

# THE HIGGS MASS, TOP PARTNERS, AND COLLIDERS: WHAT WE LOSE BY LETTING GO

Timothy Cohen

Princeton University  
Institute for Advanced Study

Fermilab Colloquium

March 4, 2015

# HIGGSDEPENDENCE DAY!

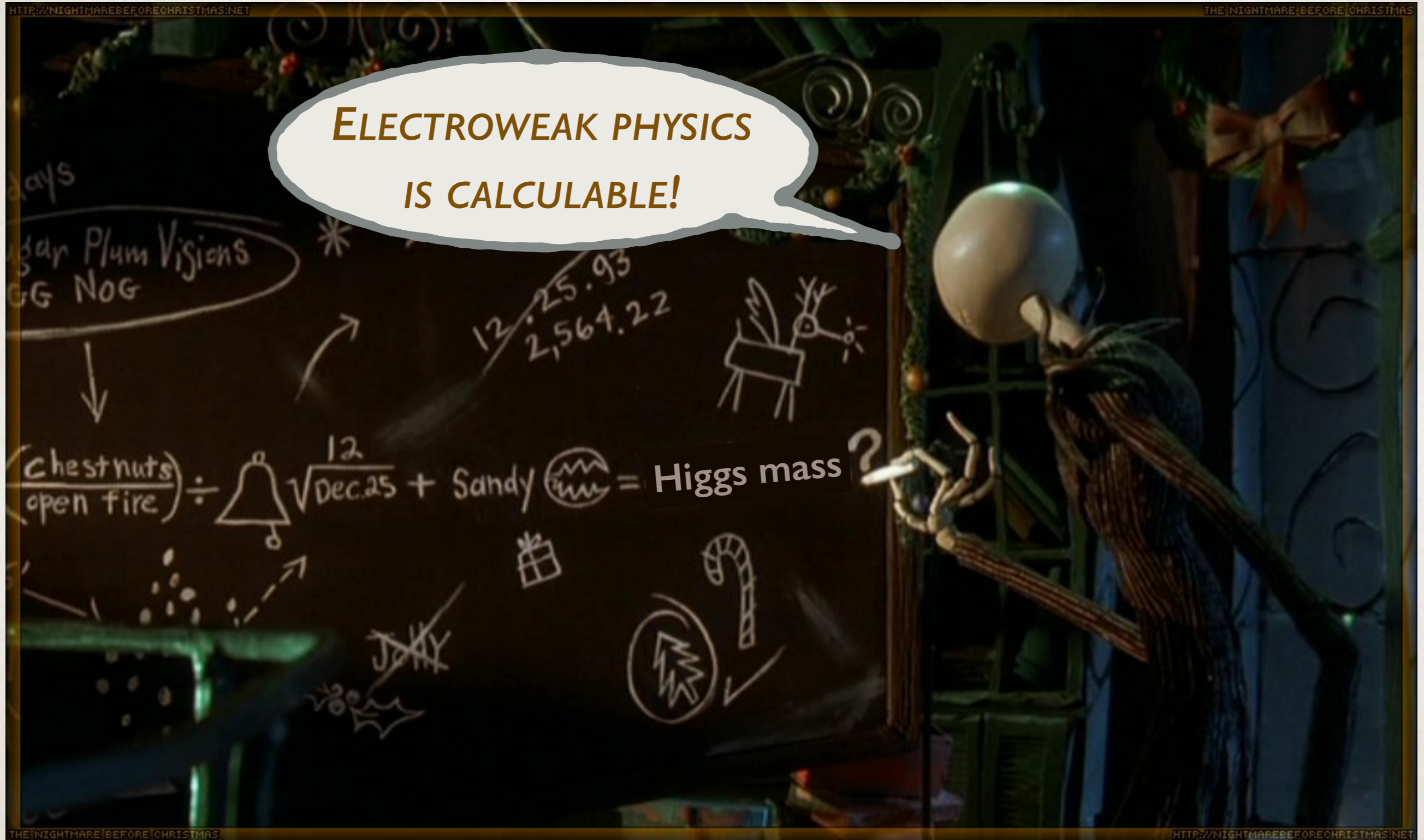


Most important lessons:

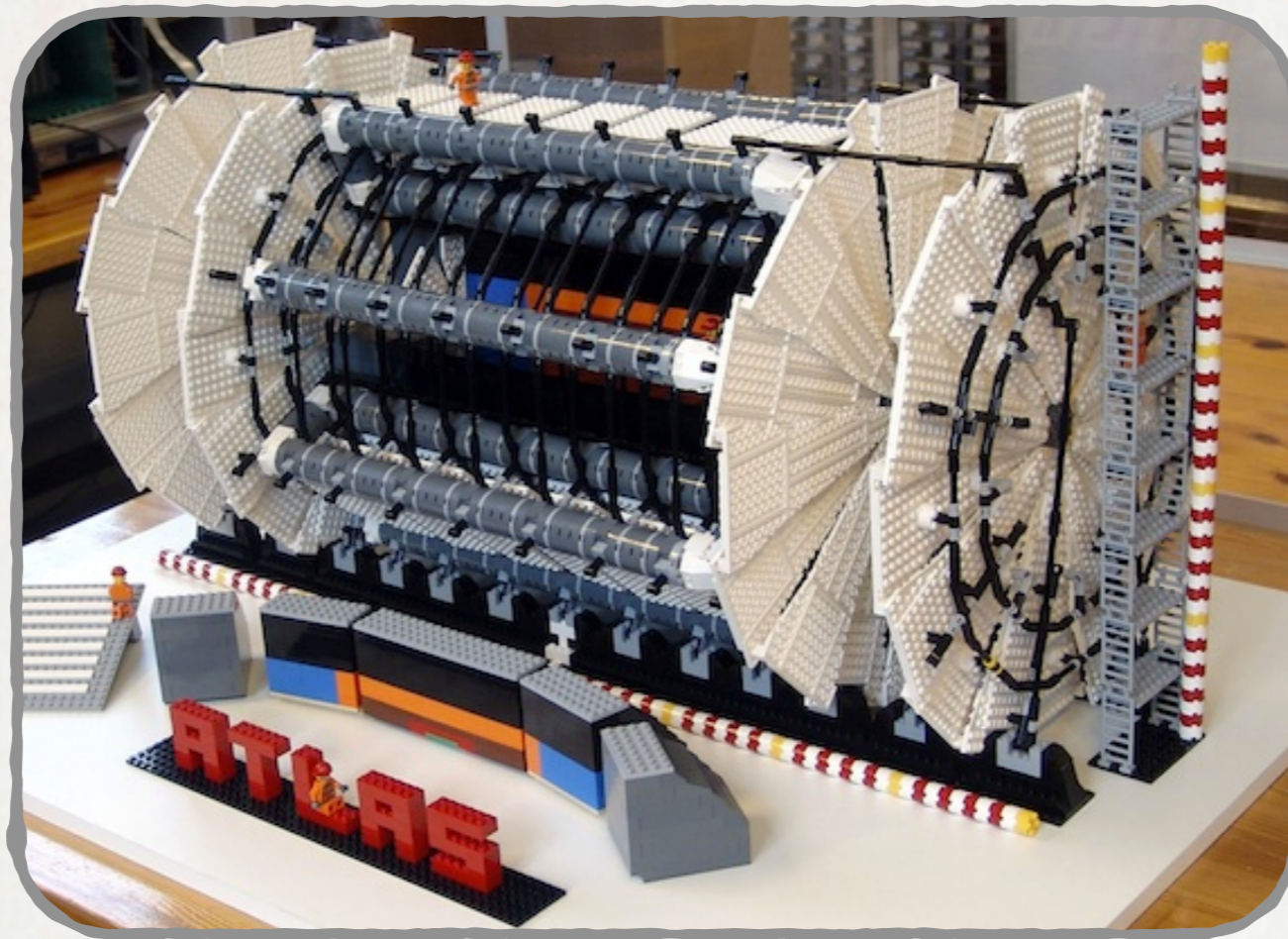
- Higgs boson exists!
- Weakly coupled!!

# HIGGSDEPENDENCE DAY!

**ELECTROWEAK PHYSICS  
IS CALCULABLE!**



# THE STANDARD MODEL IS A MODEL



Set the low energy Higgs mass and forget it.

Make predictions.

No conflict with experiment (so far).

# BRIEF SIDEBAR

# DIMENSIONFUL SCALES

## Relativity:

Space  $\equiv$  Time

Convert meters to seconds with speed of light:  $c$ .

Mass  $\equiv$  Energy

Convert mass to energy with speed of light:  $c$ .

# DIMENSIONFUL SCALES

## Relativity:

Space  $\equiv$  Time

Convert meters to seconds with speed of light:  $c$ .

Mass  $\equiv$  Energy

Convert mass to energy with speed of light:  $c$ .

---

## Quantum Mechanics:

Energy  $\equiv$  1/Time

Convert energy to time with Planck's constant:  $\hbar$ .

# DIMENSIONFUL SCALES

## Relativity:

Space  $\equiv$  Time

Convert meters to seconds with speed of light:  $c$ .

Mass  $\equiv$  Energy

Convert mass to energy with speed of light:  $c$ .

## Quantum Mechanics:

Energy  $\equiv$  1/Time

Convert energy to time with Planck's constant:  $\hbar$ .

Large mass scales are like short distance scales.

The LHC is a giant microscope!



END SIDEBAR

# REDUCTIONISM

What does it mean to have a *theory* for the Higgs mass?

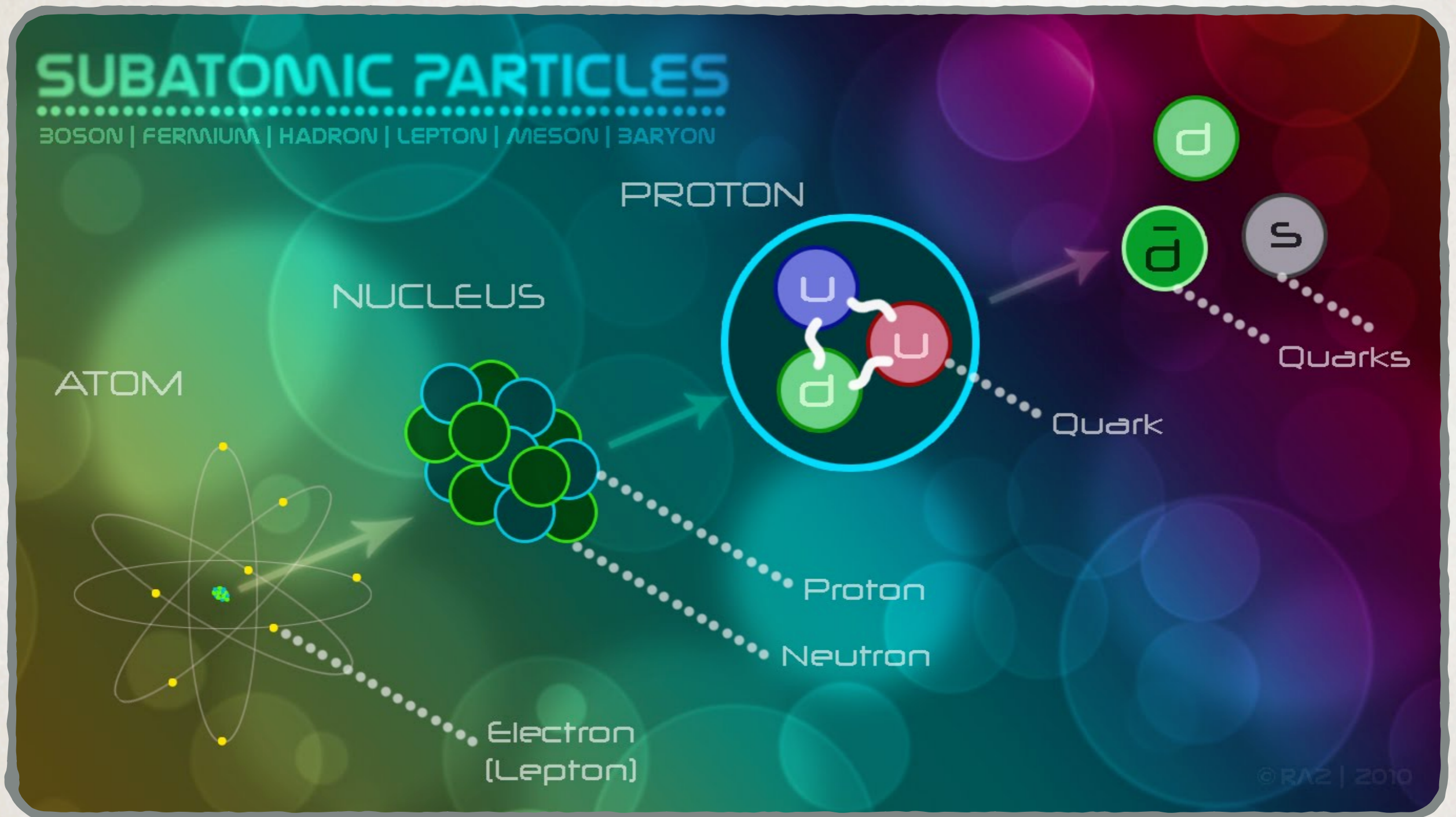
Higgs mass is a function of well defined (finite) inputs.  
The Standard Model encompassed by larger framework.

***Naively: Higgs mass quadratically sensitive to new mass scales.***

**Historical precedent for reductionism:**

Underlies progress in fundamental physics.  
New frameworks encompass the old, giving “reasons”.

# REDUCTIONISM



# NEW SCALES?

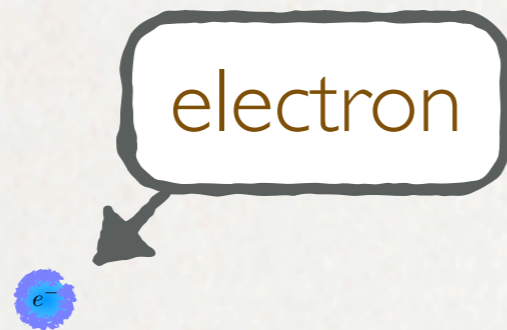
Evidence for physics beyond the Standard Model:

- dark matter
- matter/anti-matter asymmetry
- neutrino masses
- gravity

Very likely new dimensionful scale exists.

Want to protect the Higgs from physics at high energy scales.

# ELECTRON SELF-ENERGY



# ELECTRON SELF-ENERGY



CLASSICAL

$$\Delta m_e \sim \frac{e^2}{r_e}$$

“electron  
radius”

$$\frac{\Delta m_e}{m_e} \sim 10^9 \left( \frac{r_{\text{Planck}}}{r_e} \right)$$

# ELECTRON SELF-ENERGY

CLASSICAL

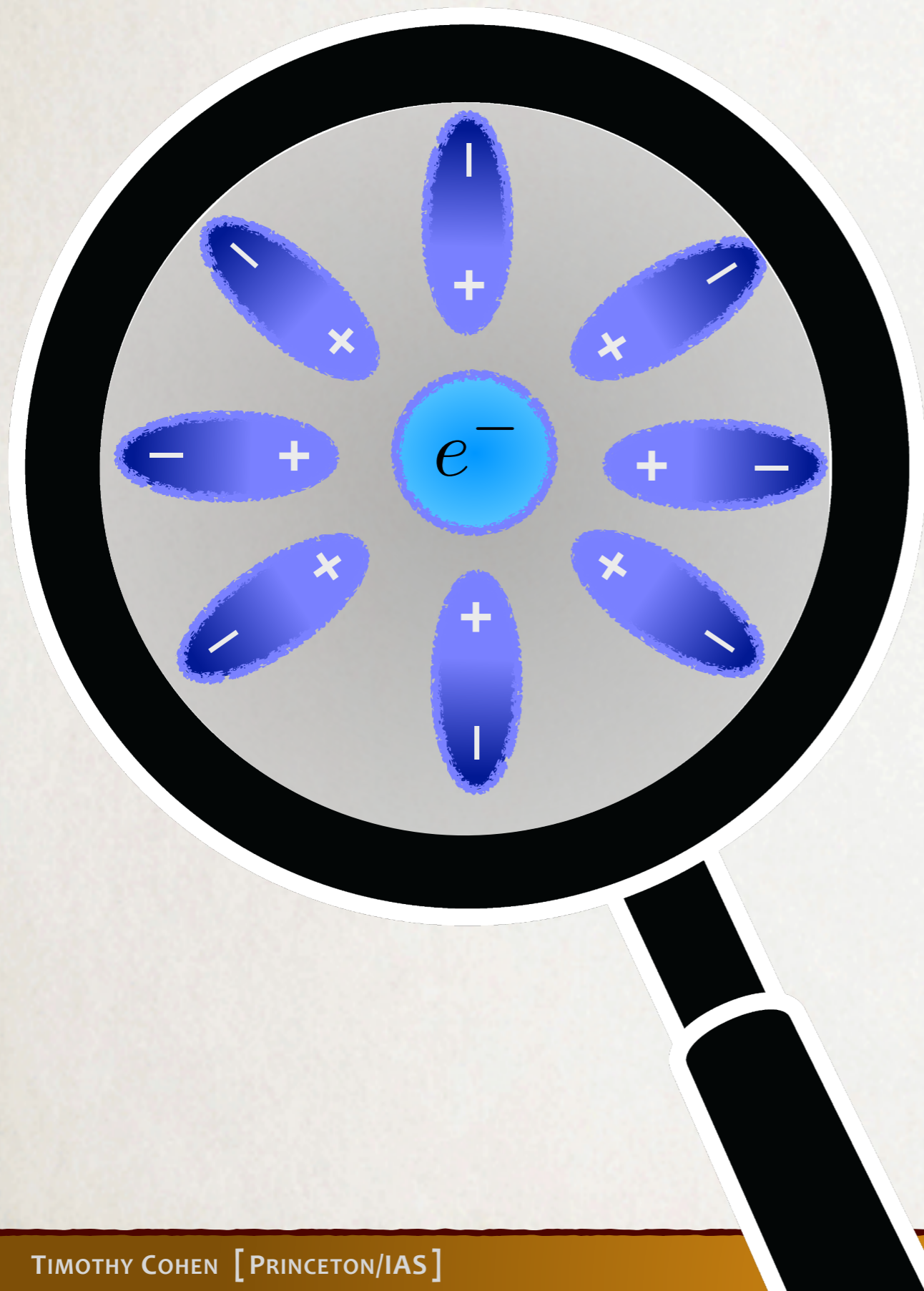
$$\Delta m_e \sim \frac{e^2}{r_e}$$

“electron  
radius”

$$\frac{\Delta m_e}{m_e} \sim 10^9 \left( \frac{r_{\text{Planck}}}{r_e} \right)$$



# ELECTRON SELF-ENERGY



CLASSICAL

$$\Delta m_e \sim \frac{e^2}{r_e}$$

“electron radius”

$$\frac{\Delta m_e}{m_e} \sim 10^9 \left( \frac{1}{e} \frac{\hbar}{k} \right)$$

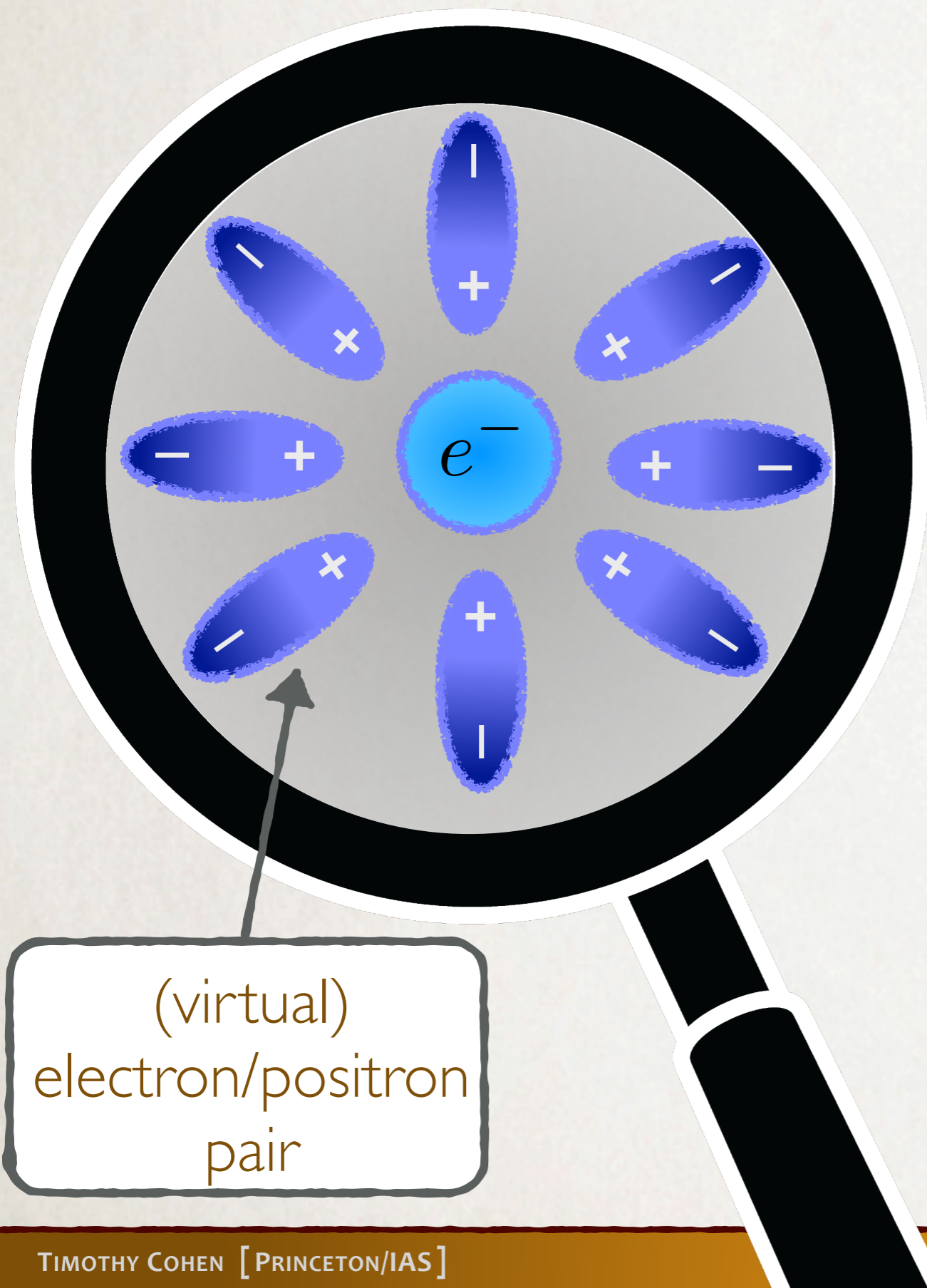
## QUANTUM

$$\Delta m_e \sim m_e \frac{e^2}{16\pi^2} \log(m_e r_e)$$

$$\frac{\Delta m_e}{m_e} \sim 10^{-2}$$



# ELECTRON SELF-ENERGY



(virtual)  
electron/positron  
pair

~~CLASSICAL~~

$$\Delta m_e \sim \frac{e^2}{r_e}$$

“electron radius”

$$\frac{\Delta m_e}{m_e} \sim 10^9 \left( \frac{1}{e} \frac{\hbar c}{e} \right)$$

## QUANTUM

$$\Delta m_e \sim m_e \frac{e^2}{16\pi^2} \log(m_e r_e)$$

$$\frac{\Delta m_e}{m_e} \sim 10^{-2}$$

**Requires new particle!**

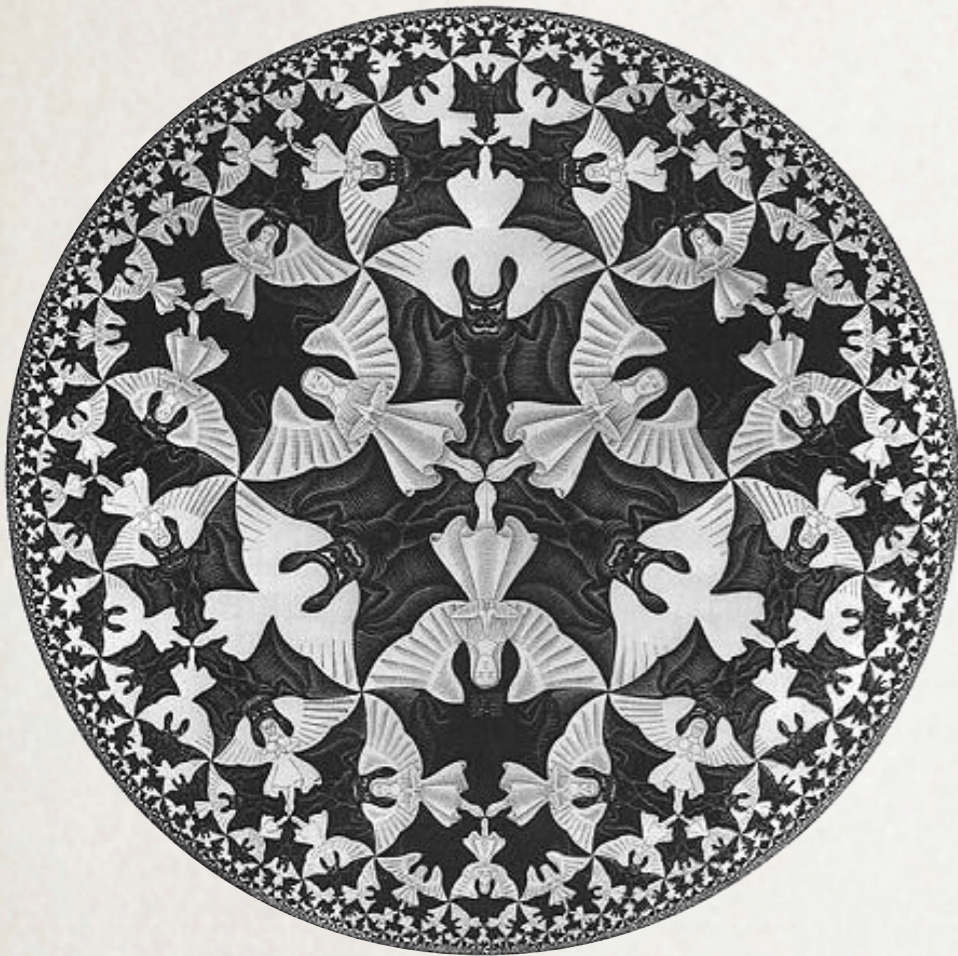
# OUTLINE

- I. The Higgs and its Potential
- II. Supersymmetry: An Example Theory
- III. SUSY Naturalness Confronts Experiment
- IV. Alternative Theories
- V. Summary

# I. THE HIGGS AND ITS POTENTIAL

# CHARGED SCALAR

Introduce a charged scalar state:  $\phi$ .



“ $\phi$  carries a charge”

is equivalent to

$$\phi \rightarrow \phi' = e^{i q_\phi \xi} \phi$$

under a symmetry (gauge) transformation.

“Charge is conserved”

is equivalent to

Lagrangian  $\mathcal{L}$  invariant under transformation.

$\phi$  and  $\phi^*$  have opposite charge.

# SCALAR MASS

Scalar mass always phase rotation invariant:

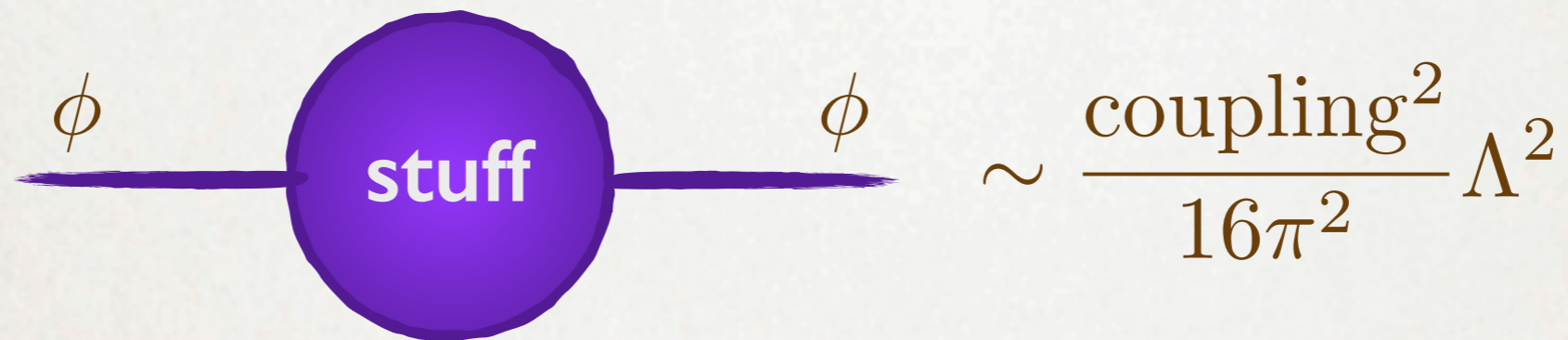
$$\mathcal{L} \supset m^2 |\phi|^2$$

No phase rotation can forbid mass.

***Anything can happen in quantum mechanics!***

Implication:

If  $\phi$  interacts, quantum corrections can generate a mass.



# SCALAR MASS

Scalar mass always phase rotation invariant:

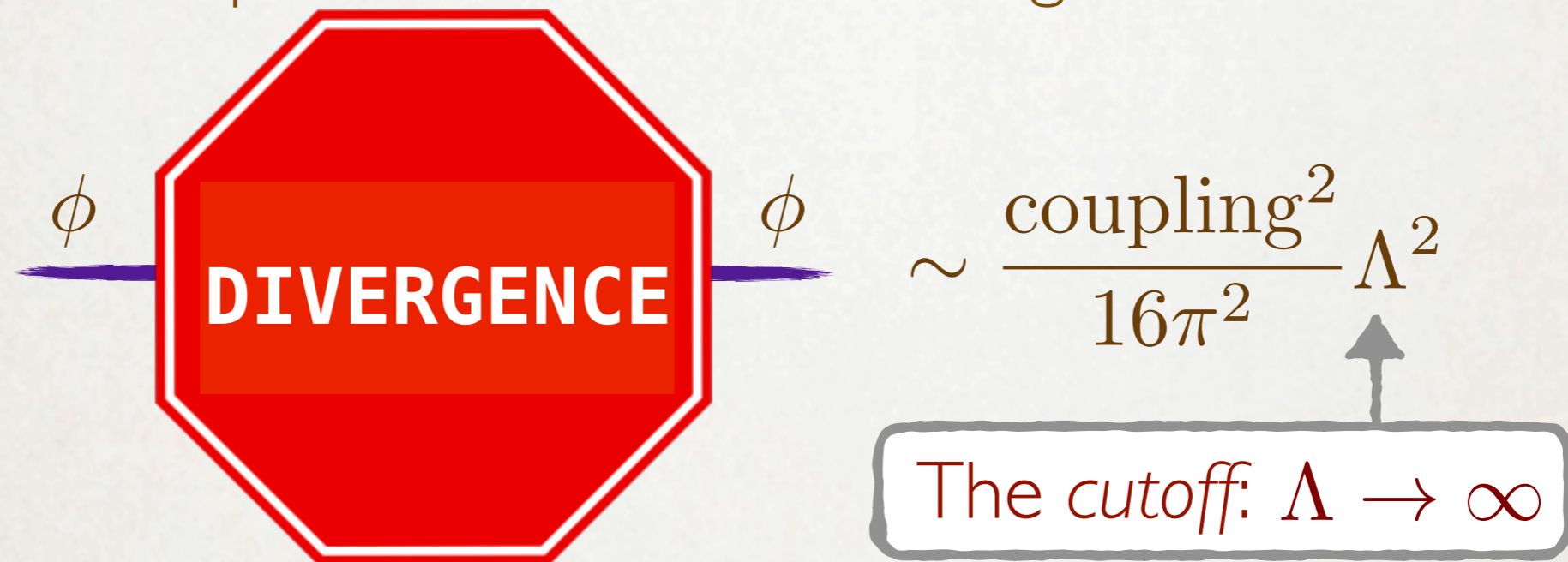
$$\mathcal{L} \supset m^2 |\phi|^2$$

No phase rotation can forbid mass.

**Anything can happen in quantum mechanics!**

Implication:

If  $\phi$  interacts, quantum corrections can generate a mass.



Problem for Higgs boson!

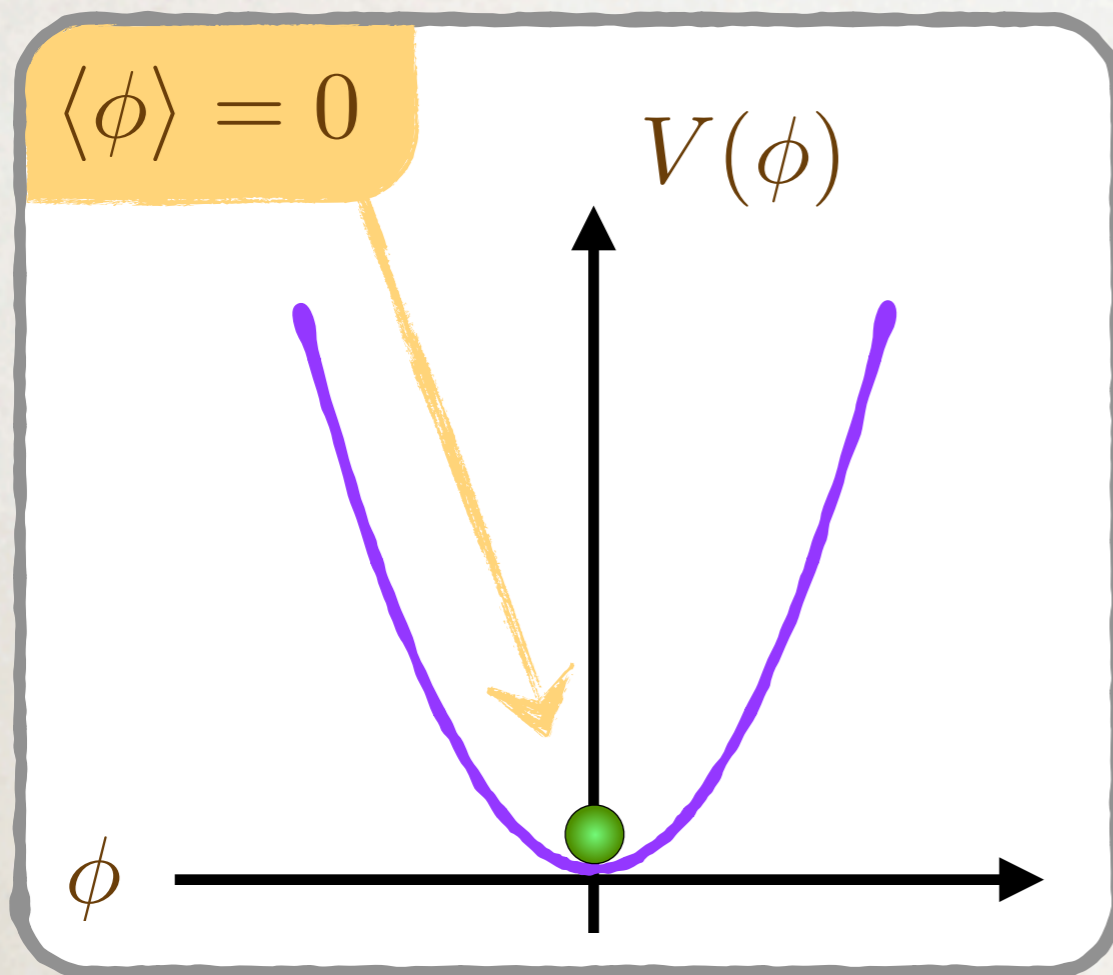
# SCALAR VEV

Add quartic for scalar (phase rotation invariant):

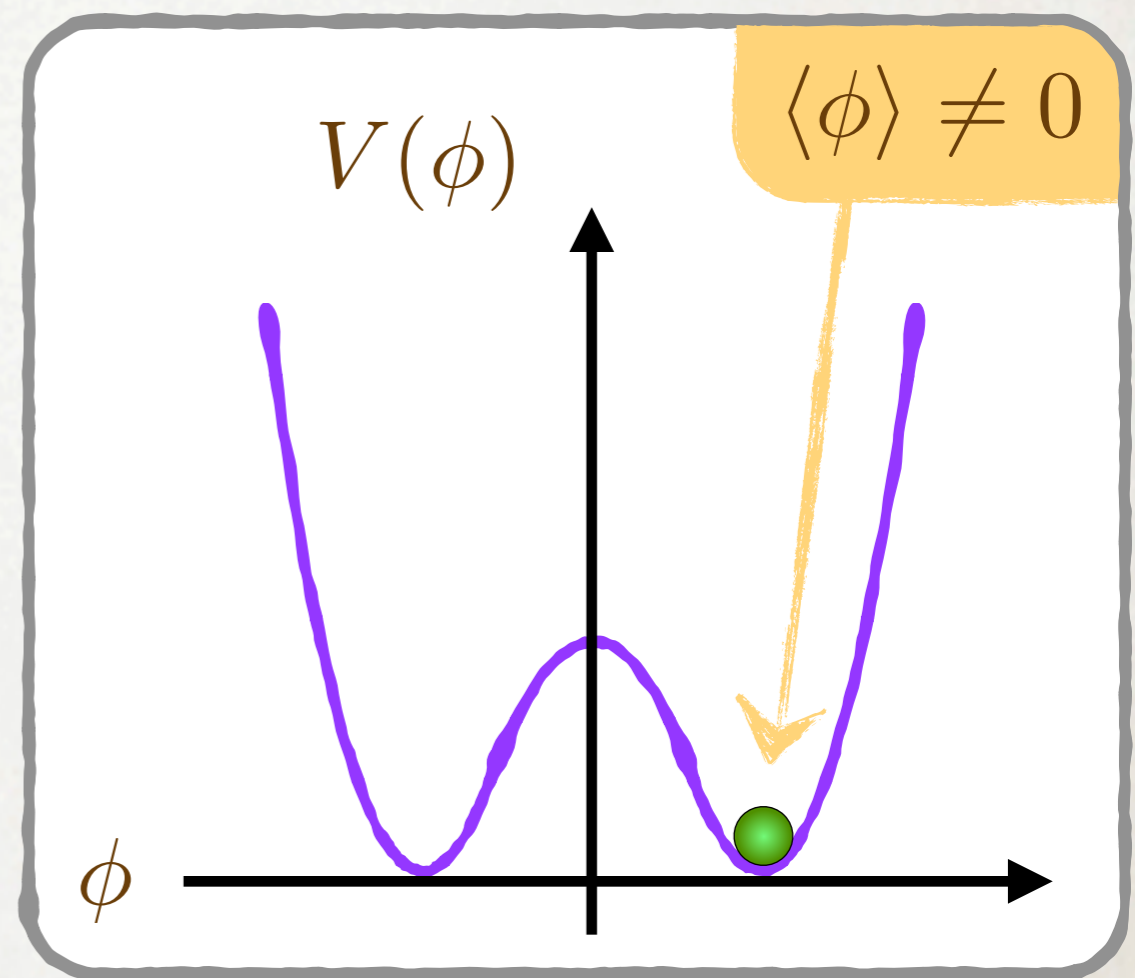
$$\mathcal{L} \supset \underbrace{m^2 |\phi|^2 + \frac{\lambda}{4} |\phi|^4}$$

The “potential”  $V(\phi)$ .

$$m^2 > 0$$

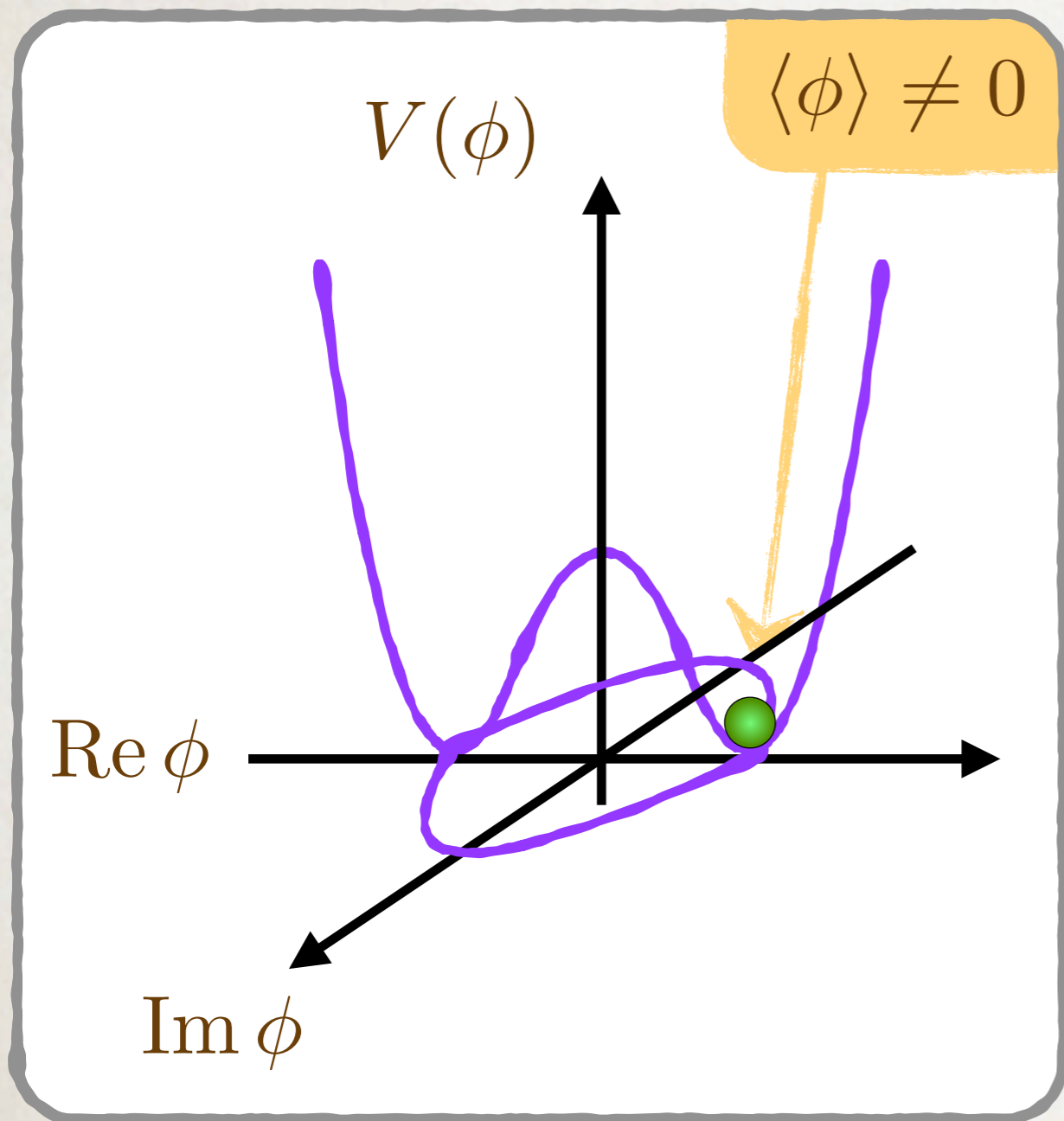


$$m^2 < 0$$



# VACUUM EXPECTATION VALUE

$\langle \phi \rangle$  : vacuum expectation value (vev).



Vacuum “sees” the phase.

The vev spontaneously breaks the symmetry.

This is how the Higgs vev breaks electroweak symmetry.



# THE HIGGS BOSON

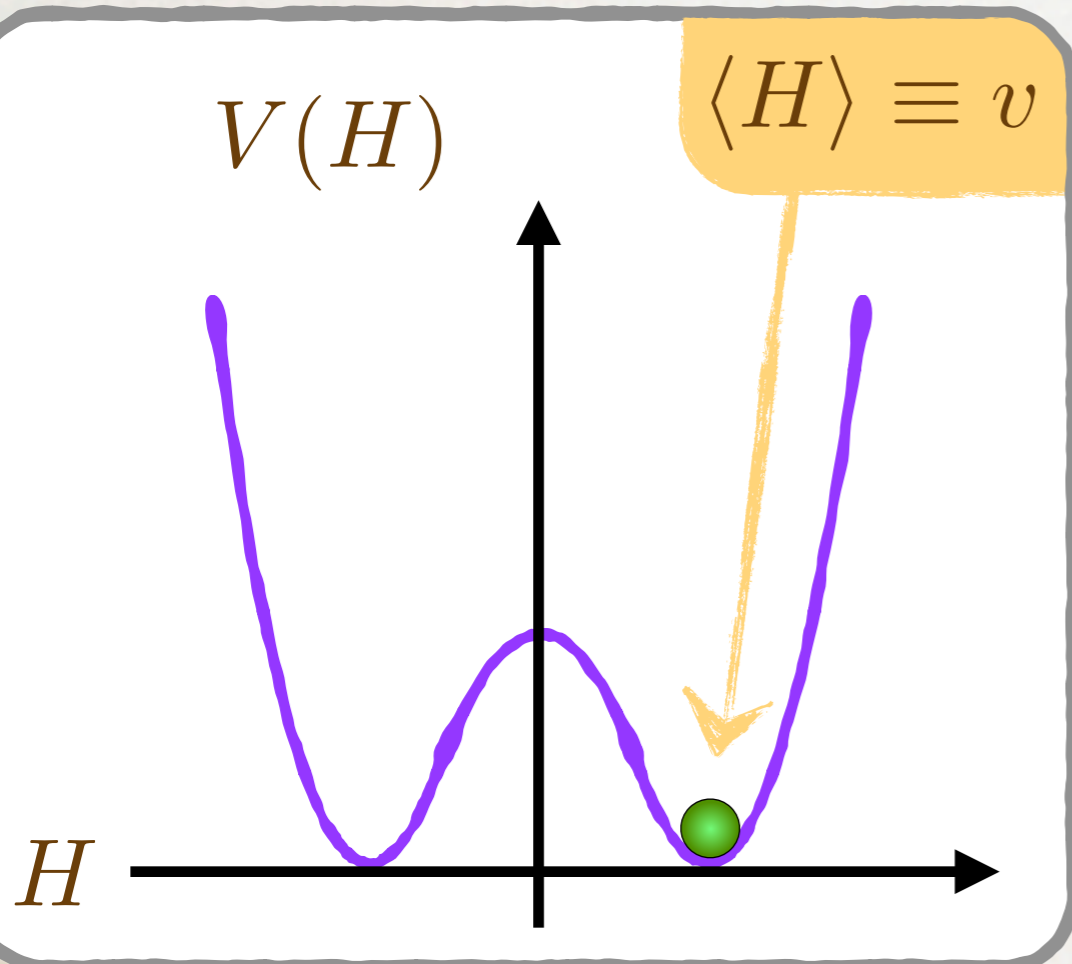
$$\phi \rightarrow H$$

# HIGGS POTENTIAL

$$V(v, H) = -\frac{\mu^2}{2} |v + H|^2 + \frac{\lambda_H}{16} |v + H|^4$$

Solving  $\frac{\partial V(v, 0)}{\partial v} = 0 \Rightarrow v^2 = \frac{2\mu^2}{\lambda_H}$

Solving  $\frac{\partial^2 V(v, 0)}{\partial v^2} = m_H^2 \Rightarrow m_H^2 = 2\mu^2$



$W^\pm$  mass  $\Rightarrow v \simeq 246$  GeV

Higgs mass  $\Rightarrow m_H \simeq 125$  GeV

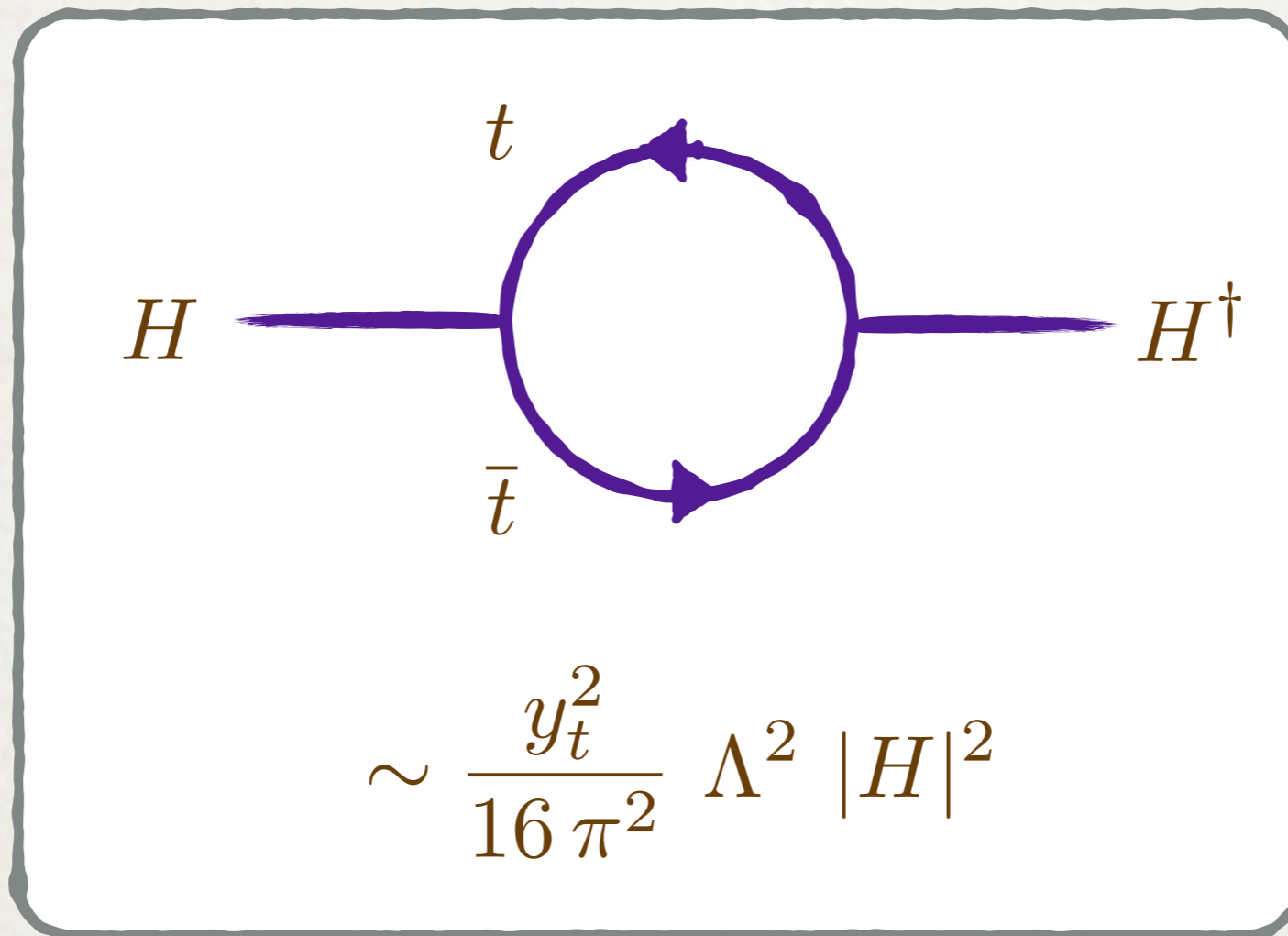
Yields  $\lambda_H \simeq 0.26$   
 $\mu \simeq 88$  GeV

**Value for all SM parameters known!**

# HIGGS MASS CORRECTIONS

$$\mathcal{L} \supset y_t H \bar{t} t \quad \Rightarrow \quad m_t = \frac{y_t}{\sqrt{2}} \langle H \rangle$$

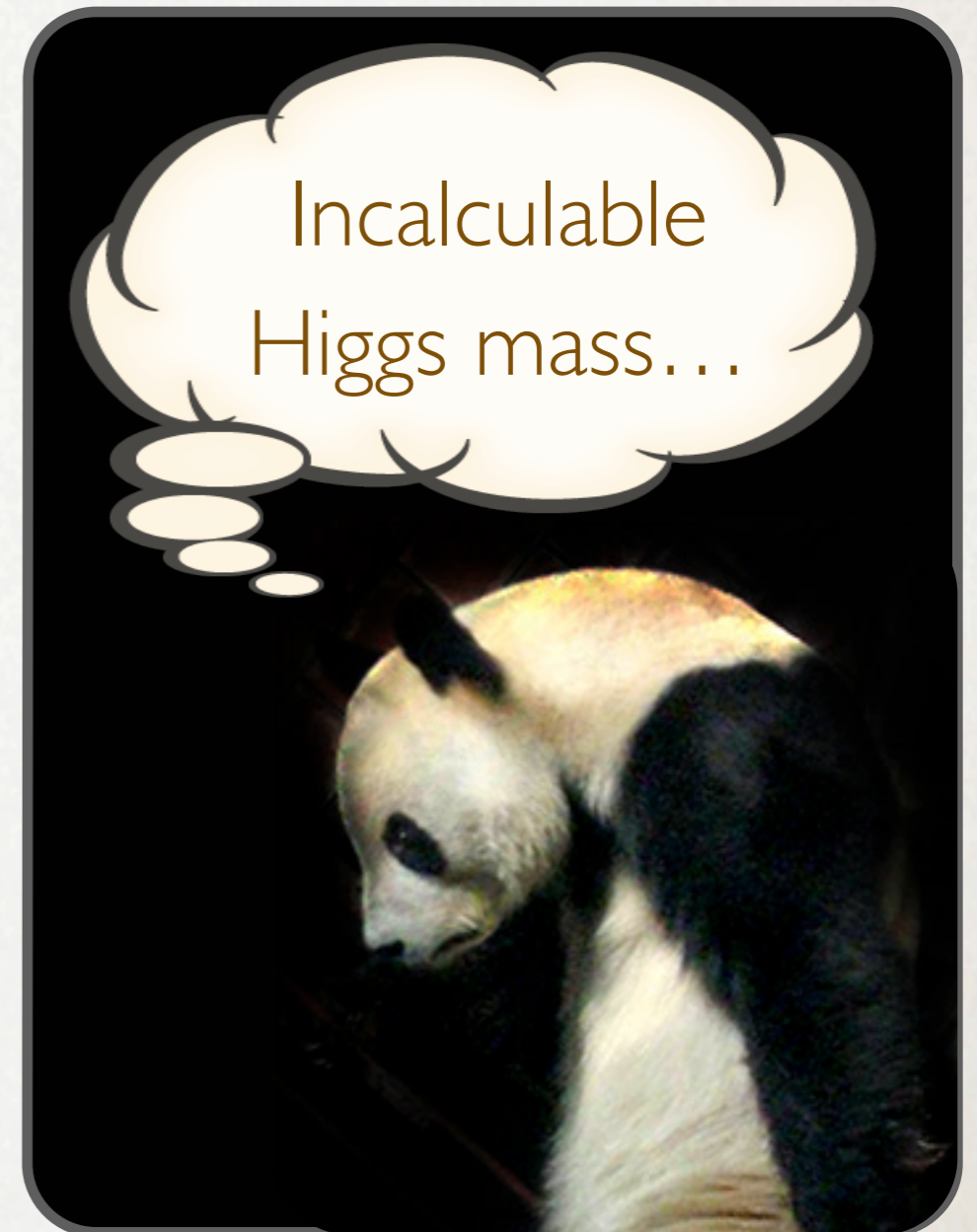
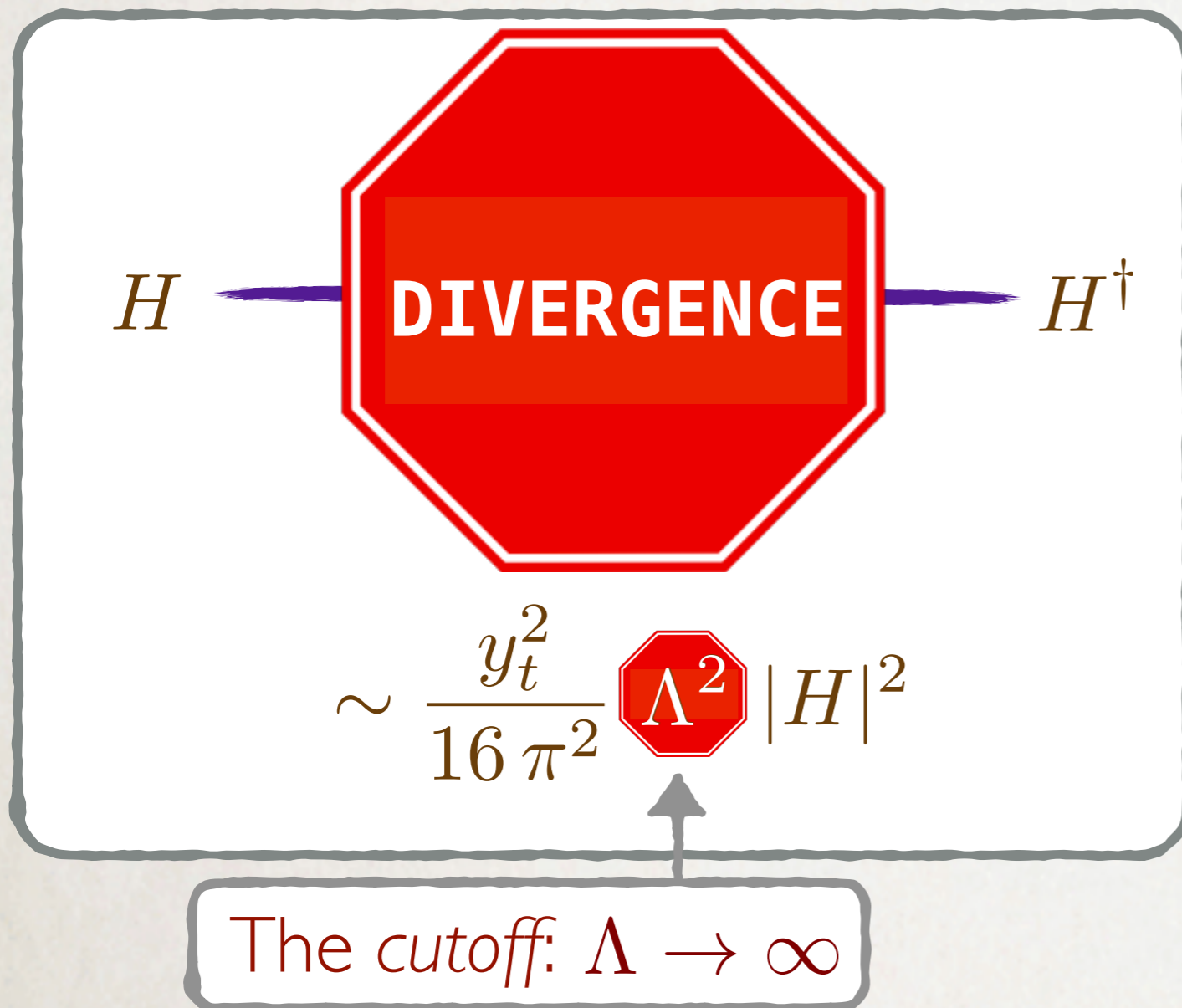
Top quark is heaviest particle; has strongest coupling to Higgs.



# HIGGS MASS CORRECTIONS

$$\mathcal{L} \supset y_t H \bar{t} t \quad \Rightarrow \quad m_t = \frac{y_t}{\sqrt{2}} \langle H \rangle$$

Top quark is heaviest particle; has strongest coupling to Higgs.



# TAME THE CUTOFF



To make sense of the model,  
“renormalize.”

$$m_{\text{finite}}^2 = m_{\text{bare}}^2 + m_{\text{loop}}^2$$

Diagram illustrating the renormalization process:

- $\mu^2$  (free parameter in Lagrangian)
- $\frac{y_t}{16\pi^2} \Lambda^2 + \mu^2$  (intermediate result)
- $-\frac{y_t}{16\pi^2} \Lambda^2$  (computed from loops)

Arrows indicate the flow of information:  $\mu^2$  and  $-\frac{y_t}{16\pi^2} \Lambda^2$  combine to form  $\frac{y_t}{16\pi^2} \Lambda^2 + \mu^2$ , which is then identified as  $m_{\text{finite}}^2$ . The term  $-\frac{y_t}{16\pi^2} \Lambda^2$  is identified as  $m_{\text{loop}}^2$ .

Recall:  $\mu^2 = (88 \text{ GeV})^2$  extracted from experiment.

IS THERE A PHYSICAL  
INTERPRETATION OF THE CUTOFF?!?

WILL SEE THAT THIS REQUIRES A  
*THEORY OF THE HIGGS MASS.*

## II. SUPERSYMMETRY: AN EXAMPLE THEORY



# FERMION MASS

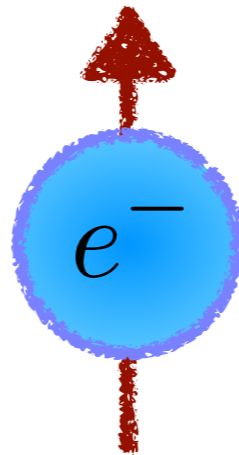
Fermions have spin  $1/2$ .

Is my fermion left or right handed?

# FERMION MASS

Fermions have spin  $1/2$ .  
Is my fermion left or right handed?

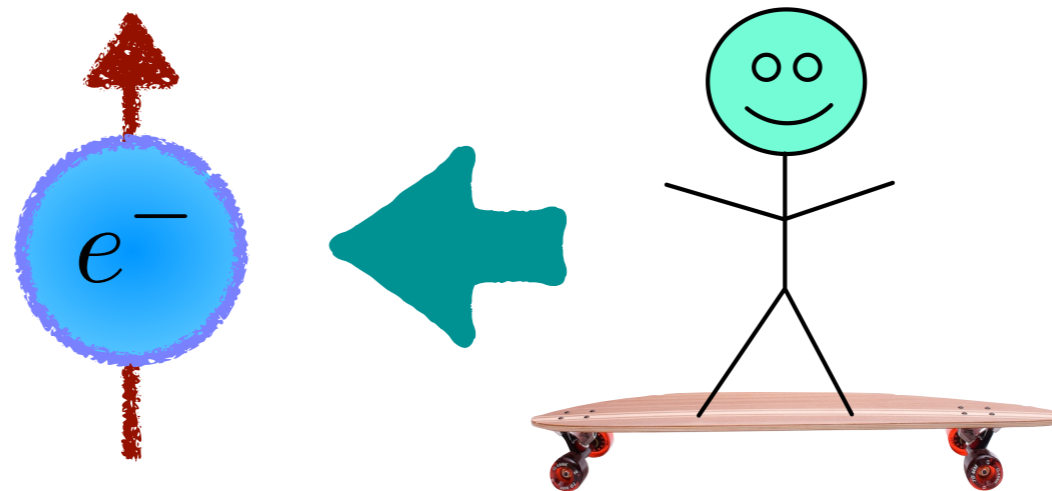
Take an electron spinning “up”.



# FERMION MASS

Fermions have spin  $1/2$ .  
Is my fermion left or right handed?

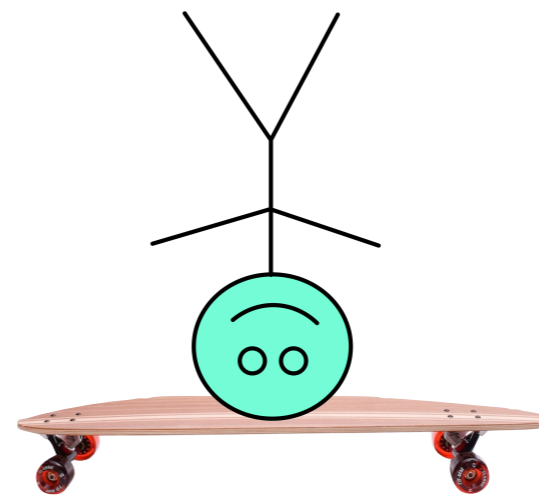
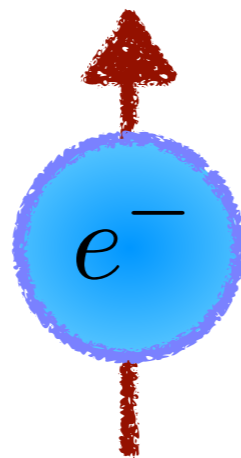
Move into its rest frame.



# FERMION MASS

Fermions have spin  $1/2$ .  
Is my fermion left or right handed?

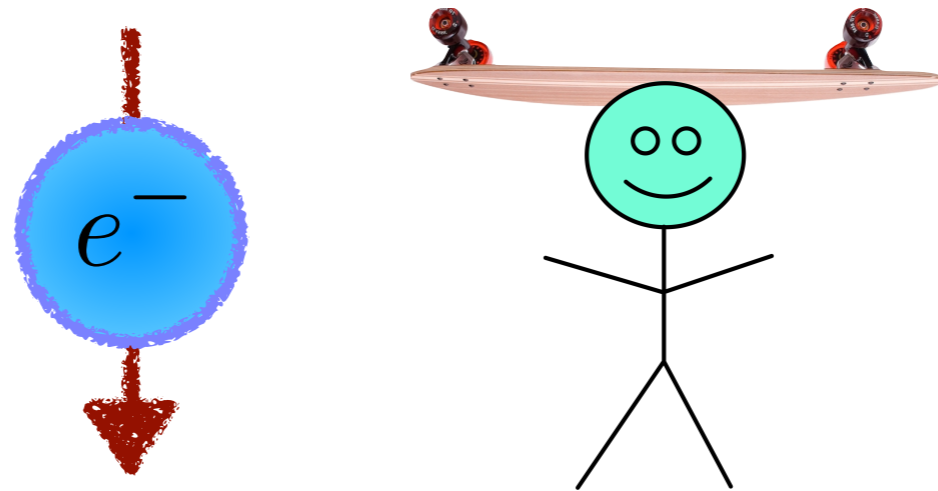
Stand on your head.



# FERMION MASS

Fermions have spin  $1/2$ .  
Is my fermion left or right handed?

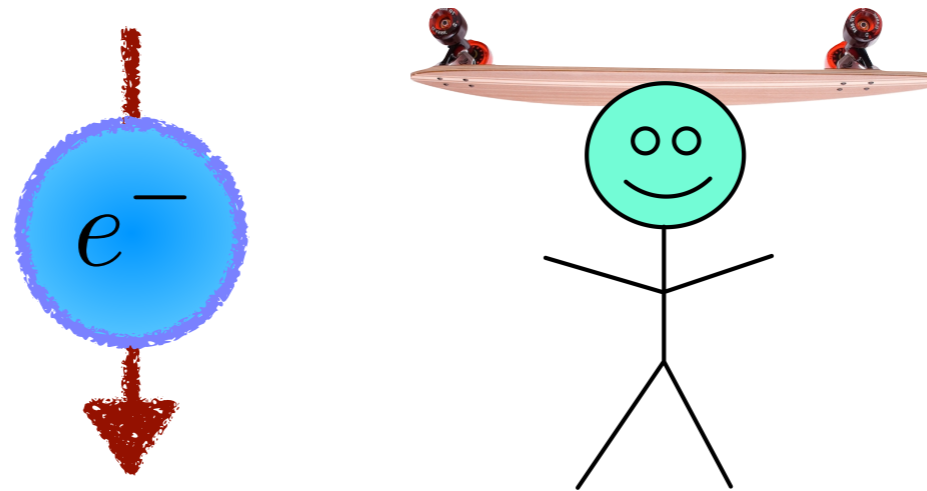
From your point of view,  
chirality is flipped.



# FERMION MASS

Fermions have spin 1/2.  
Is my fermion left or right handed?

From your point of view,  
chirality is flipped.



Going to rest frame is crucial step:  
flipping chirality requires mass.

$$\mathcal{L} \supset m \psi_L \psi_R$$

Fermion mass mixes left and right chiralities.

# FERMION MASS

Recall, scalar mass always symmetric under phase rotation.

$$\mathcal{L} \supset m_\phi |\phi|^2$$

What about fermion mass?

$$\mathcal{L} \supset m_\psi \psi_L \psi_R$$

# FERMION MASS

Recall, scalar mass always symmetric under phase rotation.

$$\mathcal{L} \supset m_\phi |\phi|^2$$

What about fermion mass?

$$\mathcal{L} \supset m_\psi \psi_L \psi_R$$

“Chiral” rotation:

$$\psi_R \rightarrow \psi'_R = e^{i\zeta} \psi_R$$

$$\psi_L \rightarrow \psi'_L = e^{i\zeta} \psi_L$$

$$\mathcal{L} \rightarrow \mathcal{L}' \subset e^{2i\xi} m_\psi \psi_L \psi_R$$



# FERMION MASS

Recall, scalar mass always symmetric under phase rotation.

$$\mathcal{L} \supset m_\phi |\phi|^2$$

What about fermion mass?

$$\mathcal{L} \supset m_\psi \psi_L \psi_R$$

“Chiral” rotation:

$$\psi_R \rightarrow \psi'_R = e^{i\zeta} \psi_R$$

$$\psi_L \rightarrow \psi'_L = e^{i\zeta} \psi_L$$

$$\mathcal{L} \rightarrow \mathcal{L}' \subset e^{2i\xi} m_\psi \psi_L \psi_R$$

Chiral symmetry can forbid fermion mass.

Fermion mass correction proportional to  $m_\psi$ .

**Cutoff does not infect fermion masses!**

# “CHIRALITY” FOR SCALARS



Scalars inherit chirality of partner fermions.

Calculability of fermion masses inherited by scalars.

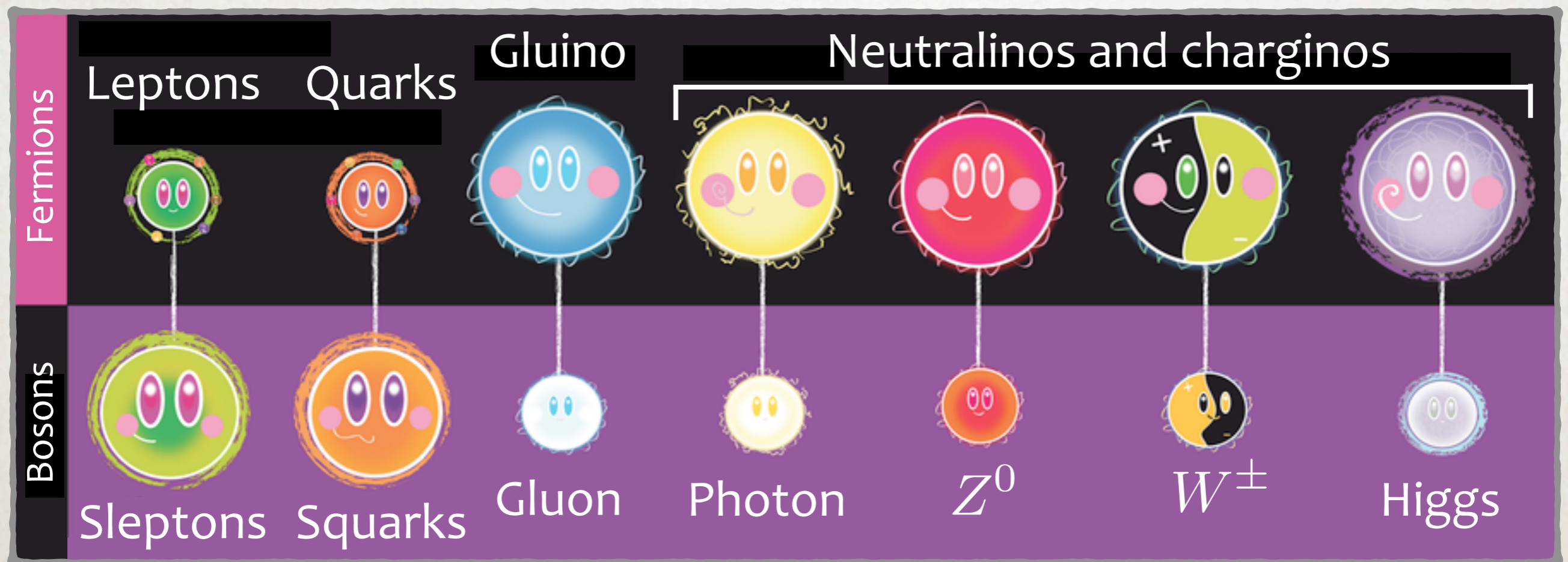
# A NEW KIND OF SYMMETRY

## Gauge invariance:

Phase rotation compensated by states of opposite charge:  
electron needed positron.

## Supersymmetry:

“Rotation” compensated by states of different spin



# THE TOP GETS A PARTNER

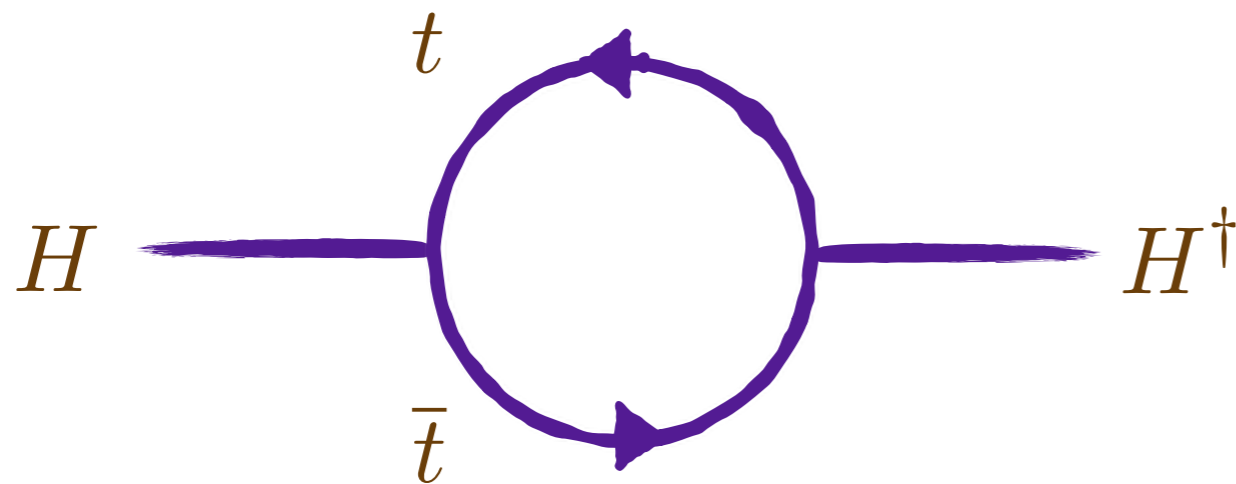


Top is a massive fermion; has  $L$  and  $R$  chirality.  
Consistency requires introducing two stops:  $\tilde{t}_L$  and  $\tilde{t}_R$ .  
Free parameters: two masses and a mixing angle.

Top coupling to the Higgs:  $y_t$ .

Strength of top partner coupling to the Higgs set by  $y_t$ .

# HIGGS MASS CORRECTIONS

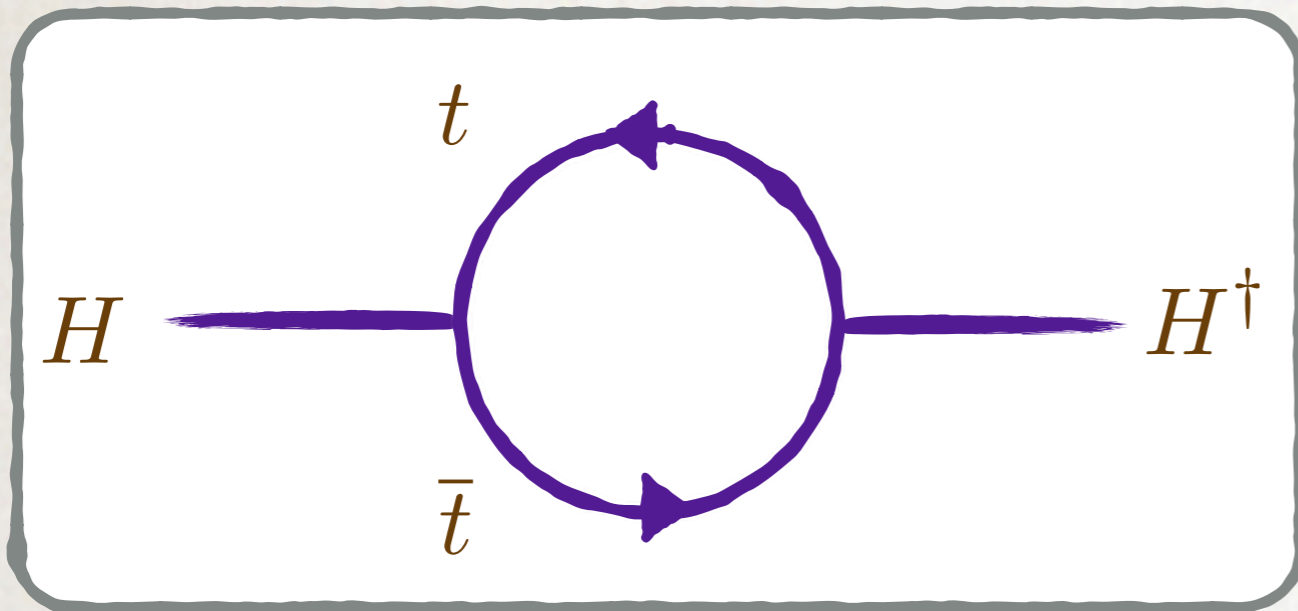


$$\sim \frac{y_t^2}{16\pi^2} \Lambda^2 |H|^2$$

Incalculable  
Higgs mass...

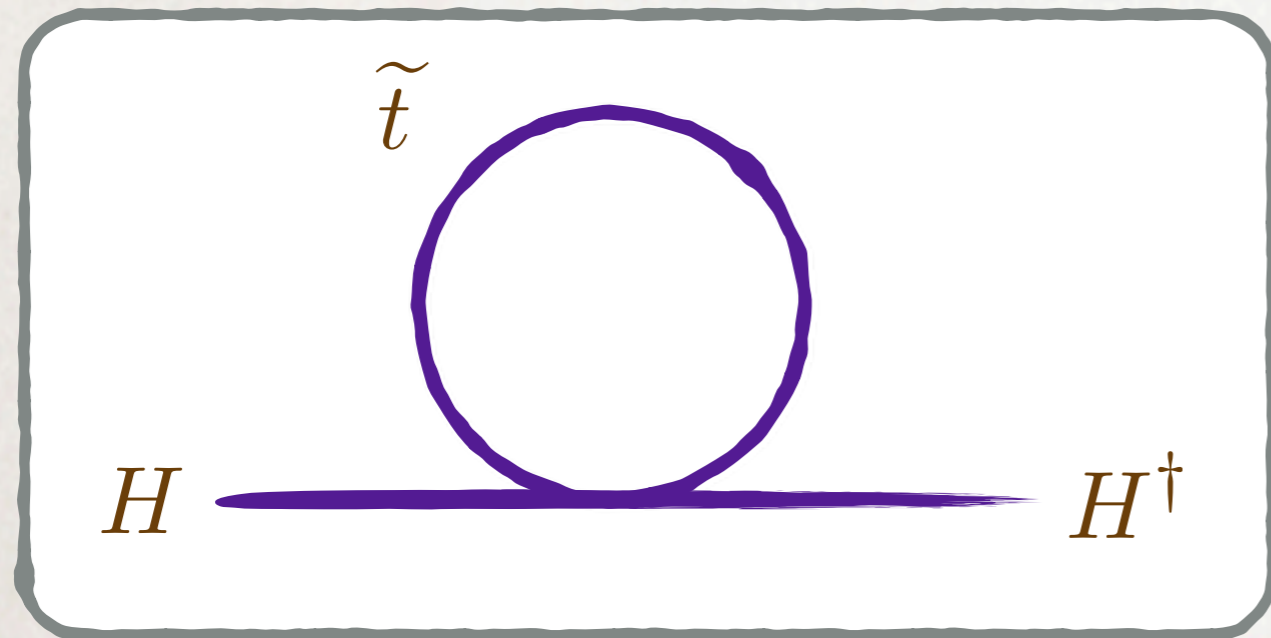


# HIGGS MASS CORRECTIONS



$$\sim \frac{y_t^2}{16\pi^2} \Lambda^2 |H|^2$$

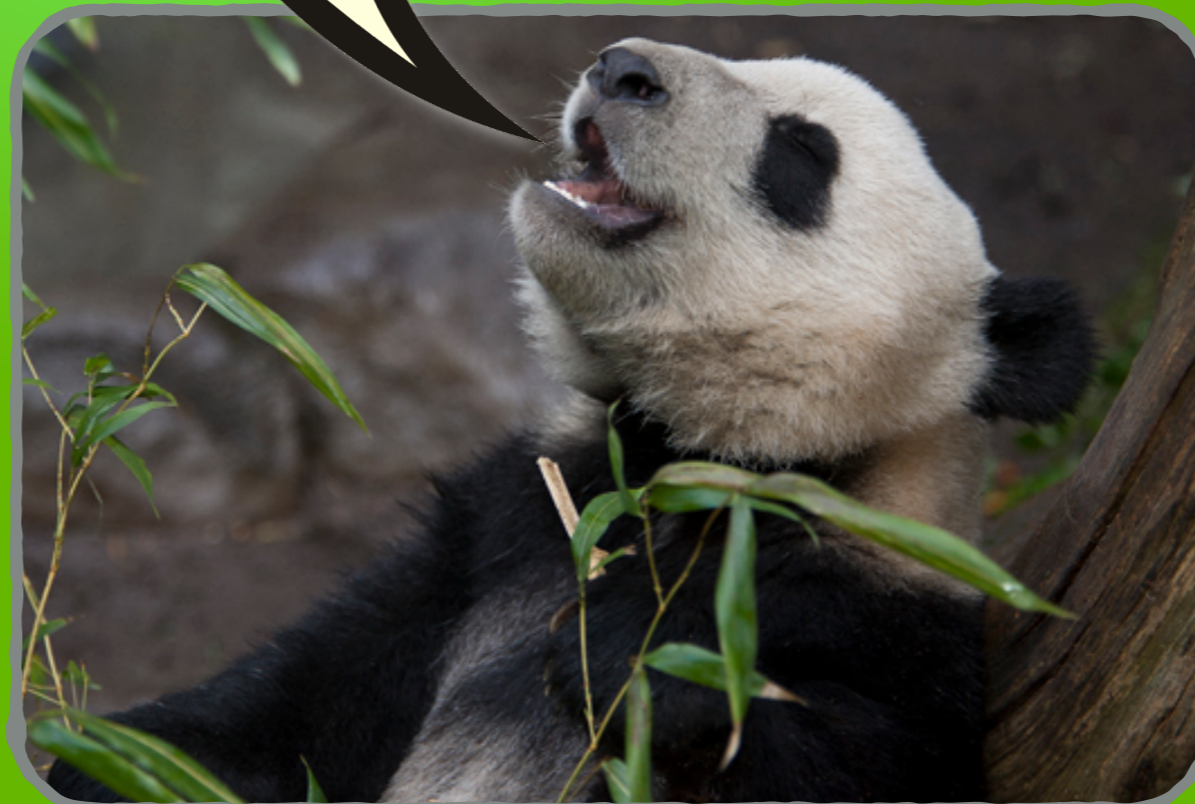
+



$$\sim -\frac{y_t^2}{16\pi^2} \left( \Lambda^2 + c m_{\tilde{t}}^2 \right) |H|^2$$

# HIGGS MASS CORRECTIONS

*Calculable  
Higgs mass!*



# SUSY HIGGS POTENTIAL\*

“Naturalness in SUSY”

$$V \simeq \left( m_H^2 - \frac{3}{4\pi^2} y_t^2 m_{\tilde{t}}^2 \right) |H|^2$$

Calculable Higgs mass.

Stop mass give physical meaning to quadratic divergence.

Heavier stops imply larger cancellation.

***Stop mass below ~1 TeV extremely plausible!***

\* in a simplified limit



# SUSY HIGGS POTENTIAL\*

$$V \simeq \left( m_H^2 - \frac{3}{4\pi^2} y_t^2 m_{\tilde{t}}^2 \right) |H|^2 + \left( \frac{g^2 + g'^2}{8} + \frac{3}{8\pi^2} y_t^4 \log \frac{m_{\tilde{t}}}{m_t} \right) |H|^4$$

Calculable Higgs mass.

Stop mass give physical meaning to quadratic divergence.

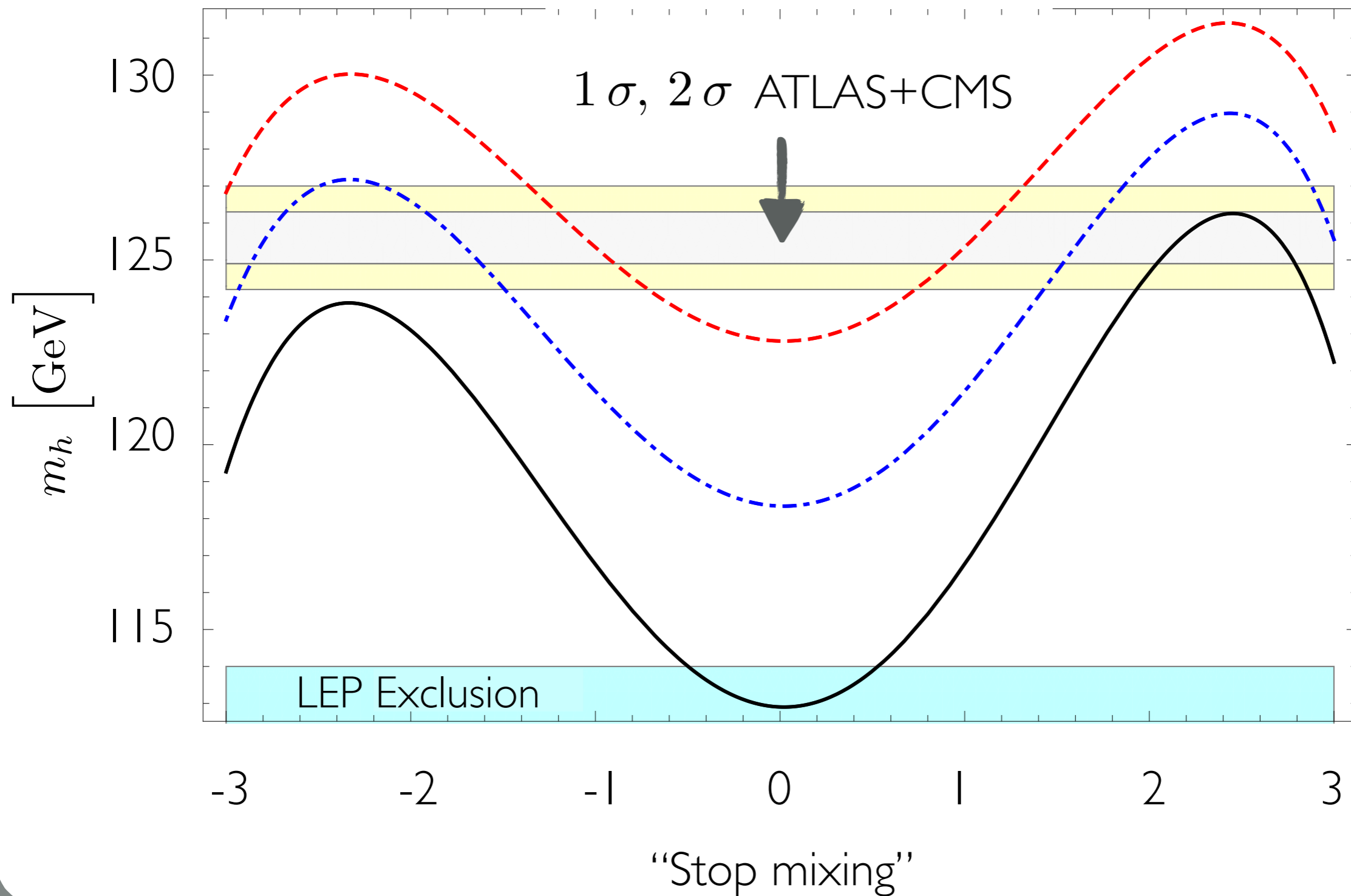
Heavier stops imply larger cancellation.

***Stop mass below ~1 TeV extremely plausible!***

Heavier stops yield larger quartic.

\* in a simplified limit

# HIGGS MASS IN MSSM

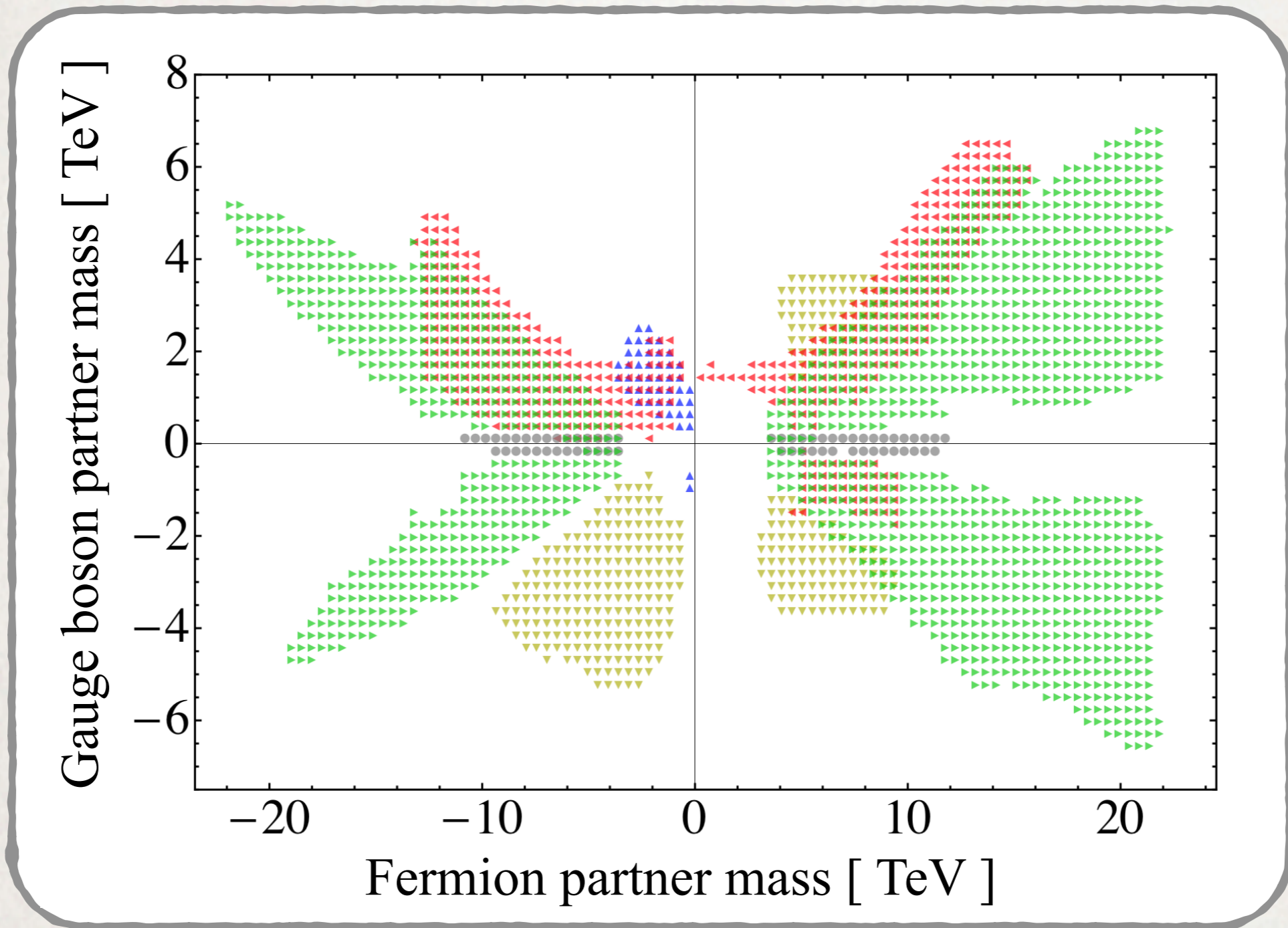


$m_{\tilde{\tau}}$   
4 TeV  
2 TeV  
1 TeV

Draper, Lee, Wagner [arXiv:1312.5743]

# SUSY PARAMETER SPACE

Highly simplified assumption for inputs (the CMSSM).



Boundaries from requiring Higgs mass and relic density of dark matter.

[TC, Wacker \[arXiv:1305.2914\]](https://arxiv.org/abs/1305.2914)

# III. SUSY NATURALNESS CONFRONTS EXPERIMENT

# MINIMAL INGREDIENTS

Biggest contributions to Higgs mass, ordered by size:

$$m_H^2 = m_{\text{Higgsino}}^2 + m_{\text{stop}}^2 + m_{\text{gluino}}^2 + \dots$$



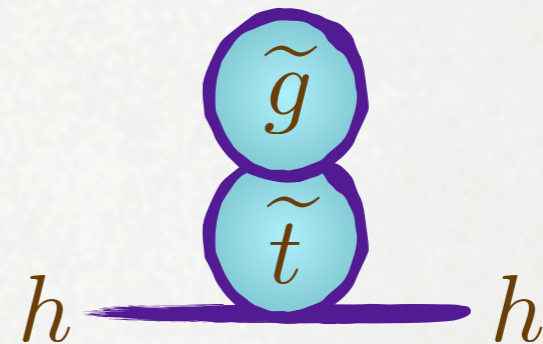
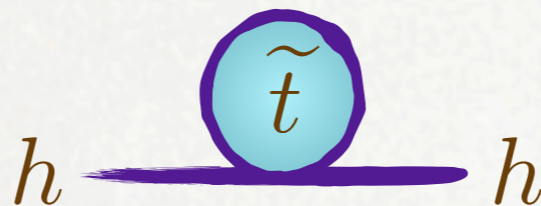
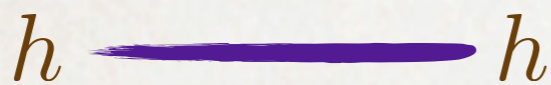
tree-level



one-loop



two-loop



Minimal spectrum



$\tilde{g}$

1 TeV



$\tilde{t}$



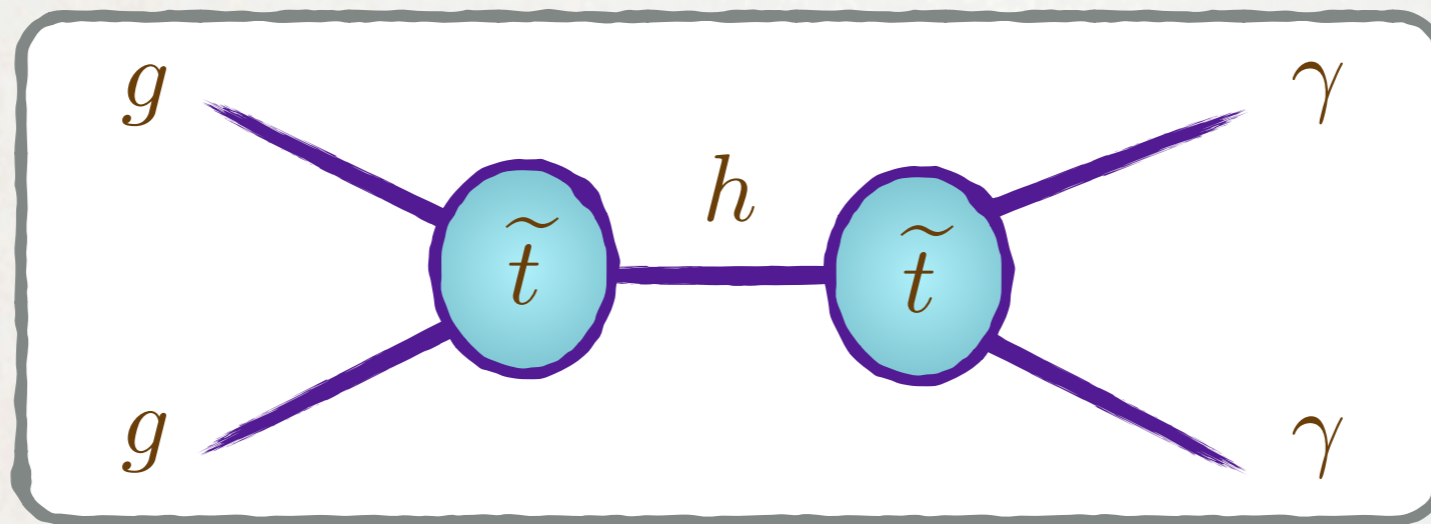
$\chi$

100 GeV

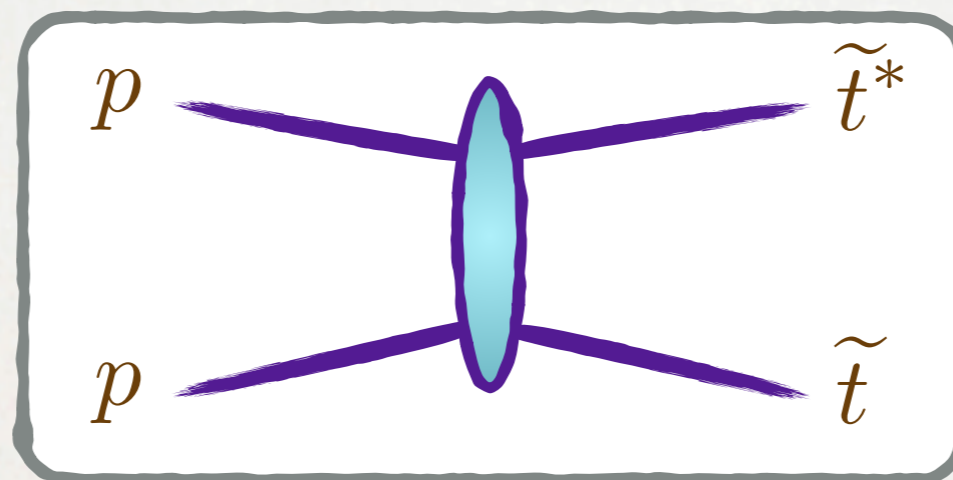
# NATURAL THEORIES HAVE OBSERVABLES

Existence of top partners  $\Rightarrow$  physical observables!

Loop corrections to Higgs properties.

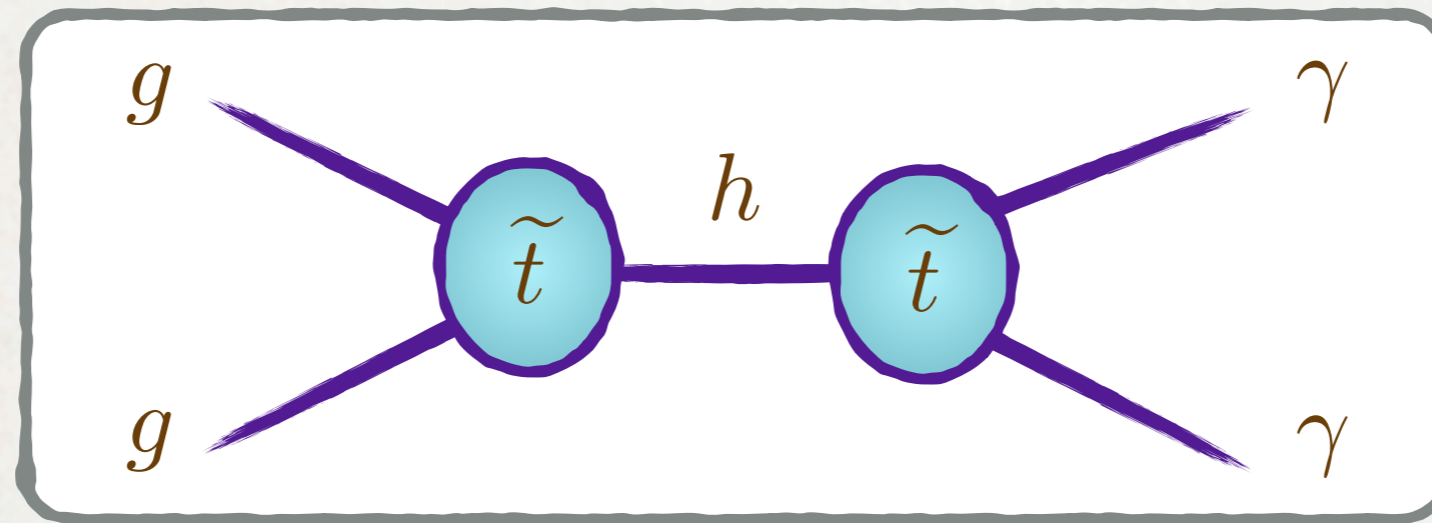


Direct production in proton collisions.



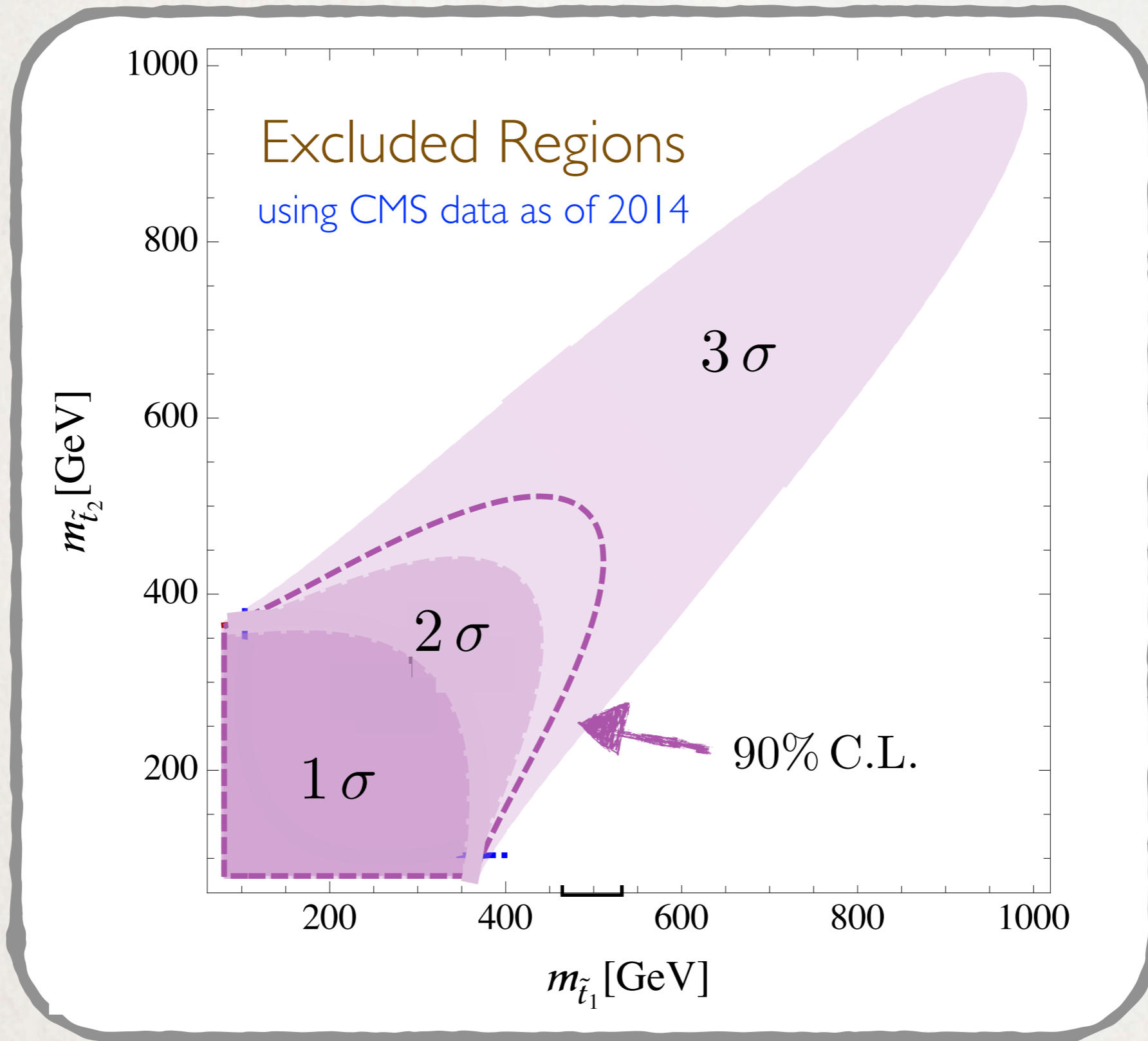
# HIGGS COUPLINGS

Modification to Higgs production and decay.



Yields bounds independent of stop decay modes.

# CONSTRAINTS

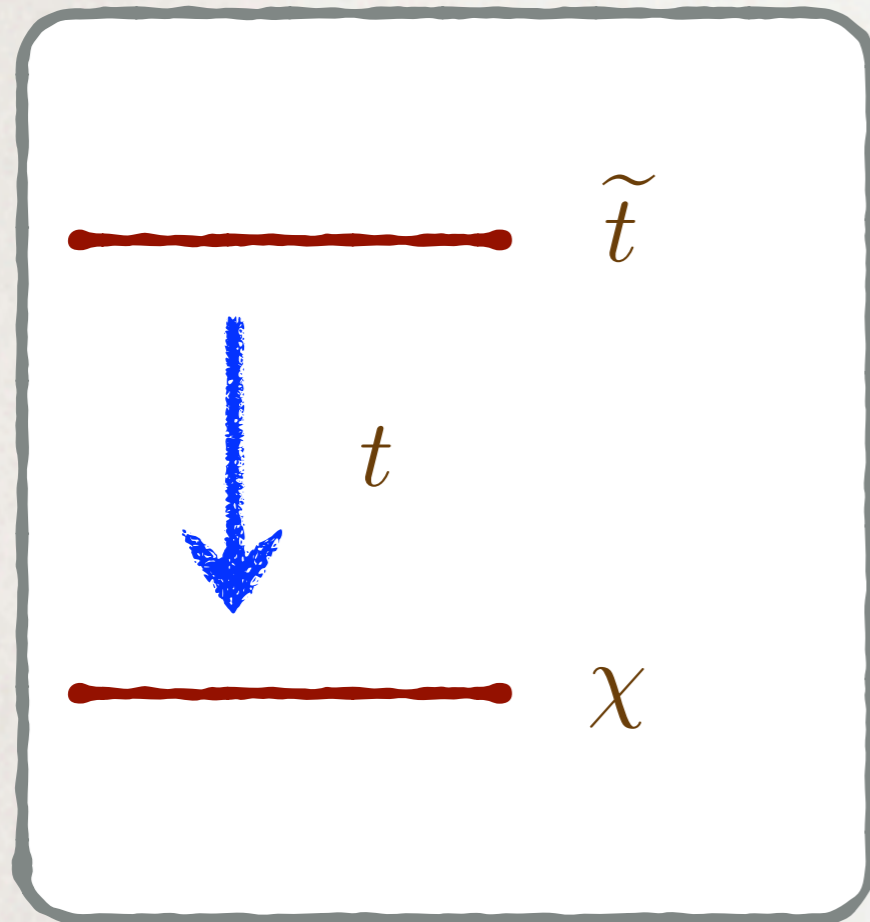


Reece, Fan [arXiv:1401.7671]

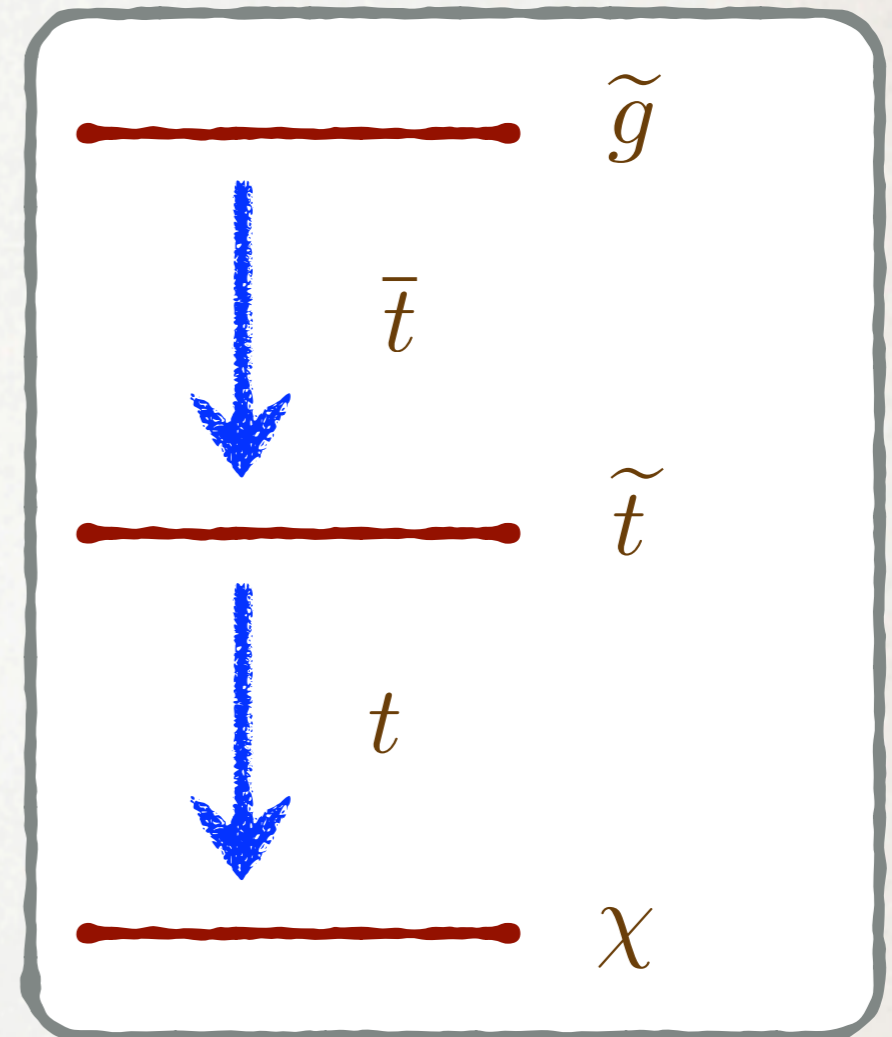


# SIMPLIFIED MODEL FOR PROTON COLLIDERS

Stop-neutralino

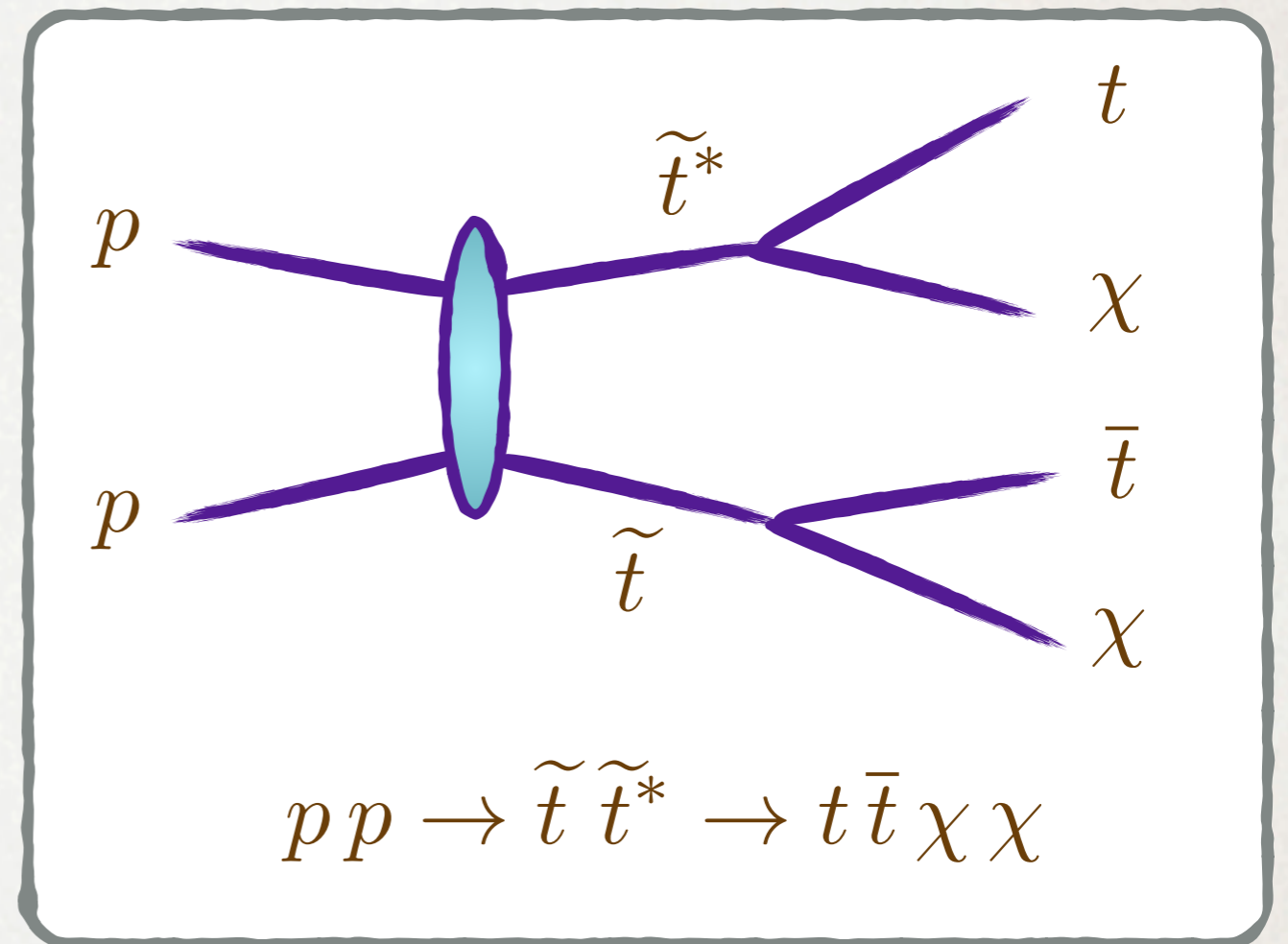
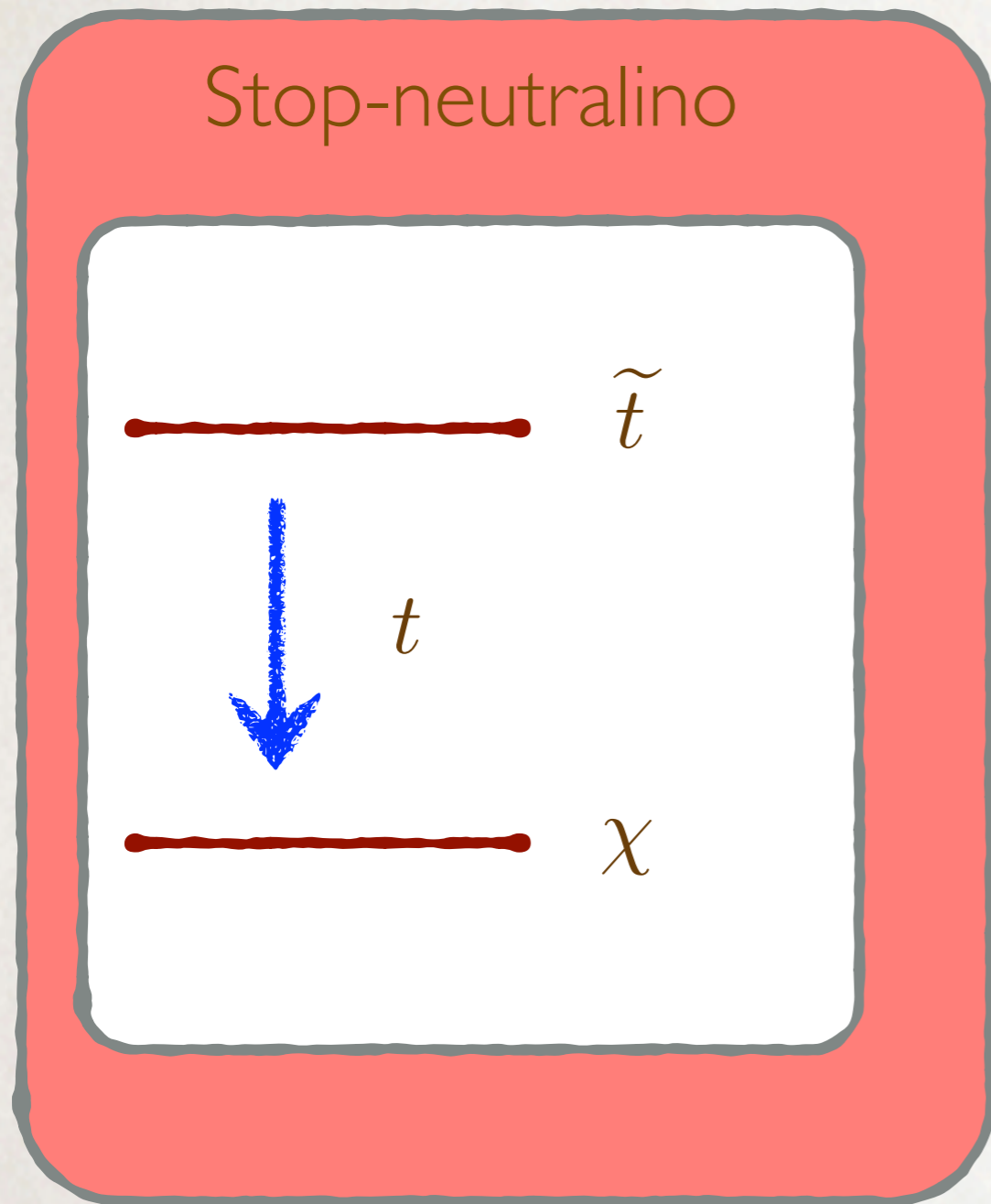


Glauino-stop-neutralino



Essig, Izaguirre, Kaplan, Wacker [arXiv:1110.6443]; Papucci, Ruderman, Weiler [arXiv:1110.6926]

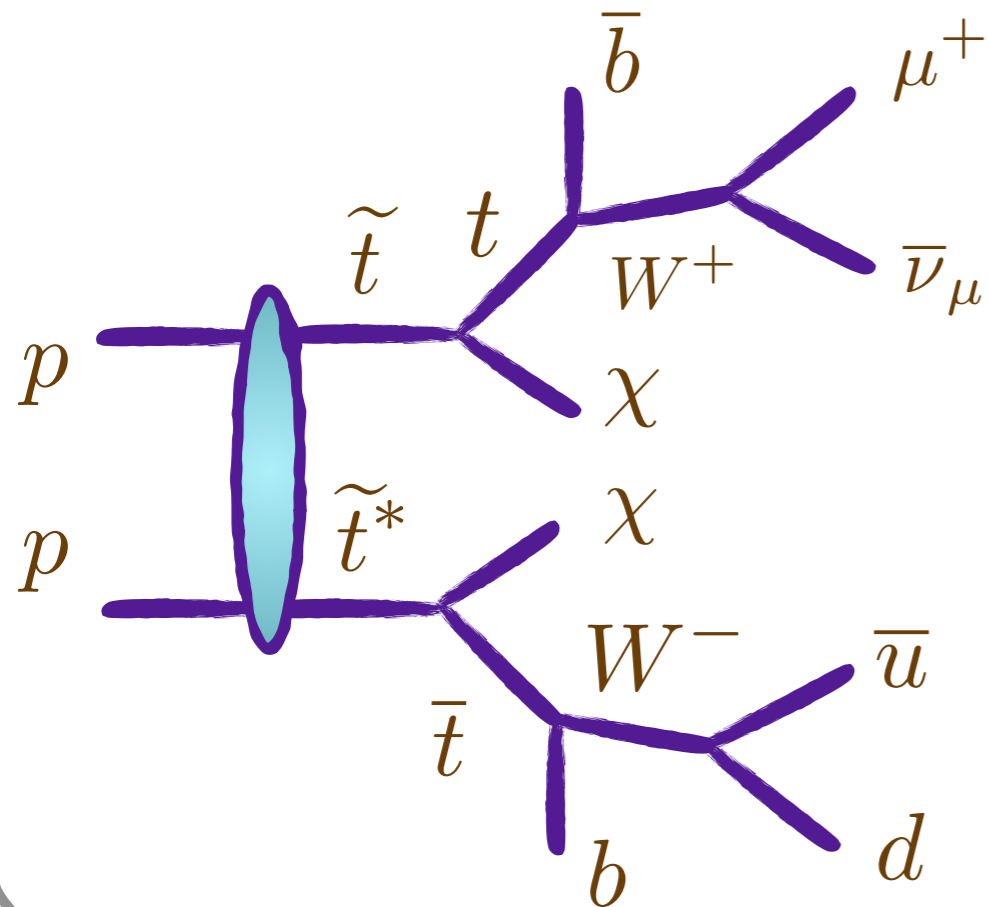
# SIMPLIFIED MODEL FOR PROTON COLLIDERS



Essig, Izaguirre, Kaplan, Wacker [arXiv:1110.6443]; Papucci, Ruderman, Weiler [arXiv:1110.6926]

# SEARCH STRATEGY

## MAIN REQUIREMENTS

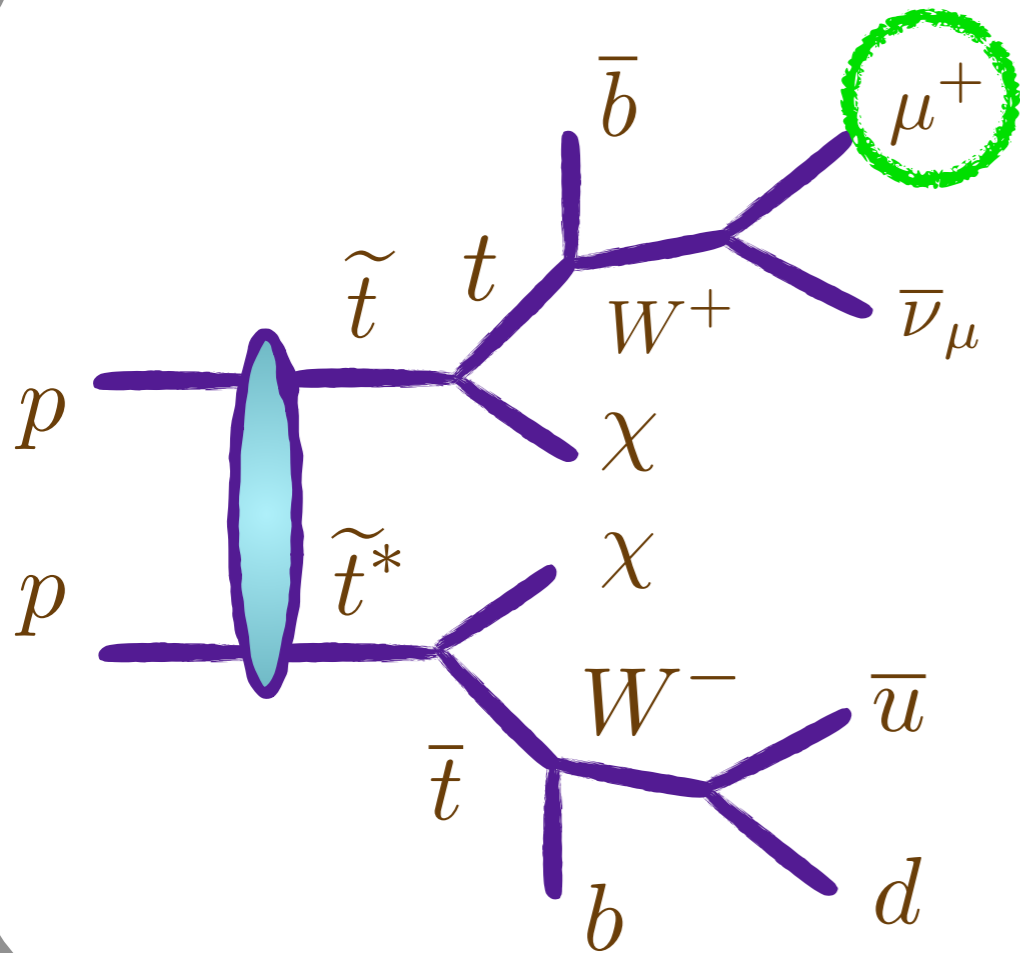


ATLAS [arXiv:1407.0583];  
see also CMS [arXiv:1502.00300]

# SEARCH STRATEGY

## MAIN REQUIREMENTS

- 1 lepton

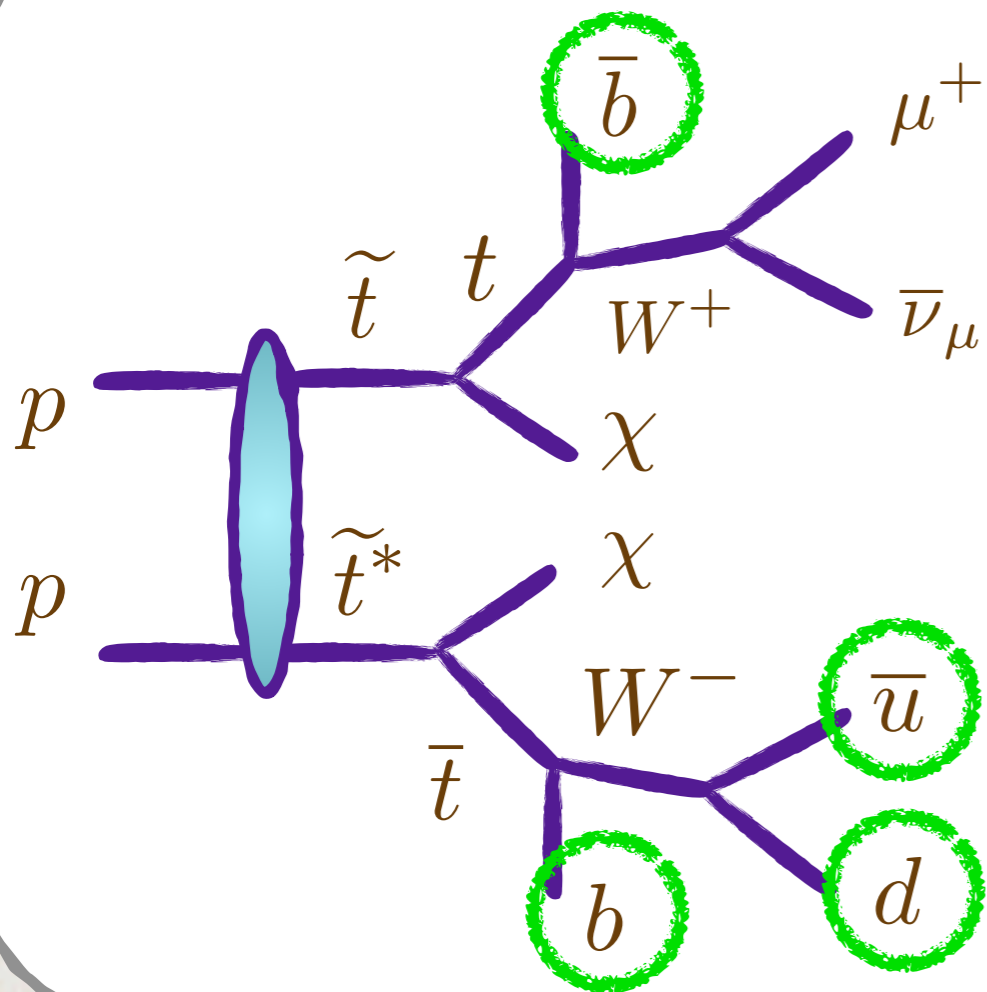


ATLAS [arXiv:1407.0583];  
see also CMS [arXiv:1502.00300]

# SEARCH STRATEGY

## MAIN REQUIREMENTS

- 1 lepton
- $\geq 4$  jets

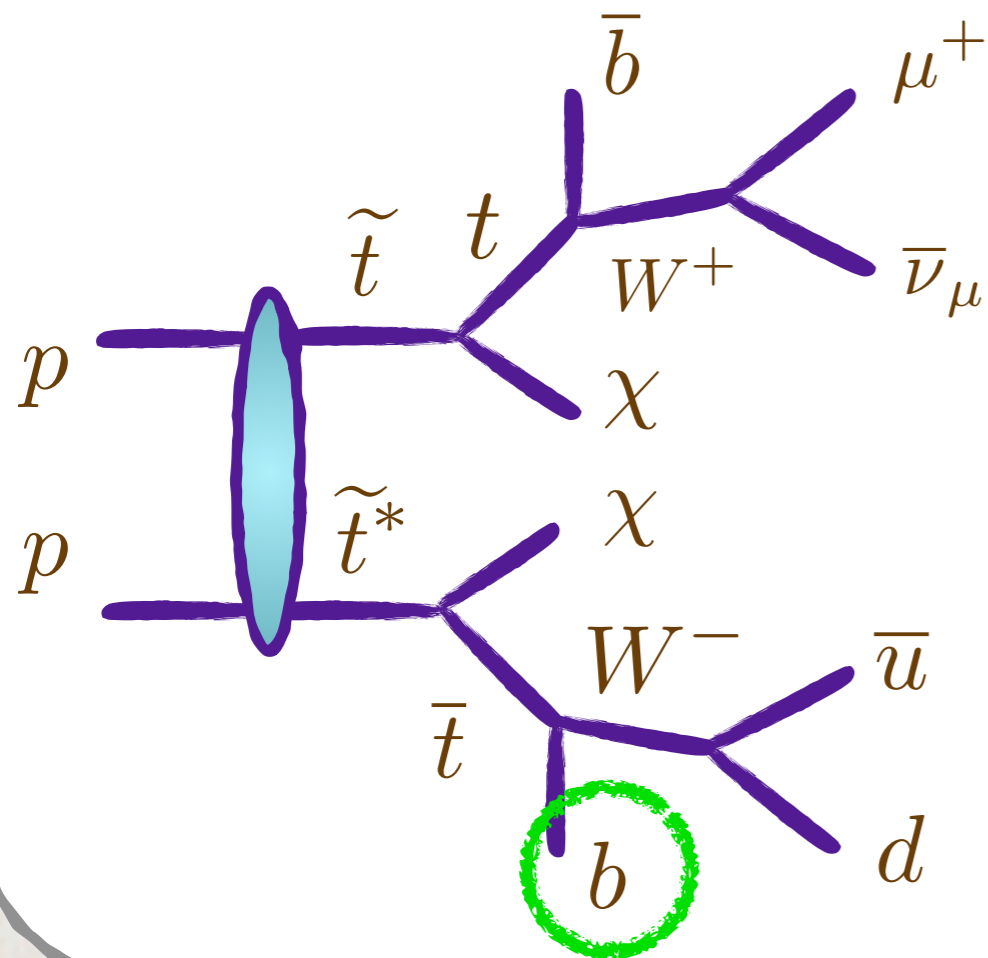


ATLAS [arXiv:1407.0583];  
see also CMS [arXiv:1502.00300]

# SEARCH STRATEGY

## MAIN REQUIREMENTS

- 1 lepton
- $\geq 4$  jets
- $\geq 1$   $b$ -jet(s)

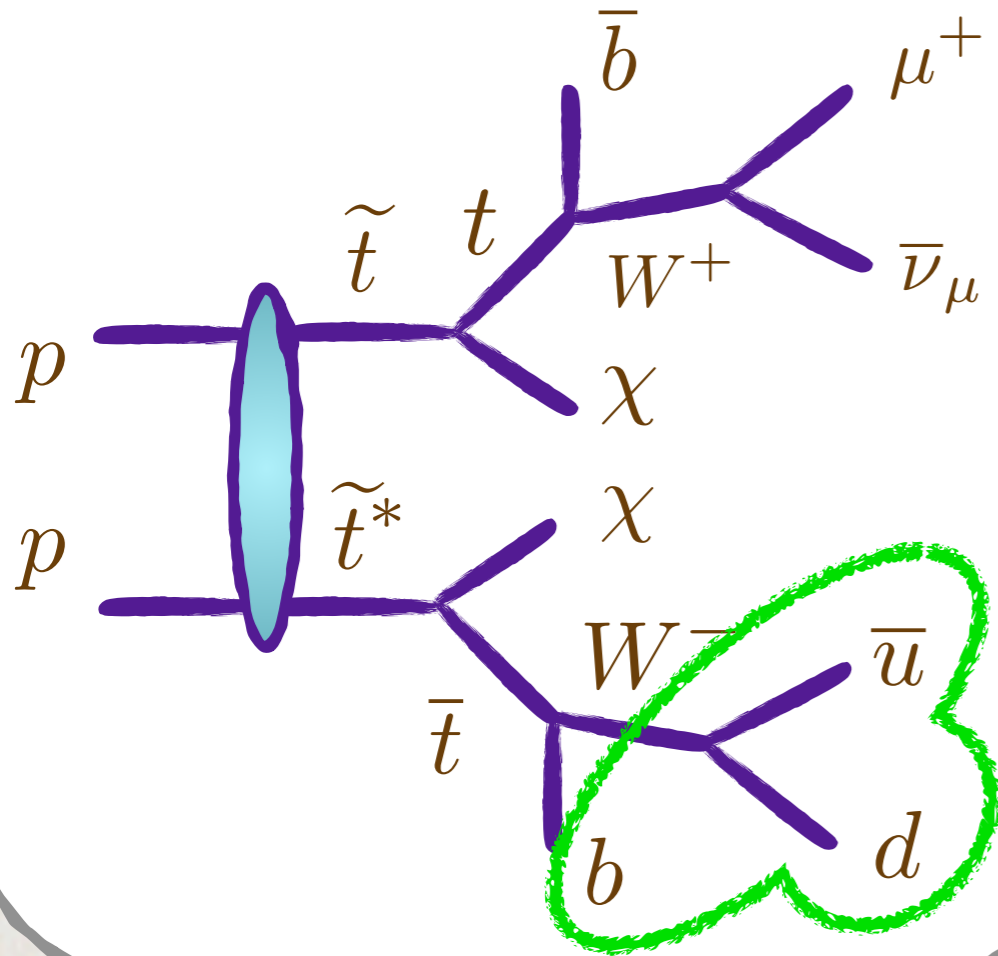


ATLAS [arXiv:1407.0583];  
see also CMS [arXiv:1502.00300]

# SEARCH STRATEGY

## MAIN REQUIREMENTS

- 1 lepton
- $\geq 4$  jets
- $\geq 1$   $b$ -jet(s)
- One “hadronic top”

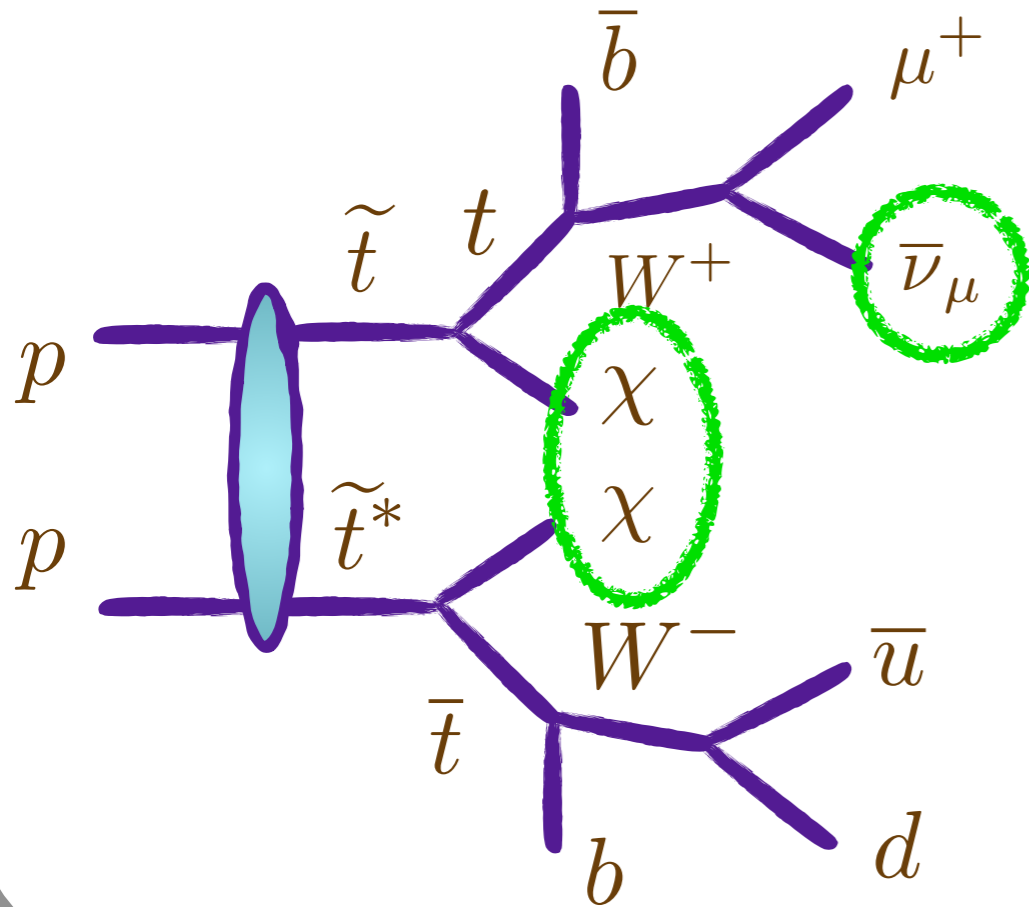


ATLAS [arXiv:1407.0583];  
see also CMS [arXiv:1502.00300]

# SEARCH STRATEGY

## MAIN REQUIREMENTS

- 1 lepton
- $\geq 4$  jets
- $\geq 1$   $b$ -jet(s)
- One “hadronic top”
- Few hundred GeV missing energy

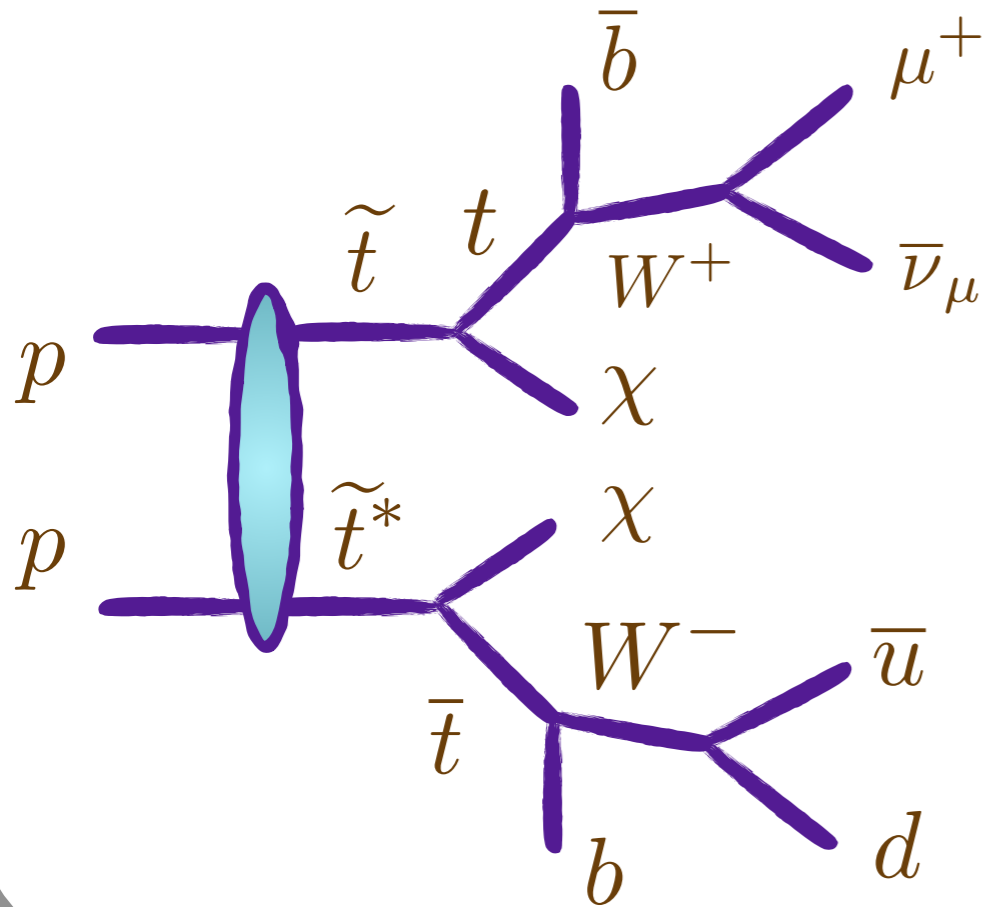


ATLAS [arXiv:1407.0583];  
see also CMS [arXiv:1502.00300]



# SEARCH STRATEGY

## MAIN REQUIREMENTS



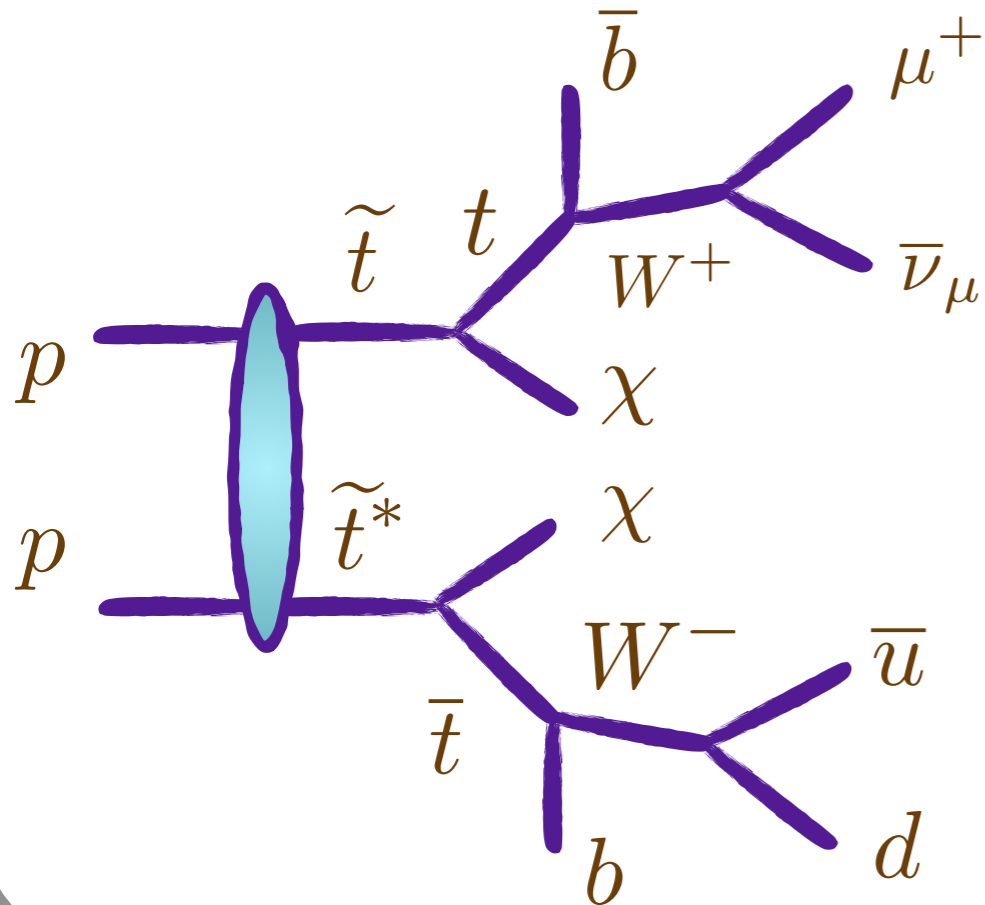
- 1 lepton
- $\geq 4$  jets
- $\geq 1$   $b$ -jet(s)
- One “hadronic top”
- Few hundred GeV missing energy
- Few hundred GeV transverse mass
- ...

ATLAS [arXiv:1407.0583];  
see also CMS [arXiv:1502.00300]

# SEARCH STRATEGY

## MAIN REQUIREMENTS

- 1 lepton
- $\geq 4$  jets
- $\geq 1$   $b$ -jet(s)
- One “hadronic top”
- Few hundred GeV missing energy
- Few hundred GeV transverse mass
- ...

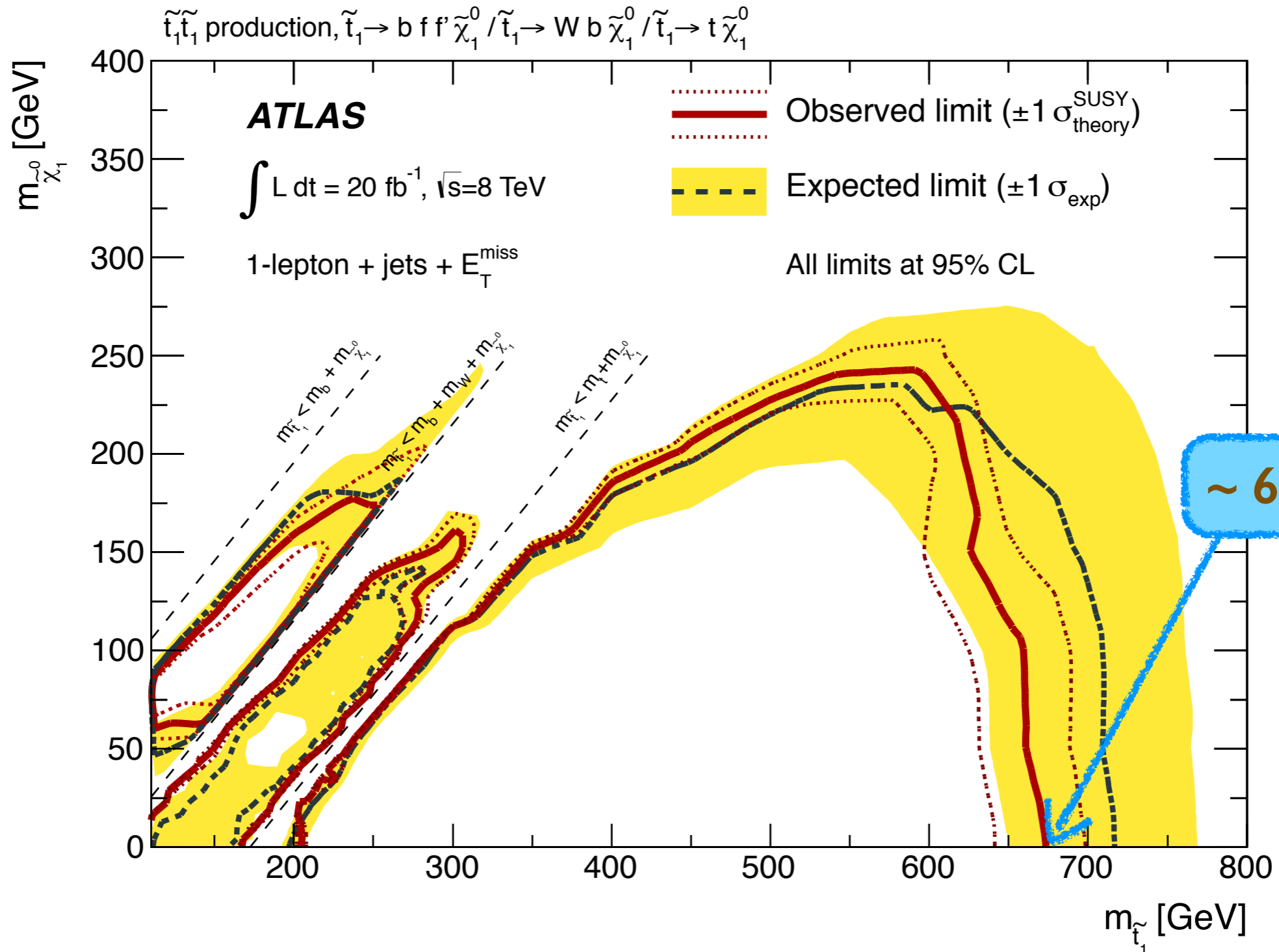


ATLAS [arXiv:1407.0583];  
see also CMS [arXiv:1502.00300]

## DOMINANT BACKGROUNDS

$t\bar{t}$ ,  $W$  + jets,  $t\bar{t}$  +  $W/Z$ , ...

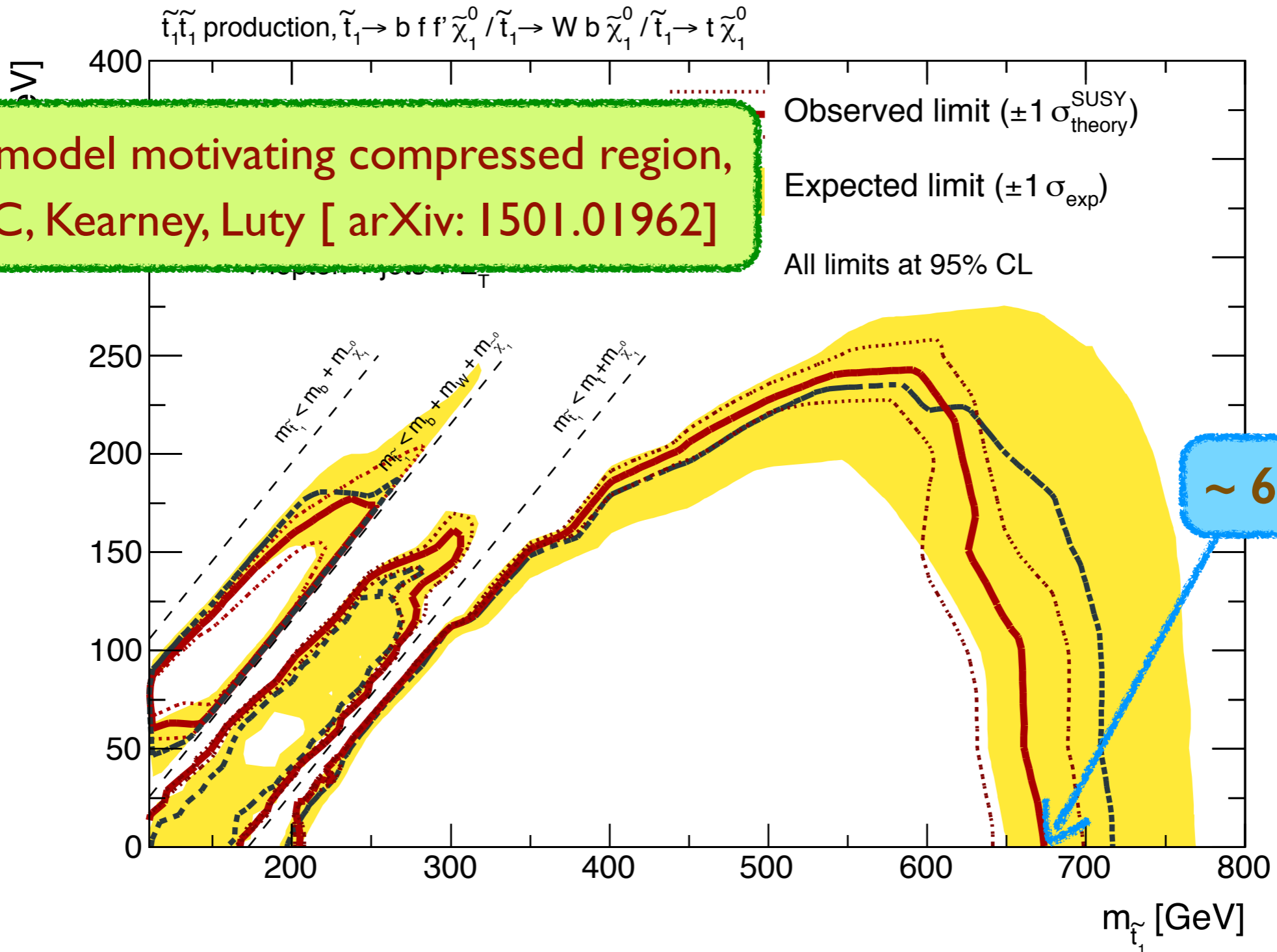
# 1-LEPTON RESULTS



ATLAS [arXiv:1407.0583]; see also CMS [arXiv:1502.00300]

# 1-LEPTON RESULTS

For a model motivating compressed region, see TC, Kearney, Luty [ arXiv: 1501.01962 ]

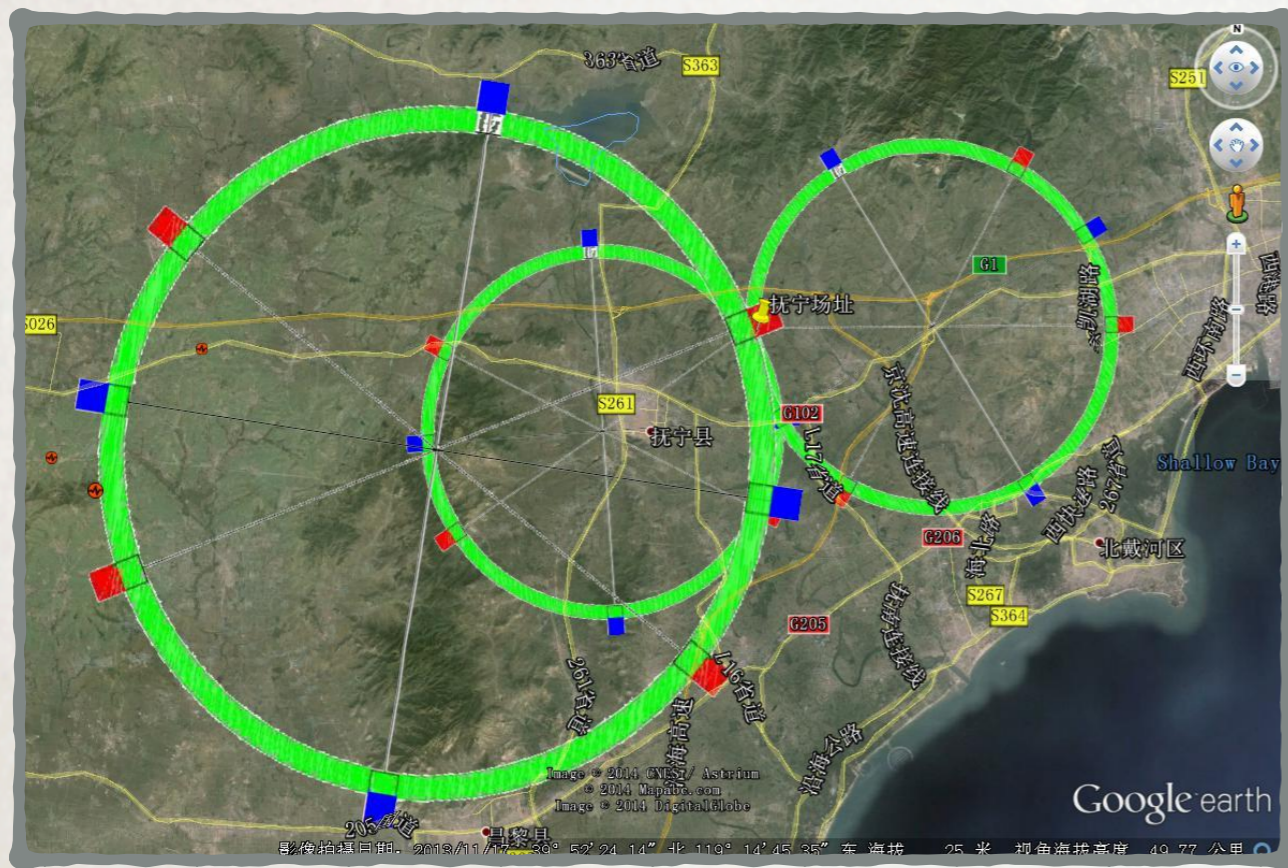


ATLAS [arXiv:1407.0583]; see also CMS [arXiv:1502.00300]

# FUTURE COLLIDERS

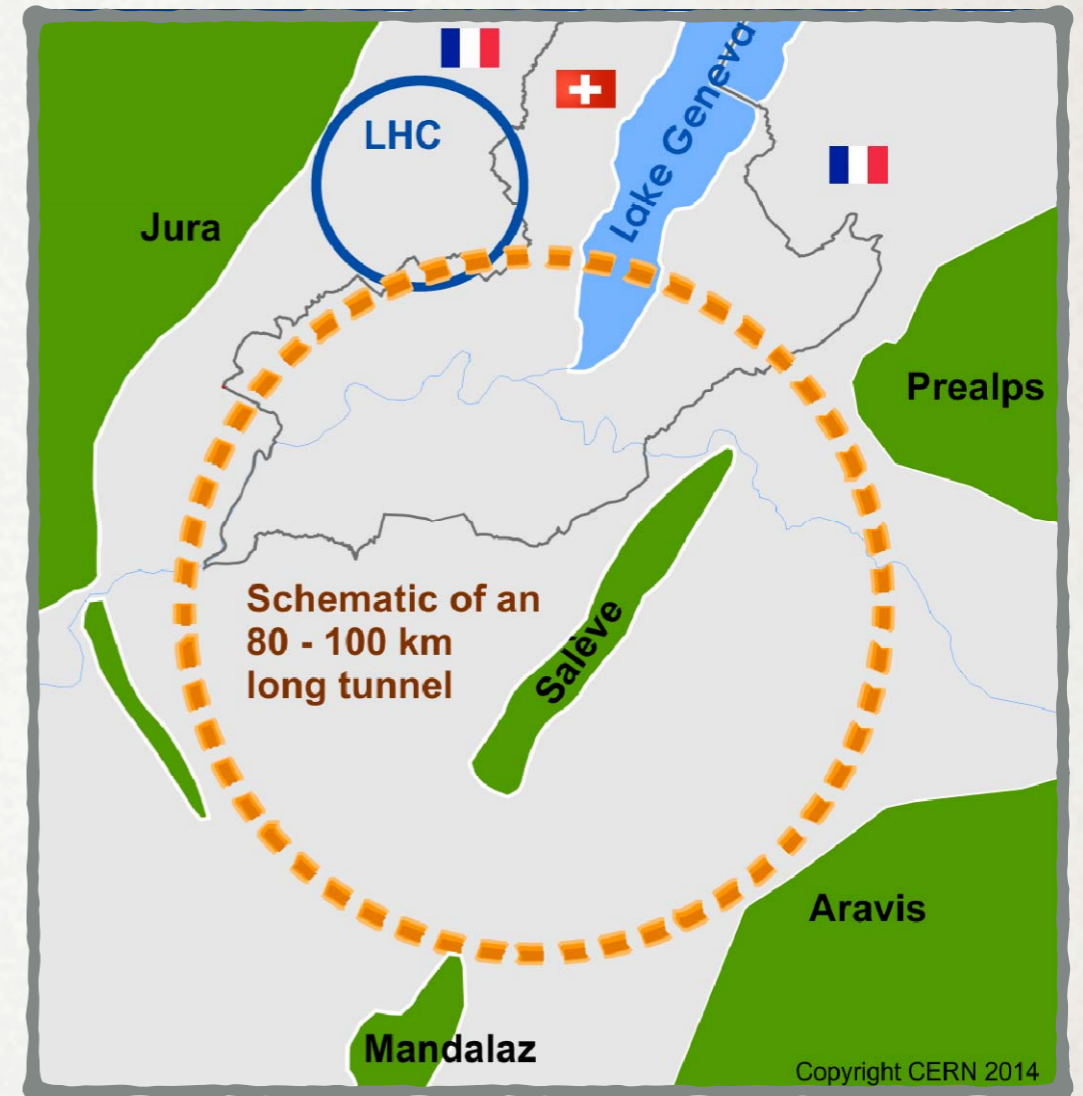
Higgs factory and  
100 TeV proton collider?

IHEP in China?



Talk by X. Lou [Aspen Future Colliders Conference, 2015]

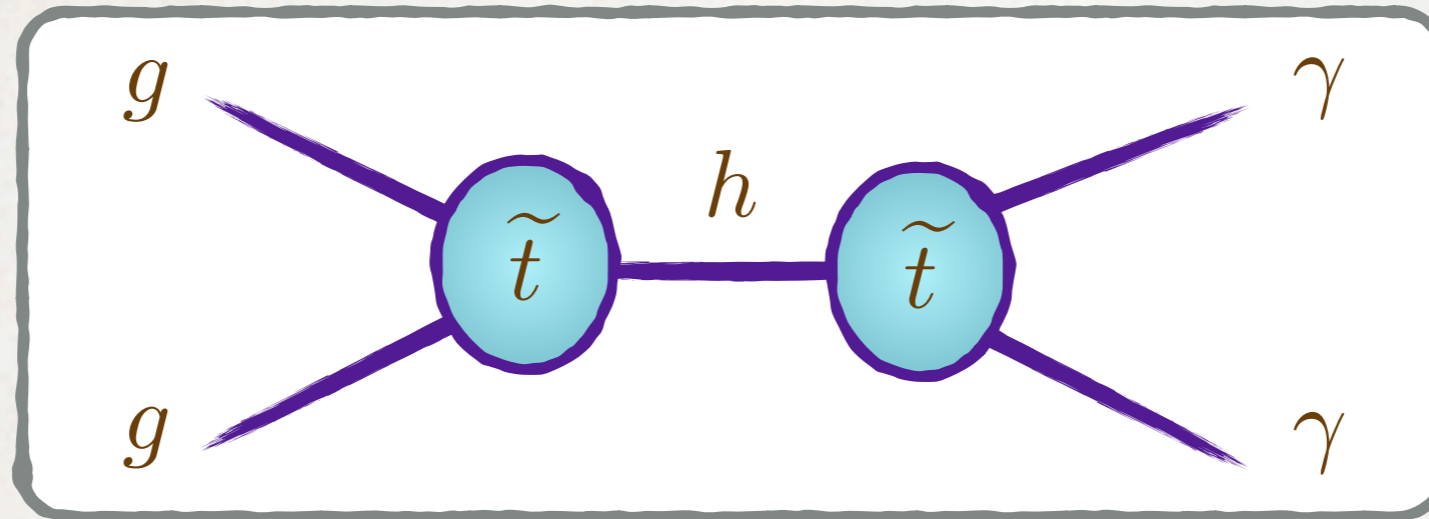
CERN in Europe?



Talk by M. Benedikt [Aspen Future Colliders Conference, 2015]

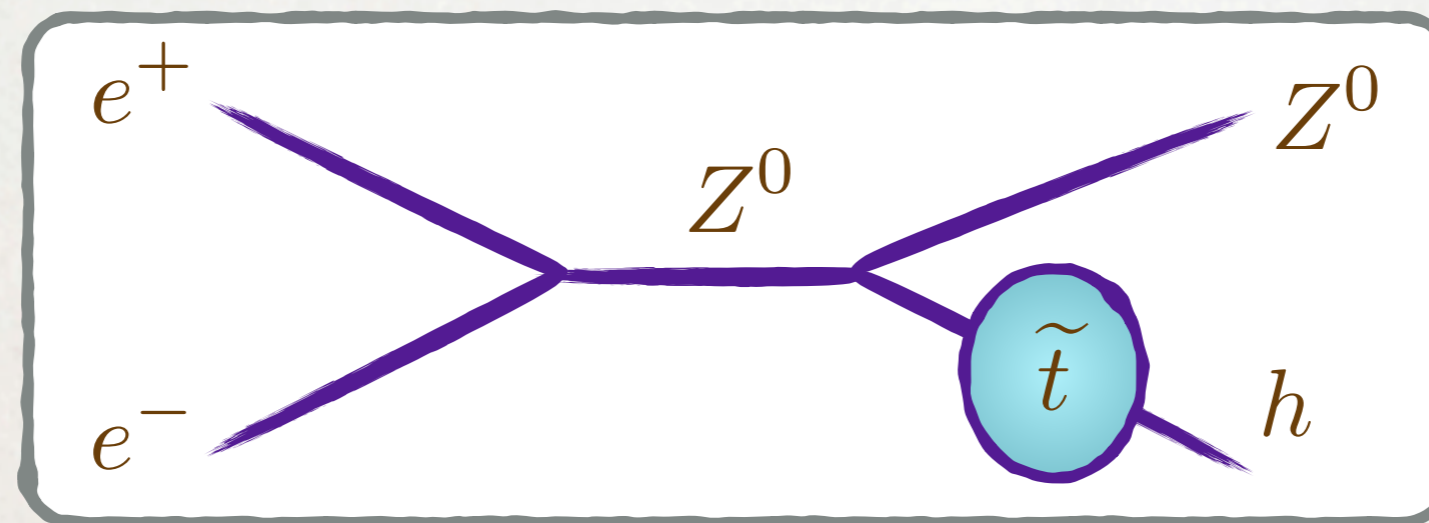
# MORE HIGGS MEASUREMENTS

Previous indirect probe required top partner be colored and charged.



Completely model independent probe:

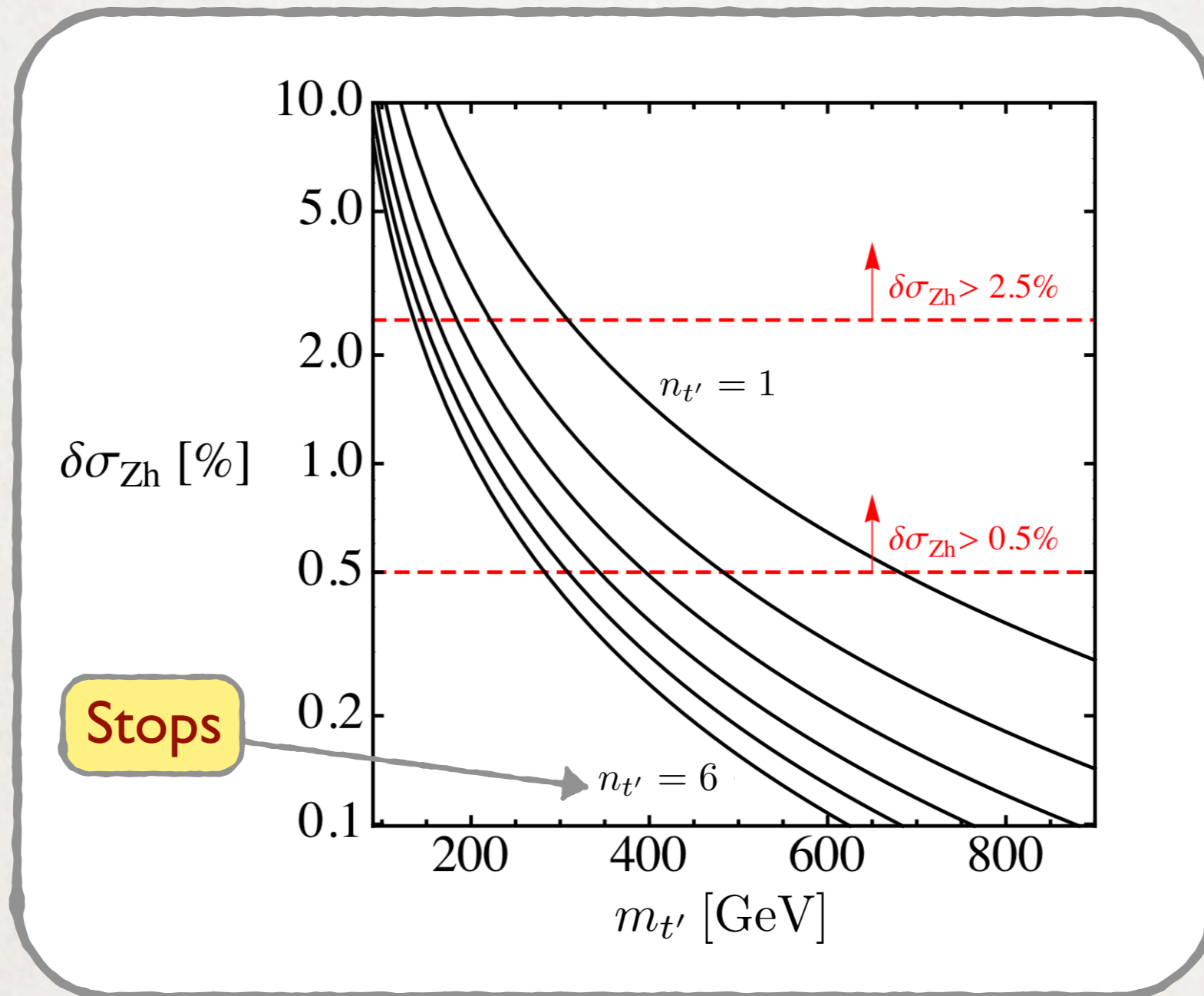
Modification to  $Z^0 - h$  associated production cross section  $\delta\sigma_{Z^0 h}$ .



Craig, Englert, McCullough [arXiv:1305.5251]

# AT A FUTURE LEPTON COLLIDER

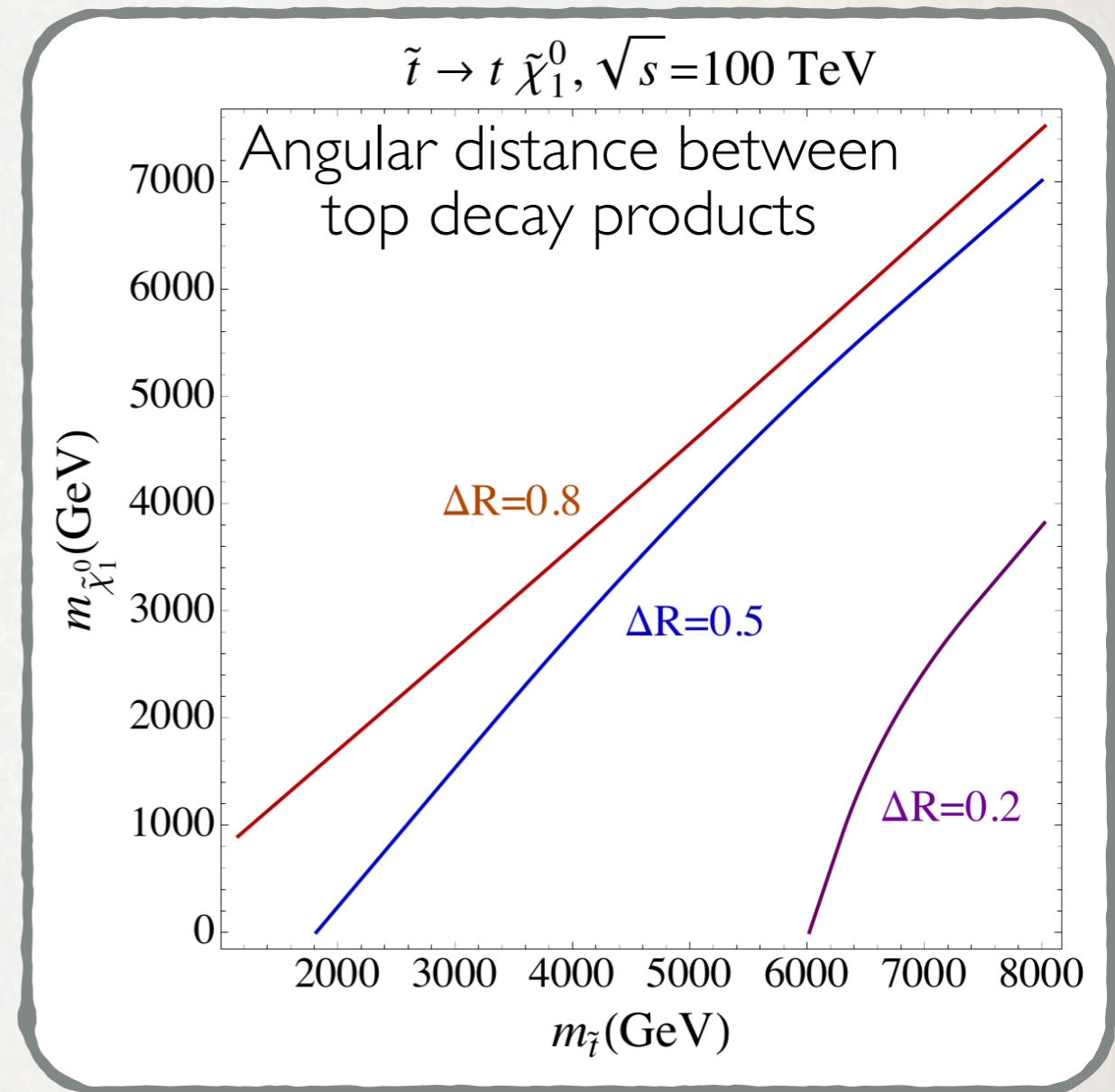
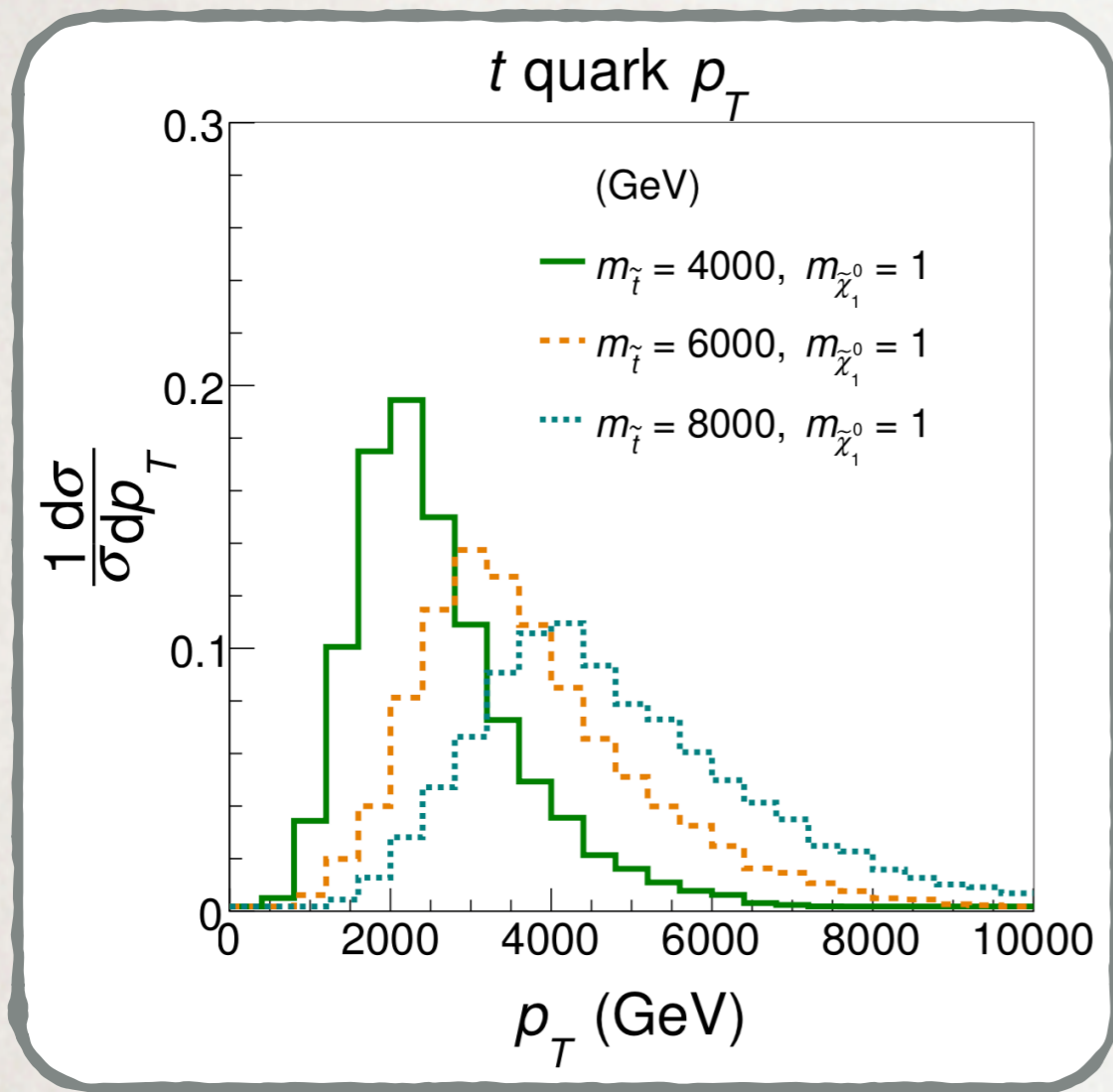
For a generic top partner  $t'$ ,  
with number of degrees of freedom  $n_{t'}$ .



Craig, Englert, McCullough [arXiv:1305.5251]

# STOP DECAYS AT A 100 TEV COLLIDER

“Top is the new bottom.”

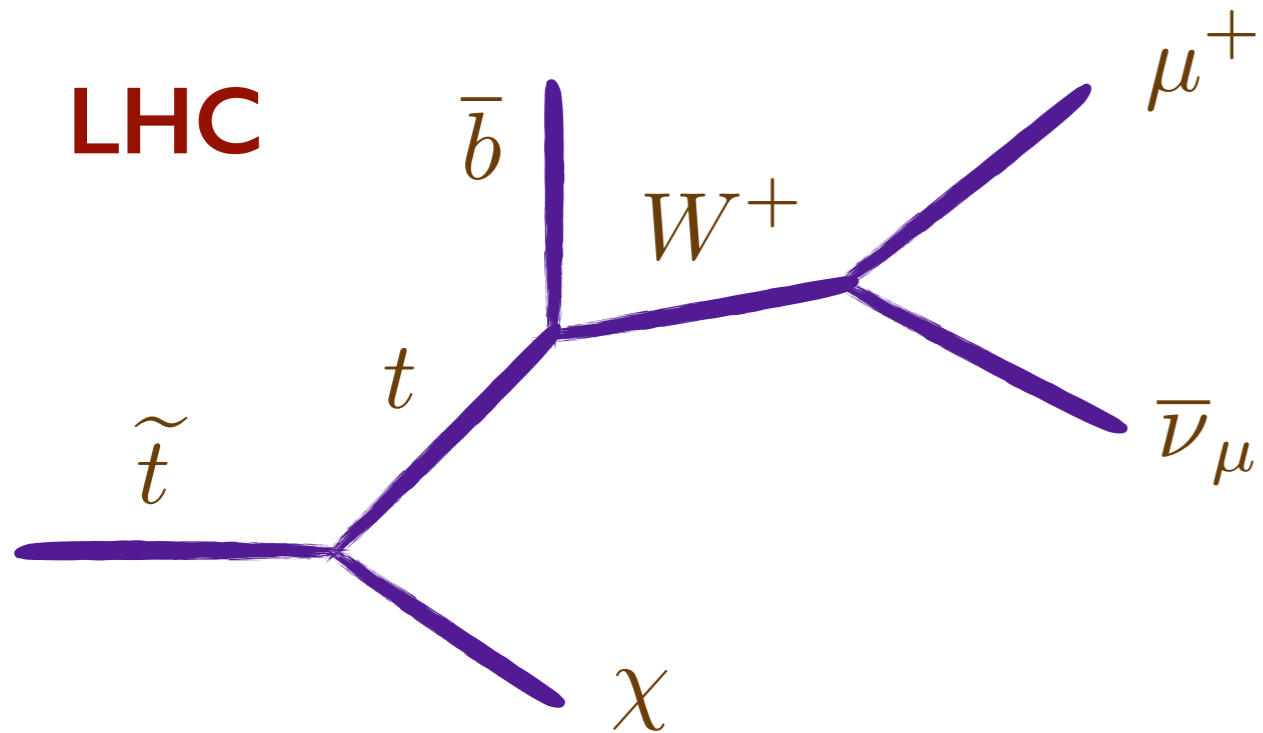


TC, D'Agnolo, Hance, Lou, Wacker [arXiv:1406.4512]



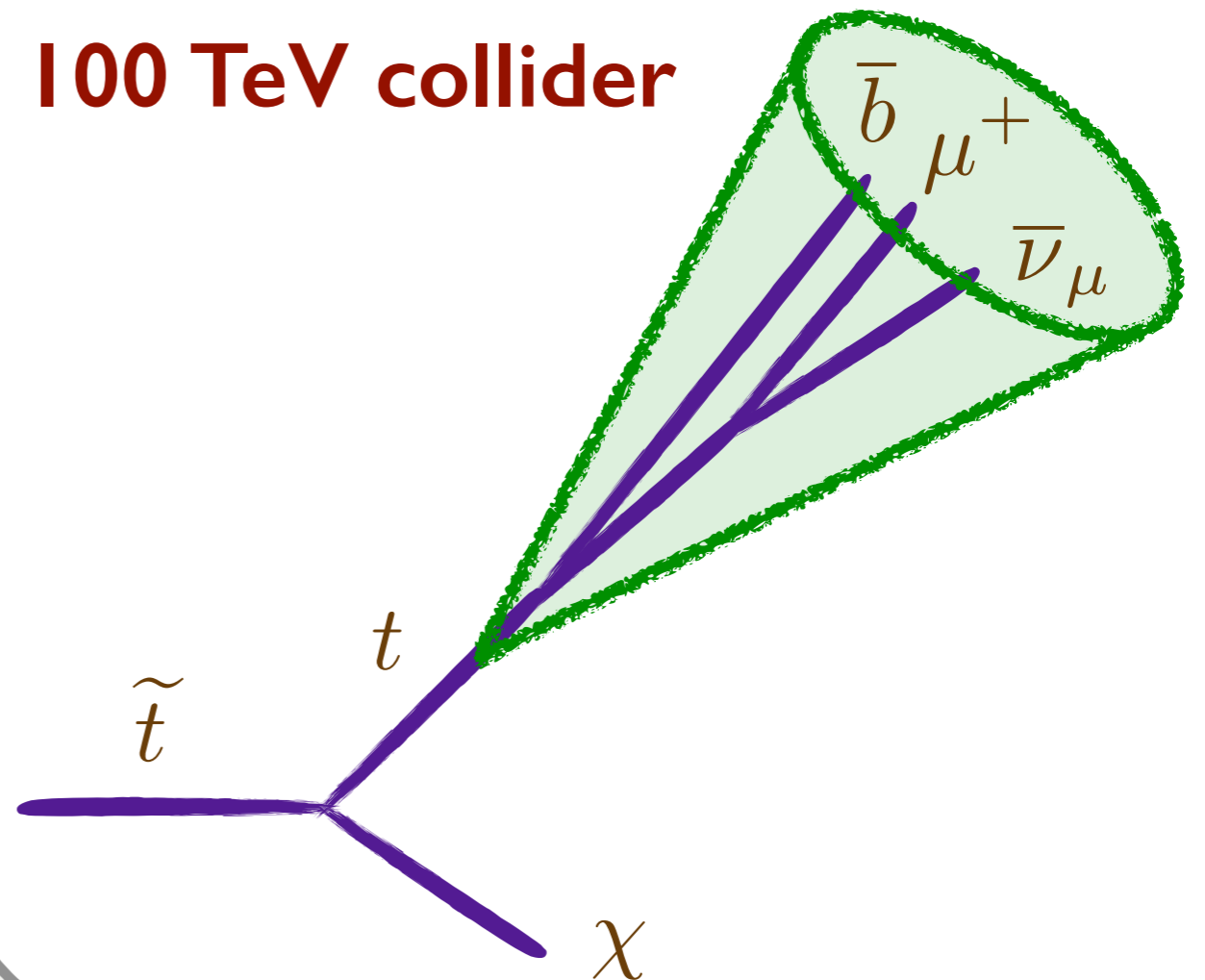
# SUPER BOOSTED TOP TAGGING

Require isolated lepton.



Require muon inside a jet.

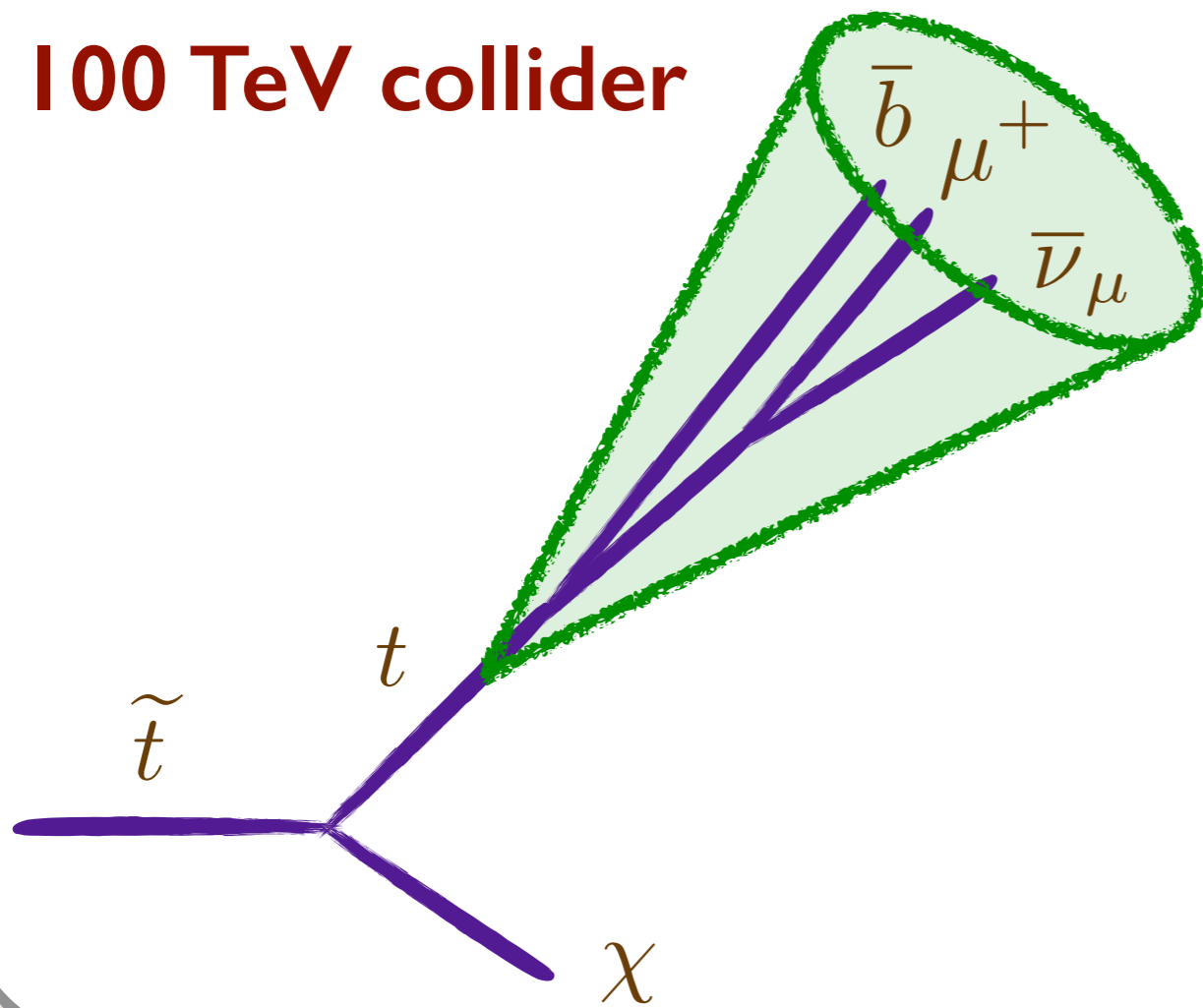
**100 TeV collider**



# SUPER BOOSTED TOP TAGGING

Require muon inside a jet.

**100 TeV collider**



## MAIN REQUIREMENTS

- $\geq 2$  jets
- $\geq 1$  muon inside a jet
- 0 isolated leptons
- Few TeV missing energy
- ...

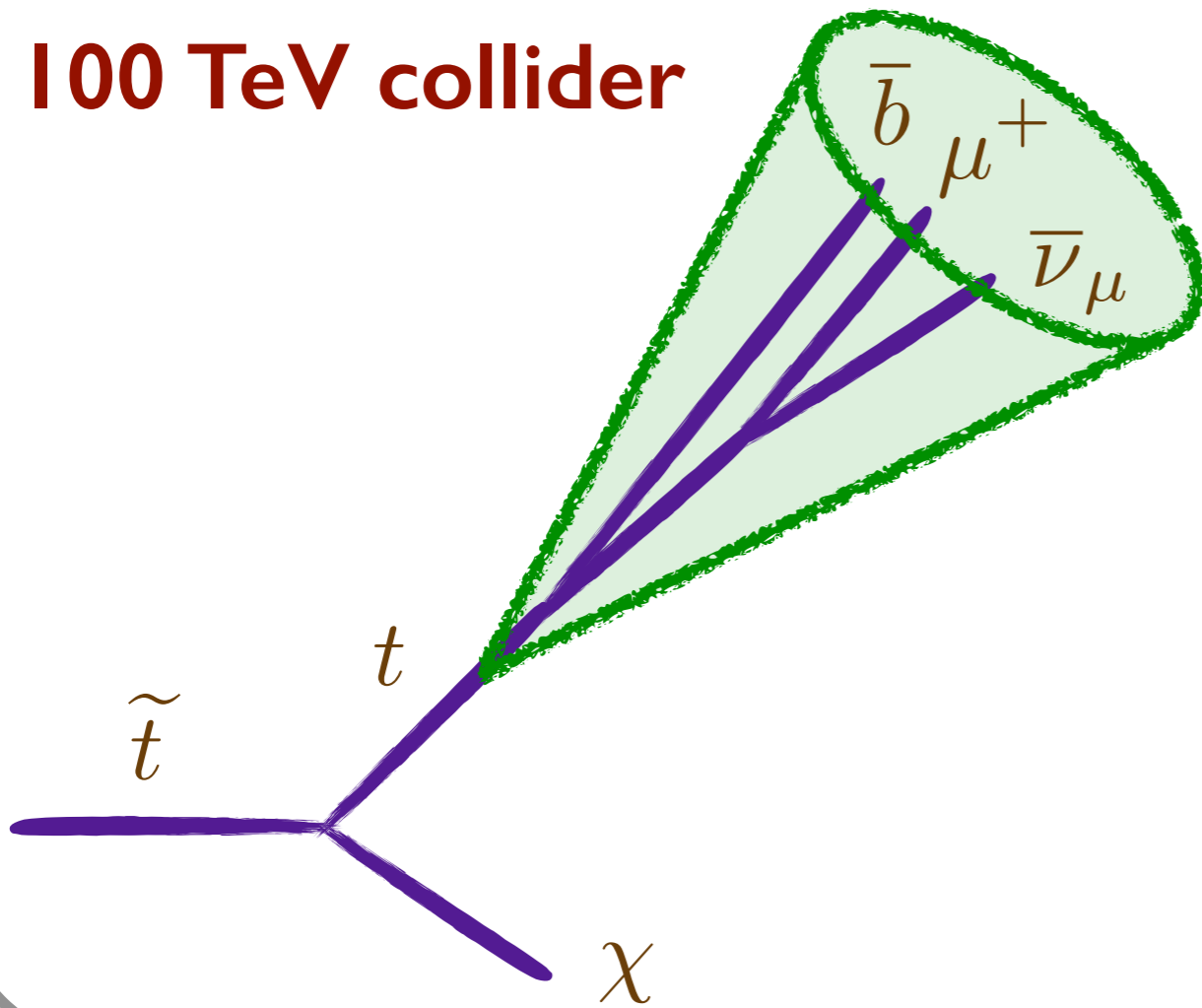
## DOMINANT BACKGROUND

$$t\bar{t} + W/Z, \dots$$

# SUPER BOOSTED TOP TAGGING

Require muon inside a jet.

**100 TeV collider**



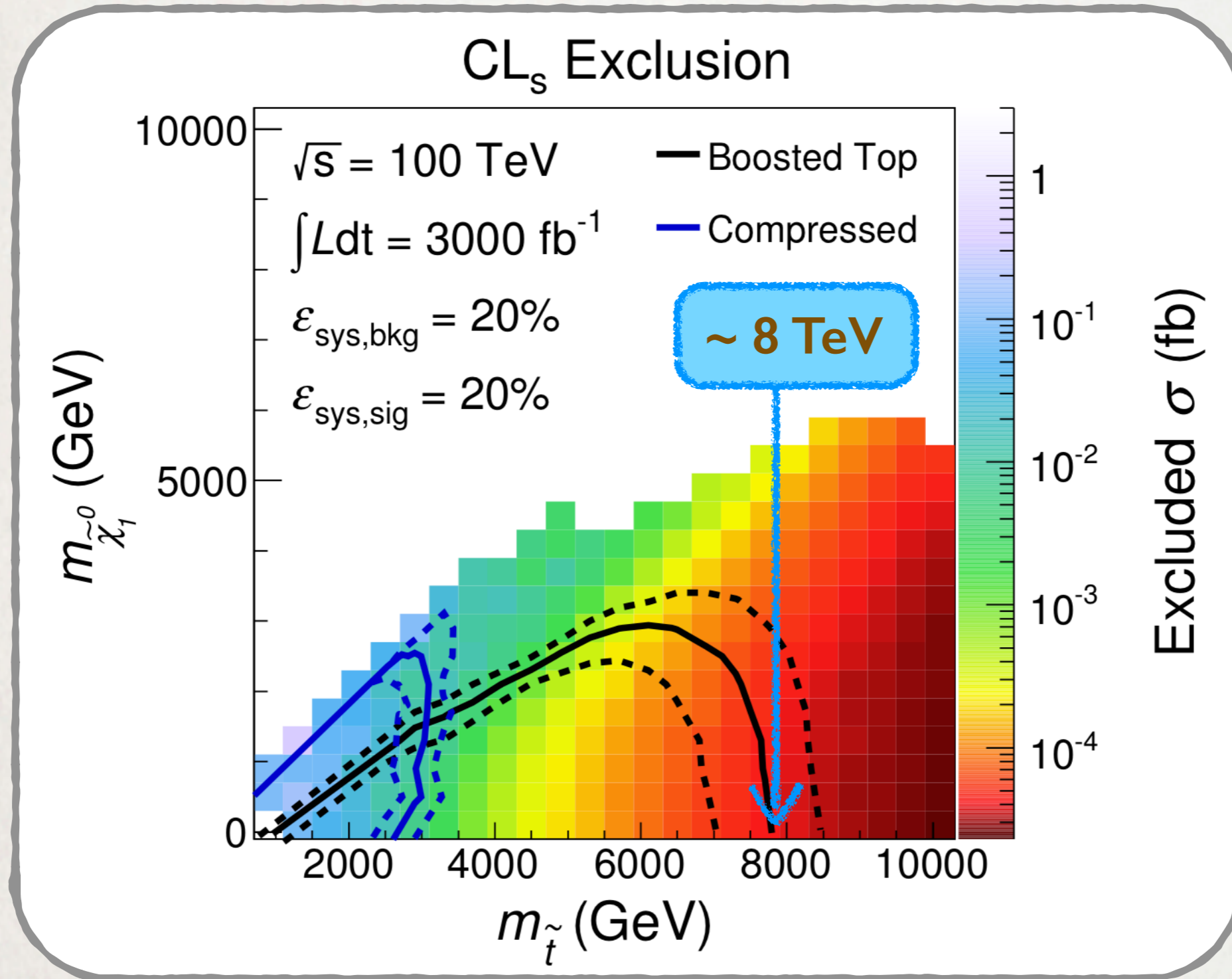
## MAIN REQUIREMENTS

- $\geq 2$  jets
- $\geq 1$  muon inside a jet
- 0 isolated leptons
- Few **TeV** missing energy
- ...

## DOMINANT BACKGROUND

$$t\bar{t} + W/Z, \dots$$

# PROJECTED LIMITS



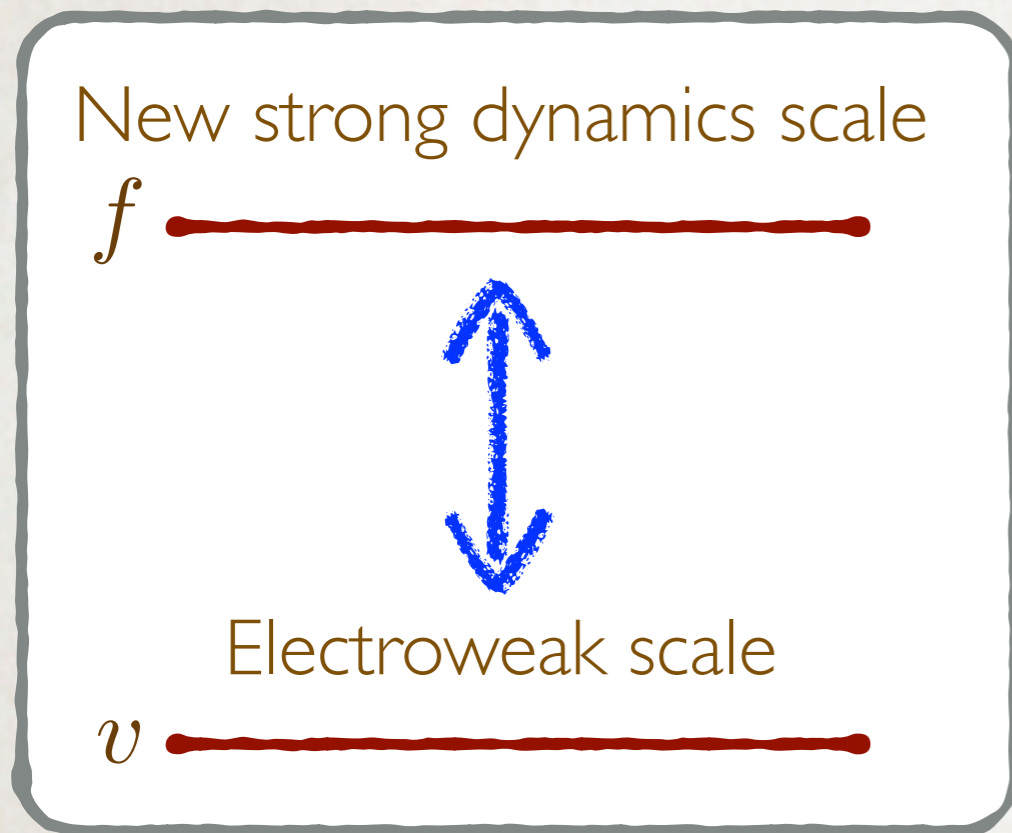
TC, D'Agnolo, Hance, Lou, Wacker [arXiv:1406.4512]

# ALTERNATIVE THEORIES

# COMPOSITE HIGGS

What if the Higgs were **not** an elementary scalar?

Requires new strong dynamics.

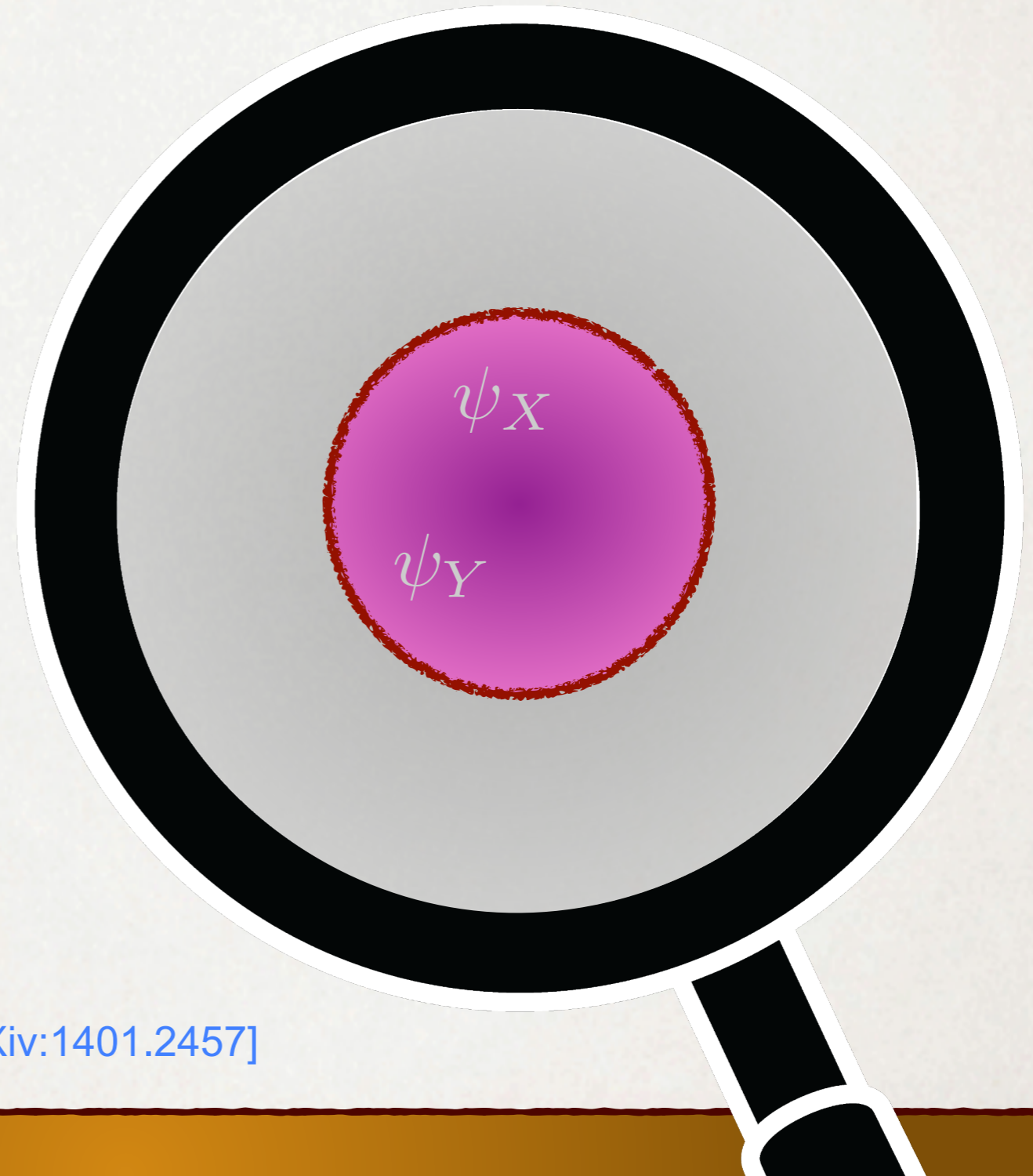
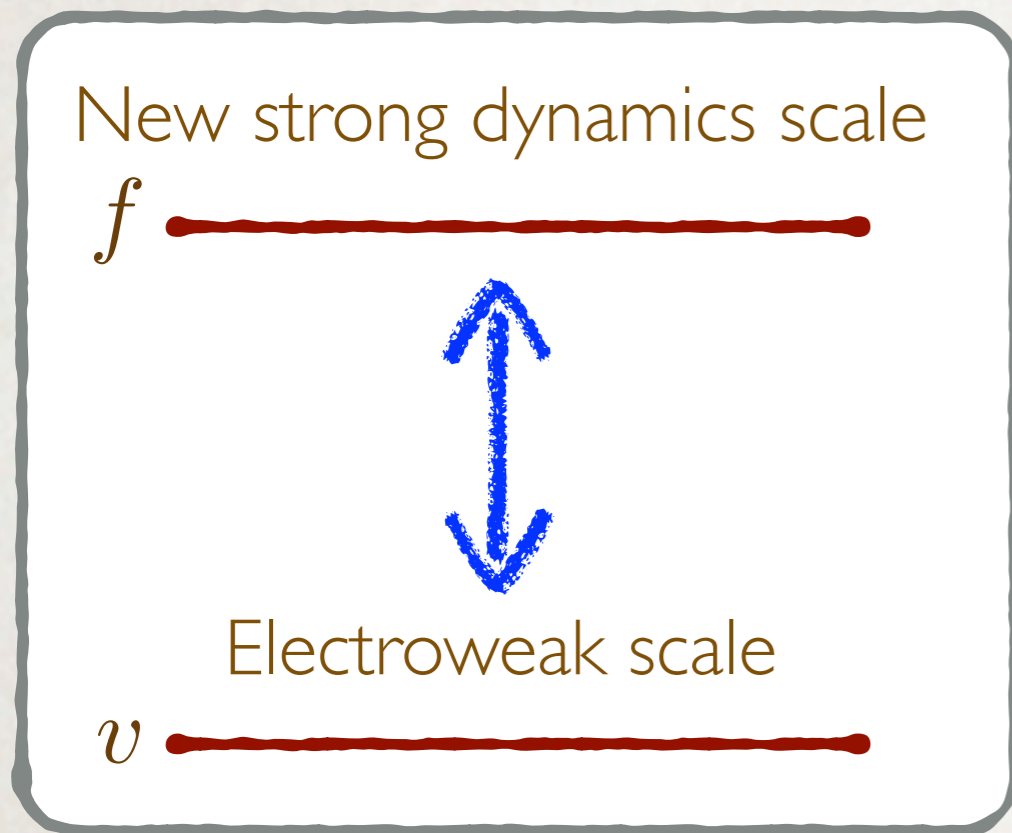


Kaplan, Georgi [1984];  
Kaplan, Georgi, Dimopoulos [1984]; ...  
for a recent review: Bellazzini, Csaki, Serra [arXiv:1401.2457]

# COMPOSITE HIGGS

What if the Higgs were **not** an elementary scalar?

Requires new strong dynamics.

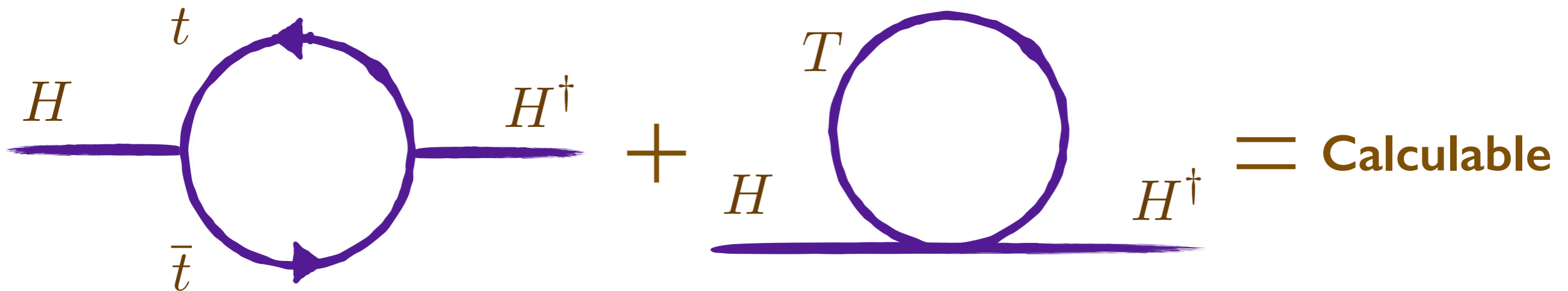


Kaplan, Georgi [1984];  
Kaplan, Georgi, Dimopoulos [1984]; ...  
for a recent review: Bellazzini, Csaki, Serra [arXiv:1401.2457]

# FERMIONIC TOP PARTNERS

Calculable requires new fermions,  $T$ .

Quadratic divergences canceled by fermionic top partner loops.



## OBSERVABLES

- Search for top partners:

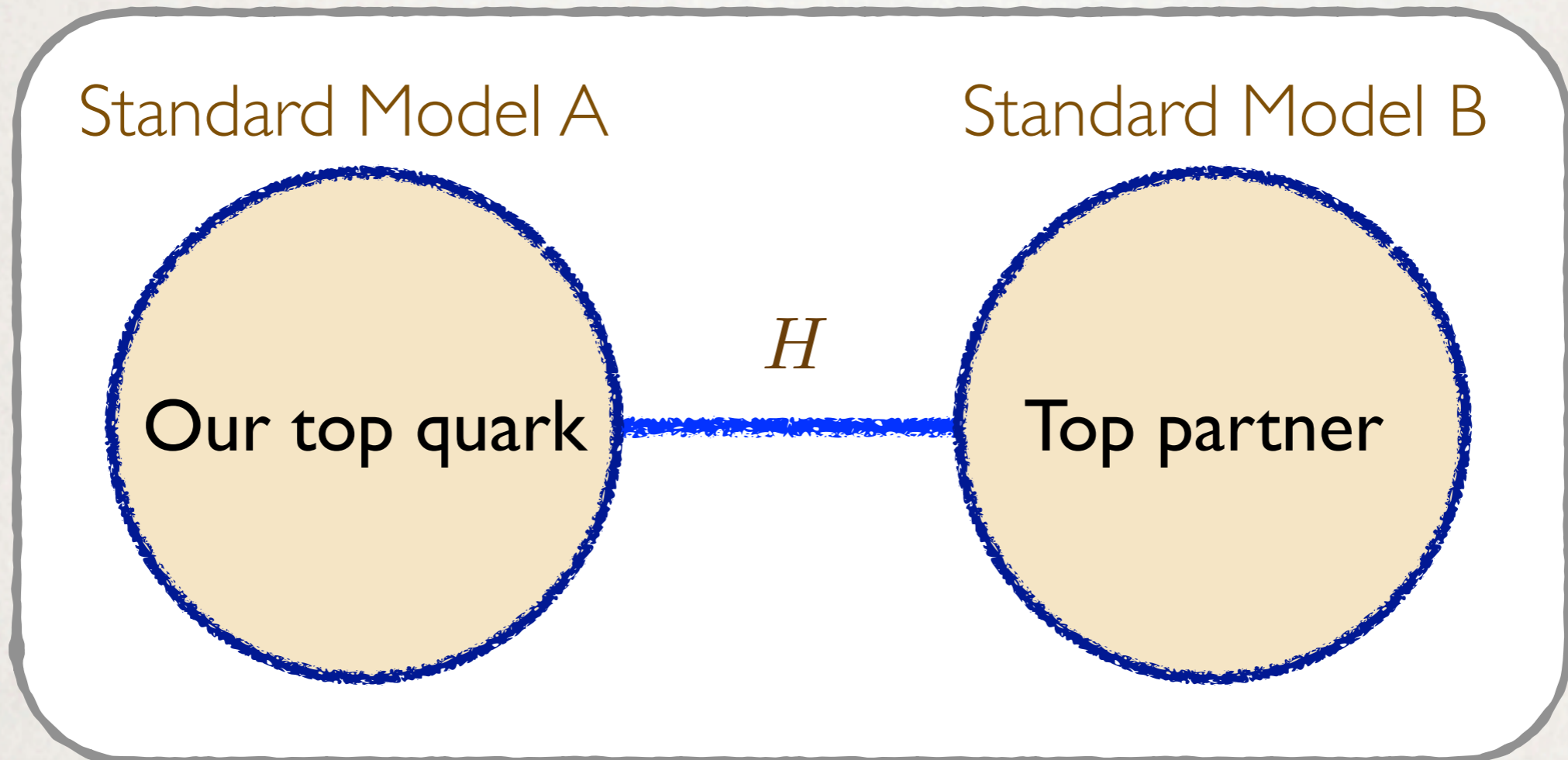
$$T \rightarrow t + Z^0, T \rightarrow t + H, T \rightarrow \bar{b} + W^+, \dots$$

- Modified Higgs properties set by  $v/f$ .



# NEUTRAL NATURALNESS

Do top partners have to be colored?



## TWO OPTIONS

- Fermionic neutral top partners: Twin Higgs  
[Chacko, Goh, Harnik \[arXiv:hep-ph/0506256\]](#)
- Scalar neutral top partners: Folded Supersymmetry  
[Burdman, Chacko, Goh, Harnik \[arXiv:hep-ph/0609152\]](#)

# SUMMARY

# SUMMARY

Reductionism:

Want “theory” of the Higgs potential.

Loops of top partners  
render Higgs mass calculable.

**MANY MANIFESTATIONS:**

- Supersymmetry: stops (scalars)
- Composite Higgs:  $T$  (fermions)
- SM  $\times$  SM: neutral scalars or fermions

**TESTABLE CONSEQUENCES:**

- Direct production at colliders
- Modification of Higgs properties

