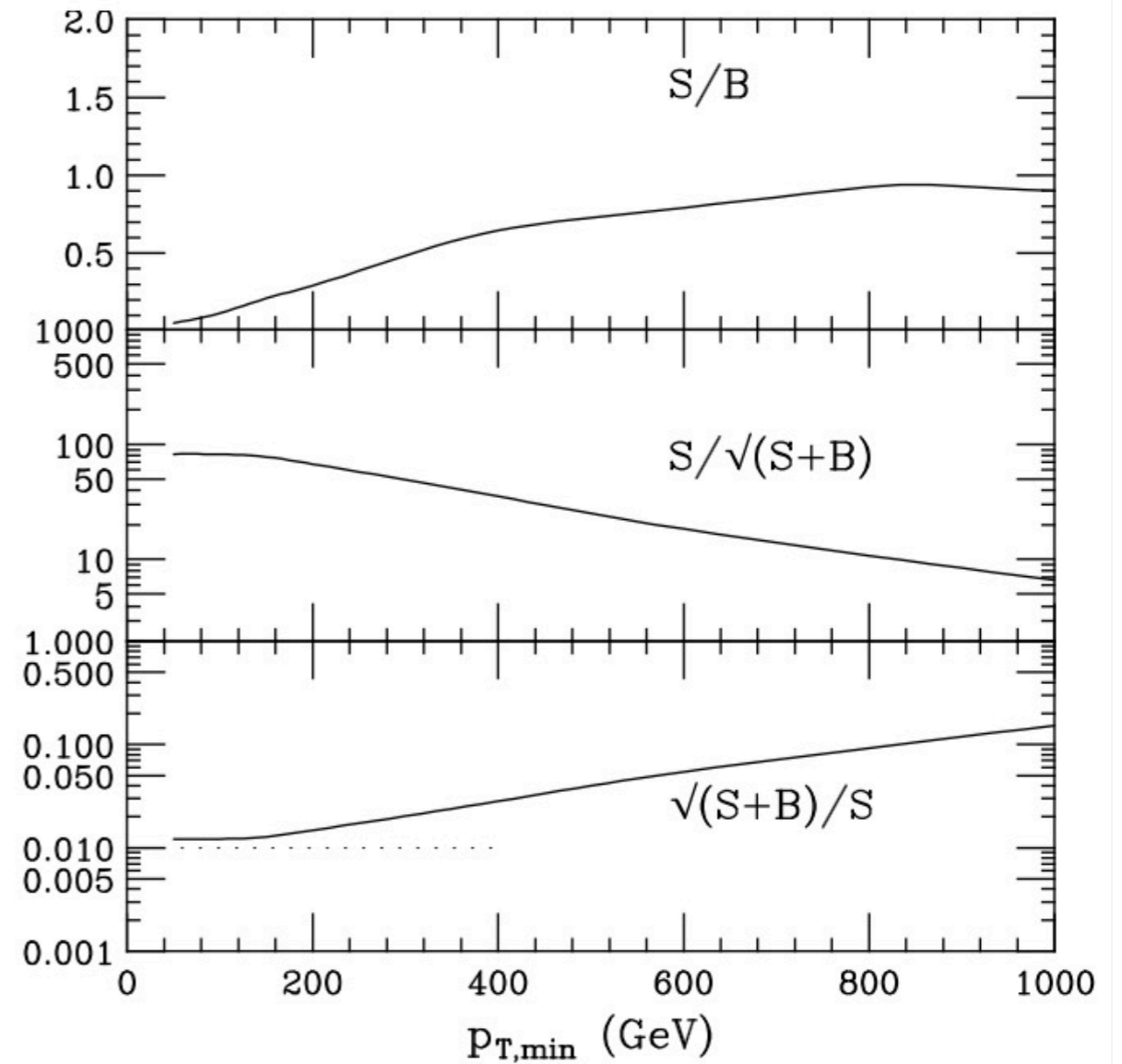
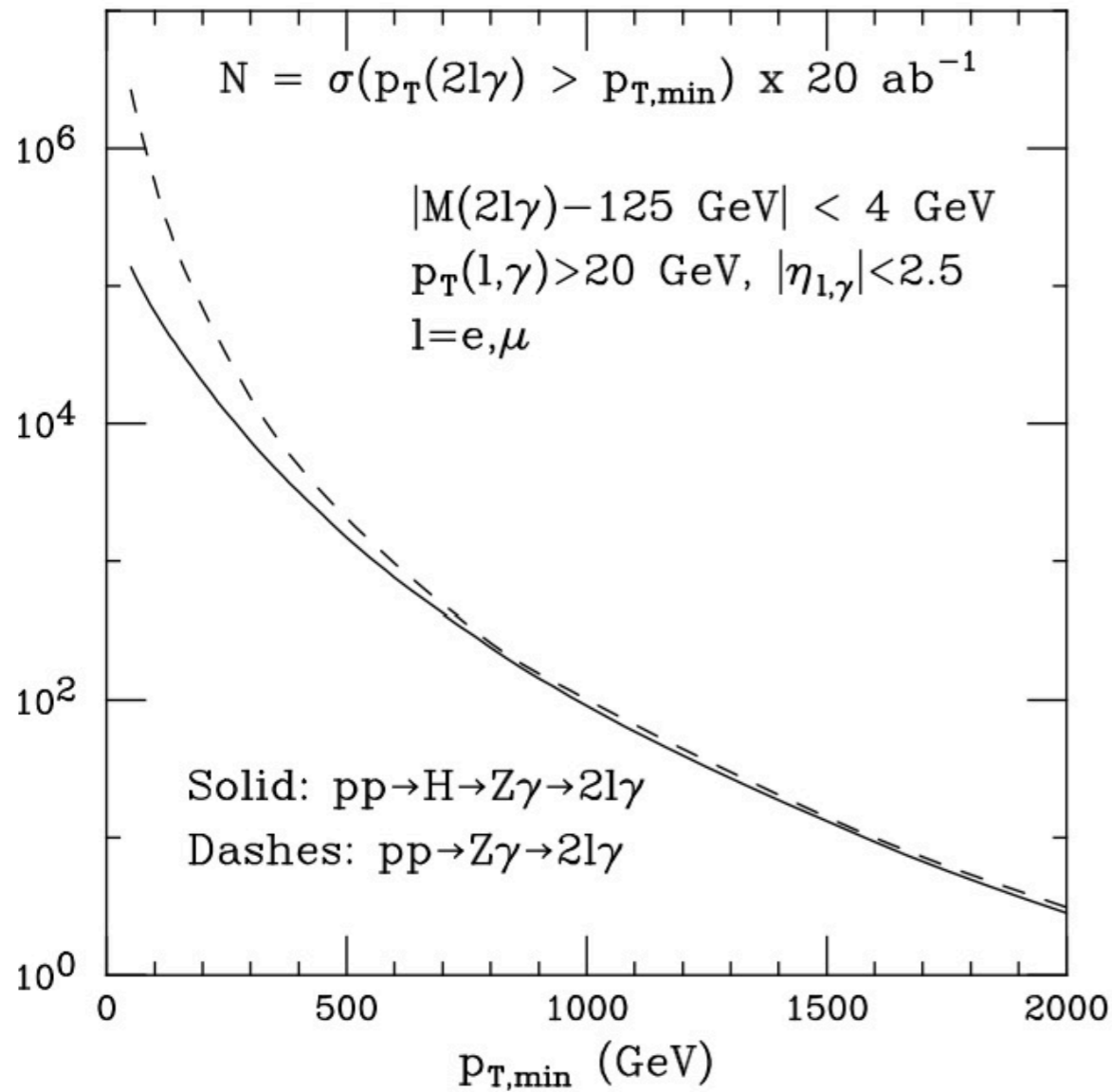
$gg \rightarrow H \rightarrow \mu\mu$ at large p_T



- Stat reach $\sim 1\%$ at $p_T \sim 100$ GeV
- Exptl systematics on $BR(\mu\mu)/BR(\gamma\gamma)$?
(use same fiducial selection to remove H modeling syst's)

$p_{T,min}$ (GeV)	δ_{stat}
100	1%
500	10%

$gg \rightarrow H \rightarrow Z\gamma \rightarrow \ell\ell\gamma$ at large p_T

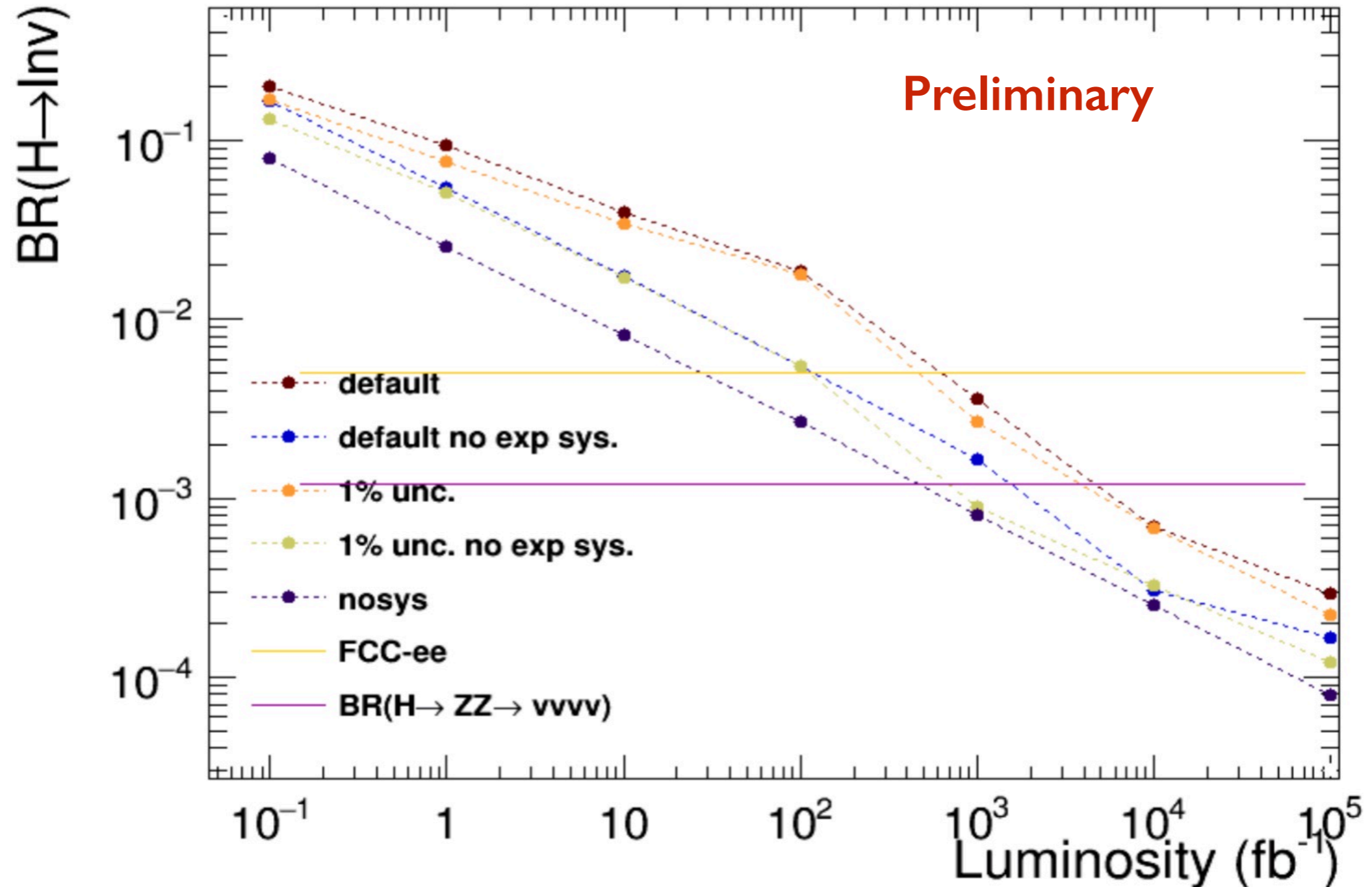


- $S/B \rightarrow 1$ at large p_T
- Stat reach $\sim 1\%$ at $p_T \sim 100 \text{ GeV}$
- Exptl systematics on $BR(Z\gamma)/BR(\gamma\gamma)$?

$p_{T,\min}$ (GeV)	δ_{stat}
100	1%
900	10%

BR(H→inv) in H+X production at large p_T(H)

Constrain bg pt spectrum from Z→vv to the % level using NNLO QCD/EW to relate to measured Z→ee, W and γ spectra



SM sensitivity with 1ab⁻¹, can reach few x 10⁻⁴ with 30ab⁻¹

Higgs couplings @ FCC

g_{HXY}	ee [240+350 (4IP)]	pp [100 TeV] 30ab ⁻¹	ep [60GeV/50TeV], 1ab ⁻¹
ZZ	0.15%	under study	
WW	0.19%		
bb	0.42%		0.2%
cc	0.71%		1.8%
gg	0.80%		
ττ	0.54%		
μμ	6.2%		<1%
γγ	1.5%		<0.5%
Zγ			<1%
tt	~13%		1%
HH	~30%	3.5%	under study
uu,dd	H->ργ, under study		
ss	H->φγ, under study		
BR _{inv}	< 0.45%	< 0.1%	
Γ _{tot}	1%		

- detailed study, stat+syst
- rather detailed, stat only (understood/limited/negligible theory syst)
- parton level S and B (from ratios, negligible TH syst, small exp syst)
- very preliminary estimates of exp/th syst (not stat-limited)

One should not underestimate the value of FCC-hh standalone precise “ratios-of-BRs” measurements:

- independent of $\alpha_S, m_b, m_c, \Gamma_{inv}$ systematics
- sensitive to BSM effects that typically influence BRs in different ways. Eg

$$\text{BR}(H \rightarrow \gamma\gamma) / \text{BR}(H \rightarrow ZZ^*)$$

loop-level

tree-level

$$\text{BR}(H \rightarrow \mu\mu) / \text{BR}(H \rightarrow ZZ^*)$$

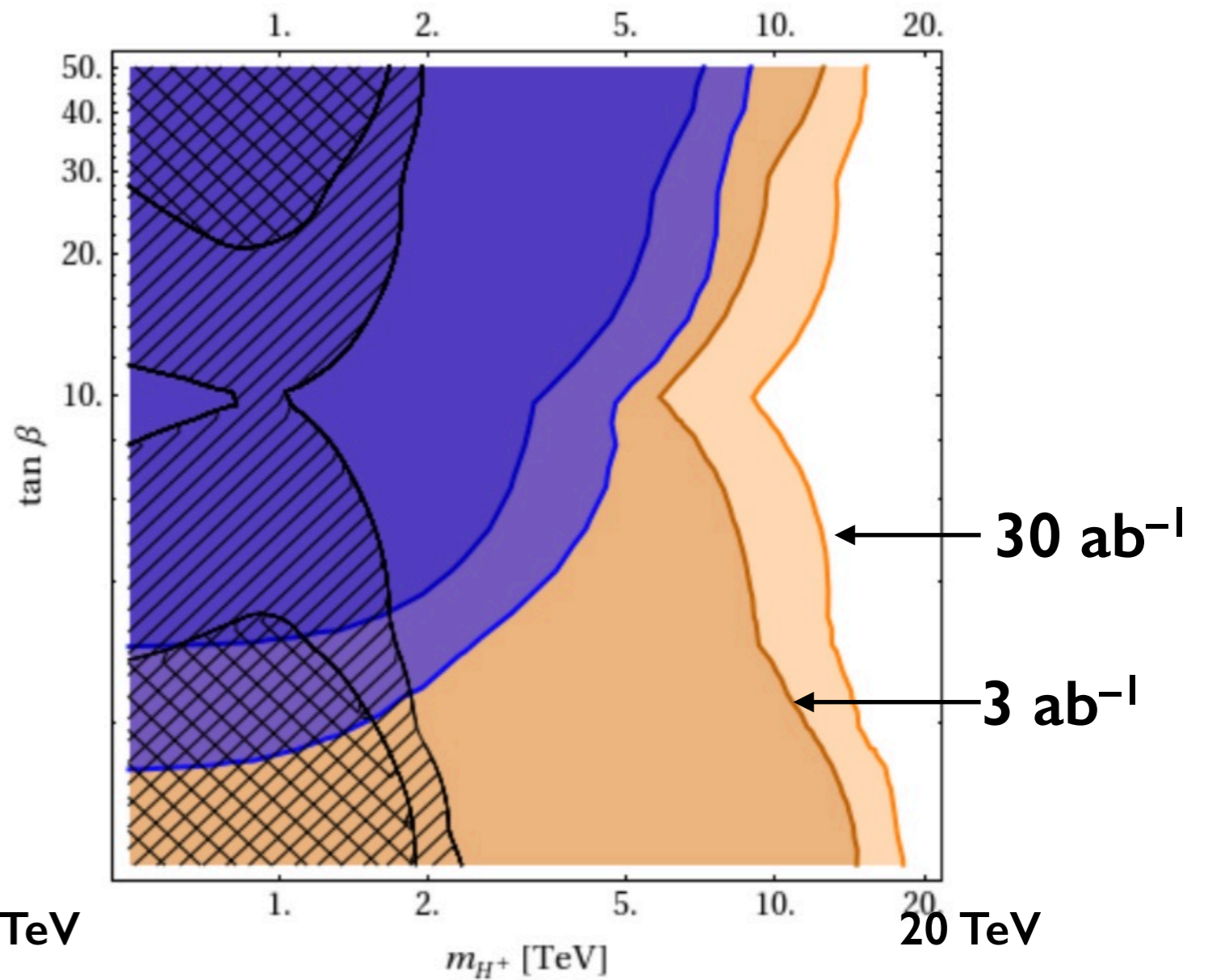
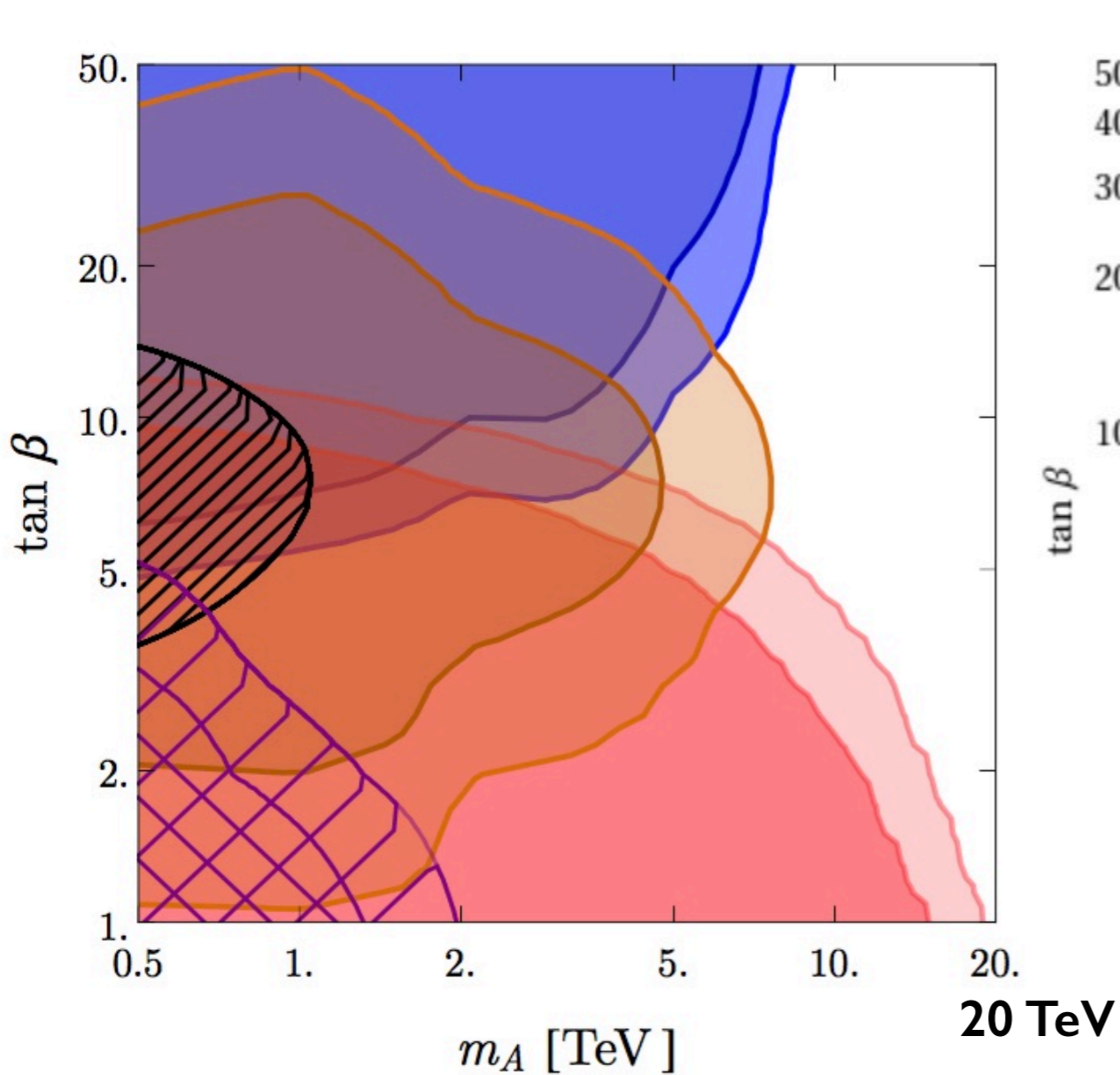
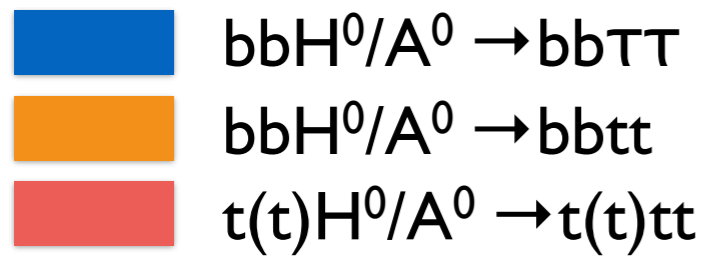
2nd gen'n Yukawa

gauge coupling

$$\text{BR}(H \rightarrow \gamma\gamma) / \text{BR}(H \rightarrow Z\gamma)$$

different EW charges in the loops of the two procs

MSSM Higgs @ 100 TeV



N. Craig, J. Hajer, Y.-Y. Li, T. Liu, H. Zhang,
arXiv:1605.08744

J. Hajer, Y.-Y. Li, T. Liu, and J. F. H. Shiu,
arXiv:1504.07617

Minimal stealthy model for a strong EW phase transition: the most challenging scenario for discovery

$$V_0 = -\mu^2 |H|^2 + \lambda |H|^4 +$$

$$\frac{1}{2}\mu_S^2 S^2 + \lambda_{HS} |H|^2 S^2 + \frac{1}{4}\lambda_S S^4$$

Unmixed SM+Singlet.
No exotic H decay, no H-S mixing,
no EWPO, ...

Two regions with strong EWPT

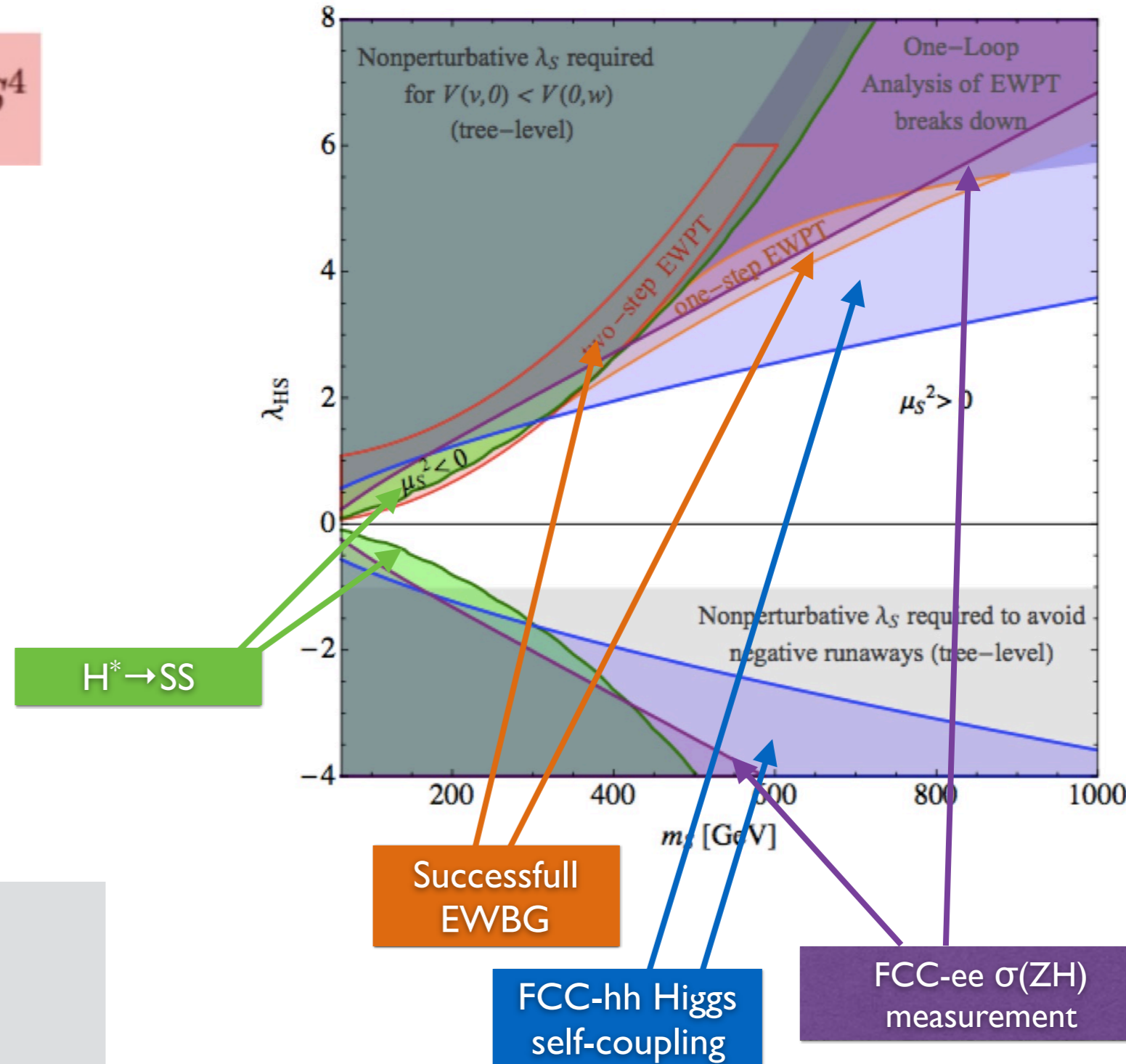
Only Higgs Portal signatures:

$h^* \rightarrow SS$ direct production

Higgs cubic coupling

$\sigma(Zh)$ deviation ($> 0.6\%$ @ TLEP)

Curtin, Meade, Yu, arXiv:1409.0005



⇒ Appearance of first “no-lose”
arguments for classes of compelling
scenarios of new physics

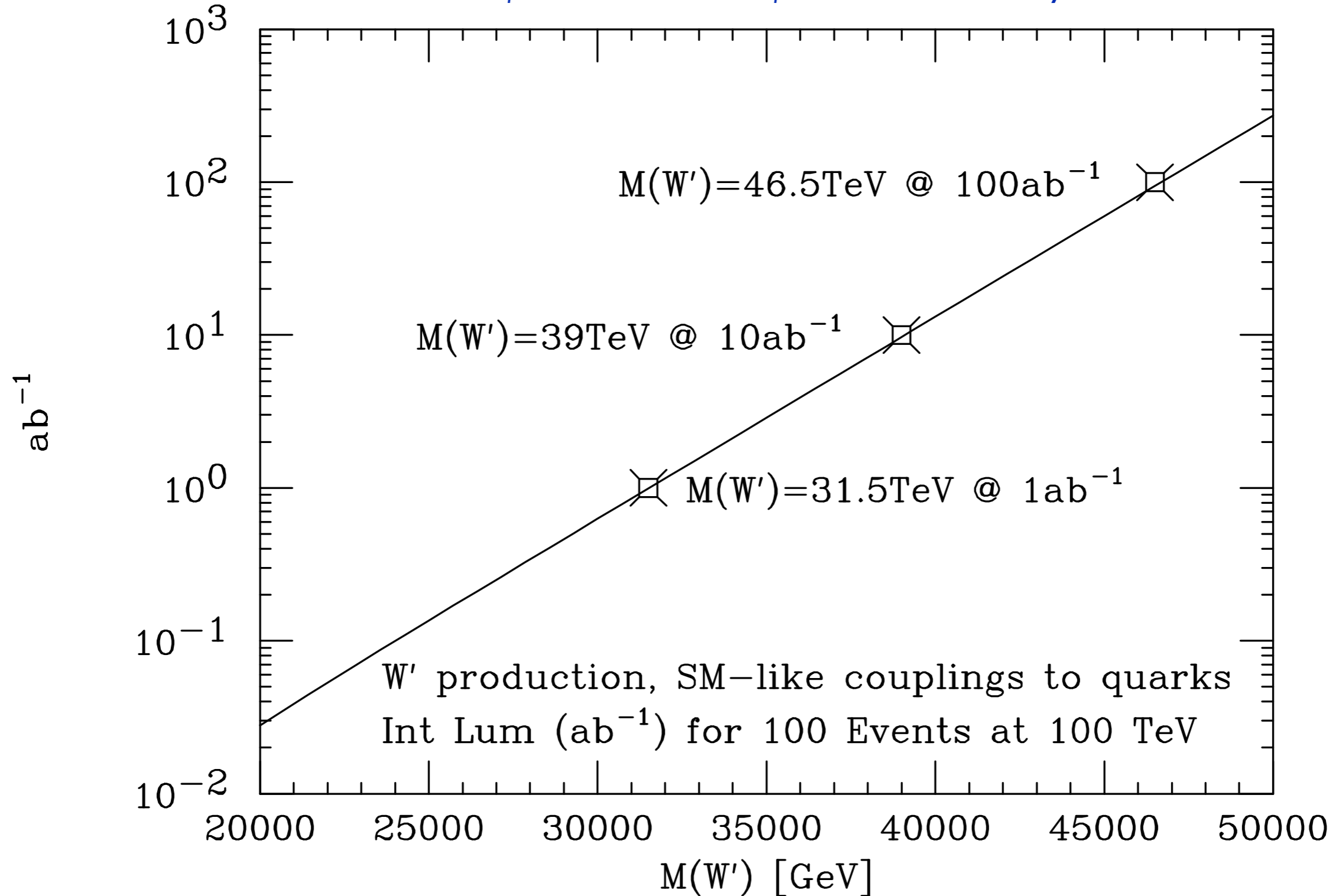
Direct discovery potential at the highest masses

at high mass, the reach of FCC-hh searches for BSM phenomena like Z' , W' , SUSY, LQs, top partners, etc.etc. scales trivially by $\sim 5-7$, depending on total luminosity ...

New gauge bosons: discovery reach

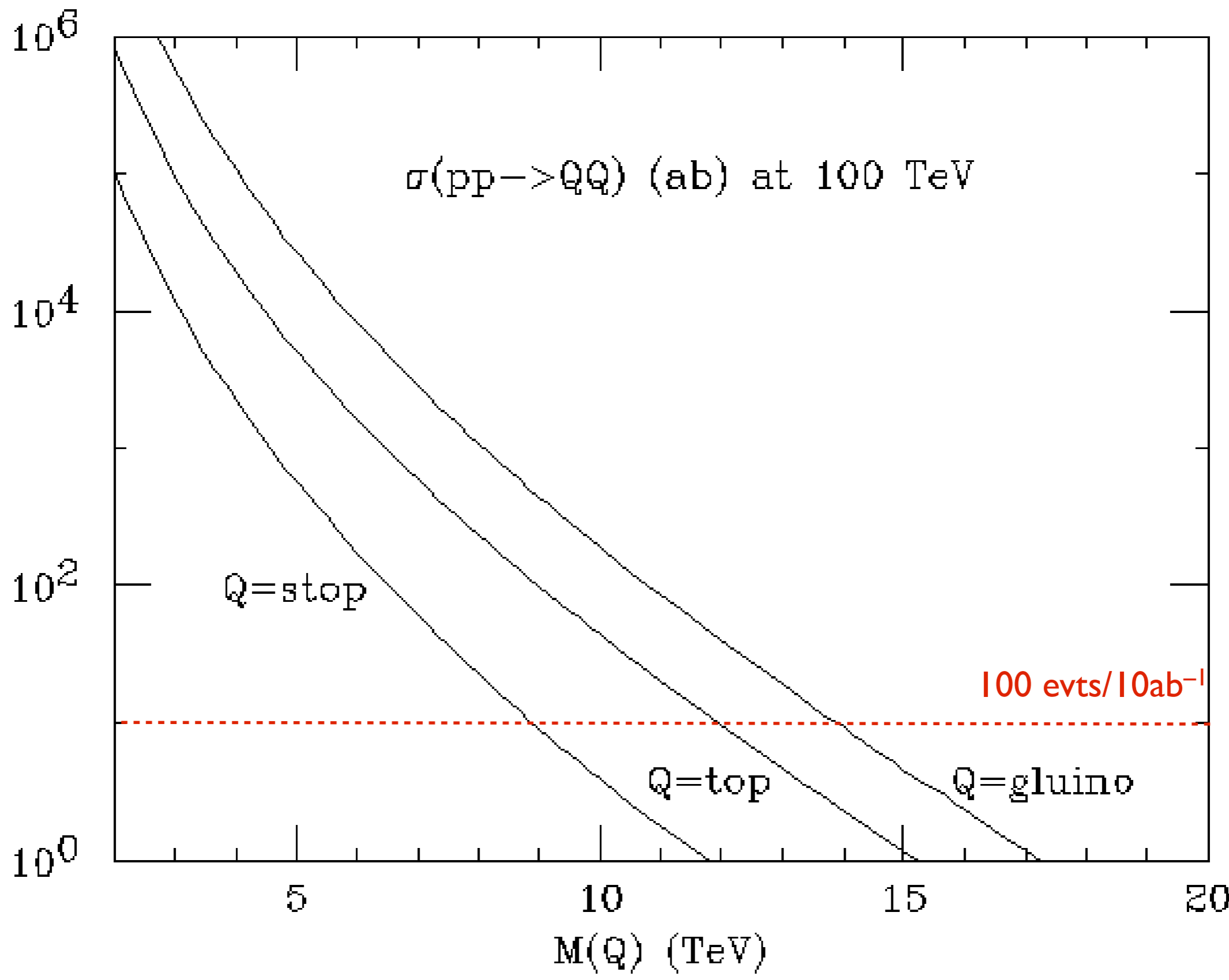
Example: W' with SM-like couplings

NB For SM-like Z' , $\sigma_{Z'} BR_{lept} \sim 0.1 \times \sigma_{W'} BR_{lept}$, \Rightarrow rescale lum by ~ 10



At $L=O(\text{ab}^{-1})$, Lum $\times 10 \Rightarrow \sim M + 7 \text{ TeV}$

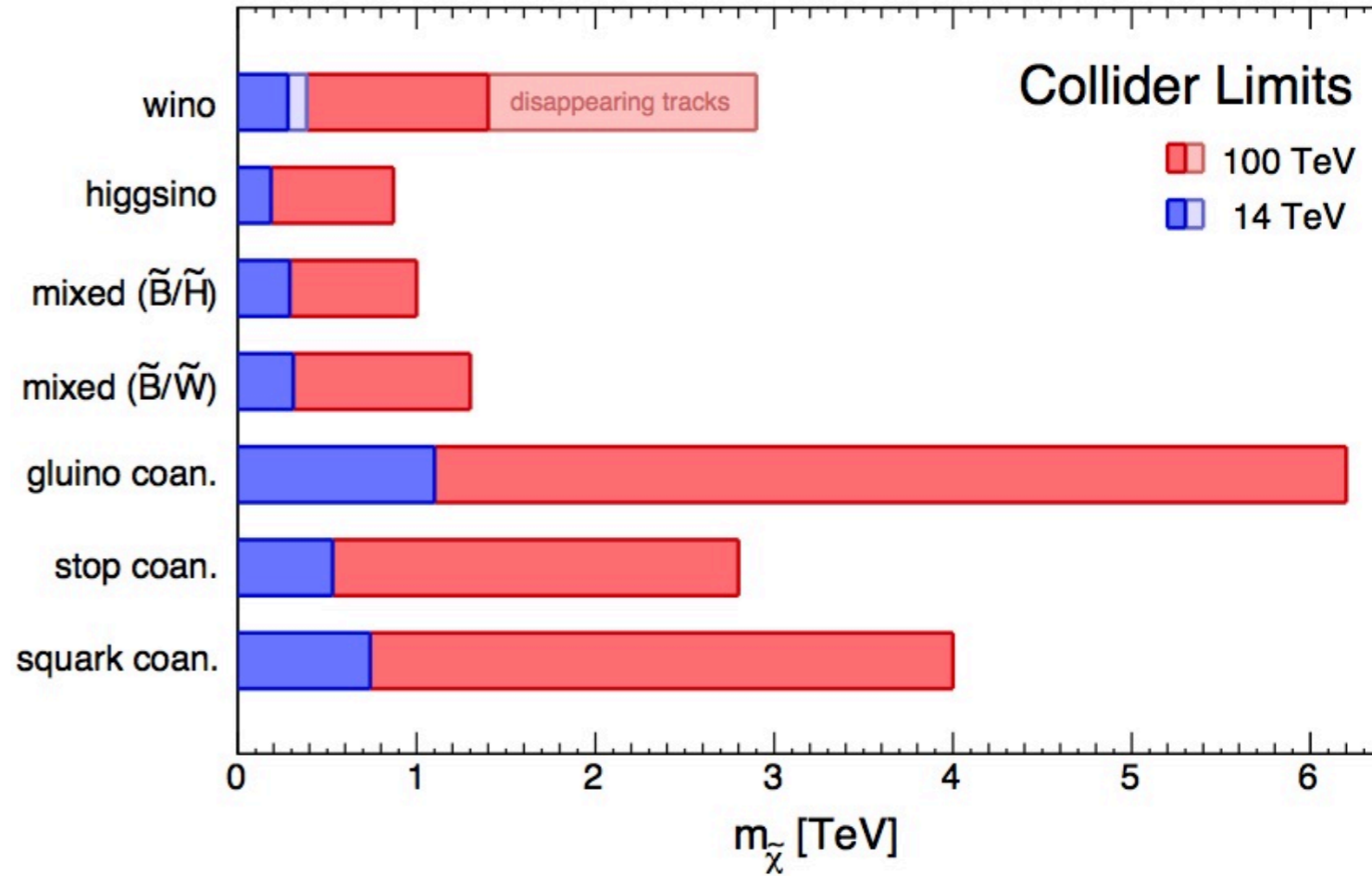
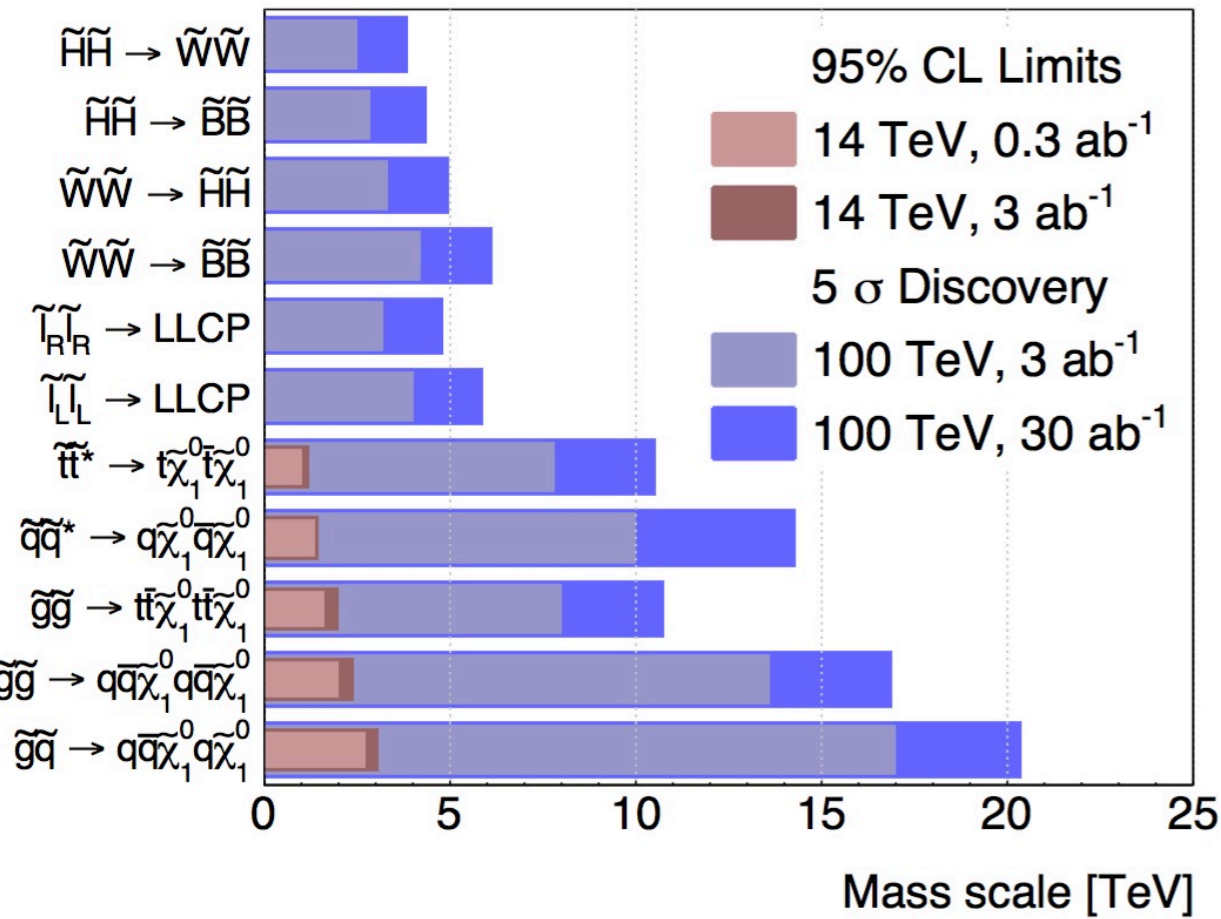
Discovery reach for pair production of strongly-interacting particles



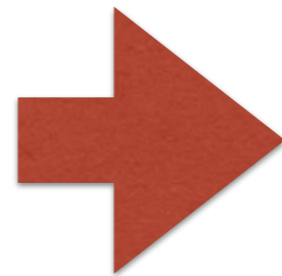
Dark Matter

- DM could be explained by BSM models that would leave no signature at any future collider (e.g. axions).
- More in general, no experiment can guarantee an answer to the question "what is DM?"
- Scenarios in which DM is a WIMP are however compelling and theoretically justified
- **We would like to understand whether a future collider can answer more specific questions, such as:**
 - do WIMPS contribute to DM?
 - can WIMPS, detectable in direct and indirect (DM annihilation) experiments, be discovered at future colliders? Is there sensitivity to the explicit detection of DM-SM mediators?
 - what are the opportunities w.r.t. new DM scenarios (e.g. interacting DM, asymmetric DM,)?

SUSY and DM reach at 100 TeV



$$M_{\text{WIMP}} \leq 1.8 \text{ TeV} \left(\frac{g^2}{0.3} \right)$$



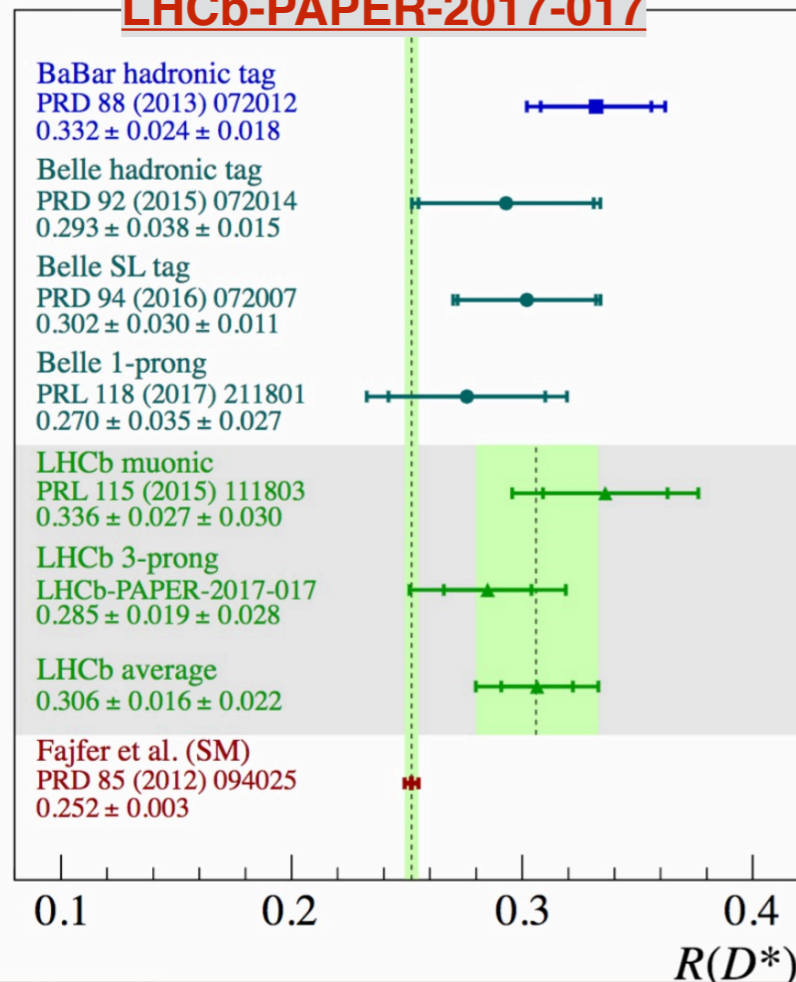
possibility to find (or rule out) thermal WIMP DM candidates

Flavour anomalies at LHC & Bfact's

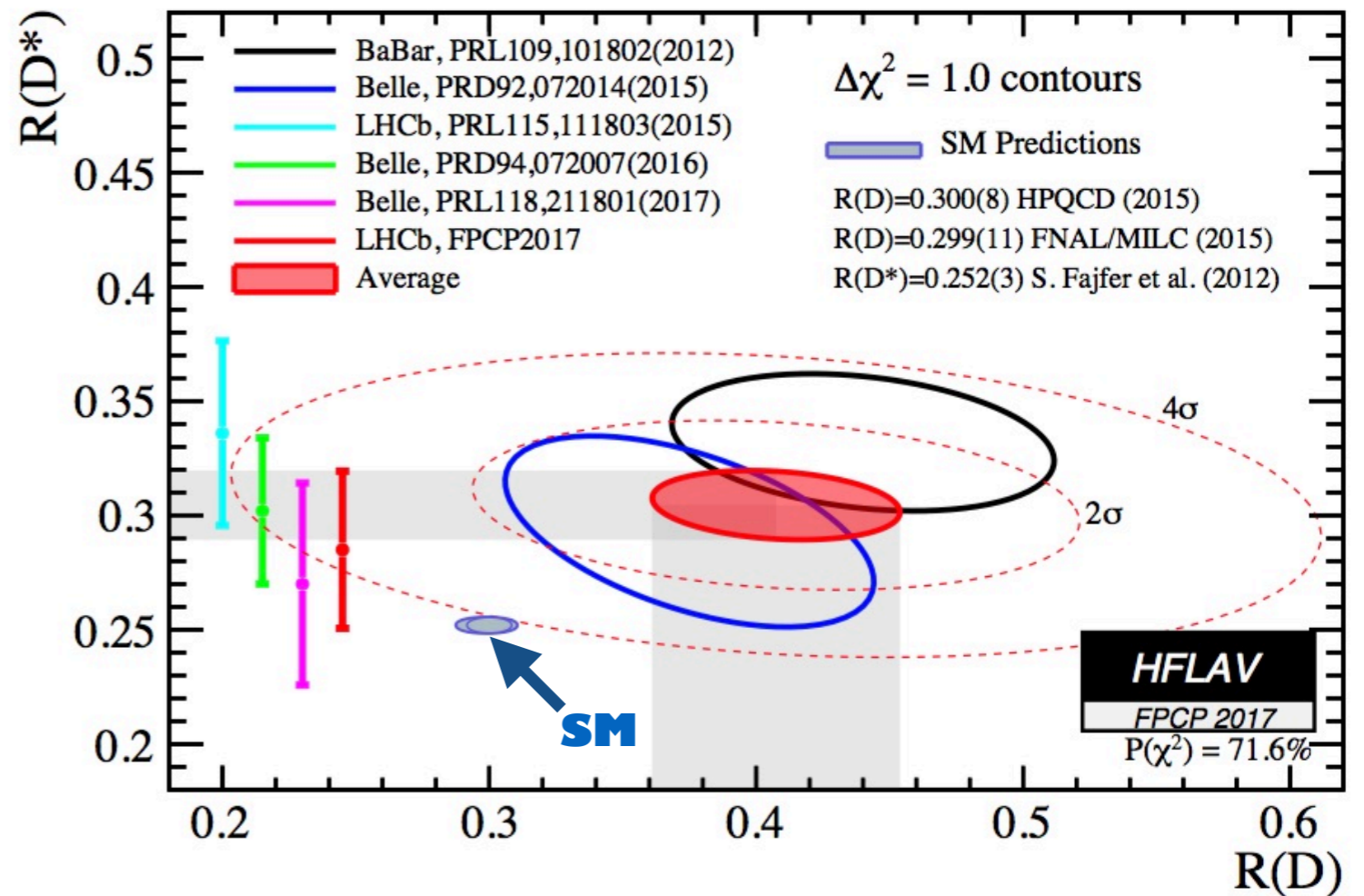
$b \rightarrow c \ell \nu$

$$R(D^{(*)}) = \frac{BR(B \rightarrow D^{(*)} \tau \nu)}{BR(B \rightarrow D^{(*)} \mu \nu)}$$

LHCb-PAPER-2017-017



Overall combination of R(D) and R(D*) is 4.1σ from SM



$b \rightarrow s \ell \ell$

$$R_{K^{(*)}} = \frac{BR(B \rightarrow K^{(*)} \mu \mu)}{BR(B \rightarrow K^{(*)} e e)}$$

$m_{\mu\mu}$ [mass range]	SM	Exp.
R_K [1-6]	1.00 ± 0.01	$0.745_{-0.074}^{+0.090} \pm 0.036$
R_{K^*} [1.1-6]	1.00 ± 0.01	$0.685_{-0.069}^{+0.113} \pm 0.047$
R_{K^*} [0.045,1.1]	0.91 ± 0.03	$0.660_{-0.070}^{+0.110} \pm 0.024$

LHCb, PRL 113 (2014) 151601, arXiv:1705.05802

Remarks

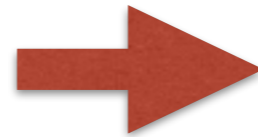
The above observables are theoretically robust: small and reliable uncertainties

Other anomalies at the $2\text{-}3\sigma$ level exist, but subject to less robust estimates of QCD uncertainties

Statistics still plays a dominant role (esp for R_K). More data will also allow use of new final states with independent exptl systematics ... eg

LHCb-PAPER-2017-035, to appear

$$R_{J/\psi} = \frac{BR(B_c \rightarrow J/\psi\tau\nu)}{BR(B_c \rightarrow J/\psi\mu\nu)}$$



$$R(J/\psi) = 0.71 \pm 0.17 \pm 0.18$$

(about 2σ from SM)

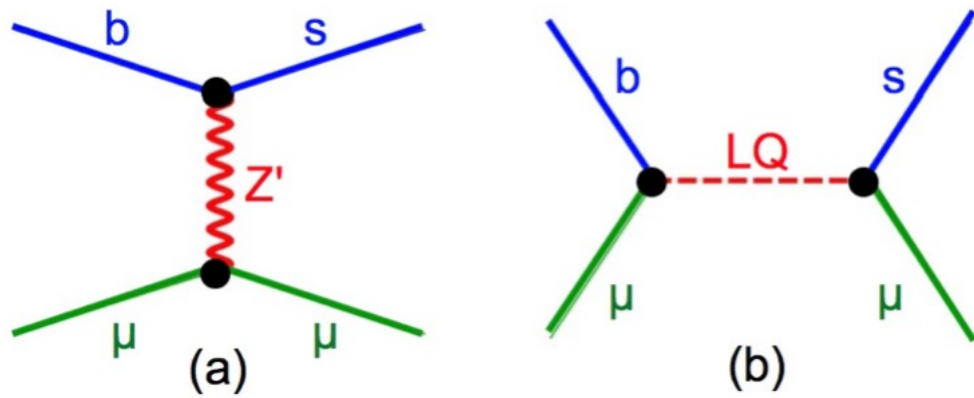
The fact that SM deviations of this type, variety and size are phenomenologically acceptable, gives a sign of how little we still know about “what’s out there” at the TeV scale, and our openness towards surprises (see also the story of the 750 GeV $\gamma\gamma$ resonance)

Example of EFT interpretation of R_K

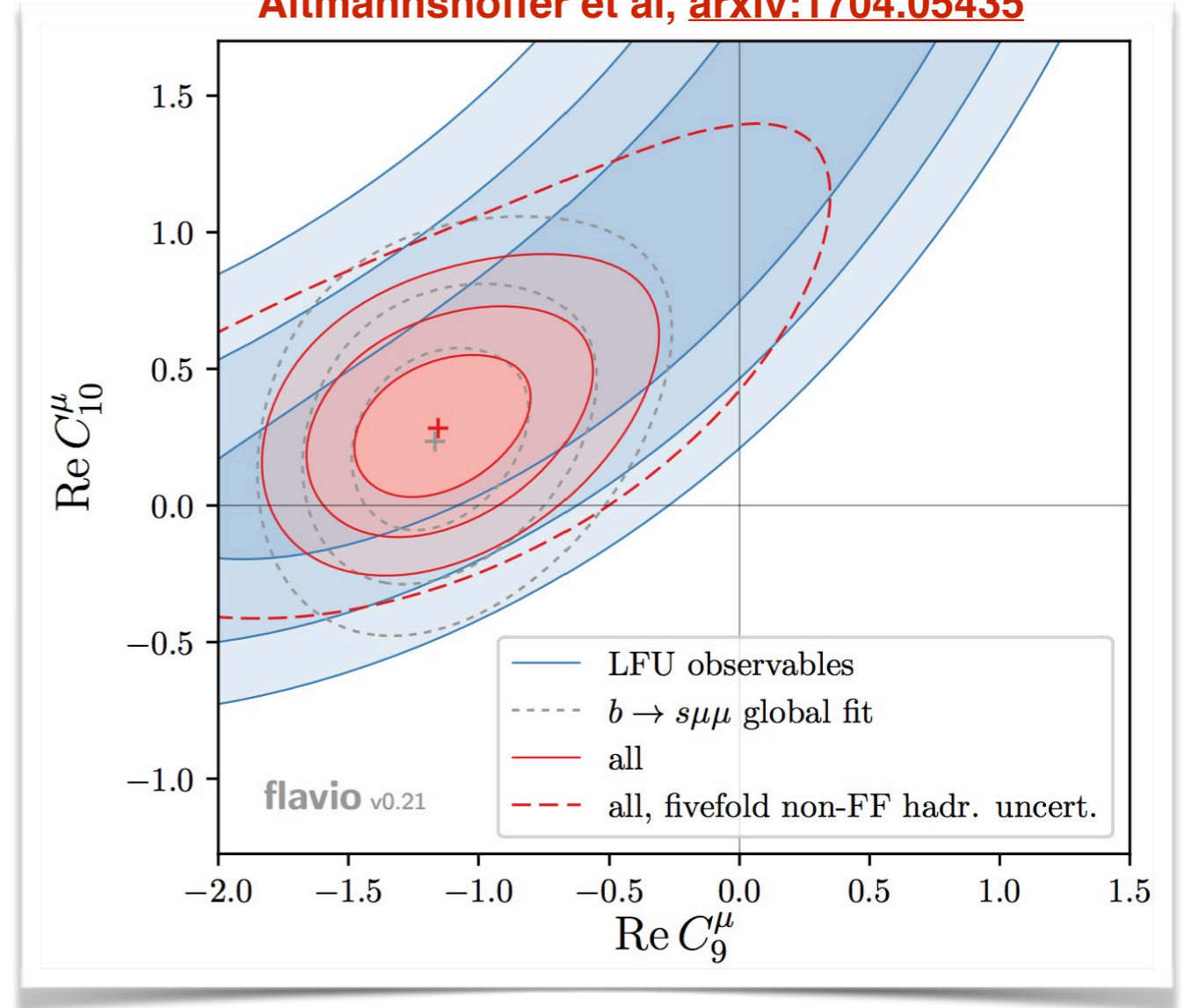
$$O_9^l = (\bar{s}\gamma_\mu P_L b)(\bar{l}\gamma^\mu l),$$

$$O_{10}^l = (\bar{s}\gamma_\mu P_L b)(\bar{l}\gamma^\mu \gamma_5 l)$$

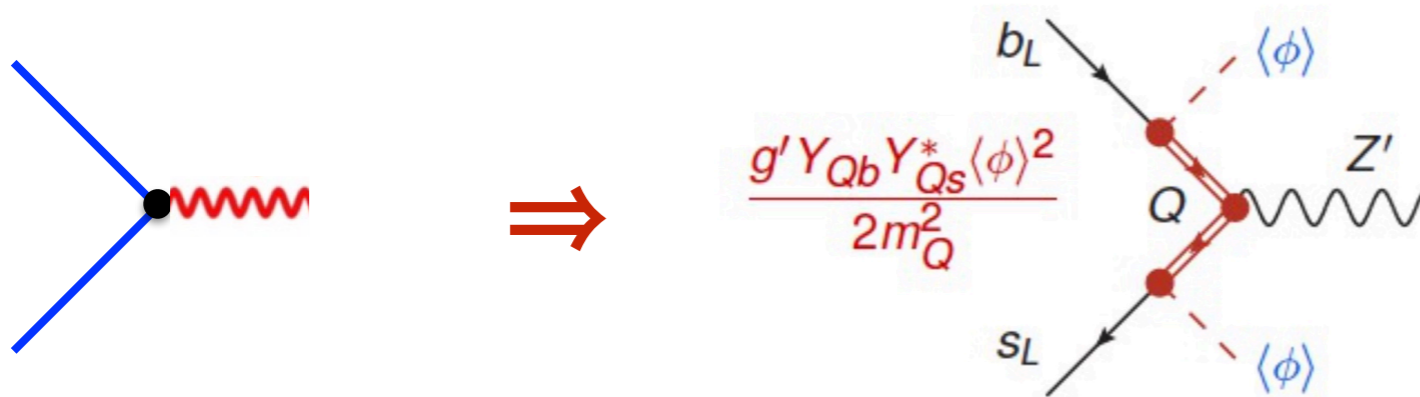
Possible explicit realizations:



Altmannshofer et al, arxiv:1704.05435



where, e.g. ,



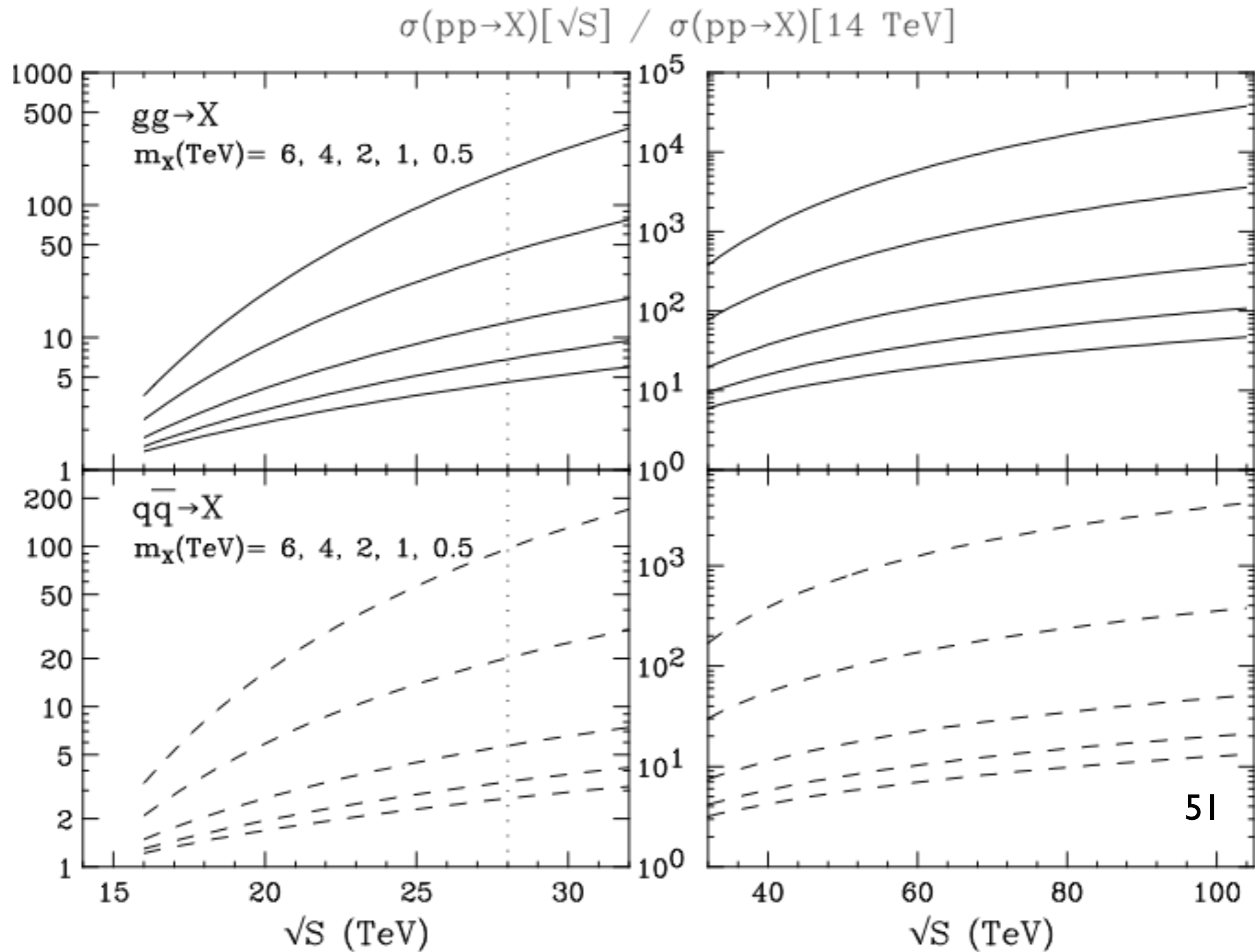
Upper limits on Z' and Leptoquark masses are model-dependent, and constrained also by other low-energy flavour phenomenology, but typically lie in the range of $1 \rightarrow O(10)$ TeV
 \Rightarrow if anomalies confirmed, we may want a no-lose theorem to identify the next facility!

100 TeV ?

200 TeV ?

28 TeV ?

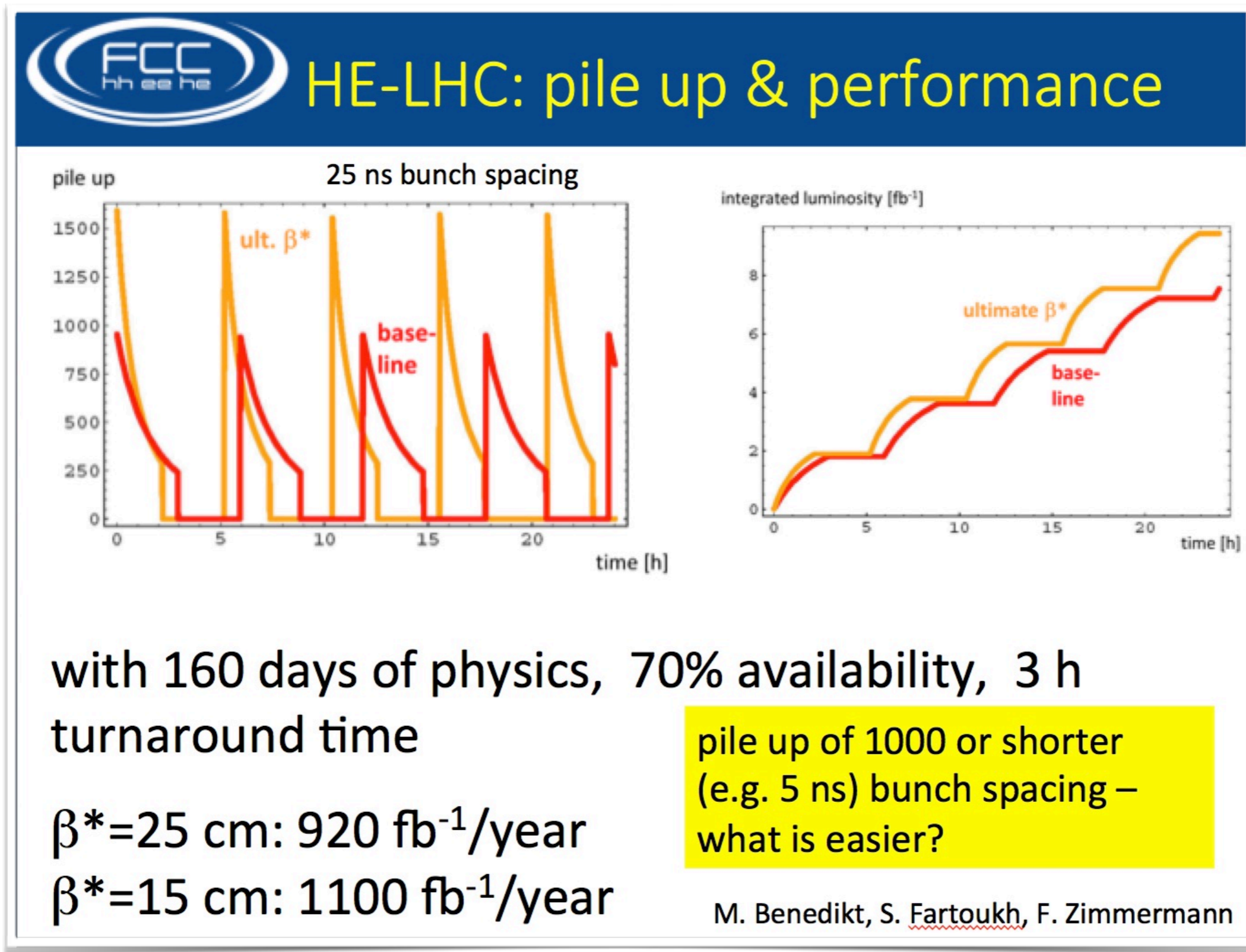
Evolution, with beam energy, of scenarios with the discovery of a new particle at the LHC



Possible questions/options

- If $m_X \sim 6$ TeV in the gg channel, rate grows x 200 @28 TeV:
 - Do we wait 40 yrs to go to pp@100TeV, or fast-track 28 TeV in the LHC tunnel?
 - Do we need 100 TeV, or 50 is enough ($\sigma_{100}/\sigma_{14} \sim 4 \cdot 10^4$, $\sigma_{50}/\sigma_{14} \sim 4 \cdot 10^3$) ?
 - ... and the answers may depend on whether we expect partners of X at masses $\gtrsim 2m_X$ (\Rightarrow 28 TeV would be *insufficient* ...)
- If $m_X \sim 0.5$ TeV in the qqbar channel, rate grows x10 @100 TeV:
 - Do we go to 100 TeV, or push by x10 $\int L$ at LHC?
 - Do we build CLIC?
- etc.etc.

HE-LHC (27 TeV), prelim performance estimates



=> O(15 ab⁻¹) over 15-20 years

Systematics studies* of the full physics potential at O(28) TeV, with O(15 ab⁻¹), need to be carried out

* except for straightfwd mass-reach extrapolations from LHC

E.g. HH at 28 TeV (back of the envelope)

$$\sigma_{HH}(28 \text{ TeV})/\sigma_{HH}(14 \text{ TeV}) \sim 4 \quad \text{Lum}(28) \sim 4 \text{ Lum}(14 \text{ TeV})$$

$$\Rightarrow N_{HH}(28) \sim 16 N_{HH}(14)$$

$$\Rightarrow \delta\lambda_{HHH}(28) \sim \delta\lambda_{HHH}(\text{HL-LHC}) / 4 \sim 10\%$$

Expect to carry out an overall evaluation of the physics potential during 2018 (in the context of the HL-LHC Physics workshop, <https://indico.cern.ch/event/647676>)

What does the HE-LHC entail?

- **Necessary:**
 - empty the tunnel (more time & \$s than removing LEP)
 - full replacement of the magnets (today's cost $\sim 4 \times$ LHC. First prototypes in ~ 2026)
 - upgrade of RF, cryogenics, collimation, beam dumps, ...
 - **Very likely:**
 - major upgrade of SPS, if need to inject at $O(1 \text{ TeV})$ (magnets, RF, transfer lines, cryo if SC, ...)
 - major overhaul of detectors (radiation damage after HL-LHC, use of new technologies)
- => it's like building the LHC ex-novo
- very unlikely to be cheaper ...
 - ... but not incompatible with a \sim constant CERN budget
 - nevertheless feasibility to be proven (eg magnets bigger than LHC's: will they fit in the tunnel ??)

Snapshots of the status of the FCC studies



progress - civil engineering studies

Review panel – Decision to focus on 100 km tunnel

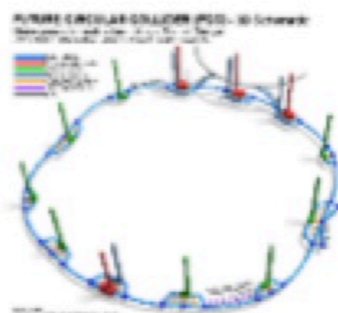


FCC week 2016 in Rome:

- Single and double tunnel
- Inclined access tunnels
- hh and ee requirements



- Revised layout for realisation studies
- Naming convention



Cost and schedule study ongoing with 2 consultants



- Cost & schedule estimates
- Inclined access shafts assessment
- Tunnel and shaft cross-section designs



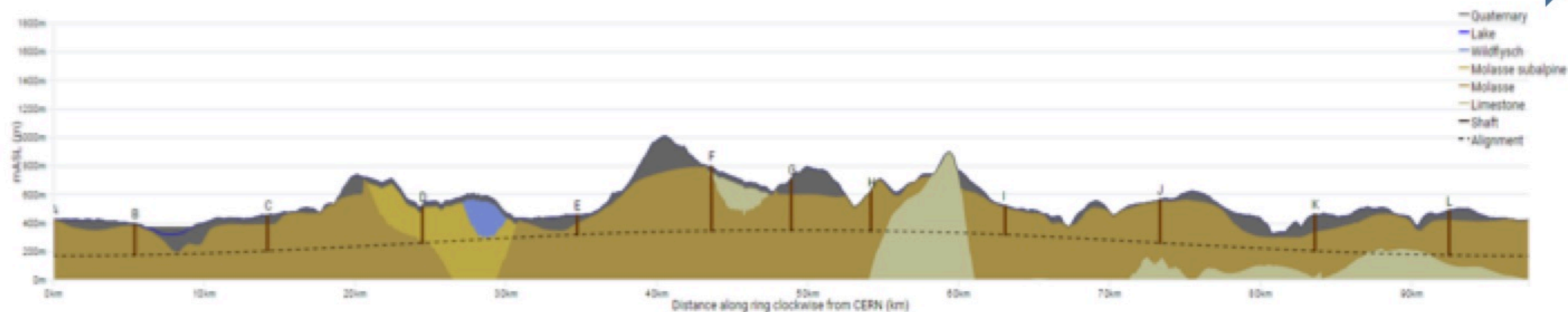
Nov. 2015

Apr. 2016

Aug. 2016

Sept. 2016

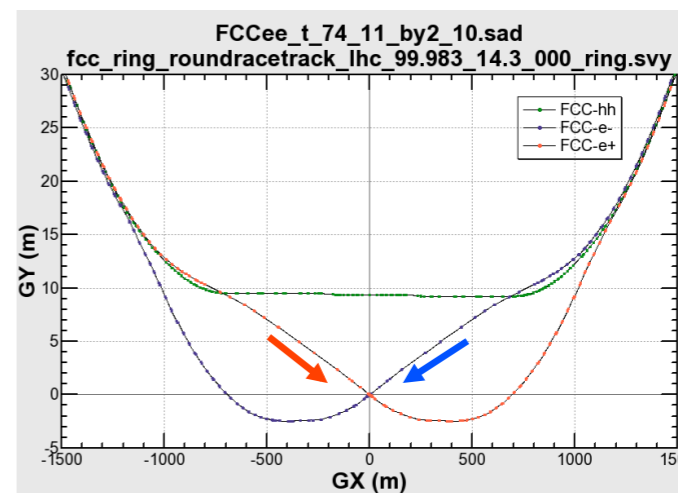
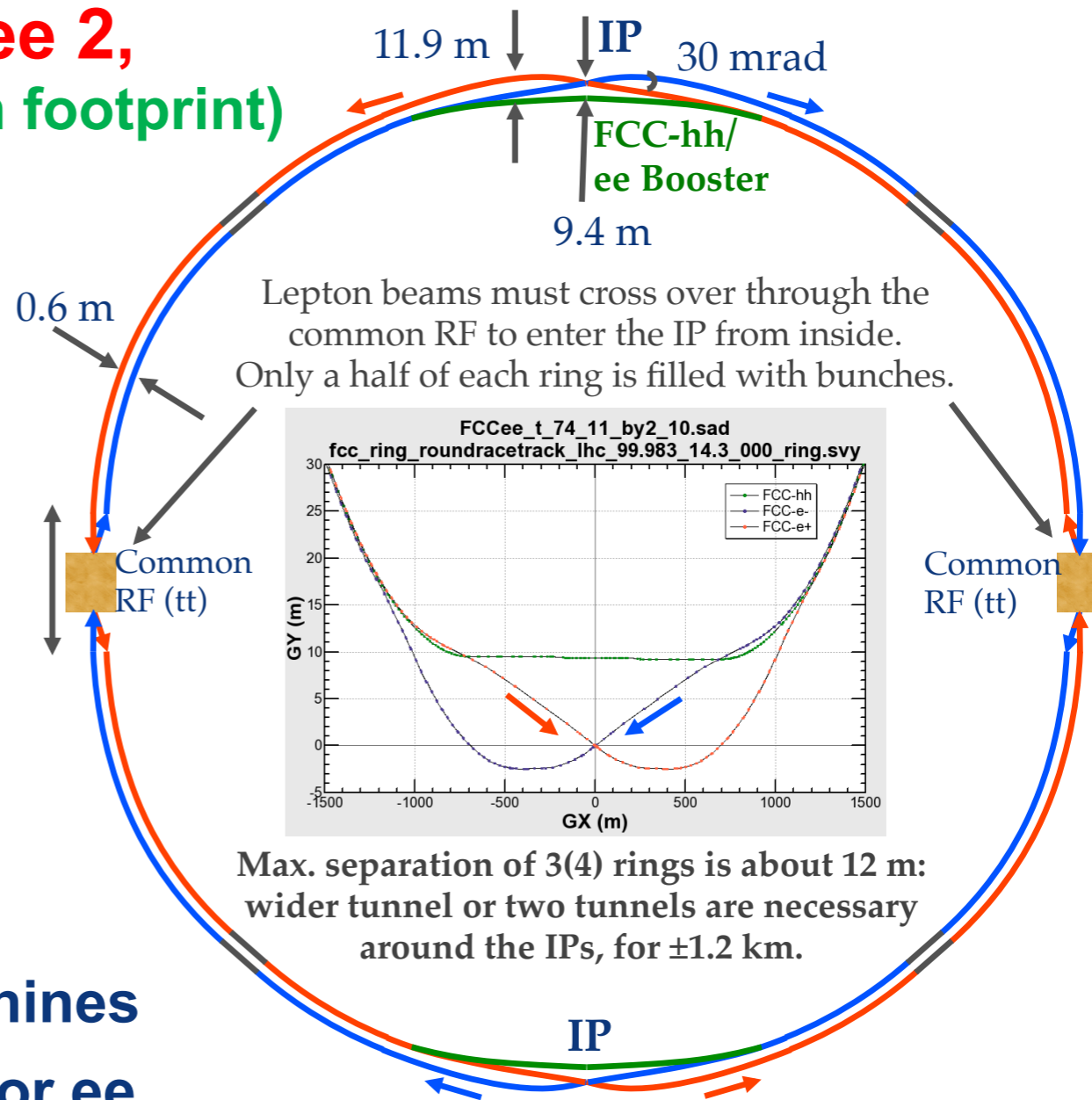
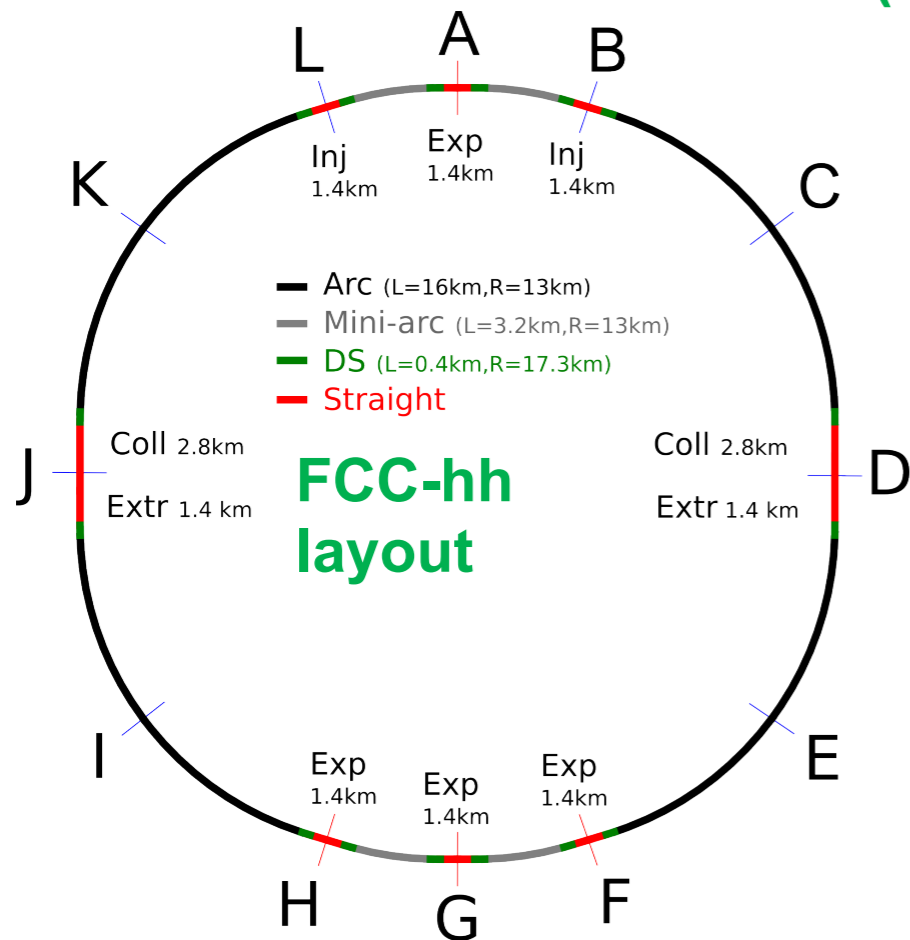
Dec. 2016



Future Circular Collider Study
 Michael Benedikt
 FCC Physics Workshop, CERN, 16 January 2017

FCC-ee 1, FCC-ee 2,

FCC-ee booster (FCC-hh footprint)



Max. separation of 3(4) rings is about 12 m:
wider tunnel or two tunnels are necessary
around the IPs, for ± 1.2 km.

- 2 main IPs in A, G for both machines
- asymmetric IR optic/geometry for ee to limit synchrotron radiation to detector

Injector options:

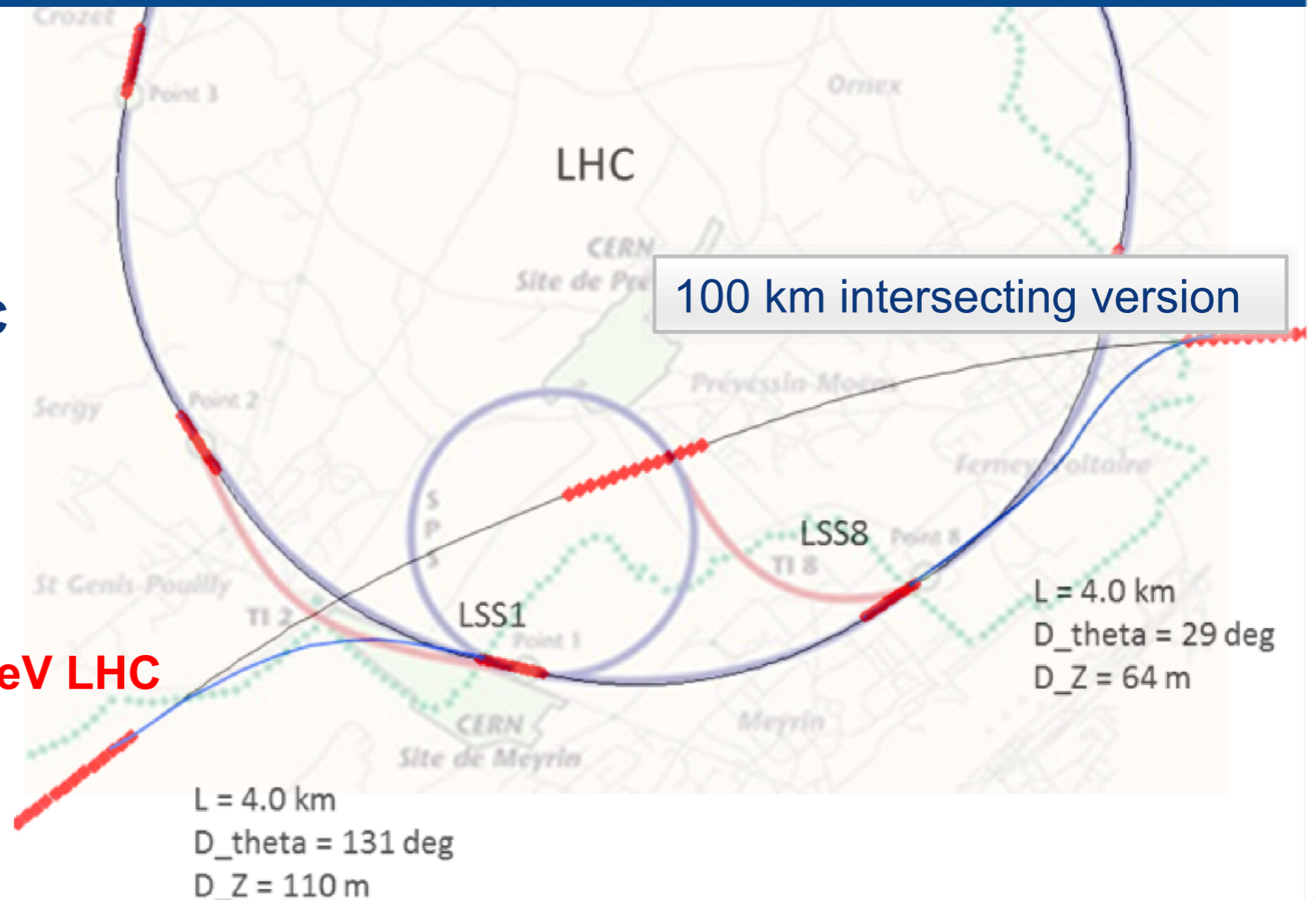
- SPS → LHC → FCC
- SPS/SPS_{upgrade} → FCC

Current baseline:

- **Injection energy 3.3 TeV LHC**

Alternative option:

- **Injection around 1.5 TeV**
- SPS_{upgrade} could be based on fast-cycling SC magnets, 6-7T, ~ 1T/s ramp





FCC-pp collider parameters



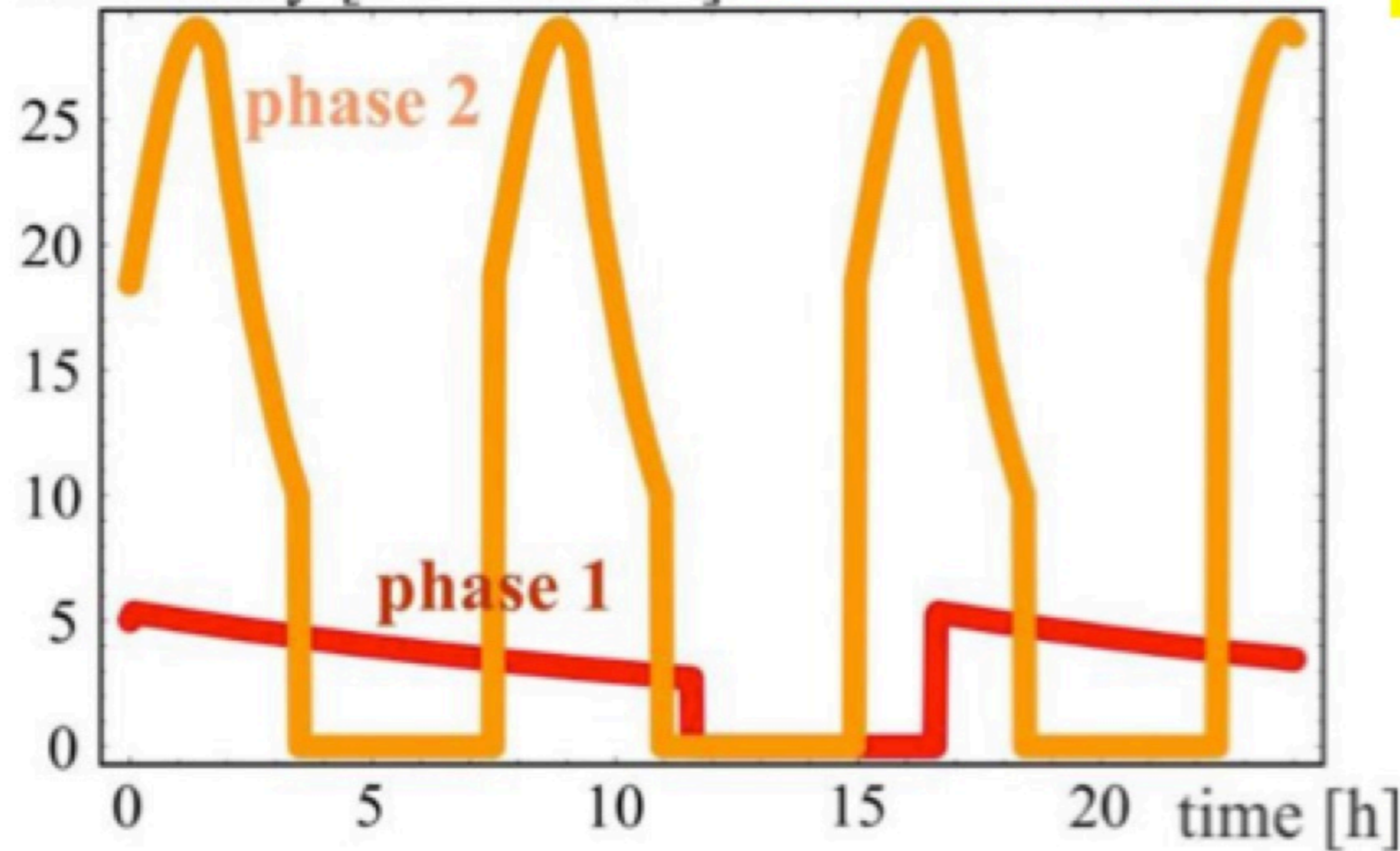
parameter	FCC-hh		HE-LHC	HL-LHC	LHC
collision energy cms [TeV]	100		27	14	14
dipole field [T]	16		16	8.33	8.33
circumference [km]	97.75		26.7	26.7	26.7
beam current [A]	0.5		1.12	1.12	0.58
bunch intensity [10^{11}]	1	1 (0.2)	2.2 (0.44)	2.2	1.15
bunch spacing [ns]	25	25 (5)	25 (5)	25	25
synchr. rad. power / ring [kW]	2400		101	7.3	3.6
SR power / length [W/m/ap.]	28.4		4.6	0.33	0.17
long. emit. damping time [h]	0.54		1.8	12.9	12.9
beta* [m]	1.1	0.3	0.25	0.20	0.55
normalized emittance [μm]	2.2 (0.4)		2.5 (0.5)	2.5	3.75
peak luminosity [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$]	5	30	25	5	1
events/bunch crossing	170	1k (200)	~800 (160)	135	27
stored energy/beam [GJ]	8.4		1.3	0.7	0.36



luminosity evolution over 24 h

luminosity [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$] radiation damping: $\tau \sim 1 \text{ h}$

PRST-AB 18, 101002 (2015)



for both phases:

beam current 0.5 A, unchanged!

total synchrotron radiation power $\sim 5 \text{ MW}$.

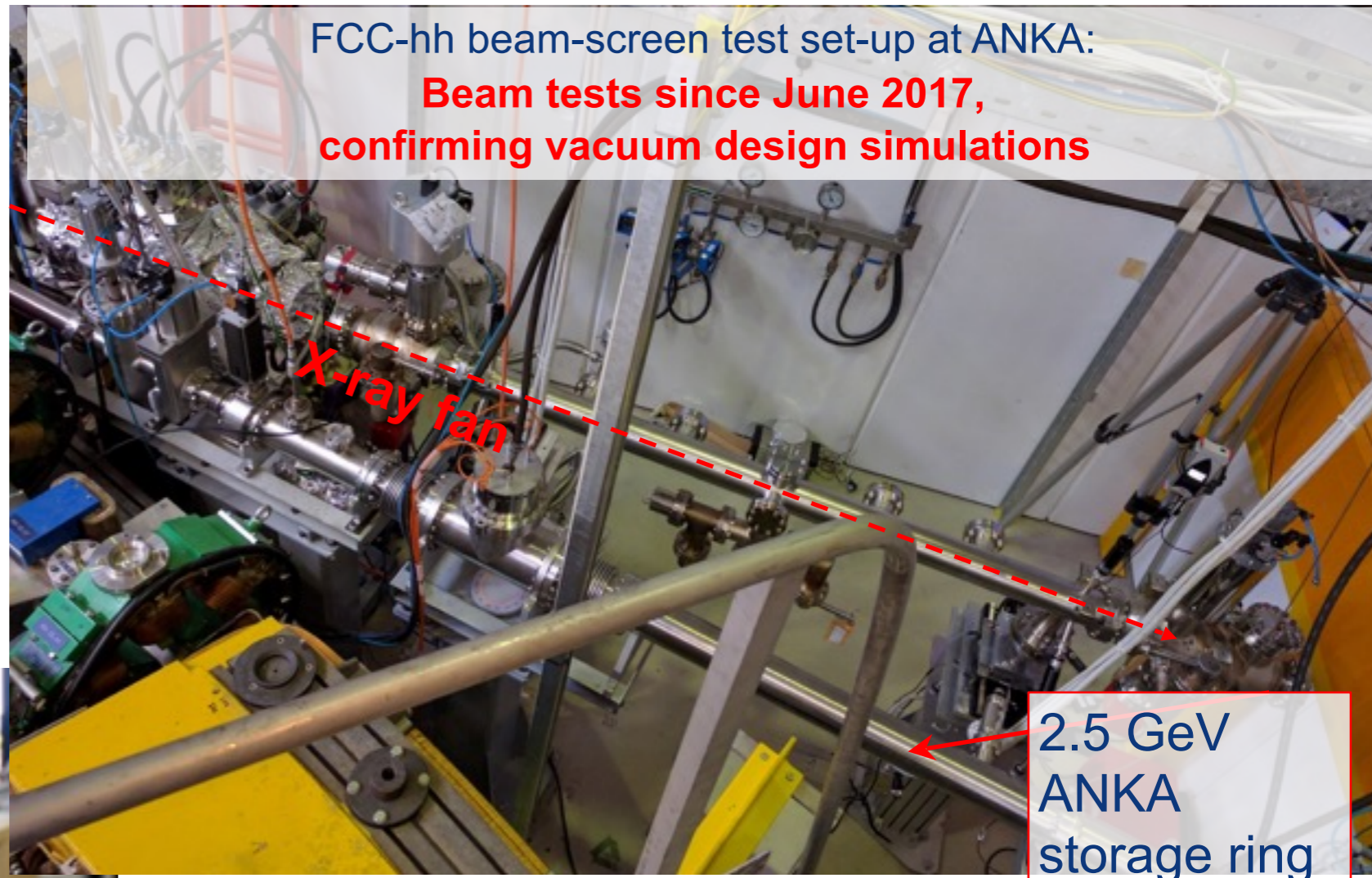
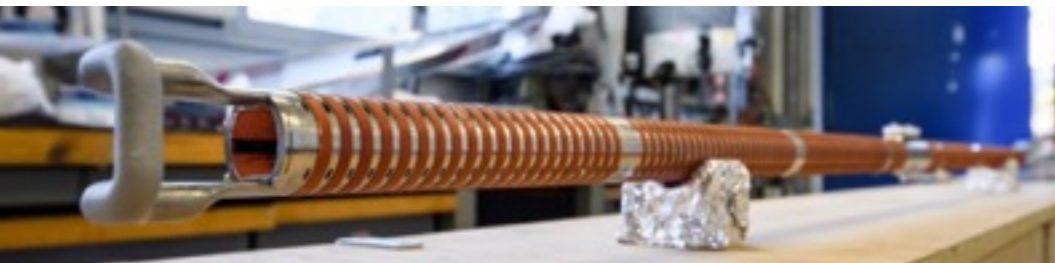
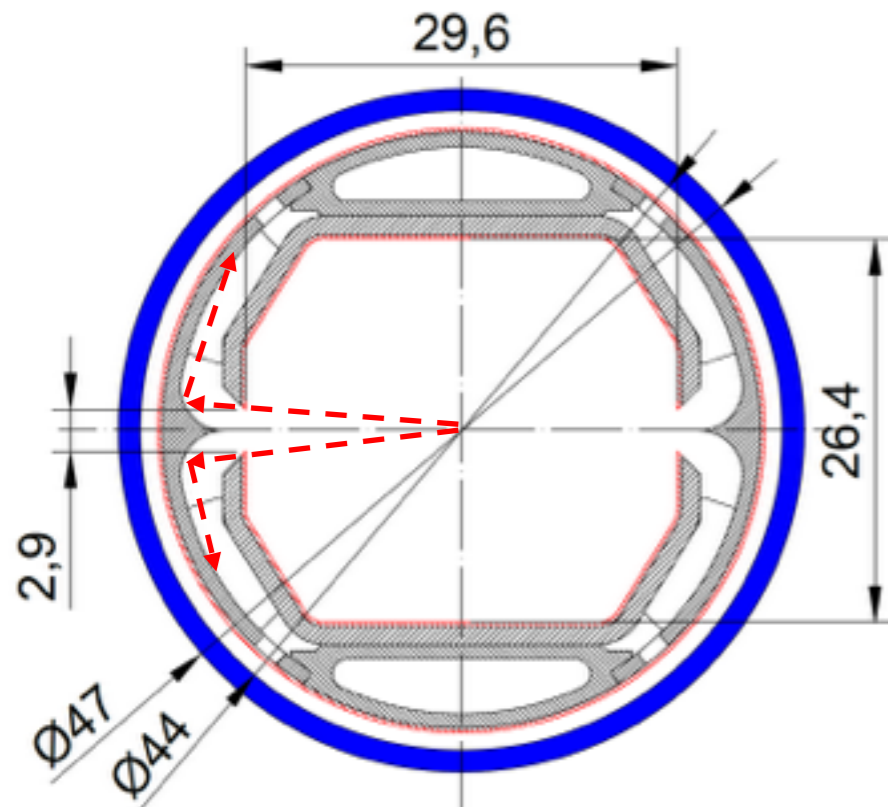
phase 1: $\beta^* = 1.1 \text{ m}$, $\xi_{\text{tot}} = 0.01$, $t_{\text{ta}} = 5 \text{ h}$, $250 \text{ fb}^{-1} / \text{year}$

phase 2: $\beta^* = 0.3 \text{ m}$, $\xi_{\text{tot}} = 0.03$, $t_{\text{ta}} = 4 \text{ h}$, $1000 \text{ fb}^{-1} / \text{year}$



FCC-hh cryogenic beam vacuum system

- **Synchrotron radiation** ($\sim 30 \text{ W/m/beam}$ (@16 T field) (LHC $< 0.2 \text{ W/m}$) $\sim 5 \text{ MW}$ total load in arcs
- **Absorption of synchrotron radiation at $\sim 50 \text{ K}$ for cryogenic efficiency (5 MW \rightarrow 100 MW cryoplant)**
- Provision of beam vacuum, suppression of photo-electrons, electron cloud effect, impedance, etc.



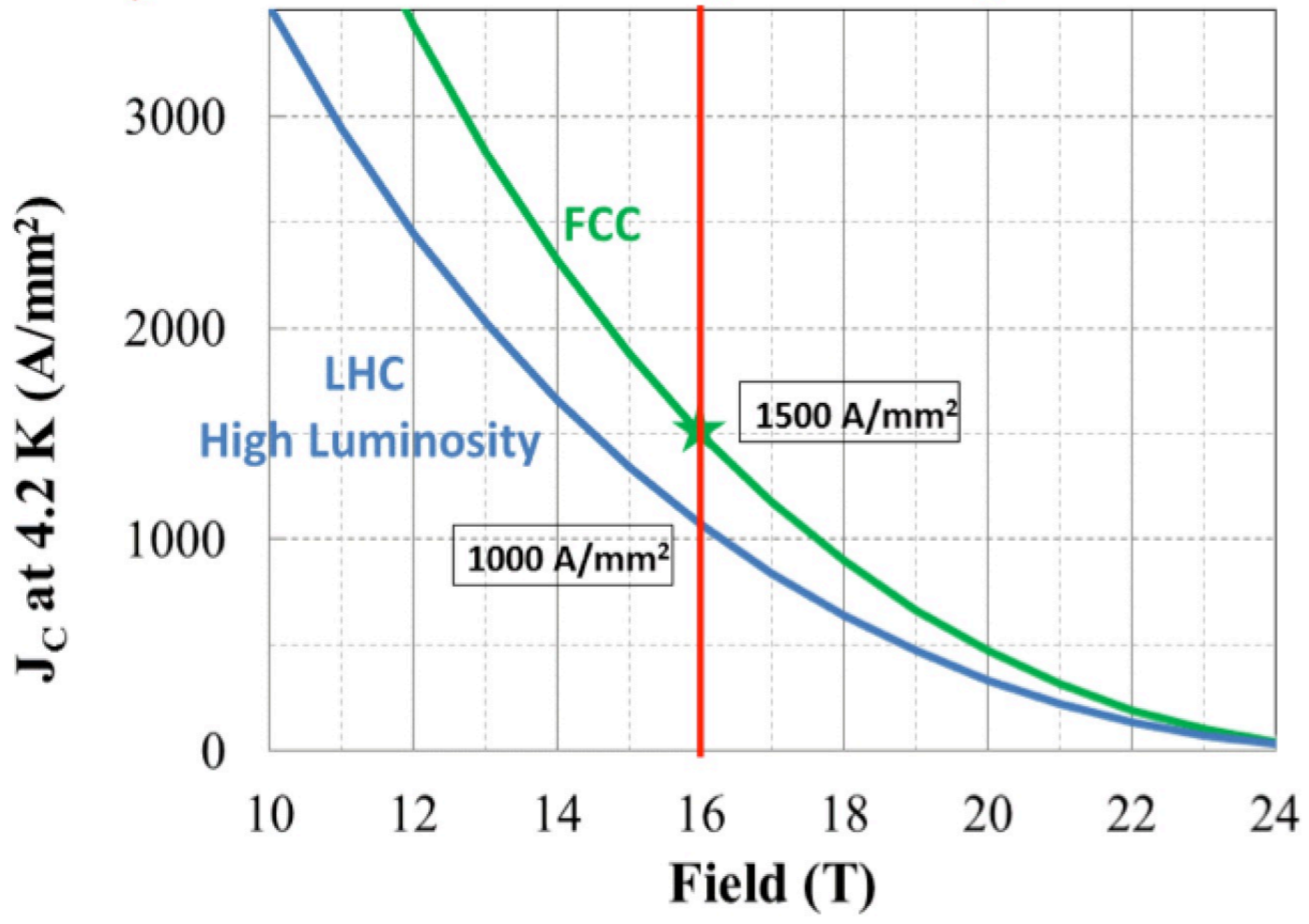
FCC-hh beam-screen test set-up at ANKA:
Beam tests since June 2017,
confirming vacuum design simulations

2.5 GeV
ANKA
storage ring



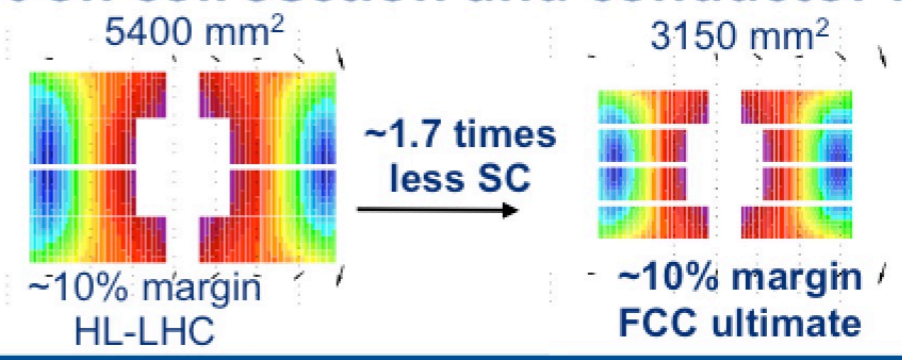
Nb₃Sn conductor development program

Nb₃Sn is one of the key cost & performance factors for FCC-hh / HE-LHC



- Main development goals:**
- J_c increase (16T, 4.2K) > 1500 A/mm² i.e. 50% increase wrt HL-LHC wire
 - Reference wire diameter 1 mm
 - Potentials for large-scale production and cost reduction

Impact on coil section and conductor mass





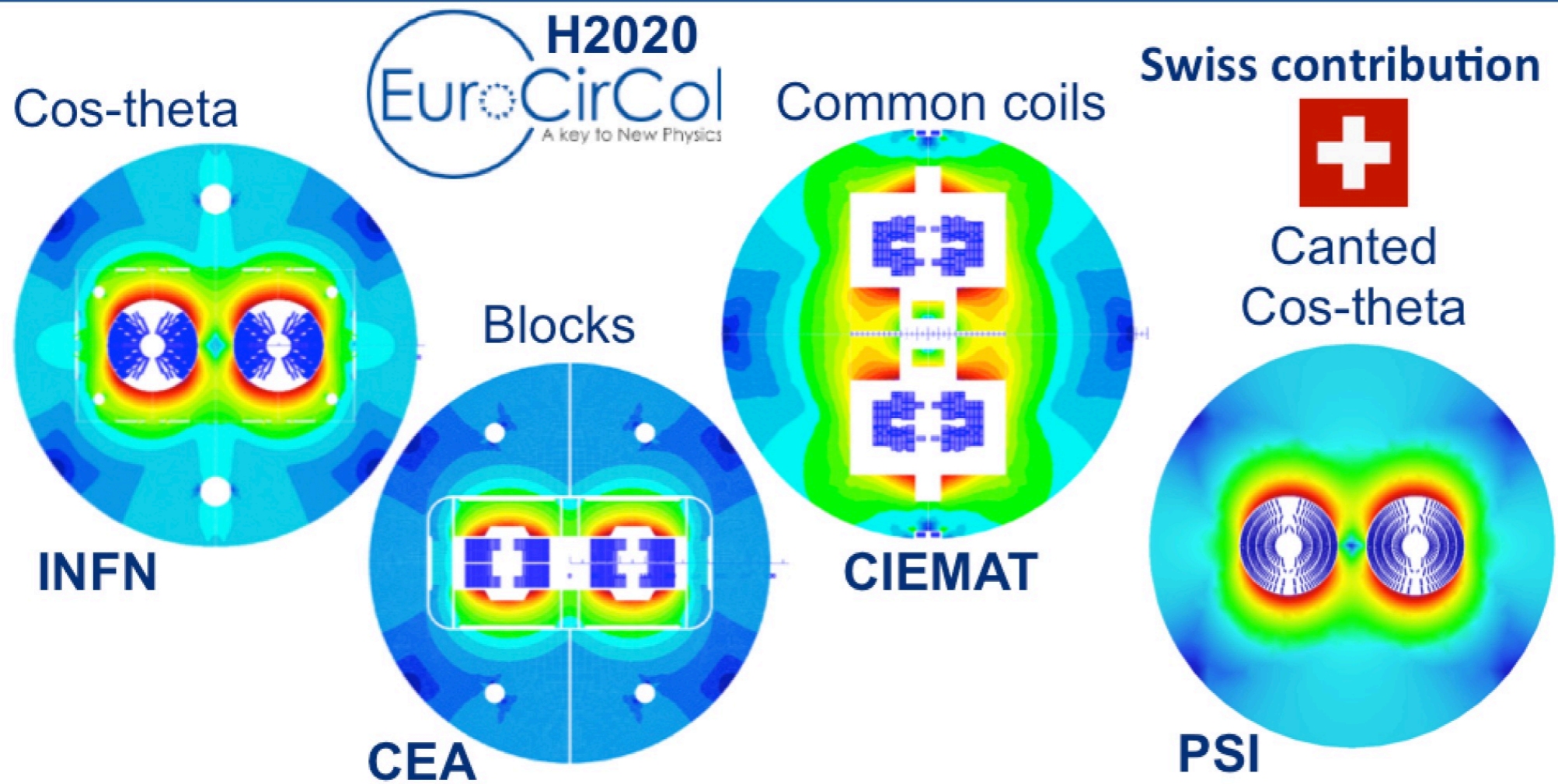
Collaborations FCC Nb₃Sn program

Established worldwide activities for Nb₃Sn development:

- **Procurement of state-of-the-art conductor for prototyping:**
 - **Bruker-OST** – **European/US**
- **Stimulation of conductor development with regional industry:**
 - **CERN/KEK** – **Japanese** contribution. Japanese **industry** (JASTEC, Furukawa, SH Copper) and laboratories (Tohoku Univ. and NIMS).
 - **CERN/Bochvar High-technology Research Inst.** – **Russian** contribution. Russian **industry** (TVEL) and laboratories
 - **CERN/KAT** – **Korean** industrial contribution
- **Characterization of conductor & research with universities:**
 - **Europe: Technical Univ. Vienna, Geneva University, University of Twente**
 - **Applied Superconductivity Centre at Florida State University**



16 T dipole design activities and options



The U.S. Magnet Development Program Plan

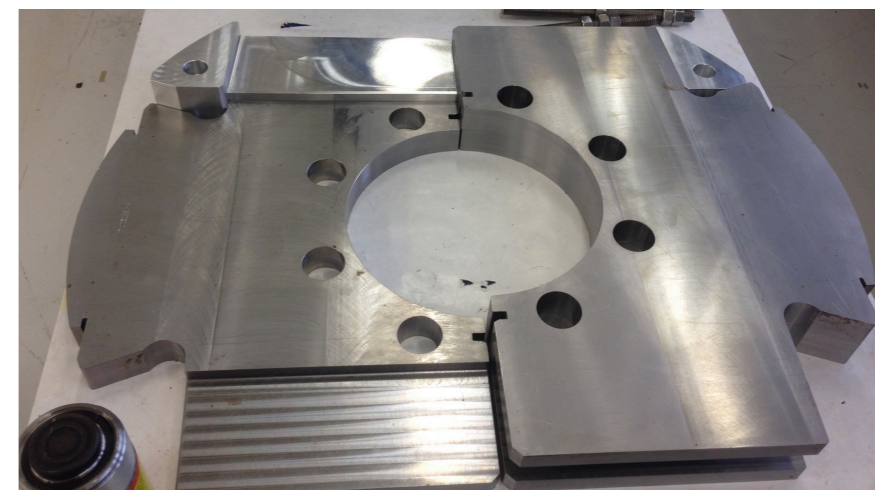
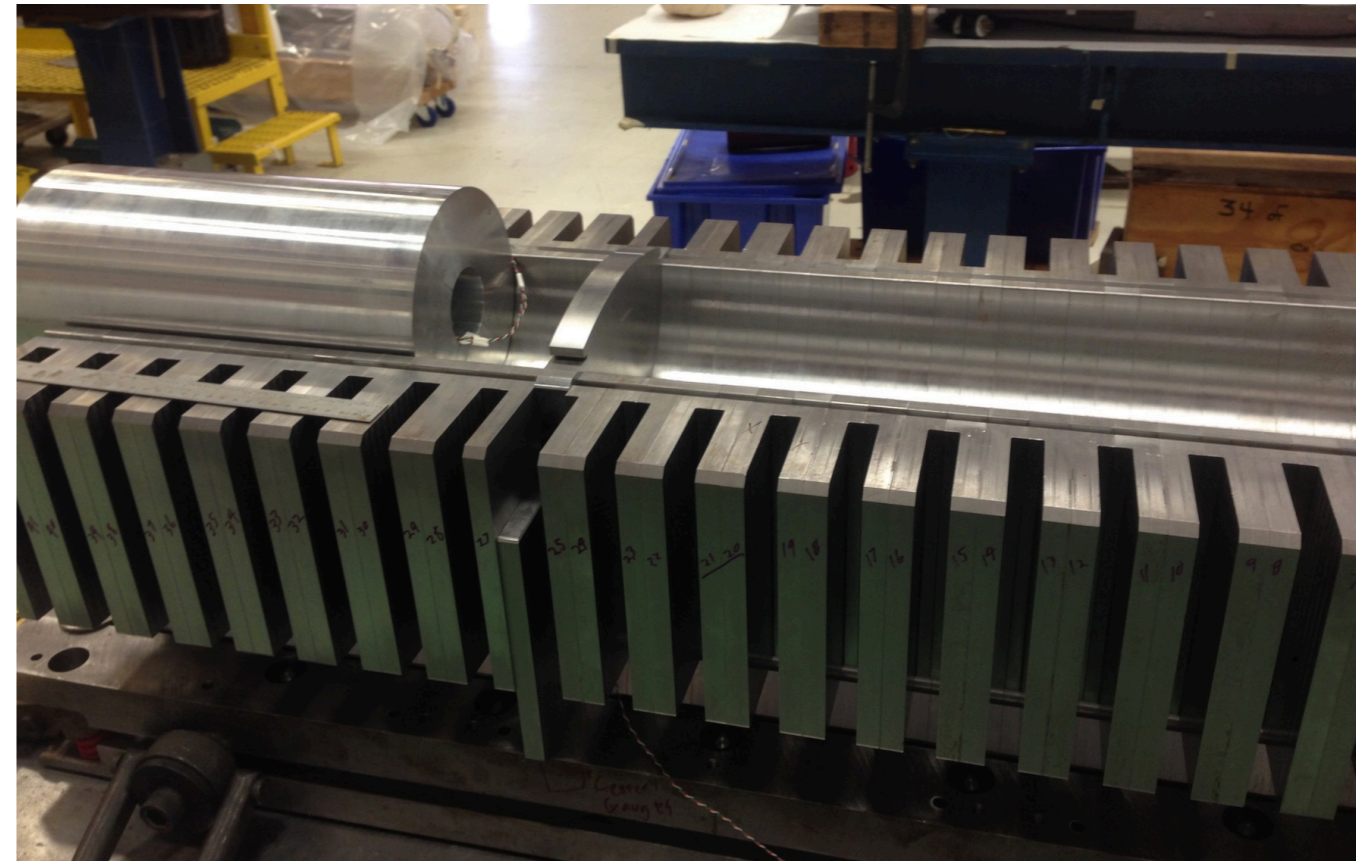
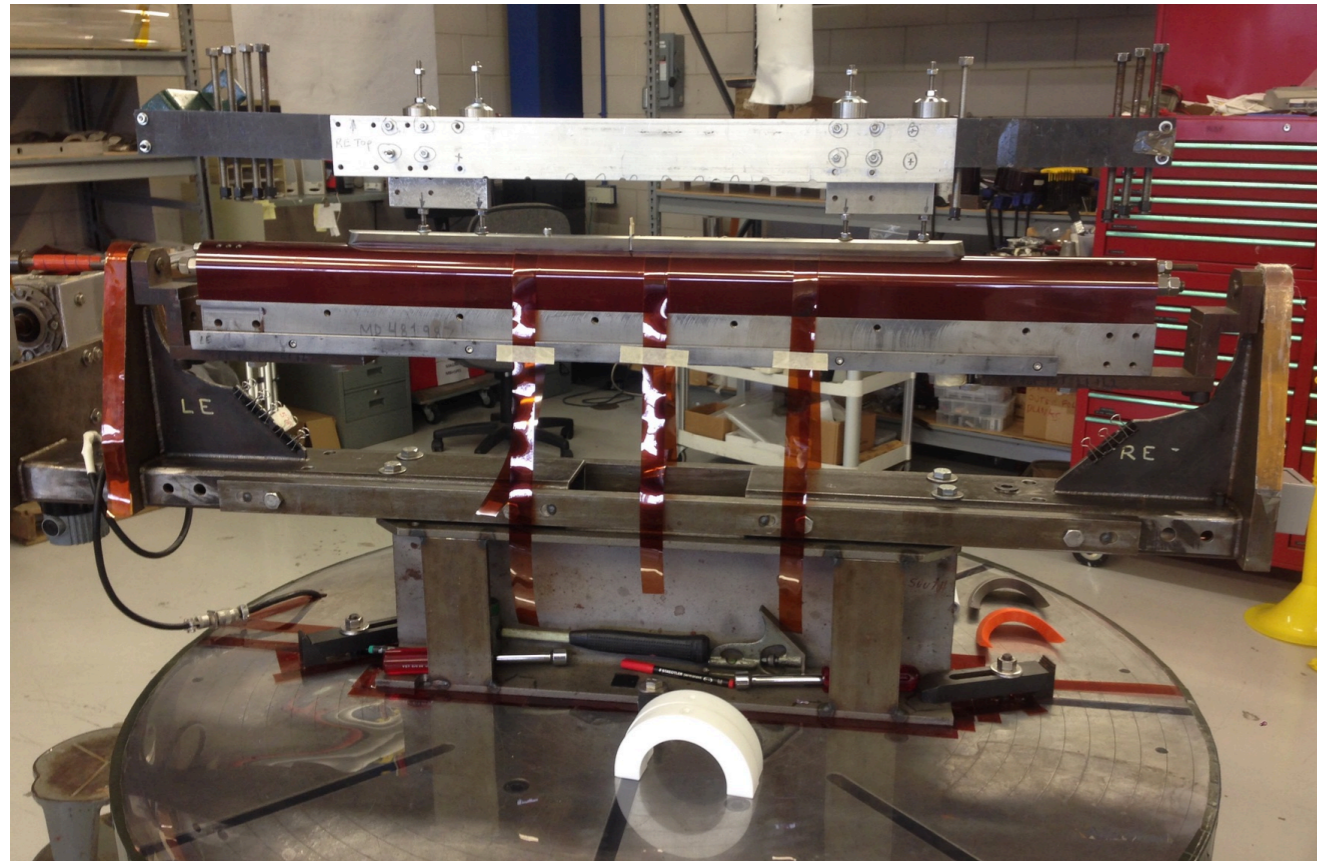
Intercepting ribs
Conductor
Spar
Shrinking Al tube

LBNL
FNAL

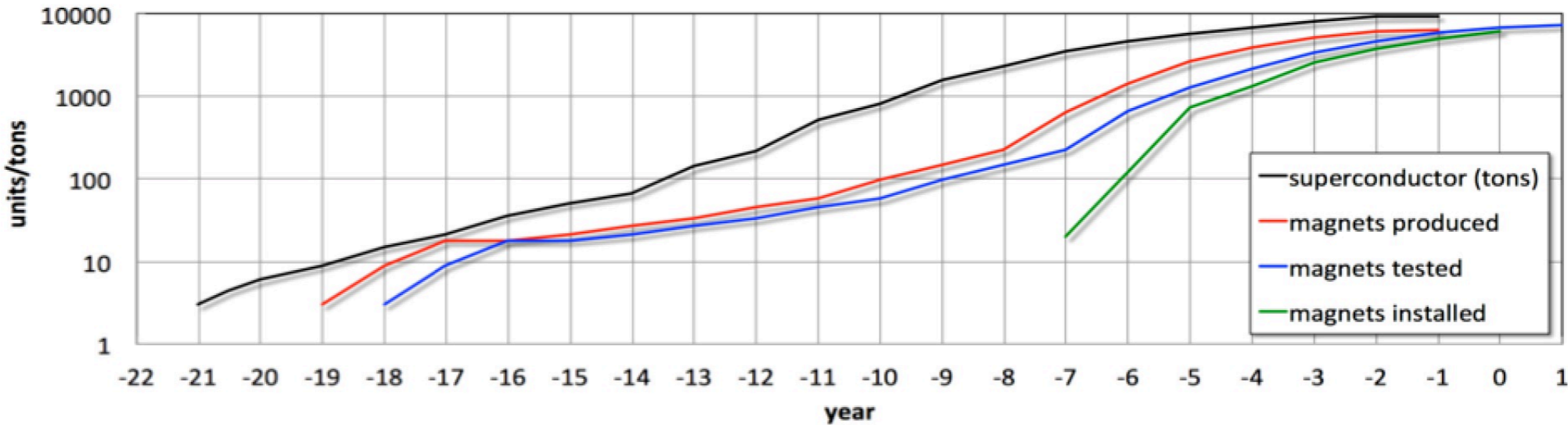
S. A. Gourlay, S. O. Prustemon
Lawrence Berkeley National Laboratory
Berkeley, CA 94720
A. V. Zlobin, L. Cooley
Fermi National Accelerator Laboratory

Short model magnets (1.5 m lengths) will be built from 2017 - 2021

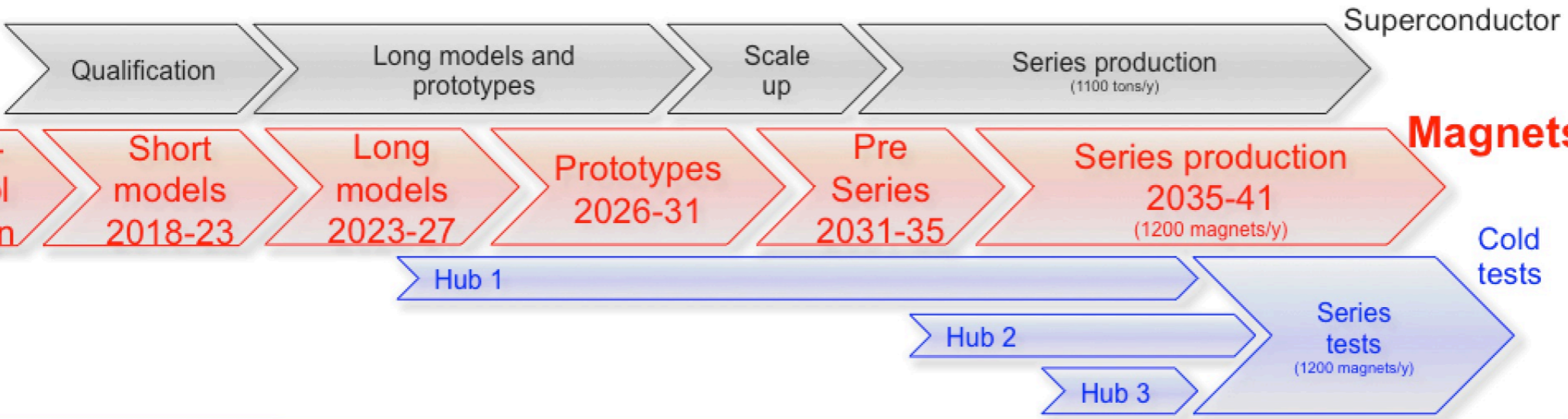
15T dipole prototyping at FNAL (60mm aperture, L=1m)



16 T magnet R&D schedule



Total duration of magnet program:
~20 years



Would follow on HL-LHC Nb_3Sn program with long models with industry from 2023/24

HE-LHC integration aspects

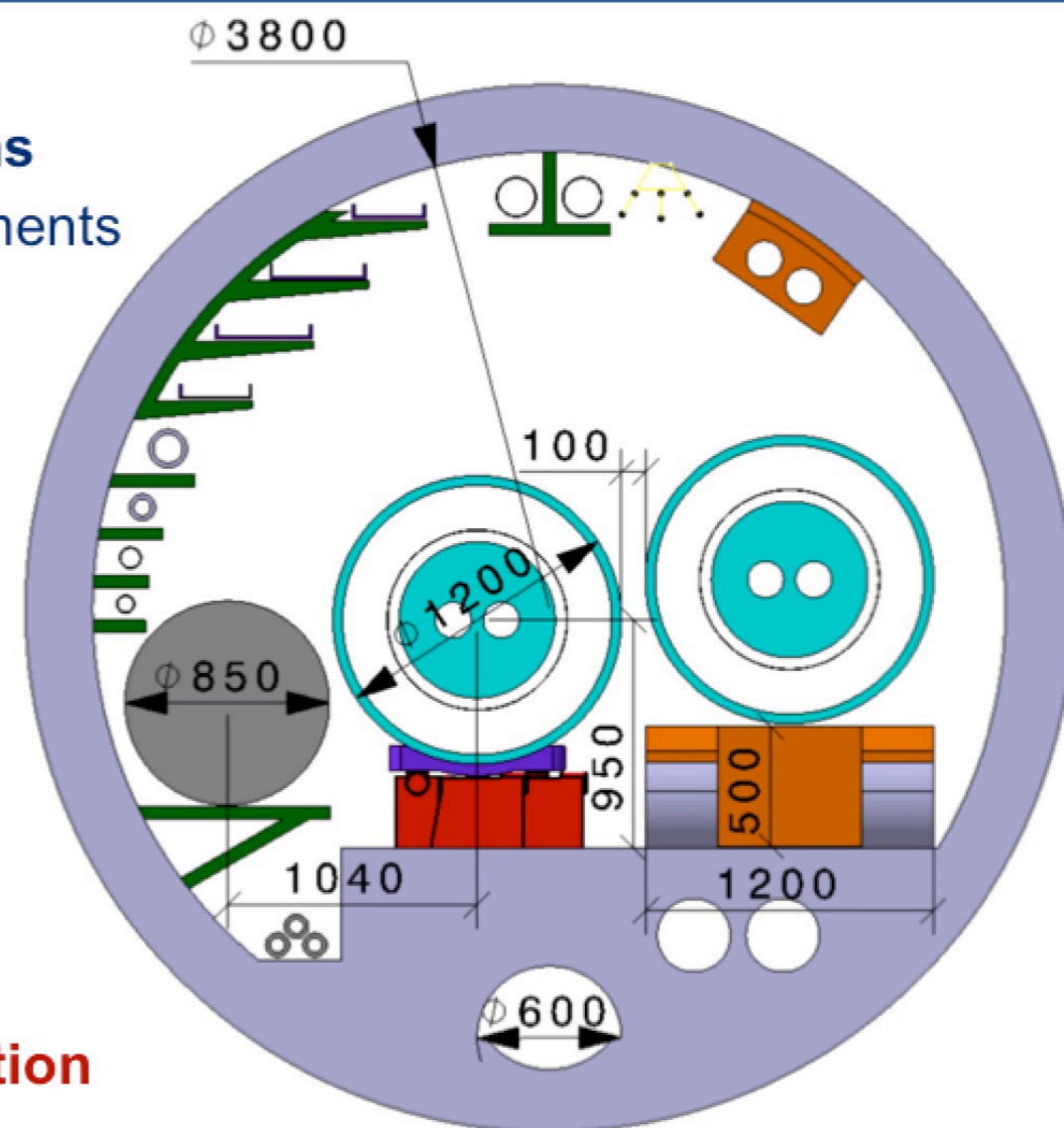
Working hypothesis for HE LHC design:

No major CE modifications on machine tunnel and caverns

- Similar geometry and layout as LHC machine and experiments
- **Maximum magnet cryostat external diameter compatible with LHC tunnel ~1200 mm**
- Classical 16 T cryostat design based on LHC approach gives ~1500 mm diameter!

Strategy: develop a single 16 T magnet, compatible with both HE LHC and FCC-hh requirements:

- Allow stray-field and/or cryostat as return-yoke
 - Optimization of inter-beam distance (compactness)
- **Smaller diam. also relevant for FCC-hh cost optimization**

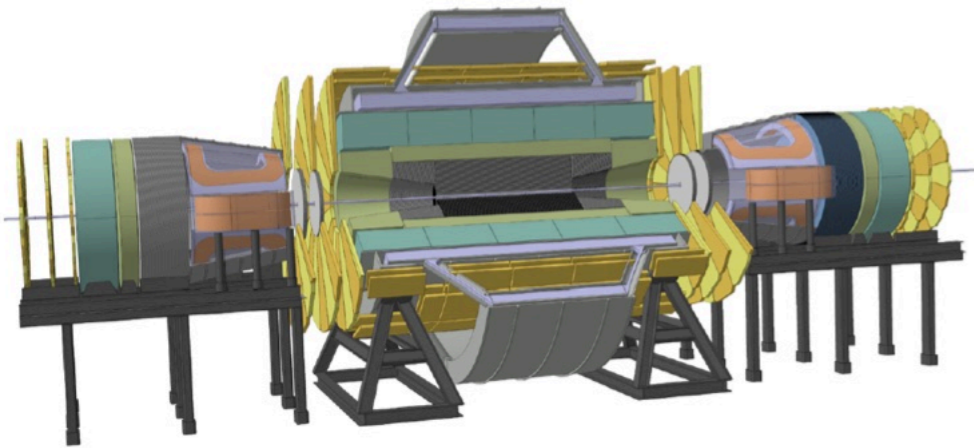


Reference detector

earlier design

6 T, 12 m bore solenoid, 10 Tm dipoles, shielding coil

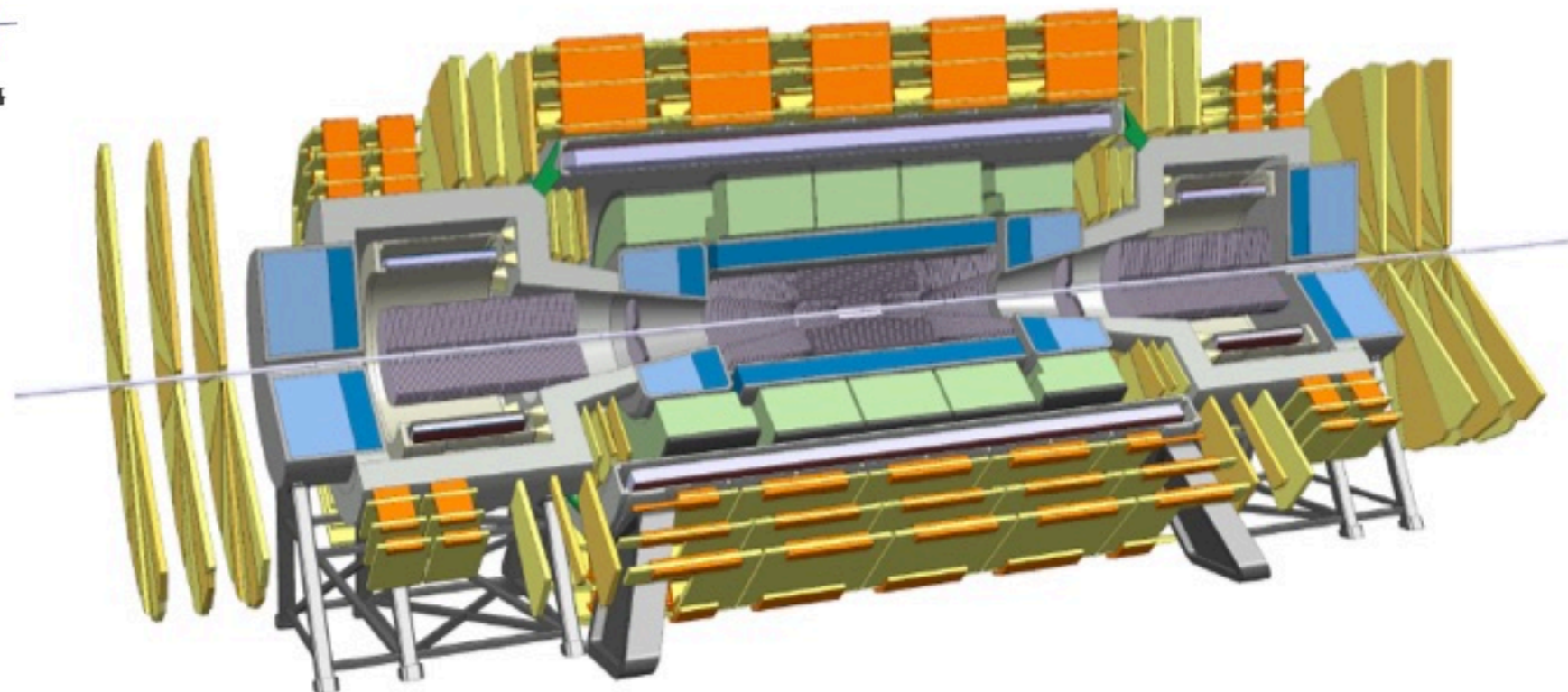
- 65 GJ stored energy
- 28 m diameter
- >30 m shaft
- multi billion project



current design

4 T, 10 m bore solenoid, 4 T forward solenoids, no shielding coil

- 14 GJ stored energy
- rotational symmetry for tracking!
- 20 m diameter (~ ATLAS)
- 15 m shaft
- ~1 billion project



latest $l^* = 40$ m

W. Riegler et al.

- **Detector design group leader: Werner Riegler**
 - Indico site of mtgs: <http://indico.cern.ch/category/8920/>
 - join the mailing list
- **Physics Simulation subgroup leaders: Heather Gray & Filip Moortgat**
 - Indico site of mtgs: <http://indico.cern.ch/category/6067/>
 - join the mailing list
- Monthly mtgs of each group, if interested register to the mailing lists

Final remarks

- FCC-hh physics studies today focus on exploring possible scenarios, assessing the physics potential, defining benchmarks for the accelerator and detector design and performance, in order to better inform the discussions that will take place when the time for decisions comes...
- The interplay of the three colliders (ee, eh and hh) is crucial to the full exploitation of the FCC physics potential
- The physics case of a 100 TeV collider is very clear as a long-term goal for the field, simply because no other proposed or foreseeable project can have direct sensitivity to such large mass scales.
- Nevertheless, the precise route followed to get there must take account of the fuller picture, to reflect the future data (and the impact they will have on the theoretical thinking) from the LHC, as well as other current and future experiments in areas ranging from flavour physics to searches for dark matter, axions, ALPs,